



# Insights from $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$ substitution on the mechanism of O-O bond formation in photosystem II

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Received: 20 June 2024 / Accepted: 6 August 2024  
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## Abstract

In recent years, there has been a steady interest in unraveling the intricate mechanistic details of water oxidation mechanism in photosynthesis. Despite the substantial progress made over several decades, a comprehensive understanding of the precise kinetics underlying O-O bond formation and subsequent evolution remains elusive. However, it is well-established that the oxygen evolving complex (OEC), specifically the  $\text{CaMn}_4\text{O}_5$  cluster, plays a crucial role in O-O bond formation, undergoing a series of four oxidative events as it progresses through the S-states of the Kok cycle. To gain further insights into the OEC, researchers have explored the substitution of the  $\text{Ca}^{2+}$  cofactor with strontium (Sr), the sole atomic replacement capable of retaining oxygen-evolving activity. Empirical investigations utilizing spectroscopic techniques such as XAS, XRD, EPR, FTIR, and XANES have been conducted to probe the structural consequences of  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution. In parallel, the development of DFT and QM/MM computational models has explored different oxidation and protonation states, as well as variations in ligand coordination at the catalytic center involving amino acid residues. In this review, we critically evaluate and integrate these computational and spectroscopic approaches, focusing on the structural and mechanistic implications of  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution in PS II. We contribute DFT modelling and simulate EXAFS Fourier transforms of Sr-substituted OEC, analyzing promising structures of the  $\text{S}_3$  state. Through the combination of computational modeling and spectroscopic investigations, valuable insights have been gained, developing a deeper understanding of the photosynthetic process.

**Keywords** Natural photosynthesis · Density functional theory · X-ray absorption spectroscopy · Electron paramagnetic spectroscopy · Water oxidation catalysis · Strontium

## Abbreviations

|       |                                          |
|-------|------------------------------------------|
| AC    | Acetate Ion, $\text{CH}_3\text{COO}^-$   |
| AO    | Atomic Orbital                           |
| APS   | Advanced Photon Source                   |
| CV    | Cyclic Voltammetry                       |
| CW    | Continuous Wave                          |
| DFT   | Density Functional Theory                |
| ENDOR | Electron-Nuclear Double Resonance        |
| EPR   | Electron Paramagnetic Resonance          |
| ESEEM | Electron Spin-Echo Envelope Modulation   |
| eV    | Electron-Volts                           |
| EXAFS | Extended X-Ray Absorption Fine Structure |

|                  |                                         |
|------------------|-----------------------------------------|
| FTIR             | Fourier-Transform Infrared Spectroscopy |
| LDAO             | Lauryldimethylamine-N-Oxide             |
| LHC              | Light-Harvesting Complex                |
| MLS              | Multiline Signal                        |
| MS               | Mass Spectroscopy                       |
| NMR              | Nuclear Magnetic Resonance              |
| $\text{O}_2$     | Molecular Oxygen                        |
| OEC              | Oxygen Evolving Complex                 |
| O-O              | Oxygen-Oxygen bond                      |
| PCET             | Proton-Coupled Electron Transfer        |
| PS I             | Photosystem I                           |
| PS II            | Photosystem II                          |
| PY               | Pyridine                                |
| RC               | Radical Coupling                        |
| RR               | Resonance Raman                         |
| Tyr <sub>Z</sub> | Tyrosine Z                              |
| UV               | Ultraviolet                             |
| WNA              | Water Nucleophilic Attack               |

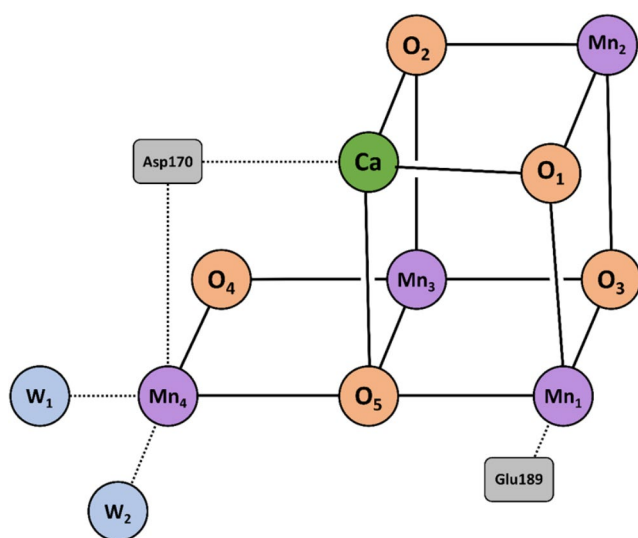
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|       |                                      |
|-------|--------------------------------------|
| WOC   | Water Oxidation Catalyst             |
| WOR   | Water Oxidation Reaction             |
| XANES | X-Ray Absorption Near Edge Structure |
| XAS   | X-ray Absorption Spectroscopy        |
| XC    | Exchange Correlation                 |
| XES   | X-ray Emission Spectroscopy          |

## Introduction

The oxygen-evolving activity of Photosystem II (PS II) protein plays a crucial role in sustaining the biosphere. Photosynthesis not only serves as the means to harvest light but also enables the conversion of sunlight into the chemical bonds of organic molecules – the fundamental building blocks of all living organisms – as well as oil, gas, and coal. PS II stands out as the sole known natural protein capable of evolving molecular oxygen from water. The photosynthetic process involves light capture and harvesting, as well as a complex electron transfer chain (Vinyard et al. 2013; Shevela et al. 2023). Reviews of the protein components in the photosynthetic process exist in the literature (Müh and Zouni 2020; Shevela et al. 2023). Central to the photosynthetic process is the oxygen evolving complex (OEC), which comprises an inorganic tetramanganese cluster,  $\text{CaMn}_4\text{O}_5$ , and serves as the driving force for the water-splitting reaction. The OEC undergoes a series of transformations referred to as the “Kok cycle,” which involves five states,



**Fig. 1** The distorted “chair-like” cartoon image of the OEC, for use in establishing notation for atoms in the OEC, as well as identifying amino acids explicitly referenced in this work. W refers to a coordinating water molecule. The coordination of certain ligating amino acid residues is shown.  $\text{Mn}_4$  is the “dangler” manganese in some models. Sr replaces Ca in  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  models, resulting in a slightly distorted structure.  $\text{Mn}_1$ ,  $\text{Mn}_2$ ,  $\text{Mn}_3$ , and  $\text{Mn}_4$  are also referred to  $\text{Mn}_D$ ,  $\text{Mn}_C$ ,  $\text{Mn}_B$ , and  $\text{Mn}_A$ , respectively

starting from the dark-adapted state ( $S_1$ ) and progressing through intermediates  $S_1$ - $S_0$  (Kok et al. 1970). Over the past four decades, noteworthy progress has been made in understanding the structure-kinetics relationship in the OEC (Cox and Messinger 2013; Yano and Yachandra 2014; Klauss et al. 2015; Vogt et al. 2015b; Shevela et al. 2023). However, the critical step involving the formation of the O—O bond remains inadequately characterized. It is believed that O—O bond formation occurs on a microsecond to millisecond timescale during the  $S_3$ -to- $S_0$  step of the catalytic cycle, ultimately resulting in oxygen evolution, in large part due to the multitude of UV-Vis, and EPR, and other spectroscopy-based studies conducted in the 1980s and 1990s (Vrettos et al. 2001; McEvoy and Brudvig 2006); recent time-resolved works followed, yielding greater insight into the precise kinetics timescales, especially in this fast  $S_3$ -to- $S_0$  transition (Chukhutsina et al. 2019; Lubitz et al. 2019; Ibrahim et al. 2020; Poddar et al. 2022; Bhowmick et al. 2023; Greife et al. 2023). Despite ongoing efforts, directly monitoring this transient process has proven challenging, leaving many details yet to be fully elucidated.

The discovery of calcium as an indispensable cofactor in oxygen evolution marked a significant breakthrough in understanding the intricate structure of the oxygen-evolving complex (OEC), discovered independently in the same year by the Babcock/Yocum and Murata laboratories (Ghanotakis et al. 1984a; Miyao and Murata 1984). A single calcium ion is contained within each OEC unit – Fig. 1 provides a cartoon guide for notation of the atoms in the OEC, as well as of some ligating amino acids and water/oxo groups. Subsequent investigations demonstrated the proximity of this metal to the tetramanganese cluster, as evidenced by the detection (Boussac et al. 1989; Sivaraja et al. 1989) and subsequent characterization (Baumgarten et al. 1990; Boussac et al. 1990; Ono and Inoue 1990; Tso et al. 1991) of a long-lived modified electron paramagnetic resonance (EPR) multiline signal generated in the  $S_2$  state of the calcium-depleted OEC. Further experimental techniques such as X-ray absorption spectroscopy (XAS) and pulsed EPR provided additional evidence supporting the presence of  $\text{Ca}^{2+}$  within the OEC; please refer to the EPR and XAS sections. Under normal catalytic conditions,  $\text{Tyr}_Z^*$  oxidizes the  $S_2$  state to generate the  $S_3\text{Tyr}_Z$  state; however, calcium depletion prevents this reaction from occurring, obstructing the S-state cycle at the  $S_2\text{Tyr}_Z^*$  state. The underlying mechanism of this effect was initially suggested to be primarily electrostatic in nature. Given its vital role in driving S-state transitions and shaping the OEC structure, the calcium ion’s significance has been extensively explored for several decades (Miqyass et al. 2008; Yocum 2008). A comparative analysis between OECs incorporating strontium (Sr) as a substitute for calcium and the native Ca OECs was performed. This work

presents an account of spectroscopic and computational investigations conducted on Sr-substituted PS II, highlighting significant disparities in the structure, function, and stability of OEC activity when compared to native (Ca) PS II. We contribute to the already discussed work with comparison of simulated EXAFS spectra produced from DFT-optimized structures of the  $S_3$  state in CaOEC and SrOEC against experimentally derived EXAFS of deconvoluted  $S_3$  states for both native and SrOEC.

## Preparation of SrOEC

Early methods for achieving Sr-PS II include Sr reconstitution of  $Ca^{2+}$ -depleted PS II centers. Ghanotakis et al. reported a protocol for inhibiting the oxygen-evolving activity of PS II (Ghanotakis et al. 1984). Purified PS II complexes were exposed to a salt treatment, releasing two water-soluble polypeptides (23 and 17 kDa) (Ghanotakis et al. 1984a, b). Such procedures to remove  $Ca^{2+}$  were found to lead to inhibition of a small proportion of PS II centers (Shen and Katoh 1991; Boussac and Rutherford 1995). Following Ca depletion, reconstitution with  $Sr^{2+}$  can be achieved with incubation in a Sr-based buffer (Boussac et al. 2000, 2004) or simply by cyanobacterial growth in a Sr-based buffer. Reports in the mid-1980s discover that following NaCl washing,  $Ca^{2+}$  release is a light-dependent process (Dekker et al. 1984; Miyao and Murata 1986). In 2004, Boussac et al. reported a method for isolating the  $Sr^{2+}/Ca^{2+}$ -containing thylakoid and PS II core complexes formed from *Thermosynechococcus elongatus* based on previously established methods for thylakoid preparation, with minor modifications (Sugiura and Inoue 1999). Other similar methods for growth of *T. elongatus* with Sr substitution have been reported (Suzuki et al. 2006; Gates et al. 2016).

## Functional effects of Sr substitution

In 1984, the laboratory of Babcock and Yocum demonstrated that the addition of small quantities of Ca cation restores oxygen-evolving activity in PS II prepared in high concentrations of NaCl; however, they note that  $Sr^{2+}$ , alongside other cations, is also able to restore lost oxygen-evolving activity, though not as completely (Ghanotakis et al. 1984). Later works by Boussac and Rutherford reported similar reductions in oxygen-evolution rates with Sr substitution preparations, alongside EPR analysis suggesting a slowing of the  $S_3$ -to- $S_0$  transition (Boussac and Rutherford 1988, 2000; Boussac et al. 1992).

An EPR-based study (Boussac et al. 2004), addressed later in greater detail, demonstrates that Sr PS II has increased stability of the  $S_3$  state, alongside increased rates of electron donation from Tyr<sub>D</sub> and an increased rate of oxidation from  $S_0 \rightarrow S_1$  by Tyr<sub>D</sub>. Slowed kinetics consistent with the rate of  $S_3Tyr_Z^-$  to  $S_0Tyr_Z$  transition suggests Tyr<sub>Z</sub> redox acts as a rate-limiting step in Ca/Sr enzymes. Later work suggests this slowdown is a function of entropic change with Sr substitution perturbing the hydrogen bonding network about Tyr<sub>Z</sub> (Rappaport et al. 2011); this is later supported with water-exchange experiments performed by (Nilsson et al. 2014), highlighting the need for a highly ordered hydrogen-bonding water network for low-energy O-O bond formation.

Ishida et al. (Ishida et al. 2008) studied the roles and importance of the  $Ca^{2+}$  and  $Cl^-$  cofactors in oxygen-evolving activity of *T. elongatus* with  $Sr^{2+}$  and  $Br^-$  substitution, respectively; this study yielded 4 different enzymes: Ca/Cl, Ca/Br, Sr/Cl, and Sr/Br. Four groups of wild type cells were grown under previously reported conditions (Sugiura et al. 2004) and were uniquely supplemented with 0.8mM YX<sub>2</sub> buffer (Y = Ca or Sr; X = Cl or Br). It was determined that the effects of  $Ca^{2+}$  to  $Sr^{2+}$  and  $Cl^-$  to  $Br^-$  substitution are additive, and the effect of  $Ca^{2+}$  to  $Sr^{2+}$  substitution is greater than the effect of the  $Cl^-$  to  $Br^-$  substitution. This is due to the perturbation of the OEC cubane structure with the larger ionic radius of  $Sr^{2+}$  compared to  $Ca^{2+}$ . Absorption changes at 292 nm were observed for each of the four enzymes, after flashes in a series 400 ms apart – enabling reduction of the Tyr<sub>Z</sub> radical by the  $Mn_4Ca/Sr$  cluster. Neither  $Sr^{2+}$  nor  $Br^-$  substitution lead to worsened turnover rate in the S-state cycle under the flashing conditions. This contrasts with the substantial decrease in oxygen evolution activity under steady-state conditions. Thus, Ishida et al. conclude a decreased efficiency of the cycle results due to an intrinsic change in kinetics in the cycle, as opposed to an increase in the number of inactive PS II centers.

Alongside the crystal structure on Sr-substituted PS II derived from *T. vulcanus*, Koua et al. analyzed the oxygen-evolving activity of the Sr-substituted photosynthetic protein (Koua et al. 2013), with an oxygen evolving rate of 46% the rate of native PS II. A rate of 3750  $\mu\text{mol O}_2/\text{mg chl a/hr}$  is reported for non-crystallized Ca-PS II, whereas Sr-substituted activity is reduced by a third (63% activity).

Gates et al. (Gates et al. 2016) perform oxygen-evolution studies on whole cells of *T. elongatus* replaced with strontium. They report lower rates of oxygen release, in line with previous reports. They highlight a higher Arrhenius activation barrier for  $O_2$  evolution associated with Sr-substituted *T. elongatus* cells compared to the Ca-based native. They observe that for PS II-WO (Sr), flux through the  $Q_AQ_B$  acceptor gate acts as the limiting process for turnover rate

*in vivo*. As such, they introduce benzoquinone derivatives binding the  $Q_B$  site and removing this kinetic inefficiency. This results in a 31% *higher* yield for Sr-substituted PS II-WOC compared to the native Ca. This effect is particularly observed at low light conditions, with fewer backward transitions and a longer lifetime of the unstable  $S_3$  state, posited to be an effect of the stabilization of the Sr-WOC compared to the photoactive Tyr<sub>z</sub>. At higher temperature ( $\sim 45^\circ\text{C}$ ), the greater activity of PS II centers in Sr-grown cells is responded with a reduced number of active PS II-WOC/Chl, resulting in comparable oxygen-evolution activity per cell.

## Structural and spectroscopic analysis of SrOEC

### XRD

Recent reviews (Cox et al. 2020; Simon et al. 2023) highlight the progress in crystallographic methods on PS II with the improvement of free-electron laser sources. In 2011, Umena et al. acquired a 1.9 Å structure (Umena et al. 2011) of purified PS II core complexes from thermophilic cyanobacterium, *T. vulcanus*. This resolution is an improvement upon prior solved structures on PS II, including work on closely related *T. elongatus*. The electron density mapping from Umena et al. (1.9 Å structure) located the metal atoms of the cubane  $\text{Mn}_4\text{CaO}_5$  cluster alongside all ligands. This structure identified more than 1300 water molecules in each photosystem II monomer, identifying extensive hydrogen-bonding networks able to serve as channels for protons,  $\text{H}_2\text{O}$ , or oxygen molecules. Subsequently, in 2013, Koua et al. published a 2.1 Å structure of Sr-substituted PS II derived from *T. vulcanus* and compared the structure to the

1.9 Å structure of Ca PS II obtained years prior. The occupancy of Sr atoms in the structure was  $\sim 0.7$  and resulted in a lower electron density of  $\text{Sr}^{2+}$  compared to that of the Mn ions, consistent with prior reports of reduced  $\text{Sr}^{2+}$  occupation in  $\text{Sr}^{2+}$  PS II (Strickler et al. 2005; Boussac et al. 2008; Yano et al. 2011). Coupled to the lower occupancy of the  $\text{Sr}^{2+}$  ion is the lowered occupancy of the W3 and W4 water molecules, themselves associated with Ca/Sr, though the accuracy of their position in the mapping is unaffected. Subsequent XFEL-based structures on Ca PS II (Kern et al. 2018; Suga et al. 2019) obtained with low X-ray induced damage have shown similar structure as reported by Umena et al.

Mn-Mn and Sr-Mn distances of key XRD structures are summarized in Tables 1 and S1. Table 1 shows that the relative positions of the four Mn ions are similar in both Ca and Sr PS II. The average  $\text{Mn}_1\text{-Mn}_4$  distance, 0.1 Å longer for Sr, and the average  $\text{Mn}_3\text{-Mn}_4$  distance, 0.1 Å shorter for Sr, differ from the native  $\text{Mn}_4\text{CaO}_5$  cluster. Please note that these distances might also be affected by variable X-ray dose in two XRD experiments which can result in variable content of photo-reduced Mn centers. Photo-reduction of Mn centers typically increase Mn-Mn distances. DFT modeling (see below and in Table S2) shown elongation of some Mn-Mn distances in Sr based OEC models in agreement with XRD results. More substantial than the slight variations in locations of the Mn ions is the positioning of the  $\text{Sr}^{2+}$  ion, located more toward the outside of the cubane. The displacements of the Ca and Sr ions in the monomers of each structure, due to conformational changes with Sr substitution average to be 0.3 Å. The resulting distorted Sr-cubane structure yielded Sr-Mn distances  $\sim 0.2$  Å greater than analogous distances for native Ca-Mn distances. These elongations may be in part attributed to a larger ionic radius

**Table 1** Distances between Ca/Sr-Mn, Mn-Mn, and Ca/Sr-O atoms. The distances are taken as average of the two monomers. All structures are of the dark-stable  $S_1$  state, either from XRD or XFEL methods

|                                   | 1.9 Å, Ca PS II<br>(Umena et al. 2011)* | 2.0 Å, Ca PS II XFEL<br>(Kern et al. 2018) | 2.05 Å Ca PS II<br>XFEL<br>(Suga et al. 2019) | 2.1 Å, Sr PS II<br>(Koua et al. 2013)* |
|-----------------------------------|-----------------------------------------|--------------------------------------------|-----------------------------------------------|----------------------------------------|
| Ca/Sr – Mn <sub>1</sub>           | 3.52                                    | 3.43                                       | 3.61                                          | 3.55                                   |
| Ca/Sr – Mn <sub>2</sub>           | 3.34                                    | 3.38                                       | 3.42                                          | 3.50                                   |
| Ca/Sr – Mn <sub>3</sub>           | 3.42                                    | 3.51                                       | 3.40                                          | 3.64                                   |
| Ca/Sr – Mn <sub>4</sub>           | 3.80                                    | 3.83                                       | 3.75                                          | 4.00                                   |
| Mn <sub>1</sub> – Mn <sub>2</sub> | 2.82                                    | 2.78                                       | 2.59                                          | 2.78                                   |
| Mn <sub>1</sub> – Mn <sub>3</sub> | 3.28                                    | 3.24                                       | 3.16                                          | 3.34                                   |
| Mn <sub>1</sub> – Mn <sub>4</sub> | 4.95                                    | 4.85                                       | 4.96                                          | 5.06                                   |
| Mn <sub>2</sub> – Mn <sub>3</sub> | 2.90                                    | 2.85                                       | 2.72                                          | 2.92                                   |
| Mn <sub>2</sub> – Mn <sub>4</sub> | 5.41                                    | 5.20                                       | 5.27                                          | 5.42                                   |
| Mn <sub>3</sub> – Mn <sub>4</sub> | 2.93                                    | 2.74                                       | 2.89                                          | 2.88                                   |
| Ca/Sr – O <sub>1</sub>            | 2.41                                    | 2.30                                       | 2.58                                          | 2.44                                   |
| Ca/Sr – O <sub>2</sub>            | 2.51                                    | 2.48                                       | 2.64                                          | 2.66                                   |
| Ca/Sr – O <sub>3</sub>            | 3.91                                    | 3.78                                       | 4.16                                          | 3.99                                   |
| Ca/Sr – O <sub>4</sub>            | 3.98                                    | 4.19                                       | 3.82                                          | 4.03                                   |
| Ca/Sr – O <sub>5</sub>            | 2.50                                    | 2.56                                       | 2.59                                          | 2.59                                   |

\* Please note that these distances might also be affected by variable X-ray dose in two XRD experiments which can result in variable content of photo-reduced Mn centers. Photo-reduction of Mn centers typically increase Mn-Mn distances

of  $\text{Sr}^{2+}$  compared to  $\text{Ca}^{2+}$  (0.99 Å to 1.12 Å, with  $\text{Ca} \rightarrow \text{Sr}$  substitution). Thorough analysis of the differences in oxo-bridging oxygens is hampered by the lower electron densities of the oxygens compared to Mn and Ca/Sr and by the reduced occupancy of the  $\text{Sr}^{2+}$  ions. These factors yielded higher experimental errors for these distances. Given that the Mn-O distances between both structures differed on average by 0.1–0.2 Å and that some of these bonds are shorter while others are longer in the Sr-based structure, at least some of these differences can be attributed to experimental error.

## XAS

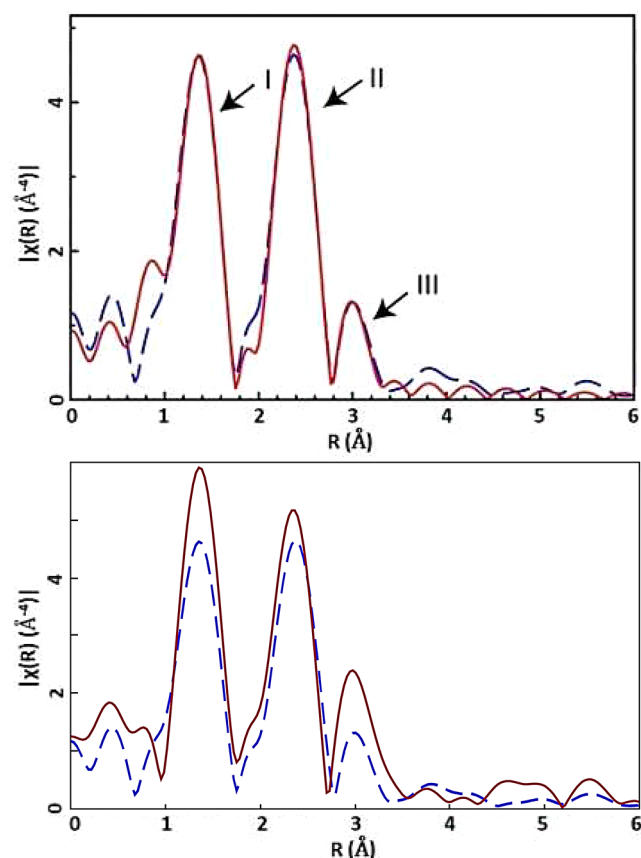
While early XRD-derived structures were able to locate the electron density associated with the  $\text{Mn}_4\text{Ca}$  cluster within the PS II complex, they were limited by resolution (3.2–3.5 Å) (Zouni et al. 2001; Kamiya and Shen 2003; Biesiadka et

al. 2004; Ferreira et al. 2004). As such, detailed structures critically depended on input from a variety of spectroscopic techniques to acquire an accurate determination of positions of the Mn/Ca and bridging/terminal ligand atoms. Early Mn (Yano et al. 2005b) and Ca (Cinco et al. 2002) K-edge XANES and EXAFS studies were able to more precisely deduce Mn-Mn/Ca/ligand distances due to the then-limiting levels of X-ray doses used in XRD. PS II crystals of the  $\text{Mn}_4\text{Ca}$  cluster were highly susceptible to radiation damage, and with the lower X-ray doses required for EXAFS/XANES experiments, any radiation damage could be precisely monitored and controlled. Many works (Dau et al. 2004; Yano et al. 2005a; Grabolle et al. 2006) warn of the potential to reduce  $\text{Mn}^{\text{III/IV}}$  to  $\text{Mn}^{\text{II}}$  at sufficiently high X-ray doses and temperatures.

Cinco et al. performed a Ca EXAFS-based study (Cinco et al. 2002) on native PS II membrane fragments, prepared as BBY particles from spinach. Fourier transforms of their Ca EXAFS from active (Chelex-treated samples) revealed an inner coordination shell, described by 5–6 O atoms at 2.4 Å. Another significant peak in their Fourier transforms was fitted to a Ca-Mn coordination at 3.4 Å. Note that these samples are characterized by 2 Ca/PS II, due to the presence of one calcium in both the OEC (Mn cluster) and the light-harvesting complex II (LHC II). Recently Ca valence-to-core X-ray emission spectroscopy was proposed as an additional tool for analysis of the Ca ligand environments (Mathe et al. 2019). In 2005, Yano et al. through Mn EXAFS, refined existing structures of the Mn-OEC. They revealed three short Mn-Mn distances, between 2.7 Å and 2.8 Å and resulting, the same number of di- $\mu$ -oxo-bridging units as well; prior works (Hasegawa et al. 1999; Dau et al. 2001; Robblee et al. 2002) were not able to determine whether there are two or three di- $\mu$ -oxo-bridged Mn-Mn moieties, due in part to the inherent error of EXAFS methods in determining the number of Mn-Mn vectors (Yachandra 1995). These insights provided much-needed clear limitations on the proposed structures based on prior spectroscopic and diffraction data for native PS II.

One EXAFS study on native PS II at room-temperature was performed by Davis et al. in 2015 (Davis and Pushkar 2015). Their single shell analysis of the first coordination sphere (Mn-O interactions) fit the data to three Mn-O interactions at 1.9 Å; later fits refined this to two Mn-O interactions at 1.82 Å and one Mn-O interaction at 1.99 Å. Analysis of peak III of the Fourier transform – shown in Fig. 2 – is challenging due to a low intensity and contributions from multiple absorber-backscatterers: Mn, Ca, and C.

Having considered EXAFS of native PS II, we can consider the work done on Sr-substituted PS II. Pushkar et al. (Pushkar et al. 2008) perform an EXAFS-based study on *T. elongatus* prepared as CaPS II and exchanged Sr PS II.



**Fig. 2** a) Overlays of Fourier transforms,  $k = 4\text{--}11.45 \text{ \AA}^{-1}$ , of EXAFS data for native OEC from spinach prepared in the dark-stable S1 state. Red: Fourier transform of room-temperature EXAFS collected by Davis et al. 2015 at APS (Davis and Pushkar 2015). Blue: EXAFS fit produced by Davis et al. revealing Mn coordination environment in agreement with prior cryogenic XAS and XRD structures. Characteristic peaks I, II, and III are indicated with arrows. b) SrOEC data from *T. elongatus* (Pushkar et al. 2008), red, are compared to the same CaOEC data from spinach, blue. Increase in the peak III intensity is due to stronger Mn-Sr scattering

Samples are advanced through the S-state cycle via consecutive laser flashes followed by immediate freezing in liquid nitrogen (Messinger et al. 2001). Mn K-edge EXAFS spectra are obtained for both native PS II and the Sr-substituted PS II for subsequent flashes, corresponding to structures in the  $S_1$ ,  $S_2$ ,  $S_3$ , and  $S_0$  states, cyclically. Mn K-edge EXAFS data ensure the integrity of the samples and the absence of radiation damage to the core structure of the  $Mn_4$  core. Fourier transforms (FT) of the Mn EXAFS data reveal three major peaks: I, characteristic of the 1.8 Å Mn-oxo bridging ligand distances, II, characteristic to 2–3 di- $\mu$ -oxo-bridged Mn-Mn interactions  $\sim 2.7$  Å, and III, characteristic of one mono- $\mu$ -oxo-bridged Mn-Mn and from Mn-Sr interactions. While EXAFS of the  $S_1$ ,  $S_2$ , and  $S_0$  states are similar, peaks II and III are attenuated and lowered in intensity for the  $S_3$  state. This suggests structural changes in the OEC in the  $S_3$  state. Figure 2 shows Fourier transforms of CaOEC showing the three characteristic peaks.

Sr EXAFS are more informative and can more clearly resolve Sr-Mn interactions in all S states. Sr XANES spectra (Pushkar et al. 2008) reveal differences in the S-state spectra and the spectrum of hydroxylamine-inhibited (HYD) “inactive” Sr-PS II, revealing a different first coordination environment for Sr in active PS II and from free ion  $Sr^{2+}$ ; FT of the HYD spectra reveal no peak II, associated with Sr-Mn interactions. The analogous peak I for Sr EXAFS suggests one shell of 7–8 oxygen atoms  $\sim 2.5$  Å, while the FT peak II corresponds to two different Sr-Mn interactions for all S states. Sr-Mn interactions are fitted at  $\sim 3.5$  Å and  $\sim 4.0$  Å for all S-states in Sr. Prior Ca EXAFS of native (Cinco et al. 2004) PS II and Sr EXAFS of Sr-reactivated PS II membranes (Cinco et al. 1998) indicated proximity of the Ca at 3.4 Å and Sr at 3.5 Å to the Mn cluster in the  $S_1$  dark-stable state of the OEC. The work of Pushkar et al. unambiguously demonstrated that Sr is proximate to the OEC in all S states with significant changes in Sr(Ca)-Mn interactions with sequential flashes – as the enzyme proceeds through the catalytic cycle.

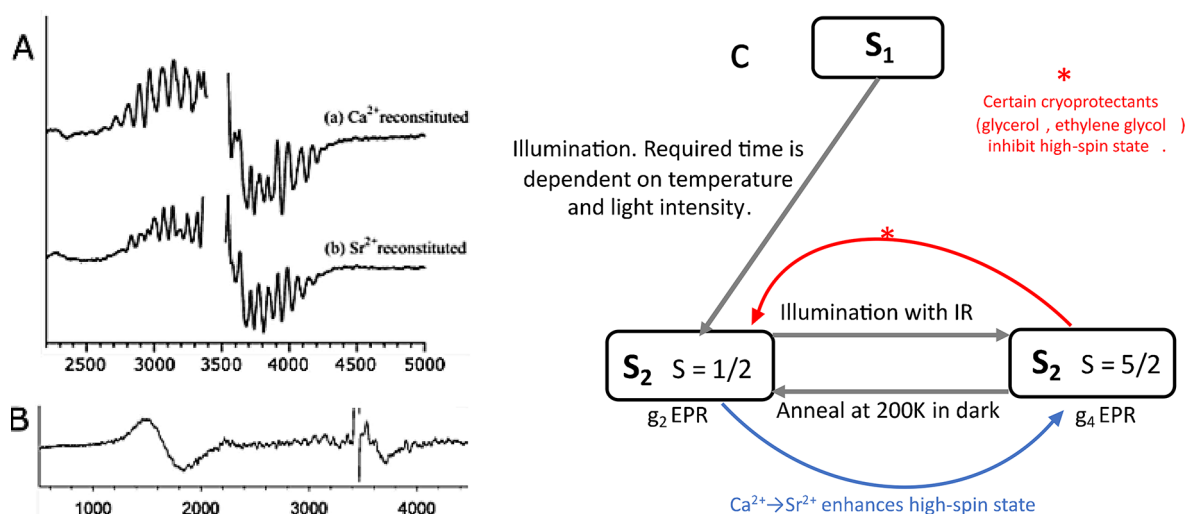
Now addressing XANES spectra, we revisit the work done on Sr-substituted PS II. Pushkar et al. (2008) perform an EXAFS-based study on *T. elongatus* prepared with native Ca and biosynthetically exchanged Sr-PS II preparations (Boussac et al. 2004). Having discussed Mn EXAFS previously, we can discuss the results of Sr XANES which can more clearly resolve Sr-Mn interactions in all S states. Sr XANES spectra (Pushkar et al. 2008) reveal differences in the S-state spectra and the spectrum of hydroxylamine-inhibited (HYD) “inactive” Sr-PS II. This indicated that the first coordination environment about Sr in active PS II differs from that of free  $Sr^{2+}$  – Fourier transforms of HYD Sr-PS II spectra did not show a peak assigned to Sr-Mn interactions for active Sr PS II. In the same work, examination of Sr PS

II Mn XANES revealed a shift in the Mn K-edge between the  $S_1$  and  $S_2$  states consistent with the oxidation of the Mn core. A much more minor shift occurs between  $S_2$  and  $S_3$ . However, with a third flash, the edge position restores to a lower energy from the  $S_3$ -to- $S_0$  spectra, expected with Mn reduction. The observation of Mn K-edge shifts from  $S_1$ - $S_3$  was further evidencing the  $Mn^{III}_2Mn^{IV}_2$   $S_1$  oxidation state and Mn-oxidation with advancement S-states, for both Ca and Sr PS II.

## EPR

When illuminated at 195 K by visible light for 5–10 min, native PS II advances to low-spin isoform of the  $S_2$  state ( $S_{total}=1/2$ , with  $Mn^{III}/Mn^{IV}$  and  $Mn^{IV}/Mn^{IV}$  antiferromagnetically coupled), characterized by the so-called multiline signal (MLS). This feature is characterized by at least eighteen partially resolved hyperfine lines near g-factor 2. There exists substantial evidence of a spin isomorphism in the  $S_2$  state, with a high-spin ( $S=5/2$ ) and a low-spin ( $S=1/2$ ) form, which have been characterized by EPR spectroscopy. The low-spin ( $S=1/2$ ) isoform of the  $S_2$  state is often obtained by illumination at 195 K, with use of a heat/IR filter; illumination times are dependent upon the intensity of the light source. The low spin to high spin transition occurs with absorption of  $\sim 820$  nm (IR) light. IR illumination at decreasing temperatures (170 K  $\rightarrow$  65 K) produces more of the high-spin state, as observed through the  $g=4.1$  signal in EPR spectroscopy (Boussac et al. 1996, 1998). The high-spin state may be converted back to the low-spin state by annealing ( $\sim 200$  K) in the dark. The presence of certain cryoprotectants, such as glycerol and ethylene glycol, have been shown to inhibit the formation of the high-spin  $S_2$  state (Zimmermann and Rutherford 1986; de Paula et al. 1987); sucrose, however, is a cryoprotectant allows for formation of either the high-spin or low-spin  $S_2$  state. It has been demonstrated that following NaCl washing,  $Sr^{2+}$ -reconstituted PS II demonstrates the high-spin  $g=4$  state is more easily reproduced than with  $Ca^{2+}$  reconstitution, when both samples are illuminated for 5 min at 200 K (Boussac and Rutherford 1988). Figure 3 summarizes the  $S_2$  state isomorphism.

Kim et al. (2004) performed X-band CW and ESEEM spectroscopy studies on  $Ca^{2+}$  and  $Sr^{2+}$  reconstituted PS II. The already characterized (Boussac and Rutherford 1988)  $S_2Tyr_Z^*$  split signal was reproduced to confirm  $Ca^{2+}$  depletion in all sample preparations. The  $Ca^{2+}$ -reconstituted sample was found to be nearly identical to the native PS II MLS, while the  $Sr^{2+}$ -reconstituted sample showed a significantly altered MLS, featuring narrower hyperfine splitting and a different intensity pattern compared to the control samples; which were reproduced (Boussac et al. 2004). Three-pulse  $^{87}Sr$  ESEEM was used to probe the Mn cluster-Sr/



**Fig. 3** Description of fundamental EPR features to PS II. **a**) A description of the effect of  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution on the multiline signal (MLS) in the  $\text{S}_2$  state, the characteristic EPR feature of the low-spin ( $S = 1/2$ ) isoform appearing near g-factor 2. Adapted from (Kim et al.

2004). While three-pulse ESEEM of the natural abundance (7% abundance  $^{87}\text{Sr}$ ) Sr sample exhibited no modulation,  $^{87}\text{Sr}$ -enriched samples revealed a clear modulation from the spin = 9/2  $^{87}\text{Sr}$  nucleus, weakly magnetically coupled to the Mn cluster. With a point dipole approximation for electron spin, they concluded a Mn-Ca/Sr distance in the 3.8–5.0 Å. This early  $^{87}\text{Sr}$  ESEEM approach provided independent evidence contemporary with XAS experiments, suggesting that the Ca/Sr binding site is close to the Mn cluster in the OEC of PS II.

Cox et al. (Cox et al. 2011) produced  $^{55}\text{Mn}$  ENDOR spectra of Ca and Sr OEC prepared in the  $\text{S}_2$  state,  $B_0 = 1260$  mT. Contrasting CW X-band and ESE Q-band spectra, the Ca and Sr OEC produce comparable results for  $^{55}\text{Mn}$ -ENDOR spectra; both display a similar total spectra breadth (60–200 MHz). Approximately six ENDOR peaks are resolved for both  $\text{Ca}^{2+}$  and  $\text{Sr}^{2+}$  variants of the OEC, with only minor changes in the line intensity of the peaks. Of the six peaks, only the first two demonstrate a minor downshift in position, by approximately 10 MHz. Comparing  $^{55}\text{Mn}$ -ENDOR spectra from *T. elongatus* with data from spinach, differences intrinsic to species have a larger effect on the spectra than does  $\text{Ca}^{2+}$ -to- $\text{Sr}^{2+}$  replacement (Kulik et al. 2005a, b). The change in CW spectra with  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  replacement is due to a decrease in the hyperfine tensor anisotropy. Some prior DFT-based models (Pantazis et al. 2009; Siegbahn 2009) were found to be consistent with the data presented by Cox et al. positing a 5-coordinate  $\text{Mn}^{\text{III}}$  as a site for substrate water binding.

The Mülheim group (Rapatskiy et al. 2012; Pérez Navarro et al. 2013; Lohmiller et al. 2014) explored the use of ELDOR-Detected NMR (EDNMR) for use to study the

2004). **b**) Depiction of the  $g = 4.1$  feature resulting from the high-spin ( $S = 5/2$ ) isoform of PS II. Adapted from (Haddy et al. 2004). **c**) A scheme for the formation and conversion between states of the high-spin and low-spin isoforms of the  $\text{S}_2$  state

PS II OEC complex. W-band  $^{17}\text{O}$ -EDNMR experiments shifted the resonant side bands out of the main EPR region. PS II (75% isotope enrichment) flashed into the  $\text{S}_2$  state, gave  $^{17}\text{O}$  resonances for three distinct oxygens coordinating the  $\text{Mn}_4\text{O}_5\text{Ca}$  cluster. The largest couplings appeared from the  $\text{O}_5$  bridge and two waters bound to Ca and  $\text{Mn}_4$ . A doublet assigned to a  $\mu$ -oxo bridge revealed sensitivity to Ca/Sr exchange ( $\sim 5\%$  change in splitting), as well as  $\text{NH}_3$  binding (displacing  $\text{W}_1$ , 30% shift). X-band CW EPR spectra of the  $\text{Sr}^{2+}$ -substituted  $\text{S}_2$  state in PS II samples with and without  $\text{H}_2^{17}\text{O}$  indicate no line broadening upon  $^{17}\text{O}$  exchange. The lack of line broadening with  $^{17}\text{O}$  exchange demonstrates that the largest  $^{17}\text{O}$  coupling represents only one exchangeable oxo bridge. A recent TR-MIMS and TR-EDNMR approach (de Lichtenberg et al. 2024) identified the central oxygen bridge,  $\text{O}_5$  as having the same exchange rate as the substrate water  $\text{W}_5$ ; these rates were also shown to be similarly affected by  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  exchange, suggesting  $\text{O}_5$  fulfills all criteria for being the substrate water.

In 2017, Guo et al. (Guo and Barry 2017) conducted a study on the effect of Sr substitution on the  $\text{Tyr}_Z^*$  decay rate and PCET by monitoring the  $\text{Tyr}_Z^*$  EPR signal intensity as a function of time, when the  $\text{S}_2$  state was trapped by illumination at 190 K. At pH 7.5, Sr-substituted PS II results in a  $\text{Tyr}_Z^*$  decay rate which is decreased relative to the decay rate for native Ca PS II. However, at pH 6.0, the  $\text{Tyr}_Z^*$  decay rate is unaffected by Sr substitution. A two-pathway model was proposed to explain these results and is consistent with evidence towards a functional interaction of an  $\text{S}_2$  protonated water cluster and  $\text{Tyr}_Z$  proton donation pathway.

In 2015, Boussac and coworkers (Boussac et al. 2015) demonstrated that  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  exchange in *T. elongatus*

perturbs proportions of OECs showing high ( $5/2$ ) versus low-spin ( $1/2$ )  $S_2$  states. Different illumination conditions – a single laser flash at room temperature versus illumination at 198 K – were shown to produce distinct EPR signals associated with the  $S_2$  state; 1 flash produces an Sr-MLS alongside an  $S=5/2$  feature at g-factor 4.9. Illumination at 198 K reveals similar levels of the MLS signals alongside a distinct  $S=5/2$  feature: at g-factor 4.3. That annealing the 198 K-illuminated sample converted the  $g\sim 4.3$  signal to the 4.98 signal suggests an equilibrium between these heterogeneous states at higher temperatures not observed at 198 K. Twice-flashed samples revealed a characteristic  $S=3$ , g-factor 9.5 feature of the  $S_3$  state alongside a reduction in the Sr-MLS feature, consistent with their miss-factor. On the other hand, once-flashed (room temperature) then 198 K-illuminated samples showed similar levels of the g-factor 9.5  $S_3$  state feature, alongside a reduction in the g-factor 4.9 feature with no change to the Sr-MLS. Similarly, as was observed in the  $S_2$  states, annealing of the “once-flashed (RT), then-illuminated (198K)” sample produced g-factor 4.9 signal, at a reduction in the Sr-MLS intensity. These results suggest that  $S=5/2$   $S_2$  state centers advanced to  $S_3$  at 198 K, while  $S=1/2$   $S_2$  states could not.

Beal et al. (Beal et al. 2018) considered the  $S_2$ ,  $S_2\text{Tyr}_Z^\bullet$ , and  $S_3$  state of the OEC for Ca and Sr-substituted forms via a BS-DFT study. Computed tyrosyl radical g-tensors and  $^1\text{H}$  hyperfine couplings for the Ca PS II form found good agreement with prior EPR experiments (Sanakis et al. 2008; Boussac et al. 2009). Difference in the calculated g-tensor for the tyrosyl radical with Sr substitution was attributed to changes in the hydrogen-bonding environment in  $S_2\text{Tyr}_Z^\bullet$  caused by the presence of Sr. Modelling of the  $S_3$ -peroxo state for Ca PS II resulted in a change in ligation for Ca in the  $S_2$ - $S_3$  transition;  $W_3$  originally bound to Ca was found to deprotonate and migrate to a vacant adjacent Mn site. However, in Sr-substituted models, this spontaneous migration was not found to occur. Beal et al. went on to suggest that this deprotonation-migration response, leading to a difference in the coordination environment of the heterocation, is a substantial factor leading to the decrease in oxygen-evolving activity for Sr-substituted PS II.

A later work by Boussac and coworkers (Boussac et al. 2018) revealed that the  $S_2^{\text{LS}}$  and  $S_2^{\text{HS}}$  states' equilibrium is pH dependent, with  $\text{pK}_a$  of 8.3 for CaOEC and 7.5 for SrOEC. Subsequent DFT study revealed Ca $\rightarrow$ Sr exchange modifies electronic structure of indirect ligands to the heterocation; several groups within the active site such as W1, W2, Asp61, His332, and His337 showed changes in electronic structure – observed by a change in charge density – consistent with modification of the  $\text{pK}_a$  with heterocation exchange. Addition of ammonia reversed the effect of high pH, leading to a detectable  $\text{NH}_3$ - $S_2^{\text{LS}}$  being present for all

studied pH values, though  $\text{NH}_3$  binding is not present in the  $S_3\text{Tyr}_Z^\bullet$  state. That  $S_2^{\text{HS}}$  is observed to advance to  $S_3$  at 198 K, while  $S_2^{\text{LS}}$  advances only at 240 K and above provided strong experimental support for theoretical studies (Bovi et al. 2013; Retegan et al. 2014; Ugur et al. 2016) suggesting  $S_2^{\text{HS}}$  is an intermediate in  $S_3$  formation. Molecular interpretation of the  $S_2^{\text{HS}}$  given by our group as an early substrate binding state with  $\text{Mn}_1\text{-OH}$  fragment agrees with this overall conclusion (Pushkar et al. 2019).

## FTIR

In an FTIR report by Chu et al. (Chu et al. 2000), a Mn-O-Mn vibrational mode in the  $S_2$  state at  $606\text{ cm}^{-1}$  was revealed and characterized. The corresponding Mn-O-Mn mode in the  $S_1$  state was identified at  $625\text{ cm}^{-1}$ . This band, through sensitive to  $^{18}\text{O}$  substitution, is unaffected by  $\text{D}_2\text{O}$  or  $^{44}\text{Ca}\rightarrow^{40}\text{Ca}$  substitution. Substitution of  $\text{Ca}^{2+}\rightarrow\text{Sr}^{2+}$  upshifted this  $S_2$   $606\text{ cm}^{-1}$  band to  $618\text{ cm}^{-1}$ , indicating that this oxo interacts with the  $\text{Ca}^{2+}/\text{Sr}^{2+}$  in the OEC. A C=O stretching mode of protonated carboxylic acid can be expected in the  $1700\text{--}1750\text{ cm}^{-1}$  region, and in the  $S_2$ -minus- $S_1$  difference spectrum of wild type *Synechocystis* PCC 6803, many such bands appear (Pokhrel and Brudvig 2014). One such band appears at  $1747\text{ cm}^{-1}$  and is sensitive to the presence of  $\text{D}_2\text{O}$ ;  $\text{Ca}^{2+}\rightarrow\text{Sr}^{2+}$  substitution results in the disappearance of this band (Strickler et al. 2005; Service et al. 2010). A work featuring flash-induced FTIR difference spectroscopy considered difference spectra of four S-state transitions –  $S_1\rightarrow S_2$ ,  $S_2\rightarrow S_3$ ,  $S_3\rightarrow S_0$ , and  $S_0\rightarrow S_1$  – on flashed core PS II complexes from spinach and *T. elongatus*, both with  $\text{Ca}^{2+}$  and with  $\text{Sr}^{2+}$  substitution (Suzuki et al. 2006). This work observed significant intensity changes in symmetric  $\text{COO}^-$  peaks at the first and third flashes, leading to the conclusion that three carboxylate ligands are perturbed by  $\text{Ca}^{2+}\rightarrow\text{Sr}^{2+}$  exchange. These data and prior mutagenesis works suggest these carboxylate groups ligate the Mn ions and are strongly coupled to the heterocation, rather than directly ligating the  $\text{Ca}^{2+}/\text{Sr}^{2+}$ . A 2017 work by Kim and Debus (Kim and Debus 2017) revealed that the  $\text{Ca}^{2+}$ -bound  $W_3$  water molecule acts as the second substrate water molecule. The deprotonated form of  $W_3$  moves adjacent to  $\text{O}_5$  with the  $S_2\rightarrow S_3$  transition. Their work, through FTIR difference spectroscopy with  $\text{Ca}^{2+}\rightarrow\text{Sr}^{2+}$  substitution, revealed a singular water molecule whose D-O-D bending mode is eliminated immediately prior to O-O bond formation. A similar FTIR work was published in 2018 by Guo et al. (Guo et al. 2018) This report included a description of two CO stretching bands ( $1503\text{ cm}^{-1}$  and  $1478\text{ cm}^{-1}$ ) present in calcium-containing PS II in the  $S_2$  state trapped by illumination. They reveal a minor upshift in the  $1503\text{ cm}^{-1}$  band to  $1507\text{ cm}^{-1}$  for Sr-substituted PS II – a band which

is greatly reduced in Ca-depleted PS II. This result is consistent with the existence of a hydrogen-bonding interaction network between the Ca site and the Tyr<sub>Z</sub><sup>•</sup>.

## Computational modelling

### SrOEC

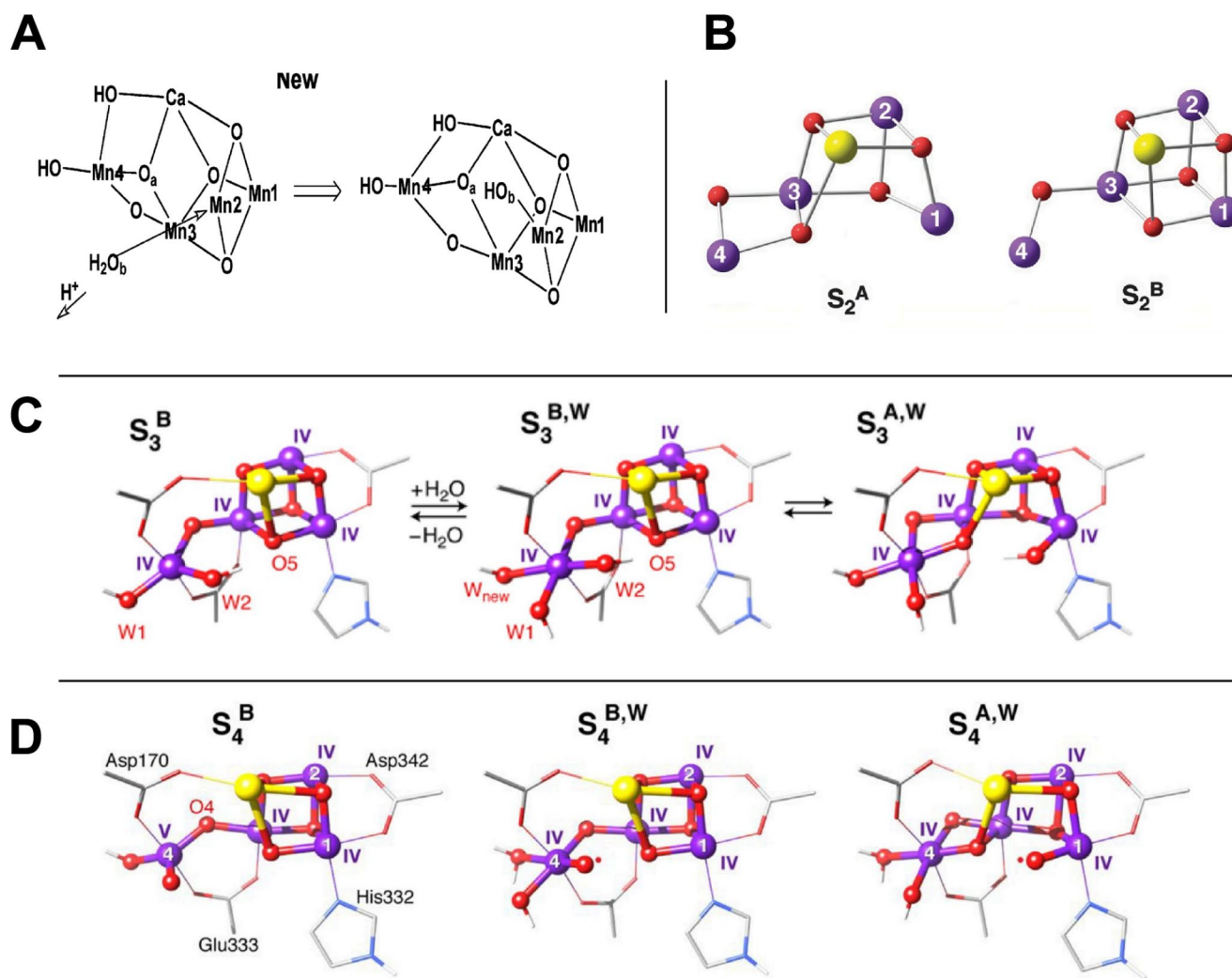
Many older computational models of the Sr-substituted OEC were based on the early 3.5 Å XRD structure produced by Ferreira et al. (Ferreira et al. 2004). The 2004 structure was an improvement over the first X-ray structure, obtained at 3.8 Å resolution for PS II isolated from the cyanobacterium *T. elongatus* by Zouni et al. (Zouni et al. 2001). A review by Gerecht et al. (Gerecht et al. 2016) discusses the role and history of Mn-Ca/Sr heterometallic compounds and their use in model study of the OEC of PS II. They note that the first Mn-Ca compound was published in 2000, followed five years later by discrete Mn-Ca and Mn-Sr compounds. As such, the emergence of higher-resolution XRD structures – both of native CaOEC and SrOEC – were vital to the development of sound computational models. Our focus will be on the more modern work that has found sufficient experimental support with spectroscopic methods.

One influential report (Siegbahn 2009) by Siegbahn drastically improved the level of detail of study of the active site by presenting a model size of 170 atoms – larger than used in previous computational studies – and found good agreement with available experimental spectroscopic evidence. The structures presented were based in part on two models: the aforementioned 3.5 Å model presented by Ferreira et al. (Ferreira et al. 2004), as well as the more recent 3.0 Å model reported by Loll et al. (Loll et al. 2005). This report, alongside prior computational work (Siegbahn 2008) on modelling the OEC, revealed that an oxygen radical preceded O-O bond formation – previous attempts to model water nucleophilic attack on Mn-oxo groups consistently led to substantial barriers (Siegbahn and Crabtree 1999). As such, the search for effective O-O bond mechanisms was conducted via exploring all possible O-O bonds in the most promising S<sub>4</sub> state structures (Siegbahn 2006). The structures produced by Siegbahn resembled a slightly distorted version of the London model (3.5 Å) (Ferreira et al. 2004; Sproviero et al. 2008). Subsequent addition of protons and electrons led to the lower S-state structures, with later optimizations achieving structures consistent with available XAS experiments, (Liang et al. 2000; Haumann et al. 2005), with shorter Mn-Mn distances to the dangling manganese (Mn<sub>4</sub>), as well as a reconstruction and rearrangement of the OEC in the S<sub>2</sub>→S<sub>3</sub> transition. The structural rearrangement, motivated by the binding of a second substrate water, occurs alongside

oxidation of Mn<sub>2</sub> from Mn<sup>III</sup> to Mn<sup>IV</sup>; with the binding of a water molecule to the dangle manganese, it becomes 7-coordinated, leading to a transfer of O<sub>5</sub> oxygen to Mn<sub>1</sub>, becoming 6-coordinated alongside its oxidation, Fig. 4a. With the loss of a proton to the protein environment, this rearrangement becomes feasible and substantially smaller than in previously reported models: Mn<sub>1</sub> moves by 0.4 Å, and all other metal atoms move by less than 0.2 Å. Despite this smaller-scale reconstruction of the OEC, it resulted in a higher energy of the S<sub>3</sub> state – already uphill – by 3 kcal/mol (0.13 eV). This energetic issue was resolved with a rearrangement of Glu189. Originally, Glu189 bridged Mn<sub>1</sub> and the calcium ion. However, its release from calcium and formation of a hydrogen bond with a water molecule allowed to bind to calcium led to more of the negative charge of Glu189 to be lent to Mn<sub>1</sub>, facilitating its oxidation and substantially lowering the energetic transition to S<sub>3</sub> by 13 kcal/mol (0.56 eV).

The 2014 computational work by Retegan and coworkers (Retegan et al. 2014) explores structures of the S<sub>2</sub>Tyr<sub>Z</sub><sup>•</sup> split signal state originating from each of the S<sub>2</sub> isoforms – open and closed cubane, Fig. 4b. They found that the oxidized S<sub>2</sub><sup>A</sup>Tyr<sub>Z</sub><sup>•</sup> and S<sub>2</sub><sup>B</sup>Tyr<sub>Z</sub><sup>•</sup> states formed prior to proton removal. The absence of valence isomers with Mn(III) oxidized rather than Tyr<sub>Z</sub> implied that chemical modification of the OEC is requisite to advance the S<sub>2</sub>-S<sub>3</sub> transition. For both structures, Tyr<sub>Z</sub> oxidation reoriented the dipole moment of the OEC, nearly coincident with the Mn<sub>4</sub>-W<sub>1</sub> bond and Asp61, suggesting that Tyr<sub>Z</sub> oxidation triggers deprotonation from W<sub>1</sub> to Asp61. Retegan et al. determined the tyrosyl radical g-tensor with the full model as well with unrelaxed subsets of the full model: Tyr<sub>Z</sub><sup>•</sup>-His with W<sub>4</sub> & HOH542, Tyr<sub>Z</sub><sup>•</sup>-His with W<sub>4</sub>, and Tyr<sub>Z</sub><sup>•</sup>-His alone. The model with the least electropositive environment about the tyrosyl (just Tyr<sub>Z</sub><sup>•</sup>-His) showed a substantial increase in the g<sub>x</sub> value to 2.0072. This suggests that the predicted g<sub>x</sub> value is highly dependent upon the presence of hydrogen-bonding interactions in which calcium-bound water molecules are involved. As such, Retegan et al. propose another role of the heterocation, in structuring water environment and optimizing the hydrogen-bonding network around Tyr<sub>Z</sub>; the ordering influence of the Ca<sup>2+</sup> fine-tunes the function of the Tyr<sub>Z</sub> residue by regulating its electron transfer properties. Rappaport (Rappaport et al. 2011) similarly claims that Ca/Sr exchange results in a change in the hydrogen-bonding network and lead to differences in entropy change in the S<sub>3</sub>Y<sub>Z</sub><sup>•</sup>→S<sub>0</sub> transition; the distribution and conformations of the substrate water molecules is critical to the catalytic rate.

A 2015 article (Vogt et al. 2015a) reported a QM/MM-based study exploring the degree of X-ray photoreduction present in the 1.9 Å CaOEC structure by Umena et al. (Umena et al. 2011) and in the 2.1 Å SrOEC structure by



**Fig. 4** Different structures relevant to computational modelling of the OEC. **a**) Proposed mechanism for  $S_2$  to  $S_3$  transition, alongside Glu189 forming a hydrogen bond with a Ca-bound water (not shown). Adapted from (Siegbahn 2009). **b**) Structures of the low-spin, open-cubane

( $S_2^A$ ) and high-spin, closed-cubane ( $S_2^B$ ) forms introduced in Pantazis et al. (Pantazis et al. 2012). Adapted from (Retegan et al. 2014). **c**) Models of the  $S_3$  and **d**)  $S_4$  states proposed by Krewald et al. Adapted from (Krewald et al. 2019)

Koua et al. (Koua et al. 2013). For both  $\text{Ca}^{2+}$  and  $\text{Sr}^{2+}$  OECs, Vogt et al. optimized structures with each cation in the resting dark-stable  $S_1$  state, as well as in a series of sequentially reduced states:  $S_0$ ,  $S_{-1}$ , and  $S_{-2}$ . Through comparison to the X-ray crystal structures and the optimized geometries – specifically of the bond lengths from the  $\text{Ca}^{2+}/\text{Sr}^{2+}$  to  $\text{O}_1$ ,  $\text{O}_2$ ,  $\text{O}_5$ ,  $\text{W}_3$ , and  $\text{W}_2$ . Their QM/MM models demonstrate that with  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution, an elongation of the heterocation bonds results, in agreement with XAS results. Based on the structural parameters studied, they suggest that the  $S_{-2}$  oxidation state of the  $\text{Mn}_4$  complex – corresponding to  $\text{Mn}_3^{\text{III}}\text{Mn}^{\text{II}}$  – is most consistent with the XRD data for both  $\text{Ca}^{2+}$  and  $\text{Sr}^{2+}$ -containing PS II. They note that their models only reproduce the  $\text{O}_5$  position when the oxygen is protonated; both  $\text{O}_4$  and  $\text{O}_5$  are coordinated in their models which best fit the XRD data for both heterocations in the  $S_2$  state.

The work of Pitari et al. (Pitari et al. 2015) in 2015 was similarly a QM/MM report inspired by the 2.1 Å XRD structure of the SrOEC; their focus was on modelling  $\text{Ca}^{2+}$ ,  $\text{Sr}^{2+}$ , and  $\text{Cd}^{2+}$  models in the  $S_2$  state. They suggested the role of the  $\text{Ca}^{2+}$  was both structural and functional, citing a difference in the computed  $\text{pK}_a$  for the water molecule bound to each of the heterocations; for  $\text{Sr}^{2+}$ , this  $\text{pK}_a$  value is 0.51 units greater than for  $\text{Ca}^{2+}$ , and the value for  $\text{Cd}^{2+}$  was 3.47 units smaller. A report by Amin et al. in 2021 (Amin et al. 2021) presents a DFT-based study with similar findings. Amin et al. employ the B3LYP hybrid functional to study Ca and Sr variants of  $S_2$  and  $S_3$  models; in particular, they examine two  $S_2$  structures:  $S_{2A}$ , with a neutral His190, and  $S_{2B}$ , where  $\text{W}_3$  is deprotonated, resulting in a  $\text{His190}^+$ . They find that the  $S_{2A} \rightarrow S_{2B}$  transition is more favorable in the CaOEC than the SrOEC for several choices of double-zeta,

polarized valence basis sets and the B3LYP functional. They also computed the  $pK_a$  values of  $W_3$  in the  $S_2$  state; the Sr water's  $pK_a$  is 4 units higher – qualitatively similar to previous literature (Pitari et al. 2015). Work by the Guidoni group (Bovi et al. 2013) suggests that the low-spin  $S_2$  state – associated with MLS – is a sequential precursor to the high-spin  $S_2$  state – associated with  $g=4.1$  signal – towards the formation of  $S_3$ , as previously hypothesized (Cox and Messinger 2013).

A 2019 computational work by Krewald and coworkers (Krewald et al. 2019) identifies three spectroscopically distinct structures of the  $S_3$  state and demonstrates that each isomer leads to a structurally distinct  $S_4$  state. Noting the hypothesized isoforms of the  $S_2$  state (open:  $S=1/2$ ,  $S_2^A$ ; closed:  $S=5/2$ ,  $S_2^B$ ) and that the  $S_2$ - $S_3$  transition may proceed with oxidation, deprotonation, and potential coordination of a water molecule, Krewald et al. posit three distinct structural models for the  $S_3$  state: (i)  $S_3^B$ , a model of five-coordinate, approximately trigonal-bipyramidal  $Mn^{IV}$  coordinated to the  $CaMn_3O_4$  cubane; (ii)  $S_3^{B,W}$ , a closed-cubane model generated by low-barrier water binding at  $Mn_4$  of  $S_3^B$  or by direct OH binding at internal site of  $Mn_4$ ; and (iii)  $S_3^{A,W}$ , a model with four six-coordinate  $Mn^{IV}$ . “W” implies a new water is bound to a manganese ion. See Fig. 4c. The  $S_3^B$  state was shown by Retegan et al. (2016) as requisite for  $S_3$  formation, since only the  $S_2^B$  component of the  $S_2$  state proceeds to  $S_3$ , as demonstrated in prior EPR study. The  $S_3^{A,W}$  state is supported by EXAFS simulations and by magnetic resonance studies; (Cox et al. 2014; Askerka et al. 2016)  $S_3^{B,W}$  connects the two. Producing  $S_4$  models from these distinct  $S_3$  models was done with consideration of the more studied  $S_2$ - $S_3$  transition; In the  $S_2$  and  $S_2Tyr_Z^\bullet$  states, the proton removed was part of  $W1$  in the open cubane  $S_2^A$ , and  $W1$  is the only Mn-bound  $H_2O$  in that state. Applying these considerations lead to the observation that the most acidic protons in  $S_3^{A,W}$  and  $S_3^{B,W}$  are on newly coordinated waters, the only doubly protonated water ligand; in  $S_3^B$ , this proton belongs to either of the hydroxyl ligands  $W1$  or  $W2$ . With the observation that following the  $S_3Tyr_Z^\bullet$  state substrates cannot exchange disfavors deprotonation of a water group in  $S_3^{A,W}$  or  $S_3^{B,W}$ . The resulting 3 models,  $S_4^B$ ,  $S_4^{A,W}$ , and  $S_4^{B,W}$  were produced (See Fig. 4d).  $S_4^{A,W}$  and  $S_4^{B,W}$  contain Mn(IV)-oxyl units –  $S_4^{A,W}$  is an open-cubane structure with oxyl attached to  $Mn_1$ , and  $S_4^{B,W}$  is a closed-cubane structure with oxyl *trans* to the  $W_1$  at (octahedral)  $Mn_4$ . These contrast to the genuine Mn(V)-oxo produced in the  $S_4^B$  model with five-coordinate, trigonal-bipyramidal-like  $Mn_4$ . Positioning the unbound water molecule in  $S_4^B$  to the second coordination sphere led to the models being energetically indistinguishable relative to the uncertainties of their computational methods. Noting that the high radical character of the oxyl groups in  $S_4^{A,W}$  and  $S_4^{B,W}$ ,

alongside the increased p/d-mixing of the Mn(V)-oxo group in  $S_4^B$  compared to  $Mn^{IV}$ -oxyl groups, Krewald suggests spectroscopic signatures of each model, were spectroscopic observation of  $S_4$  intermediates become possible. Spin density on the oxyl radical could be detectable by EPR methods such as  $O^{17}$  ENDOR, while Mn  $K_\alpha$  pre-edge intensity might reveal the  $S_4^B$  intermediate.

Allgöwer et al. (2022) performed a computational study to explore the redox-coupled protonation dynamics in various S-state transitions in photosystem II. They observed that with oxidation of  $Tyr_Z$ , conformational changes in the Asp61/Lys317 ion-pair lower the reaction barrier for water mediated proton transfer from  $W3$  ( $Ca^{2+}$ -bound) to Asp61. The deprotonation of  $W3$  substrate water results in its translation towards  $Mn1$ , in good agreement with the position identified in recent XFEL-based structures (Ibrahim et al. 2020). Optimized minimum energy reaction pathways reveal that the proton transfer from the  $Ca^{2+}$ -bound  $W3$  to Asp61 occurs across a series of five water molecules; the overall reaction barrier was shown to be  $14 \text{ kcal mol}^{-1}$ . With subsequent oxidation of the OEC, the proton transfer barrier about the water ligand sphere of the OEC is further lowered. Allgöwer et al. ultimately propose a radical-coupling mechanism towards O-O bond formation.

## Synthetic analogues

An electrochemistry-based approach by Agapie and Tsui (Tsui and Agapie 2013) investigated compounds of  $MMn_3O_4$  structure (where  $M \rightarrow Ca^{2+}, Sr^{2+}, Zn^{2+}, Sc^{3+}, Y^{3+}$ , and  $Mn^{3+}$ ). They observed that cluster reduction potentials of these compounds and of similar mixed-metal Mn-tetraoxido complexes correlate with the Lewis acidity of the apical redox-inactive metal. This relates to similar findings for heterometallic Mn-dioxo clusters they worked with prior. The further finding that  $Ca/Sr$ - $Mn_3O_4$  clusters shared similar reduction potentials, combined with the knowledge that  $Ca$  and  $Sr$  are the only atoms for a functional  $MMn_4O_5$  OEC reinforced the idea that these heterocations may have a role in modulating the reduction potential of the cluster. The relationship in findings between the tetraoxo and dioxo models suggests that the role of the Lewis acidities of the incorporated cations is held across a broadly diverse set of structural motifs. A subsequent computationally focused work by Krewald and coworkers (Krewald et al. 2016) further investigated these findings. They computed the relative redox potentials of various substituted derivatives of the OEC with cations  $Sr^{2+}$ ,  $Gd^{3+}$ ,  $Cd^{2+}$ ,  $Zn^{2+}$ ,  $Mg^{2+}$ ,  $Sc^{3+}$ ,  $Na^+$  and  $Y^{3+}$ , across two sequential transitions of the catalytic cycle. This computational approach was first validated against several  $Mn_3AO_4$  model complexes – good structural mimics of the OEC – and found good reproduction of the

empirically observed relationship between Lewis acidities of the cations and reduction potentials of the clusters. However, the correlation was not observed with the OEC substitutions. The cluster's redox potential within the enzyme was dependent upon the charge of the redox-inactive cations; in other words, the redox potential was shown to be insensitive to the identity of the substitution, hindering the stance that one primary role of the heterocation is in redox-tuning.

Computational study of synthetic OEC analogues is also worth consideration. The work of Kanady and coworkers (Kanady et al. 2014) presents a synthetic strategy and roadmap for producing accurate synthetic models of the tetramanganese oxygen evolving complex. They synthesized  $\text{Mn}^{\text{IV}}_3\text{GdO}_4$  and  $\text{Mn}^{\text{IV}}_3\text{CaO}_4$  cubanes. A procedure of sequential ligand substitutions produced substantial distortions resulting in a highly asymmetric cubane; these models, when coordinated by a fifth metal ion, were able to mirror the structure of the OEC by displaying the known cubane motif alongside a dangling transition metal. Other works led by Zhang (Li et al. 2020; Yao et al. 2021) report synthetic  $\text{Mn}_4\text{XO}_4$  clusters –  $\text{X}=\text{Ca}, \text{Y}, \text{Gd}$  – which closely resemble the OEC, closely mirroring the oxidation states of the manganese atoms and overall redox activity. Their works demonstrate the capabilities of rare-Earth metals to structurally replace Ca in neutral  $\text{Mn}_4\text{XO}_4$  clusters; this challenges the view of the role of Lewis acidity modulating the redox potential of heteronuclear-oxide clusters. These synthetic model compounds improve the capabilities to study electronic structure, spectroscopy and reactivity of the OEC. Of similar interest is the study of the closest-known catalyst analog for PS II:  $[\text{Co}_4\text{O}_4\text{Py}_4\text{Ac}_4]$  ( $\text{Py}$  = pyridine and  $\text{Ac}=\text{CH}_3\text{COO}^-$ ), recently reported upon by several groups (Nguyen et al. 2015; Smith et al. 2015; Ezhov et al. 2021).

## Simulated EXAFS analysis of $\text{S}_3$ state models

Here, we examine three computational models, each analogous to those in the preceding computational section, Fig. 5a. We refer to these three  $\text{S}_3$  state models as  $\text{S}_{3\text{Mn}=\text{O}}$ : model based on Suga et al. (2017) coordinates (PDB, 5WS6);  $\text{S}_{3\text{Mn}-\text{OH}}$ : a model with a different protonation pattern, with one proton moved from the  $\text{Mn}_4$  water to the  $\text{Mn}_1=\text{O}$  oxo-group forming a  $\text{Mn}_{1\text{D}}-\text{OH}$  moiety, and  $\text{S}_{3\text{OO}}$ : a peroxo isoform of the  $\text{S}_3$  state which was hypothesized to be forming prior to the final electron transfer to  $\text{Tyr}_Z$ . Previous work (Boussac et al. 2015) showed that the capacity of the OEC in  $\text{S}_3$  to absorb near-infrared (NIR) light to induce new EPR signals from the OEC  $\text{S}_3$  state at cryogenic temperatures suggests the possibility for a potential high-energy intermediate of the  $\text{S}_3$  state; other work by our group have demonstrated that such an intermediate may be accessible (+0.2 eV, easily

achievable by NIR light,  $\sim 1.5$  eV), while  $\text{S}_4$  generation prior to O-O bond formation does not improve energetics towards O-O bond formation (Pushkar et al. 2018).

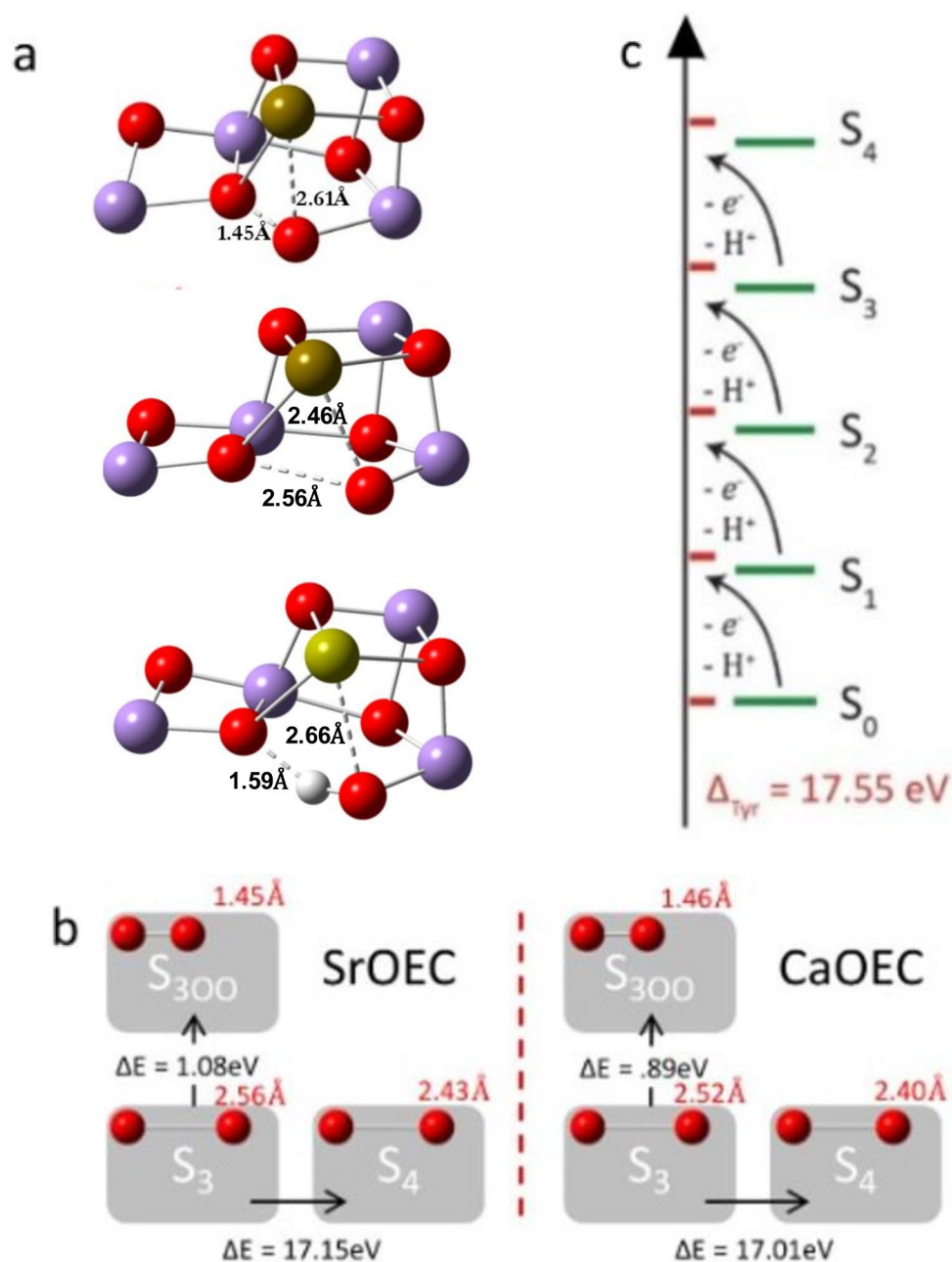
Table 2 summarizes the key supporting energies for the driving force and the energies of the S-state transitions; see Tables S3 and S4 for charge, spin, and electronic energies of each model. Note that with computation of redox processes – S-state transitions – we consider both the energies of the proton and electron upon their removal, avoiding the need to compute individual proton energies; as such, we can directly compare computed energies with the energy of the driving force,  $\text{Tyr}_Z^\bullet$  reduction with the PCET process:  $\text{Tyr}-\text{OH}=\text{Tyr}-\text{O}^\bullet + \text{e}^- + \text{H}^+$ ,  $\sim 17.55$  eV. We have to note that  $\text{S}_3$  to  $\text{S}_0$  transition, unlike other S-states transitions in Table 2 draws energy greater than available 17.55 eV from  $\text{Tyr}_Z^\bullet$  reduction. This can be due to the imperfection of the DFT model (note QM/MM might be used in the future to allow additional restrains on amino acids positions) or in the sophisticated scheme realized by PS II ligand environment where excess energy in prior  $\text{S}_1$ - $\text{S}_3$  transitions is stored in conformational or  $\text{H}^+$  uphill configurations of higher energy and later released to help the energy demanding  $\text{S}_3$  to  $\text{S}_0$  transition.

We observe that with the CaOEC Suga model (PDB, 5WS6) as a starting point and substitution of  $\text{Ca}^{2+}\rightarrow\text{Sr}^{2+}$  followed by DFT optimization – see >Table S2 for a summary of the  $\text{S}_1$  geometric differences between the CaOEC and SrOEC models – that a suitably analogous SrOEC  $\text{S}_3$  state may be found. In fact, the  $\text{S}_2\rightarrow\text{S}_3$  transition is slightly easier for SrOEC than for the CaOEC model. This energy difference  $\Delta\Delta E_{\text{Ca}\rightarrow\text{Sr}}$  is like that of the initial  $\text{S}_1\rightarrow\text{S}_2$  transition, with  $\Delta\Delta E_{\text{Ca}\rightarrow\text{Sr}}\sim 0.2$  eV. This suggests that when optimized with the BP86 functional, incorporation of water in the  $\text{S}_2\rightarrow\text{S}_3$  state transition, alongside the resulting conformational change about the ligating structure to the OEC, the SrOEC  $\text{S}_3$  state may be more stabilized relative to CaOEC. Despite the difference in  $\text{S}_3$  state energetics, the relative energies of the peroxo  $\text{S}_{3\text{OO}}$  state and the  $\text{S}_3$  state between models does not change drastically (0.89 eV for Ca and 1.08 eV for Sr, BP86 functional); the work of Yamaguchi addresses peroxo and superoxo formation models in depth (Isobe et al. 2016, 2019; Yamaguchi et al. 2022).

The prior DFT work on the  $\text{S}_{3\text{OO}}$  state revealed that the B3LYP\* functional (containing 15% Hartree-Fock exchange) resulted in a decrease in the energy difference of the peroxo state for CaOEC, a finding that is consistent with other works, notably of Isobe et al. (Isobe et al. 2016).

We do note that the geometric optimization of the SrOEC  $\text{S}_{3\text{OO}}$  state was 0.1 eV more stable in the triplet ( $\text{S}=1$ ) than quintet ( $\text{S}=2$ ) spin state. Indeed, when studied previously (Pushkar et al. 2018), the CaOEC peroxo isoform tended to converge to the  $\text{S}_3$  state when the system spin exceeded

**Fig. 5** Schematic summarizing key geometric and energetic results from our DFT study. **a)** Structures of the Sr-substituted  $S_{3OO}$  (top),  $S_{3Mn=O}$  (middle), and  $S_{3Mn-OH}$  (bottom) models. Bond distances of Sr-O<sub>6</sub> and O<sub>5</sub>-O<sub>6</sub> are indicated. **b)** Scheme comparing the energies of the  $S_3$ ,  $S_{3OO}$ , and  $S_4$  states for both CaOEC and SrOEC model, using the BP86 functional. **c)** Energy diagram of the S-state transitions with every step been modeled as a proton coupled electron transfer (PCET) to ensure a straight forward comparison with a driving force of Tyr<sub>Z</sub> redox couple. Red ticks represent the energy change due to Tyr<sub>Z</sub><sup>•</sup> reduction with the PCET process: Tyr-OH = Tyr-O<sup>•</sup> + e<sup>-</sup> + H<sup>+</sup>. Implicit treatment of protons. Adapted from (Pushkar et al. 2018)



$S = 1$ , though could converge to the peroxo isoform (O<sub>5</sub>-O<sub>6</sub> bond distance  $\sim 1.5 \text{ \AA}$ ) when optimized in the triplet state, a bond distance in good agreement with the 2017 structure of Suga et al. (1.5 Å) but disagreement with their later structure (2.0 Å). SrOEC in the  $S_3$ -state converges to a peroxo isoform for both  $S = 1$  (1.45 Å) and  $S = 2$  (1.44 Å). Heisenberg exchange Hamiltonian fitted to broken symmetry calculations was used to confirm the total spin of CaOEC DFT models for the  $S_3$  state (Pushkar et al. 2018). Such analysis for SrOEC is outside the scope of this manuscript, thus listed spin states of the SrOEC DFT models are suggestive at the moment. The  $S_3 \rightarrow S_4$  and  $S_3 \rightarrow S_0$  transitions proceed

with a similar  $\Delta\Delta E_{Ca \rightarrow Sr}$  of  $\sim 0.1 \text{ eV}$  and  $\sim 0.2 \text{ eV}$ . Figure 5 summarizes the key geometric and energetic findings of the  $S_3$  and later stages of the Kok cycle for the computed CaOEC and SrOEC models. A refinement of these SrOEC models may be possible with the use of computational methods employing some Hartree-Fock exchange, potentially resulting a further stabilization of the  $S_{3OO}$  state for SrOEC. However, the current model predicts a more uphill transition for  $S_3 \rightarrow S_{3OO}$  ( $\Delta\Delta E_{Ca \rightarrow Sr} \sim 0.2 \text{ eV}$ ) in SrOEC than in CaOEC, suggesting that access to O-O bond formation early in the  $S_3$ -to- $S_0$  transition is faster for native than SrOEC,

**Table 2** Supporting energies for driving force and energies of the S-state transitions\*. BP86 methodology, def2tzvp basis set DFT calculations. By considering the energies of both the proton and electron upon their removal,  $\Delta E$  incorporates their combined contributions, circumventing the need to estimate individual proton energies, which could be subject to significant error

|                                                                                                       | BP86 / def2tzvp Energy (eV) |                   |
|-------------------------------------------------------------------------------------------------------|-----------------------------|-------------------|
| H <sub>2</sub> O                                                                                      | -2080.665                   |                   |
| OH <sup>-</sup>                                                                                       | -2062.975                   |                   |
| O <sub>2</sub>                                                                                        | -4092.570                   |                   |
| Tyr-OH                                                                                                | -8369.840                   |                   |
| Tyr-O <sup>•</sup>                                                                                    | -8352.286                   |                   |
| Tyr-OH = Tyr-O <sup>•</sup> + e <sup>-</sup> + H <sup>+</sup>                                         | 17.55                       |                   |
|                                                                                                       | Ca OEC Model (eV)           | Sr OEC Model (eV) |
| S <sub>1</sub> = S <sub>2</sub> + H <sup>+</sup> + e <sup>-</sup>                                     | 16.33                       | 16.15             |
| S <sub>2</sub> + H <sub>2</sub> O = S <sub>3</sub> + H <sup>+</sup> + e <sup>-</sup>                  | 17.10                       | 16.88             |
| S <sub>3</sub> = S <sub>4</sub> + H <sup>+</sup> + e <sup>-</sup>                                     | 17.01                       | 17.15             |
| S <sub>3</sub> + H <sub>2</sub> O = S <sub>0</sub> + H <sup>+</sup> + e <sup>-</sup> + O <sub>2</sub> | 19.45                       | 19.66             |
| S <sub>3</sub> → S <sub>3OO</sub>                                                                     | 0.89                        | 1.08              |

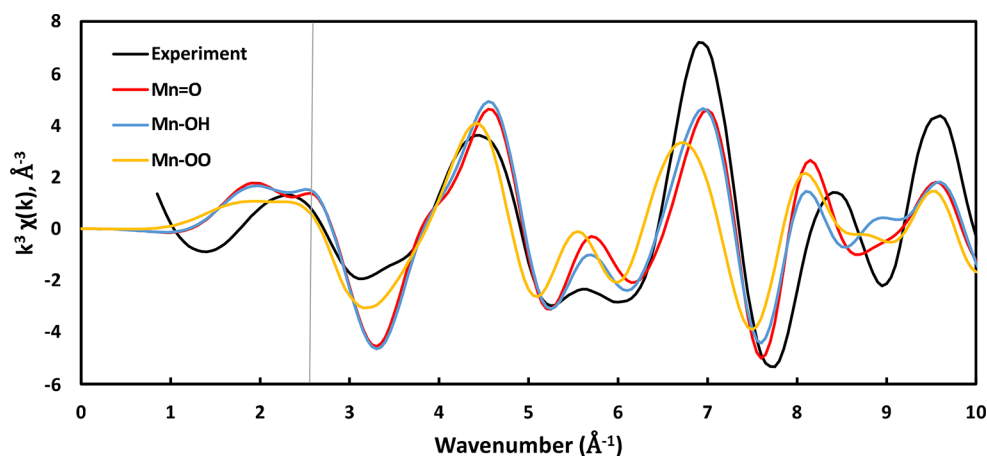
\* While release of protons to the lumen is S-state dependent, the Kok cycle was DFT modeled for OEC core with every step been a proton coupled electron transfer (PCET) to ensure a straightforward comparison with a driving force of Tyr<sub>Z</sub> redox couple

consistent with the lower rate of oxygen-evolving activity in SrOEC.

With three optimized models for CaOEC, we compute simulated Mn EXAFS – see Methods, supplementary information – for each of the models and compare against experimental data for S<sub>3</sub> state Fig. 6. MnXAS from CaOEC spinach samples was reported earlier (Pushkar et al. 2007) measured by the same techniques as in (Yano et al. 2005b). Each simulated spectrum from each computational model

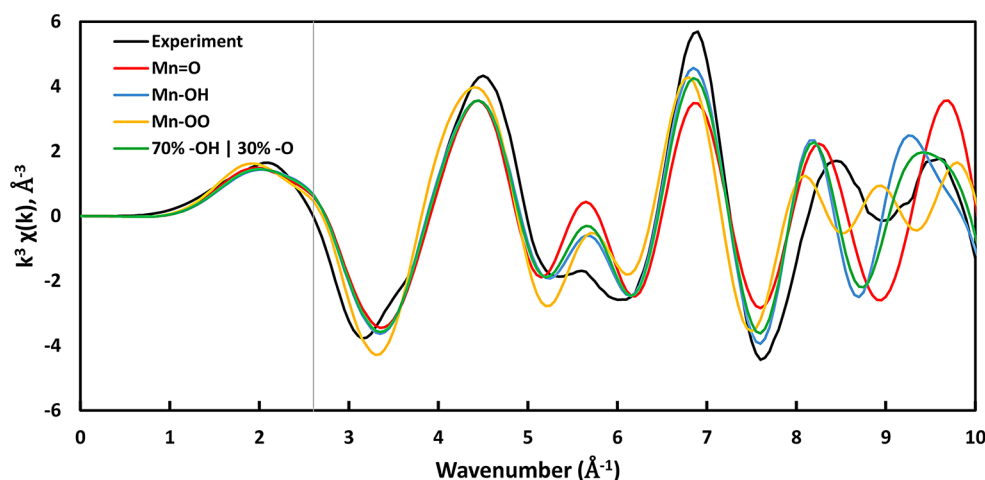
had its amplitude rescaled by a unique multiplicative constant to best fit the experimental data by least squares; see methods in supplementary information. By comparison of least-squares fitting for each computational model's simulated spectrum against the experimental spectrum, the hydroxide model, S<sub>3</sub> Mn-OH provides the best fit. Models of both Mn=O and Mn-OH protonation states on Mn<sub>1</sub> fit the experimental spectrum relatively well until wavenumber 8.5 Å<sup>-1</sup>, at which point all the models fit relatively poorly. Smoothing of the experimental spectrum (Savitzky-Golay algorithm, 10 points) does not substantially change the required vertical scaling of the simulated spectra nor the relative quality of fits and agreement with experiment, Fig. 6, S1-S2, and Table S5. The Fourier transforms of each of the computational models are overlaid against the Fourier transform in the experimental data in Figure S3. The aforementioned peaks I, II, and III are reproduced well for each of the three models, though the relative peak intensities differ slightly. This is easily explained by the selection of k-range used for the Fourier transform; a small difference in range can produce substantial changes in the R-space representation.

Similarly, the XRD structures from Suga et al. in 2017 and 2019 were evaluated and compared with experiment. The more recent 2019 structure provided a better fit than the 2017 structure, Figure S4 and Table S5. The S<sub>3OO</sub> model and the Suga 2017 XRD structure both feature a very short O<sub>5</sub>-O<sub>6</sub> bond distances ~ 1.5 Å. That these provide poor fits with the deconvoluted S<sub>3</sub> state data suggests that such a characteristic is not representative of the S<sub>3</sub>-state geometry. Despite both XRD-structure-based models fitting more poorly than DFT-optimized hydroxide models, the XRD-based structures did provide some insights. Despite the maximum scattering distance parameter being set equal in both DFT-optimized



**Fig. 6** Overlays of CaOEC Mn EXAFS spectra from spinach and the three computational models discussed. S<sub>3</sub> Mn=O: model based on Suga et al. (Suga et al. 2017) coordinates (PDB, 5WS6); S<sub>3</sub> Mn-OH: a model with a different protonation pattern, with one proton moved from the

Mn<sub>4</sub> water to the Mn<sub>1</sub>=O oxo-group forming a Mn<sub>1D</sub>-OH moiety, and S<sub>3</sub> OO: a peroxo isoform of the S<sub>3</sub> state forming prior to the final electron transfer from Tyr<sub>Z</sub>. Fitting is done in the region of 2.6 Å<sup>-1</sup>-10.0 Å<sup>-1</sup>, marked by the vertical gray line. k-weight = 3



**Fig. 7** Overlays of SrOEC Mn EXAFS spectra from experiment and the computational models. Experimental data is smoothed, deconvoluted, twice-flashed *T. elongatus* (Pushkar et al. 2008).  $S_{3\text{Mn=O}}$ : model based on Suga et al. (Suga et al. 2017) coordinates (PDB, 5WS6);  $S_{3\text{Mn-OH}}$ : a model with a different protonation pattern, with one proton moved from the  $\text{Mn}_4$  water to the  $\text{Mn}_1=\text{O}$  oxo-group forming a  $\text{Mn}_{1\text{D}}$ -

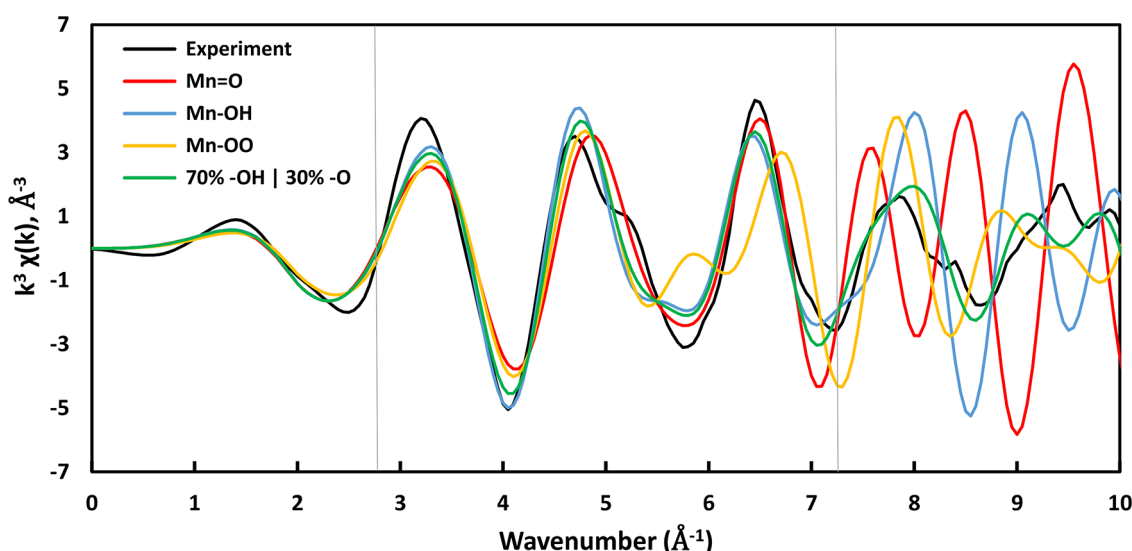
and XRD-derived structures' simulations and the much larger "model" size of the XRD structures, both approaches yielded the same number of scattering paths per Mn atom. This suggests that DFT model size ( $\sim 110$  atoms, containing the tetramanganese cluster and the first ligation sphere) is not a limiting factor in the number of scattering paths and may yield better fits against experimental EXAFS than XRD-based structures.

In a similar way, we will now consider analogous SrOEC models fitted against deconvoluted  $S_3$  state *T. elongatus* experimental Mn EXAFS (Pushkar et al. 2008). The DFT structures were optimized following  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution in the models of  $S_1$ - $S_3$  states; see methods in supplementary information. With some of the already discussed works highlighting possible heterogeneities of the  $S_3$  state, we will also consider linear combination of the  $\text{Mn}_{1\text{D}}=\text{O}$  and  $\text{Mn}_{1\text{D}}\text{-OH}$  protonation state models. Figure 7 provides overlays of each of the computational models presented. The combination of  $x\%$  hydroxide and  $(100-x)\%$   $\text{Mn}=\text{O}$  protonation state models was itself optimized by similar least-squares methodology and resulted in 70%  $\text{Mn-OH}$  and 30%  $\text{Mn}=\text{O}$  model. The improved fit with a linear combination of chemically distinct structures is consistent with recent claims of heterogeneity in the  $S_3$  state. While all three base computational models fit reasonably well at low wavenumber, the linear combination of the  $\text{Mn-OH}$  and  $\text{Mn}=\text{O}$  models produce a noticeably better fit at high wavenumber, though the  $\sim 8.5 \text{ \AA}^{-1}$  wavenumber region provides a poor fit for all simulated spectra. Of the three computational models, the hydroxide-based model provided the best least-squares fit against the experimental data up to  $k=10 \text{ \AA}^{-1}$ , and the linear combination model provided a substantially

OH moiety; a linear combination of these, 70% hydroxide model plus 30%  $\text{Mn}=\text{O}$  model; and  $S_{3\text{OO}}$ : a peroxo isoform of the  $S_3$  state forming prior to the final electron transfer from  $\text{Tyr}_Z$ . Fitting is done in the region of  $2.6 \text{ \AA}^{-1}$ - $10.0 \text{ \AA}^{-1}$ , marked by the vertical gray line.  $k\text{-weight}=3$

improved fit, Table S5. That a linear combination mostly constituted of the  $\text{Mn-OH}$  protonation model improves fitting to the experimental spectrum seems appropriate; the simulated Mn and FTs based on the  $\text{Mn-OH}$  model reproduce experimentally observed (Pushkar et al. 2008) changes in the  $S_2 \rightarrow S_3$  transition: a vanishing of peak III, Fig. 5a.

The merits of the linear combination of models EXAFS approach are more pronounced with Sr EXAFS, Fig. 8, S5-6. One immediate observation is that the individual DFT-based models demonstrate a very poor fit against the experimental Sr EXAFS above the  $\sim 8 \text{ \AA}^{-1}$ . This limited the range of  $k$ -space least-squares fitting to a maximum of  $7.25 \text{ \AA}^{-1}$ . In this range, the hydroxide model again fits better than the other two computational models, Table S5. However, by repeating the same linear combination of spectra – 70% hydroxide model and 30%  $\text{Mn}=\text{O}$  model – we find a reasonable agreement with experiment, up to wavenumber  $10 \text{ \AA}^{-1}$ . Experimental Sr EXAFS has shown that in the  $S_2 \rightarrow S_3$  transition, the Mn-Sr peak ( $\sim$  apparent distance  $\sim 2 \text{ \AA}$ ) split into two distinct peaks with slightly lower and higher apparent distances (Pushkar et al. 2008). In the simulated spectra, a similar splitting is present for the  $S_{3\text{OO}}$  and the  $S_{3\text{Mn-OH}}$  models, less so the  $S_{3\text{Mn=O}}$  structure, Figure S5b. As is shown in Fig. 7, a linear combination (mostly constituted of the  $\text{Mn-OH}$  protonation model) of  $\text{Mn-OH}$  and  $\text{Mn}=\text{O}$  spectra find good agreement with experimental  $k$ -space spectra. The fact that the linear combination of  $\text{Mn-OH}$  (which reproduces experimentally observed FT trends) and  $\text{Mn}=\text{O}$  (which reproduces these Sr FT trends more poorly) models improves the quality of the fit is likely due to a combination of two effects. One is that the linear combination of chemically distinct models may better represent a heterogeneity



**Fig. 8** Overlays of SrOEC Sr EXAFS spectra from experiment (twice-flashed *T. elongatus*, deconvoluted  $S_3$ ) and the computational models. Experimental data is smoothed, deconvoluted, twice-flashed *T. elongatus* (Pushkar et al. 2008).  $S_3$   $Mn=O$ : model based on Suga et al. (Suga et al. 2017) coordinates (PDB, 5WS6);  $S_3$   $Mn-OH$ : a model with a different protonation pattern, with one proton moved from the  $Mn_4$

water to the  $Mn_1=O$  oxo-group forming a  $Mn_{1D}-OH$  moiety; a linear combination of these, 70% hydroxide model plus 30%  $Mn=O$  model; and  $S_3$   $OO$ : a peroxo isoform of the  $S_3$  state forming prior to the final electron transfer from Tyr<sub>Z</sub>. Fitting is done in the region of  $2.8\text{ Å}^{-1}$ – $7.5\text{ Å}^{-1}$ , marked by the vertical gray lines.  $k$ -weight = 3

in the  $S_3$  state. Additionally, however, there may be effect that a linear combination of different models' Sr EXAFS may capture some of the large degree of structural disorder about Sr: the likely heterogeneity in Sr position and scattering distances; see large Debye-Waller factors for Sr EXAFS fits (Pushkar et al. 2008). Indeed, even simulated Sr EXAFS from the large Koua model (Koua et al. 2013) is unable to capture the decreased amplitude at large wavenumber when compared to an experimental  $S_2$  state spectrum, Figure S6.

Holistically, we can draw several useful conclusions from enumerated comparisons. Computational models of the CaOEC in the  $S_3$  state with a  $Mn_{1D}-OH$  moiety provide the best fits against  $S_3$ -state experimental data. DFT geometry optimizations on  $Ca^{2+} \rightarrow Sr^{2+}$  models demonstrate the same results for SrOEC fitting to experimental EXAFS. Given the similarity of CaOEC and SrOEC Mn EXAFS, that  $Mn_{1D}-OH$  remains the best-fitting model to the experimental data is unsurprising and supports the use of the computational model for SrOEC. A similar conclusion was recently achieved via DFT based analysis of high energy resolution fluorescence detected (HERFD) X-ray absorption spectra of the  $S_3$  state showing preference for the model with  $Mn_{1D}-OH$  moiety (Chrysinia et al. 2023). Turning to Sr EXAFS fitting (Ca EXAFS are non-viable due to difficulties in experimentation at low Ca K-edge energy  $\sim 4\text{ keV}$ ), we find that Sr EXAFS find poor agreement with all experimental models past  $7.5\text{ Å}^{-1}$ . This suggests possible heterogeneity in Sr position and scattering distances, resulting in greater disagreement with experimental data at higher

wavenumber. Given that XRD structures with the same  $R_{Max}$  parameter produced the same number of scattering paths as did the CaOEC DFT models, this aspect is likely not responsible for the poor agreement at high wavenumber. DFT-based models that extend beyond the second ligation sphere about the OEC, or models with appropriate structural constraints upon the first ligation sphere may better model the protein environment about the OEC. This is especially true for SrOEC, which are modelled as a perturbation of the CaOEC models derived from XRD structure. "SrOEC 2F" XRD structures may inform models able to be better fit to Sr EXAFS and gain insight into the heterocation EXAFS. However, if we consider proposals of a heterogeneity in the  $S_3$  state and apply them to the current least-squares fitting methodology applied to DFT-optimized models, we can find stronger agreement with experiment, particularly for the Sr EXAFS simulations. Given that a linear combination of different  $S_3$  models – or a heterogeneous  $S_3$  state model – with different protonation on  $Mn_{1D}$  were able to fit experiment in ways single structure modelling cannot, we cannot exclude the possibility of a heterogeneity in the  $S_3$  state. Furthermore, such a modelling approach may address the likely heterogeneity of Sr position and scattering distances in Sr-substituted specimens.

## Conclusions

The study of natural photosynthesis informs the study of artificial photosynthesis and inspires rational design of water oxidation catalysts. With this investigation into the effects of  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution in the OEC of PS II, we can hope to achieve mechanistic insights transferable to designs of artificial systems. Early hypotheses suggested that the role of the heterocation in the OEC was in altering the redox chemistry of the cluster itself; indeed, many cubane-like model compounds demonstrated altered redox chemistry with different alkaline and atomic substitutions. However, in the more complex OEC, similar findings were not observed. Further computational and electrochemical study then investigated the influence of  $\text{Ca}^{2+}$  and  $\text{Sr}^{2+}$  on water pathways and hydrogen-bonding networks. Indeed, the  $\text{Ca}^{2+}/\text{Sr}^{2+}$  ion is coordinated by two waters and forms a hydrogen-bonded network ending at the redox-active  $\text{Tyr}_Z$ . Many works have suggested the subtle 0.1–0.2 Å increase in Mn–Mn and Mn–Sr bond distances resulting from  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution, and this could very possibly disturb the hydrogen-bonding network and alter the electrochemistry of the redox active moiety; such has been suggested by both computational and EPR study. This influence on a redox-active unit then finds agreement with electrochemical studies on model cubane systems. Similar computational works explored the movement of substrate waters about the calcium atom; conformational changes about the tetramanganese cluster may then disrupt the hydrogen bonding network facilitating the substrate waters' movement. Given these insights, it is possible that the  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution is the only one which maintains sufficiently similar valence chemistry while perturbing the surrounding water and hydrogen-bonding network minimally to retain oxygen-evolving activity.

To reveal insights into the O–O bond formation chemistry of PS II, researchers have employed a variety of experimental and computational techniques, directed both at the protein and through model complexes. One means with which to probe the OEC core of the photosynthetic protein is the replacement of  $\text{Ca}^{2+}$  with  $\text{Sr}^{2+}$ , the only atom substitution able to retain (diminished) oxygen-evolving activity. These investigations encompassed native cyanobacteria and spinach variants with Ca-based systems, as well as proteins grown and reconstituted with Sr. The development of ever-higher resolution crystal structures of the OEC prepared in various stages of the Kok cycle, as well as the emergence of sophisticated DFT and QM/MM models consistent with the spectroscopic data, yield substantial insights into the mechanistic and structural effects of  $\text{Ca}^{2+} \rightarrow \text{Sr}^{2+}$  substitution. These combined studies overall do not contradict the hypothesis of O–O bond formation through  $\text{O}_5$  and  $\text{O}_6$ . Such

an interaction could occur with a peroxo-like formation in the  $\text{S}_3\text{Tyr}_Z^-$  oxidation state, prior to the final oxidation event. These collective efforts aim to develop a robust mechanistic understanding of the O–O bond formation process in PS II, Nature's finest catalyst. Moreover, these developments contribute substantially to the ongoing pursuit of designing artificial water-splitting catalysts capable of addressing the world's current energy crisis through sustainable means.

**Supplementary Information** The online version contains supplementary material available at <https://doi.org/10.1007/s11120-024-01117-2>.

**Acknowledgements** This research was supported by the NSF, CHE-2303743 (Y.P.) and by GM132024 (T32 Molecular Biophysics Training Program to G.B.) from the National Institute of General Medicinal Sciences (NIGMS). We thank Dr Alireza K. Ravari for the helpful discussion.

**Author contributions** G.B.: writing – original draft preparation; data curation; investigation; methodology. Y.P.: writing – review and editing; supervision, project administration; funding acquisition; conceptualization. All authors have read and agreed to the published version of the manuscript.

**Data availability** No datasets were generated or analysed during the current study.

## Declarations

**Competing interests** The authors declare no competing interests.

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