



OPEN

Synchronicity of the Gulf Stream path downstream of Cape Hatteras and the region of maximum wind stress curl

Ian Gifford¹, Avijit Gangopadhyay^{1,2}✉, Magdalena Andres³, Hilde Oliver³,
Glen Gawarkiewicz³ & Adrienne Silver³

The Gulf Stream, a major ocean current in the North Atlantic ocean is a key component in the global redistribution of heat and is important for marine ecosystems. Based on 27 years (1993–2019) of wind reanalysis and satellite altimetry measurements, we present observational evidence that the path of this freely meandering jet after its separation from the continental slope at Cape Hatteras, aligns with the region of maximum cyclonic vorticity of the wind stress field known as the positive vorticity pool. This synchronicity between the wind stress curl maximum region and the Gulf Stream path is observed at multiple time-scales ranging from months to decades, spanning a distance of 1500 km between 70 and 55W. The wind stress curl in the positive vorticity pool is estimated to drive persistent upward vertical velocities ranging from 5 to 17 cm day⁻¹ over its ~ 400,000 km² area; this upwelling may supply a steady source of deep nutrients to the Slope Sea region, and can explain as much as a quarter of estimated primary productivity there.

The Gulf Stream (GS) is well recognized as the northward-flowing western boundary current of the anticyclonic North Atlantic subtropical gyre. At Cape Hatteras, North Carolina the GS separates from the continental margin and becomes a free jet flowing east-northeastward. Gulf Stream path variability over the 2500-km distance downstream of Cape Hatteras has impacts that span fisheries responses¹ to atmospheric events² and the GS path is often interpreted as an indicator of climate-related changes in the Atlantic Meridional Overturning Circulation (AMOC,^{3,4}). Recent changes in the northwest Atlantic water properties and ecosystems have been linked to the variations of the GS path and to Warm and Cold Core Rings (WCRs, CCRs) shed from the meandering current^{5–14}. Andres⁵ reported that the destabilization point of the GS, the point at which the GS transitions from a stable to a vigorously meandering current, moved westward, closer to Cape Hatteras, in the late 1990s. Additionally, a recent study by Wang et al.¹⁵ reported that the Gulf Stream has moved northward west of the New England Seamount Chain (NESC; ~ 65° W) in recent years, while moving southward east of the NESC, in agreement with previous studies^{16–18}. An observation-based ring census shows that the annual number of WCR formations underwent a regime change around 2000, nearly doubling from an average of 18 rings shed per year (1980–1999) to 33 per year (2000–2017)⁷, while the CCR formations did not show such a change¹³.

Many previous studies have tried to capture the GS path variability with a simple path index. These indices apply statistical techniques to GS surface characteristics, related to the location of the current maximum or the sea surface temperature gradients¹⁹, or to GS subsurface markers, e.g., a particular isotherm at a particular depth²⁰. These surface or subsurface measurements at various discrete longitudes are then combined to obtain a single path-integrated index of the GS position at monthly or annual time scales. Chi et al.²¹ summarized a number of these different GS indices. However, the complexities associated with along-stream longitudinal variability of the often multi-valued GS path behavior—such as the different periodicities of the path to the west and east of the NESC¹⁶ or the asymmetrical movement of the GS east and west of 65° W¹⁵—are obscured in such path-integrated index representations of the GS position. Furthermore, using path-integrated indices makes it hard to identify the driver(s) of GS path variability such as those indicated above. This is borne out by different studies which indicate conflicting correlations between the GS path and the North Atlantic Oscillation (NAO, an atmospheric index related to the wind and pressure fields) over different time periods and at different time

¹Department of Physics, University of Massachusetts at Dartmouth, Dartmouth, MA 02747, USA. ²School for Marine Science and Technology, University of Massachusetts at Dartmouth, Dartmouth, MA 02747, USA. ³Woods Hole Oceanographic Institution, Woods Hole, MA 02543, USA. ✉email: avijit.gangopadhyay@umassd.edu

scales^{17,19,20,22–25}, possibly due to different representations of the GS path from different single integrated indices spanning different longitudinal domains in the different studies.

Among the potential factors affecting the mean GS path and its variability are the overlying mean and time-varying wind-stress curl fields. The basin-wide wind stress curl over the subtropical North Atlantic and the associated westward propagating Rossby waves have been linked to the formation of the western boundary current (Florida Current), its northward transport along the US eastern seaboard between Florida and North Carolina, and its separation latitude near Cape Hatteras^{26–29}. The NAO has been linked with the transport and north-south movement of the Gulf Stream path after separation [Refs. ^{19,22,30} and references therein]. In these studies, the dynamical impact of the wind stress curl during different phases of the NAO was cited as the primary factor driving the observed responses of the GS path at and beyond the separation latitude.

In the North Pacific Ocean, decadal shifts of the Kuroshio Extension, the North Pacific analog to the separated Gulf Stream, are associated with a weak (strong) transport and an unstable (stable) meandering configuration³¹. These opposing phases have been linked to the basin-wide wind-stress curl, which forces negative (positive) sea surface height (SSH) anomalies that propagate westward in the form of Rossby waves during negative (positive) phases of the North Pacific Gyre Oscillation. Furthermore, it was recently shown that the strong and stable state of the Kuroshio Extension is also associated with a strong southern recirculation gyre³². Wind stress curl anomalies over the North Atlantic may similarly affect the state (stable versus unstable path) of the GS as speculated by Kelly et al.³³. The sensitivity of the GS path to cyclonic wind stress curl was investigated by Gangopadhyay and Chao³⁴ using a 1/6-degree numerical model with climatological³⁵ and operational (ECMWF)³⁶ wind forcing. It was found that the SST-based GS north wall followed the wind-stress curl maximum over a period of 4 years (1983–1986) in the model forced by the operational fields, but not in the model forced with the coarse climatology (see Gangopadhyay and Chao³⁴, their Fig. 2). Untangling the various effects responsible for the interannual variability of the GS path downstream of the separation remains an area of active research and motivates this study.

This study investigates the relationship between GS path downstream of Cape Hatteras, and the overlying wind stress curl field. There are conflicting conceptual models about whether the Gulf Stream path after separating from coast follows the zero wind stress curl (as elucidated by linear theories) or the maximum wind stress curl north of the zero curl (as established from early non-linear numerical experiments). Our study aims to address these conflicting conceptual models from a purely observational perspective from the 27-years of satellite data of GS path and atmospheric wind field reanalysis. Details of the data sets and methods used to characterize the GS path and the wind stress curl (WSC) field are given in the Supplementary Information (SI). The following analysis of twenty-seven years (1993–2019) of a wind reanalysis product and satellite SSH data demonstrates that the path of the freely meandering GS between 70° and 55° W is synchronous with the region of maximum cyclonic vorticity of the wind stress field (known as the positive vorticity pool). This study also examines the interannual variability in this positive vorticity pool and the implications of this variability on upwelling and ecosystem responses and concludes with discussion of possible future theoretical developments based on the observed correlations.

Observation: wind stress curl spatial pattern and location of the GS free jet

To investigate the relationship between the wind and the GS, the co-variability of the WSC field over the domain 30–45° N, 80–45° W from the Japanese 55-year Reanalysis (JRA-55) is compared with variability in the path of the free GS jet from 75° W to 50° W, as inferred from the altimeter SSH data. See Materials and Methods within the Supplementary Information for the determination of the wind stress curl field from JRA-55 and for identification of the GS axis path based on the method of Andres⁵. Daily fields of both WSC and GS path are averaged over various time-scales (monthly, seasonal, annual, and interannual) and these averaged fields are then compared (e.g., Fig. 1). The synoptic variability of the wind stress curl of individual month and season for each year with the relevant GS path during 1993–2019 superimposed is presented in Figs. S1, S2 and S3, and movies S1 and S2; and the wind stress curl data are provided in Gifford et al.³⁷.

As evident in Fig. 1, the WSC field in the northwestern North Atlantic generally has two extrema associated with the mid-latitude westerlies— one negative to the south of a zero-curl line and one positive to the north of the zero-curl line. For the year shown, 1997, the location of the axis of the GS downstream of Cape Hatteras lies within the region of positive maximum wind stress curl at all time-scales examined— monthly (Fig. 1A), seasonal (Fig. 1B), and annual (Fig. 1C). This synchronicity between the GS and the WSC maximum region is also evident in the 27-year average fields (Fig. 1D). Also, Fig. S1 confirms this synchronous pattern between the WSC maximum region and the annually-averaged GS paths for each of the 27 years (1993–2019) examined here. Furthermore, Figs. S2 and S3 provide the average monthly and seasonal patterns of the curl and the corresponding GS paths, supporting their synchronicity at these timescales.

This qualitative relationship observed between the positive WSC maximum region (i.e., the ‘positive vorticity pool’) and the path of the free GS jet provides the motivation to determine a way to quantify changes in the strength and extent of the region of maximum wind stress curl within which vorticity is strongly positive (cyclonic). Identification of the ‘positive vorticity pool’ allows for characterization of interannual variability in the curl maximum amplitudes and patterns and provides a means to quantify temporal changes including the area it occupies, as well as the relationship between this positive vorticity region and the free GS jet.

As a first step to quantify interannual variability in the positive vorticity pool, the distribution of the daily WSC within the domain over the entire 27 year period is examined to evaluate the relative frequency of occurrence of different curl values within the region. The methodology to generate this wind stress curl distribution is presented in Materials and Methods within the SI and the distribution is shown in Fig. 2. Based on this distribution, a positive curl ‘threshold’ is chosen such that it exists in every year (1993–2019) within the region of

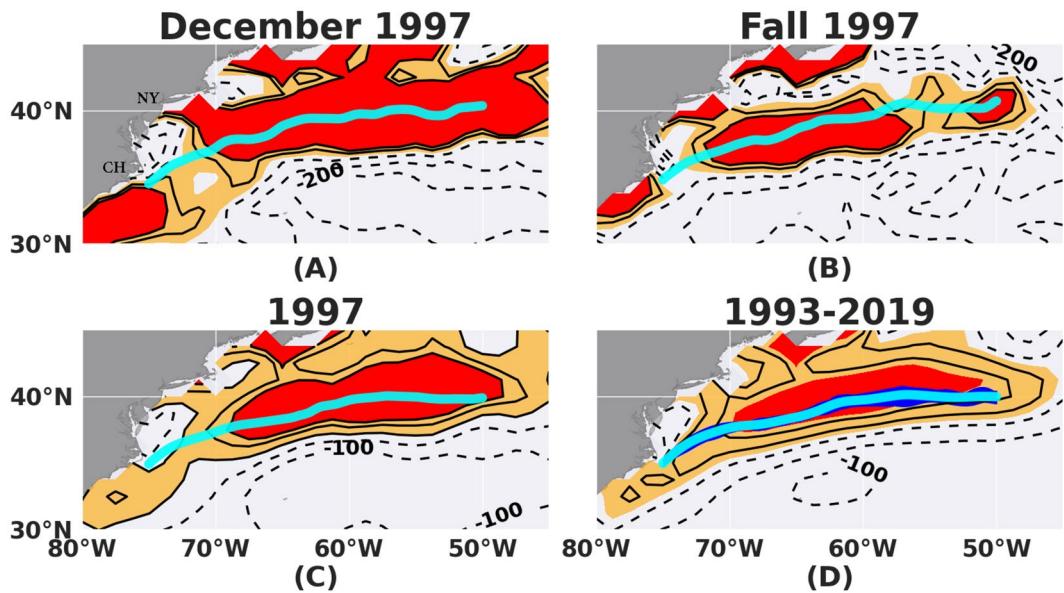


Figure 1. Multi-scale synchronicity of the region of maximum wind stress curl and the GS path. (A) Monthly example (Dec, 1997), (B) Seasonal example (Fall: OND, 1997), (C) Annual example (1997), and (D) 27-year (1993–2019) averaged WSC (contours) superimposed with corresponding averaged GS path (cyan line) from altimetry over the domain of 80° W to 45° W and 30° N to 45° N. Curl amplitude values are shown in contours of $\text{Pa}/\text{m} \times 10^{-9}$. The monthly, seasonal and annual fields are shown for the year 1997 as an example. The east coast of USA is shown in gray with Cape Hatteras (CH) near 75° W, 35° N and the city of New York (NY) near 74° W, 41° N marked for reference in panel (A). See also Fig. S1 and movie S1 in Supplementary Information (SI) for all individual years from 1993 through 2019. All monthly equivalents are shown in movie S2 in the SI. Shaded (orange and red) regions in all panels are regions of positive vorticity with the ‘positive vorticity pool’ (see text) shaded red ($\geq 80 \text{ Pa}/\text{m} \times 10^{-9}$). Yellow is $< 80 \text{ Pa}/\text{m} \times 10^{-9}$. Dashed contours denote negative curl. The blue ‘envelope’ in (D) indicates the latitudinal spread of the 1993–2019 annual GS path means calculated at each 0.1° longitude bin.

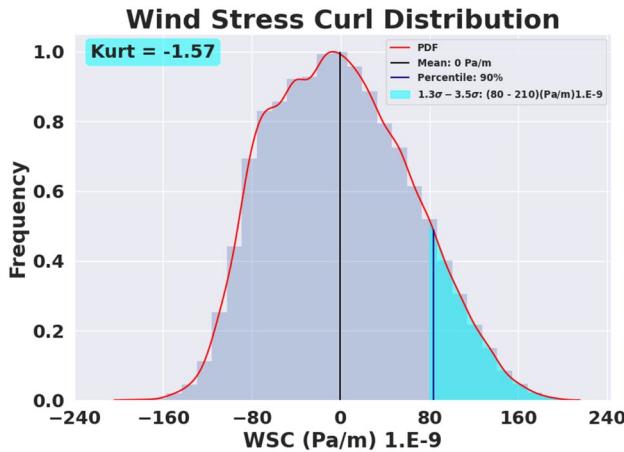


Figure 2. Distribution of wind stress curl over the North Atlantic (30–45° N, 80–45° W). This distribution is based on 27 years of wind stress data. See Methods for details. The horizontal axis extends from the most negative WSC value to the most positive WSC value. The cyan region is defined as the region of maximum wind stress curl or the positive vorticity pool. The mean is the black line located at 0 Pa/m and the blue line demonstrates the lower bound of the observed maximum as the 90 percentile.

interest and, by definition, it bounds each year’s annually-averaged positive (cyclonic) vorticity pool (i.e., the red shaded region in Fig. 1 and in Figs. S1 and S2).

Based on the curl distribution in Fig. 2, a threshold is chosen at $80 \times 10^{-9} \text{ Pa}/\text{m}$ (which is the value at 1.3 times the standard deviation, σ , of the distribution); this threshold defines a wind stress curl value above which the maximum region (i.e., the positive vorticity pool) lies for all the individual years. The region of the WSC

maximum is then defined by the area occupied by values that fall within the observed range bounded by 1.3σ to 3.5σ above the mean of the field (cyan shaded region in Fig. 2). The 1.3σ limit is chosen as it is very close to the 90th percentile of the observed curl range indicating the robustness of this choice. An excess kurtosis of -1.57 indicates the presence of outliers in the distribution. This area of the positive vorticity pool (i.e., within which the WSC exceeds the threshold value) provides a metric to determine the relationship between WSC maximum and the path of the free GS jet.

Results: variability of the wind stress curl parameters

As discussed above, the Gulf Stream path is a multi-valued curve in space and the wind-stress curl is a two-dimensional field, which for our purposes is best represented by the spatial extent (surface area) and the area-averaged value (amplitude) of the curl within a fixed-value contour such that the amplitude of curl within that area varies in different years. Thus, we choose to create the MCL (maximum curl line) and ZCL (zero curl line) metrices along with a smoother version of the Gulf Stream path between 75 and 50W, and then cast the problem with a proximity analysis as measured by 'absolute deviation' between these curves in space and time as presented throughout the analysis (Figs. 3 and 4 below).

The maximum curl line (MCL) and the zero curl line (ZCL)

Two contour lines are chosen to delineate and characterize the mean wind field – one marks the zero of the wind stress curl and the other outlines the area encompassed by the threshold contour (80×10^{-9} Pa/m) of curl, or the region of the positive vorticity pool. These contour lines are obtained from the annual curl fields mapped at 0.1° intervals (see "Materials and methods" within SI). From this we establish two separate metrics for determining the statistical significance of co-variability in the observed patterns of GS path and winds: the annual mean deviation (absolute distance) between the GS path and the zero wind stress curl line (ZCL) and the annual mean deviation of the GS path to the maximum wind stress curl line (MCL) within the region of the positive vorticity pool. Both of these are depicted in Fig. 3A,B. The ZCL is used to compare and contrast the colocation of the GS path with the MCL. The ZCL is known to lie south of the GS path after separation (see Talley³⁸—their Figs. S9.1 and S9.3; see also Seidov et al.³⁹—their Figs. 2 and 3).

The slow interannual variation in the spatial synchronicity of the GS path and the region of positive vorticity pool (quantified by the location of the MCL) is presented in Fig. 3A and is compared here with the deviation of the stream path from the ZCL. On average (zonally-averaged over the domain), for the 27-year record, the annually-averaged ZCL is $240 \text{ km} \pm 60 \text{ km}$ south of the GS jet axis (over the longitudinal domain considered, 70–55° W), while the MCL is $70 \text{ km} \pm 30 \text{ km}$ north of the GS axis (Fig. 3A). Note that the typical width of the GS is 50 km to the north and 100 km to the south of the jet axis where the maximum near surface geostrophic

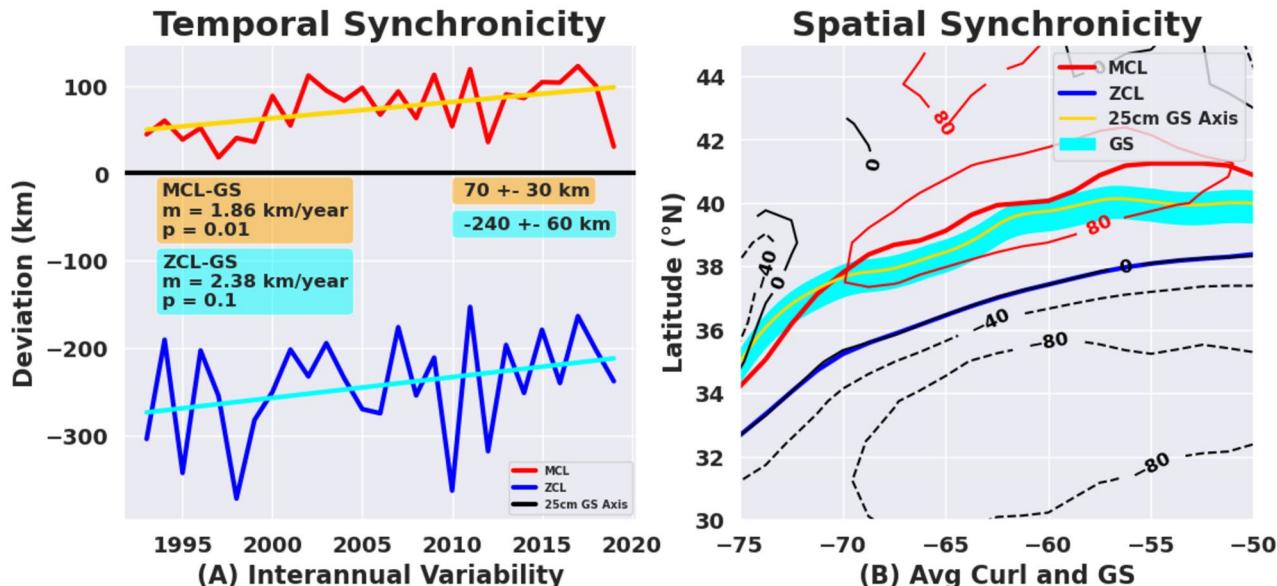


Figure 3. Temporal and spatial synchronicity of the GS path with the wind stress curl. (A) Temporal synchronicity with slow interannual variation over the 27-year study period. The average southward (northward) deviation from the GS axis to the zero curl line (ZCL, blue) and maximum curl line (MCL, red). The GS axis is situated at the $y=0$ line, with the MCL to the north (positive) and ZCL to the south (negative) of this. (B) Spatial synchronicity of the maximum curl region (bordered by the red thin line) with the Gulf Stream (cyan shaded path with the GS axis in yellow) over the 27-year period. The maximum curl line is the solid red line over the positive vorticity pool. The zero curl line is the blue line to the south, which separates the negative wind stress curl region from the positive vorticity to the north where wind-driven upwelling is expected. Similar maps for the annual and monthly averages are shown in the Movies S1 and S2, which highlight the multiscale nature of this region-wise synchronicity.

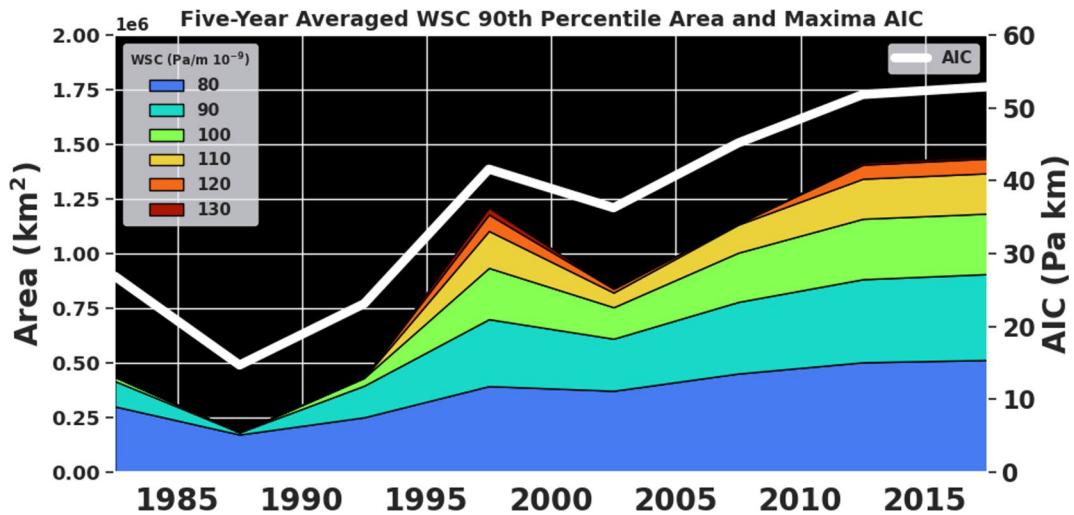


Figure 4. Temporal variability of the area of the positive vorticity pool region and the area integrated curl over this region. Five-year averaged area of WSC maximum broken down by its major contributory ranges of curl (see Gifford⁴¹ for contributions from specific maximum ranges). Units are in km^2 for the area wedges. Note how the high magnitude curl areas have been increasing over the recent pentads. Five-year averaged area integrated curl, (see Gifford⁴¹ for values) are depicted in the solid white line, whose units are in Pa km and shown along the right axis.

velocity is found⁴⁰ (see also Chi et al.²¹ -- their Fig. 1). This indicates that the maximum curl region is generally directly over the $\sim 150\text{-km}$ wide GS current between 70°W and 55°W .

The wind field and its position relative to the GS have undergone low-frequency changes over the 27-year period examined here. The MCL has shifted northward relative to the GS path by about 50 km, and the ZCL has also shifted northward (closer to the GS) by about 64 km (Fig. 3A, trend lines). While the northward trend of MCL shift is significant at 95% level, that of ZCL shift is significant only at 90% level. The variability of the ZCL shift (± 60 km) is almost twice than the variability of the MCL (± 30 km), indicating the robustness of the proximity of the GS to the MCL compared to being close to ZCL. A recent study by Seidov et al.³⁹ shows the positions of the zero lines and their relative southward separations from the GS North Wall locations of almost 200 km, which agrees with our observations.

To examine spatial patterns in the GS path and the wind field, and how they covary, the longitudinal variations of distances of the MCL and ZCL from the GS are examined at different time-scales (monthly, annually and over the whole 27-year period). The spatial synchronicity between the GS and the MCL region for the whole period is presented in Fig. 3B. The yearly and monthly variations of this spatial synchronicity are presented in the Movies S1 and S2 (SI). The GS passes through the region of maximum wind stress curl in all of the 27 years. Interestingly, the distance between the MCL and the ZCL is minimum near Cape Hatteras where the GS first separates from the western boundary and then increases eastward until about 50°W , after which the distance decreases again. This was also observed by Seidov et al.³⁹ (see their Figs. 3 and 4). The amplitude of the maximum of the wind stress curl also increases eastward, reaching a maximum around 65°W and remaining high until about 58°W (Fig. 3b). East of this region, the amplitude decreases and the contours close to form the eastern boundary of the positive vorticity pool at around 50°W . For details, see Gifford⁴¹.

Changes of the wind curl maximum region (positive vorticity pool) over 1980–2019

Next, we characterize changes in the area of the wind stress curl maximum or the ‘positive vorticity pool’ region and changes in the area-integrated curl (AIC) within this region over the study period (Fig. 4). Overall, the area and intensity of the maximum has increased over the GS region in the recent decades.

Therefore, larger portions of ocean surface area have been exposed to the WSC maxima (i.e., the positive vorticity pools) in recent years. Furthermore, contributions from larger WSC values ($\geq 120 \times 10^{-9} \text{ Pa/m}$) become prevalent in the 1995–1999 pentad and persist beyond this pentad, indicating possible changes in the forcing and upwelling within the free jet. The pentad spanning 2015–2019 shows the largest AIC value ($\sim 50 \text{ Pa km}$) over the temporal domain considered here.

Aside from changes in the spatial extent (area) of the WSC maximum, the interannual variability of the center of mass of the curl maximum (defined by the region bounded by $80 \times 10^{-9} \text{ Pa/m}$ contour) is also changing. Movement of the center of mass (COM) quantifies local changes in the position of the positive vorticity pool that may also result in changes in the region of strong upwelling. Upwelling-favorable wind stress curl has been documented to drive the upward transport of nutrient-rich deeper waters into the sunlit upper ocean, supporting elevated primary productivity near the coast^{43–45}.

Figure 5A demonstrates a coupling (see the strong linear fit with $r=0.7$) between the COM longitude and COM latitude, in that the COM is moving along a line that is not purely zonal or meridional. Furthermore, this non-unity correlation or the deviations from this line might be associated with the change of shape, size

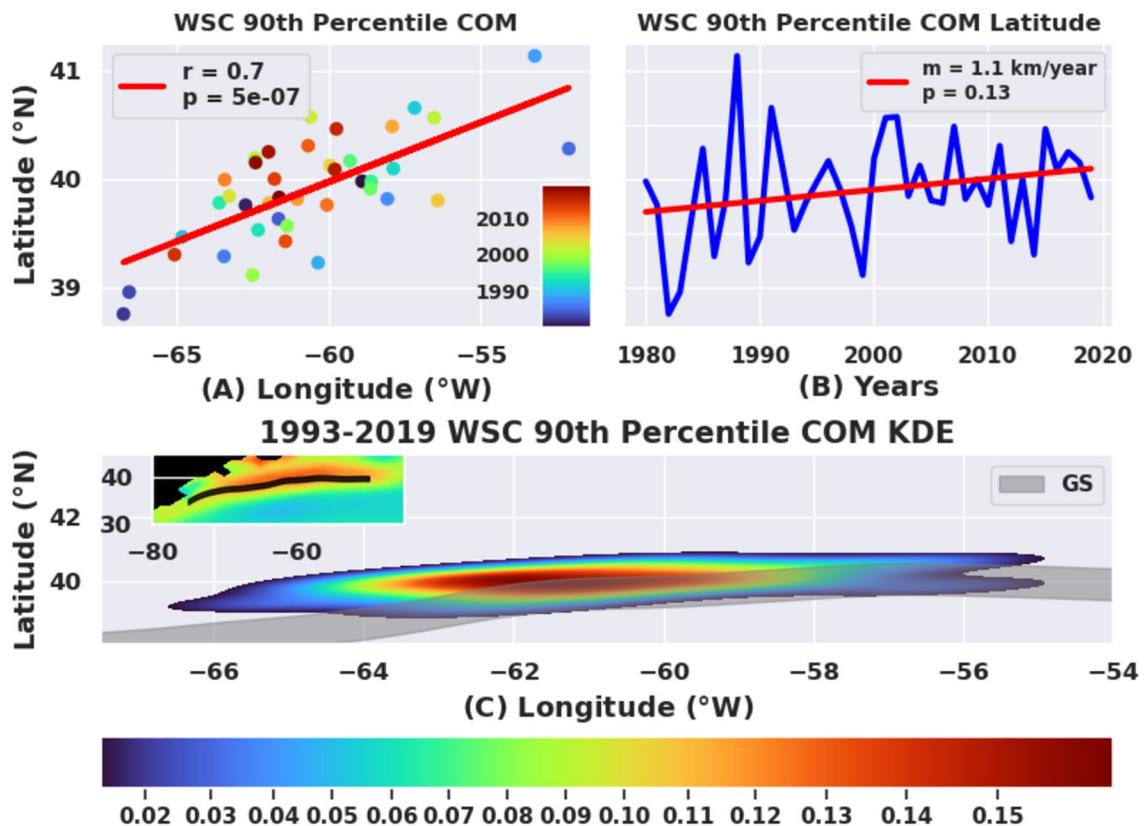


Figure 5. Spatial and temporal behavior of the center of mass (COM) of the positive vorticity pool. (A) The positive vorticity pool's center of mass (COM) longitude against latitude over forty years (1980–2019). (B) Annually averaged COM latitude from 1980 to 2019 (blue) with line of best fit (red) showing possible northward movement of the COM during this period ($p=0.13$). (C) A kernel density estimation (KDE) map of the distribution of COM location of a region bounded by the 90th percentile of the wind stress curl. The COM has been converging to a small region around 61° – 62° W primarily due to recent years which are shown in red/orange dots in panel (A).

and amplitude of the positive vorticity pool which is part of the larger atmospheric circulation system over the North Atlantic. In fact, from 1980 to 2019, the COM of the WSC maximum lower bound has shown a possible shift northward (Fig. 5B, trend line, $p=0.13$) by about 40 km and westward by about 40 km (not shown here, see Gifford^{41,42}). This indicative northwestward movement of the curl maximum is similar to that observed in the decadal analysis of the GS path movement in recent decades³⁹, providing further evidence of synchronicity between the two. Figure 5C shows that the COM has become more localized to a focused region, leading to the higher density towards the center of its kernel density estimation (KDE). KDE is a representation of the probability density estimator, when kernel smoothing is applied based on weights and used for finite data samples^{46,47}. The red and green shaded contours in Fig. 5C represent the COMs in the recent decade, indicating agreement with Seidov et al.³⁹.

Discussion

Reconciling theory with observations

The overall Gulf Stream system reaches from Florida to the Grand Banks and beyond, and different sub-regions within this system exhibit distinct controlling dynamics. These connected sub-regions comprise (1) the attached western boundary current (part of which is the Florida Current), (2) the separation region near Cape Hatteras and (3) the free jet after separation. Our results strongly indicate that variability in the path of the separated GS (i.e. the free jet) coincides with the positive maximum region of the wind stress curl at time scales from monthly, to annual, to decadal.

Western boundary intensification as observed in subtropical oceanic gyres was first explained with a simple analytical model by Stommel⁴⁸ for a homogenous and flat-bottom ocean. The essential model elements included a sinusoidal zonal wind stress (neglecting the much smaller meridional component), latitudinal variation of the Coriolis parameter (the beta term) and a bottom drag proportional to velocity. This formulation led to an elegant “single-gyre” stream function solution where the streamlines crowded along the western boundary in the presence of a non-zero Coriolis parameter (beta) to form a narrow, swift poleward-flowing western boundary current (see Vallis⁴⁹, their Fig. 14.4). By construct, the sinusoidal ($\sin \pi y/b$, where b is the domain width in the y -direction) zonal wind stress led to another sinusoidal function ($\cos \pi y/b$) to describe the associated wind stress

curl such that the zero wind stress curl lines occurred along the northern and southern boundaries of the gyre and a negative maximum curl was found along the center of the domain, where the poleward western boundary current transport was maximum.

The existence of the positive vorticity pool and the positive maximum curl within the subtropical gyres (in both the North Atlantic and North Pacific) was recognized around the same time^{50,51} from observed wind fields. Early studies of the wind-driven circulation^{49,52,53} considered linear dynamics and did not examine the dynamics within a separated boundary current like the Gulf Stream east of Cape Hatteras. Non-linear effects alter the simple picture, and the early studies were followed with “double-gyre” formulations^{54,55} and multiple numerical solutions^{52,53,56}. Later, Verron and Le Provost⁵⁷ maintained that the separation of the boundary current is largely independent of the location of the zero of the wind stress curl. A recent study covering multiple decades of observation (Figs. 3 and 4 of Seidov et al.³⁹) shows the zero of the wind stress curl follows the GS path in the “extension region” (i.e., downstream of Cape Hatteras) only around 50–40 W (i.e., east of Grand Banks). In fact, the separation latitude has been shown to depend on the balance of large-scale Ekman pumping and the boundary current’s transport while following the coast by Gangopadhyay et al.²⁷ based on earlier works by Parsons⁵⁸ and Veronis⁵⁹. This separation mechanism hypothesis was recently validated using forty years of observations⁶⁰ and was used to develop a GS separation latitude forecasting model. In contrast to these studies of the “attached” GS, our study here focuses on the path of the GS after separation—in the 75°–55° W region, east of the separation latitude, where the dynamics of the path might be better controlled by the maximum positive vorticity pool of the wind stress as previously envisaged by several investigators^{52–57}. In this study, we do not address the region further to the east of 50° W, where the mean GS path, the maximum of the WSC and the zero of the WSC, all collapse into a confined area (Figs. 3 and 4 of Seidov et al.³⁹).

The amplitude of the meandering of the GS path after separation is also linked with its separation near Cape Hatteras (75° W, 35° N). The separation of the GS from the continental slope near Cape Hatteras is governed by multiple factors that include inertial control⁶¹, time-integrated basin-wide wind stress^{26–28,58,59} and bathymetric control^{62,63}. Many studies have proposed that the path of the GS is also influenced by the southward flow of cold, fresh waters of the Labrador Current⁶⁴, dictated by the strength and location of one of the NAO’s “centers-of-action,” namely, the Icelandic low-pressure center^{65,66}. A recent study by Silver et al.⁶⁰ created a regression prediction method to forecast the GS North Wall position with 1-year lead using the Icelandic low center pressure and longitude paired with the Southern Oscillation Index (SOI). The separation of the GS from the continental margin has been discussed by Pickart and Smethie⁶⁷ in the context of the Deep Western Boundary Current, which the GS crosses near Cape Hatteras. This might play a second order role in determining the position of the GS, especially locally just downstream of the GS separation point, and may add some of the “noise” to the robust relationship found here between wind stress curl maximum and GS position—especially at interannual scales.

The variability of the path and transport of heat and mass by the GS is also linked to the variability of the AMOC since the GS carries both wind-driven and density-driven (thermohaline) components (e.g., Bryden⁶⁸) that are associated with horizontal and vertical recirculation cells, respectively. The AMOC transports warm limb waters northward within the upper ocean and colder, denser waters southward at depth. Since the GS carries all of the AMOC warm limb flow⁶⁹, understanding GS path variability as a component of AMOC, might lead to a better understanding and prediction of the variability of the overall AMOC system^{70,71}. As shown by Silver et al.⁶⁰, the Parsons–Veronis two-layer model whereby Ekman wind drift affects the GS separation latitude^{27,57,58} has worked well (i.e., has good prediction skill) for four decades on the backdrop of an actively varying AMOC. Given that most of the AMOC variability is in fact dominated by this Ekman Drift^{70–74}, it is possible that one could develop an AMOC predictability scheme using climate projections of the wind field with a focus on changes in the region of the maximum wind stress curl.

Biological implications

The interplay between the region of positive wind stress curl with the GS and its meanders and WCRs has implications for biological changes near the GS in the domain of interest. This also suggests significant implications for the New England fisheries sector as well as the region’s economy and ecological health. This relationship between GS path and the positive vorticity pool demonstrated above hints at a relationship between the free GS jet and wind induced gyres and secondary circulations. This may inform understanding of the biological productivity related to upwelling regions (found beneath positive wind stress curl) versus downwelling regions (found beneath negative wind stress curl) and the ecology of the free jet region. A positive vorticity region of the wind stress curl would typically be a favorable region for positive vertical velocities^{75,76}. Thus, according to the observational evidence presented here, the GS path and its latitudinal width of 150 km should be subjected to continuous upward motion due to the overlying positive wind stress curl. Multidecadal evolution of the wind field and shifts in the regions of wind stress curl maximum in the domain of interest (e.g., Fig. 5) suggest changes in Ekman pumping that can drive upwelling within and north of the GS. This upwelling has the potential to introduce more nutrient-rich water to the high light levels in the near-surface layer, and support biological productivity.

Using a climatology of the wind field, a first order estimate of upwelling within the maximum curl region (i.e., in the positive vorticity pool) ranges from 5 cm day⁻¹ (0.05 m day⁻¹) in July to about 17 cm day⁻¹ (0.17 m day⁻¹) in January (Fig. 6A). The springtime upwelling with a peak of about 10 cm day⁻¹ is also shown in Fig. 6B. This matches reasonably well with previous estimates of 5–20 cm day⁻¹ seasonal fluctuations from a QuickSCAT wind analysis by Risen and Chelton⁷⁷. These vertical velocities are relatively low, but they may be important biologically as the upwelling is continuous and occurs over a large area (400,000 km²) spanned by the positive maximum of the wind stress curl. If we assume a ~ 10 μM nitrate endmember at the nutricline at 50 m depth in the Slope Sea⁷⁸, daily WSC-driven nitrate flux into the euphotic zone over the course of a season can be estimated (see methods for details). Assuming an area integrated upwelling flux of about 0.35 Sv (see Fig. 6C for area

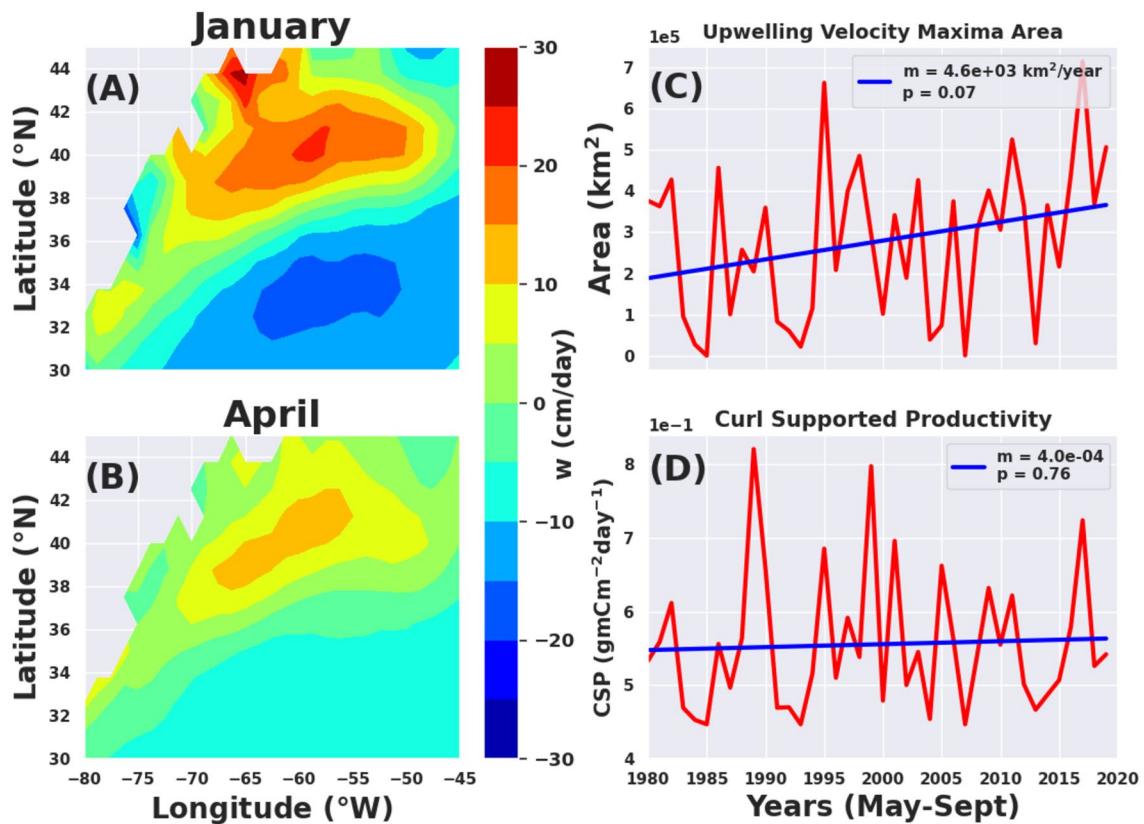


Figure 6. Validation of upwelling velocities and potential impact on primary productivity. (A) for January (maximum upwelling); (B) for April (springtime upwelling). The fields reasonably agree with the results of Risen and Chelton⁷⁷. (C) Interannual variability of the Area covered by the maximum upwelling velocity region defined by the area bounded by the threshold of 6 cm day^{-1} , which is very similar to the positive vorticity pool area. The area expansion is really high in 1995 during the nutrient limited period (May–September). (D) The inter-annual variability of potential wind-stress curl supported productivity (CSP in $\text{gm C m}^{-2} \text{ day}^{-1}$) averaged for the biologically active season of May–September over the positive vorticity pool. There is no significant trend in the CSP's interannual variation.

and Gifford⁴¹ for integrated upwelling flux) over the maximum curl region, this wind stress curl-driven upwelling would potentially provide an additional flux of $3.02 \times 10^8 \text{ mol N/day}$. This has the potential to support more than $24,000 \text{ tonnes C day}^{-1}$ of primary production over the GS and Slope Sea region. When considered over the $400,000 \text{ km}^2$ maximum wind stress curl region, this nitrate flux translates to about $60 \text{ mg C m}^{-2} \text{ day}^{-1}$ of primary productivity. The interannual variability in WSC-driven productivity over the period 1993–2019 is presented in Fig. 6D. It is curious to note that there is large interannual variability without any significant long-term trend in the curl supported productivity (Fig. 6D), despite the observed significant long-term trends in the positive vorticity pool area and the GS path discussed above. This lack of significant long-term trend in curl-supported productivity, despite the fact that there appears to be a clear and more statistically significant trend in upwelling velocity max area as shown in Fig. 6C, might be due to other factors including unresolved strength of upwelling from inadequate mixed-layer depth data, unknown vertical structure of upwelling, unknown distribution of rings and upwelling filaments in the slope sea underneath the positive vorticity pool.

This estimate is substantial in context of multiple metrics of productivity in the region. Estimates of primary productivity (PP) in the Mid Atlantic Bight from Ma and Smith⁷⁹ range from $200 \text{ mg C m}^{-2} \text{ day}^{-1}$ to about $1000 \text{ mg C m}^{-2} \text{ d}^{-1}$ depending on the location and timing (see their Figs. 2 and 3 and Tables 2–4 of Ma and Smith⁷⁹) and the type of PP models (vertically resolved, estimated from ^{14}C uptake, or estimate from surface variables). Thus, our estimate of as much as $60 \text{ mg C m}^{-2} \text{ day}^{-1}$ of primary productivity supported by wind-induced upwelling would be $\sim 10\text{--}25\%$ of the current estimates of offshore primary productivity in the continental slope region of the western North Atlantic. The wind-induced upwelling likely supports a larger proportion of productivity offshore than on the shelf, as the primary productivity in the slope and Gulf Stream region is generally lower.

Conclusions

As demonstrated here, the GS is observed to lie within the region of the wind stress curl maximum or the positive vorticity pool to the north of the zero wind stress curl. For monthly averaged values, the maximum positive vorticity line is found $70 \pm 30 \text{ km}$ north of the axis of the GS (where the standard deviation represents the along-GS variability between 70° W and 55° W), while the line of zero wind stress curl is south of the GS by about

240 ± 60 km. The whole width of the Gulf Stream (~ 150 km) falls within the region of maximum wind stress curl, i.e., the positive vorticity pool (Figs. 1, 3B,C; Figs. S1, S2, S3; and Movies S1, S2) at multiple time-scales over a 27-year period (monthly, annually and over the 27-year average). Annually, the GS is always within the region of maximum curl. Evolution of the WSC maximum shows an increase in the area and in the intensity of the maximum (increasing from 80×10^{-9} to 140×10^{-9} Pa/m); specifically a significant shift in area, AIC, and the introduction of stronger WSC values of $\geq 120 \times 10^{-9}$ Pa/m in the pentad of 1995–1999 is observed. This has important implications on the region's fundamental physics and on the chemistry and biology of North Atlantic ecosystems. The resulting wind-induced upwelling could support as much as $60 \text{ mg C m}^{-2} \text{ day}^{-1}$ of primary productivity, explaining about 10–25% of the estimated mean primary productivity in the offshore western North Atlantic.

The observational evidence presented herein, suggests that the meandering GS (or the “free jet”) after separation adjusts to the nearby positive maximum of the wind stress curl. We note that this 27-year-long observational phenomenon is robust and is inconsistent with the paradigm that the Gulf Stream path after separation follows the isopleth of zero in the wind stress curl field where the wind stress's vorticity is zero. Clearly, the dynamical framework underpinning this observed alignment remains to be explored in the future with models that might (or might not) build on the early, linear theories of (attached) western boundary currents. Note that it is possible that the GS path is not just passively responding to the overlying wind field. The GS is known to actively influence the wind field above it via impacts on surface winds, atmospheric convergences and divergences, rain bands, air-sea fluxes, etc.^{80–82}. The implication of the synchronicity found here between the wind stress curl maximum and the position of the warm GS path over the span of 1500–2500 km opens up a range of questions related to ocean-atmosphere coupling in the western North Atlantic across temporal scales spanning storms, to weather, and to climate. This observational evidence suggests the possibility of development of new ways of understanding the behavior of the GS path as a free jet after Cape Hatteras, especially in the 75° – 55° W region, which is important for studying future AMOC scenarios by more clearly delineating how Gulf Stream variability relates to changes in AMOC.

Data availability

The wind velocity data is freely available from the Japan Meteorological Agency's 55-year reanalysis (JRA-55) distribution site at <https://jra.kishou.go.jp/JRA-55/atlas/en/index.html>. The wind stress and wind stress curl data computed for this study are available in the zenodo repository at <https://zenodo.org/record/8200832>. The Gulf Stream paths were extracted from altimeter gridded product available at <https://www.copernicus.eu/en>. The monthly and annual Gulf Stream paths, and contours of zero and maximum wind stress curl are available at <https://zenodo.org/records/8217388>. The Curl distribution, upwelling fields and other datasets generated or analyzed during the study are available from the corresponding author on reasonable request. All the data are available in the links provided in the Supplementary Information and in the Data availability statement. All of the data generated or analyzed during the current study are available from the corresponding author on reasonable request.

Received: 12 March 2024; Accepted: 24 July 2024

Published online: 09 August 2024

References

1. Nye, J. A., Joyce, T. M., Kwon, Y. O. & Link, J. S. Silver hake tracks changes in Northwest Atlantic circulation. *Nat. Commun.* **2**, 1 (2011).
2. Joyce, T. M., Kwon, Y. O. & Yu, L. On the relationship between synoptic wintertime atmospheric variability and path shifts in the Gulf Stream and the Kuroshio Extension. *J. Clim.* **22**, 3177–3192 (2009).
3. Caeser, L., Rahmstorf, S., Robinson, A., Feulner, G. & Saba, V. Observed fingerprint of a weakening Atlantic Ocean overturning circulation. *Nature* **556**, 191–196 (2018).
4. Zhang, R. *et al.* A review of the role of the Atlantic meridional overturning circulation in Atlantic multidecadal variability and associated climate impacts. *Rev. Geophys.* **57**(2), 316–375 (2019).
5. Andres, M. On the recent destabilization of the Gulf Stream path downstream of Cape Hatteras. *Geophys. Res. Lett.* **43**, 9836–9842 (2016).
6. Brickman, D., Hebert, D. & Wang, Z. Mechanism for the recent ocean warming events on the Scotian Shelf of eastern Canada. *Continent. Shelf Res.* **156**, 11–22 (2018).
7. Gangopadhyay, A., Gawarkiewicz, G., Silva, E. N. S., Monim, M. & Clark, J. An observed regime shift in the formation of warm core rings from the Gulf Stream. *Sci. Rep.* **9**, 1–9 (2019).
8. Gawarkiewicz, G., Todd, R. E., Plueddemann, A. J., Andres, M. & Manning, J. P. Direct interaction between the Gulf Stream and the shelfbreak south of New England. *Sci. Rep.* **2**, 553 (2012).
9. Gawarkiewicz, G. *et al.* The changing nature of shelf-break exchange revealed by the OOI Pioneer Array. *Oceanography* **31**, 60–70 (2018).
10. Gawarkiewicz, G. *et al.* Characteristics of an advective marine heatwave in the Middle Atlantic Bight in early 2017. *Front. Mar. Sci.* **6**, 6 (2019).
11. Mills, K. E., Pershing, A. J., Sheehan, T. F. & Mountain, D. Climate and ecosystem linkages explain widespread declines in North American Atlantic Salmon Populations. *Global Change Biol.* **19**, 3046–3061 (2013).
12. Pershing, A. J. *et al.* Slow adaptation in the face of rapid warming leads to collapse of the Gulf of Maine cod fishery. *Science* **350**, 809–812 (2015).
13. Silver, A., Gangopadhyay, G., Gawarkiewicz, G., Silva, E. N. S. & Clark, J. Interannual and seasonal asymmetries in Gulf Stream Ring Formations from 1980 to 2019. *Sci. Rep.* **11**, 1 (2021).
14. Silver, A., Gangopadhyay, A., Gawarkiewicz, G., Fratantonio, P. & Clark, J. Increased gulf stream warm core ring formations contributes to an observed increase in salinity maximum intrusions on the Northeast Shelf. *Sci. Rep.* **13**(1), 1–9 (2023).
15. Wang, Z., Yang, J., Johnson, C. & DeTracey, B. Changes in Deep Ocean Contribute to a “See-Sawing” Gulf Stream Path. *Geophys. Res. Lett.* **49**, 21 (2022).

16. Gangopadhyay, A., Chaudhuri, A. H. & Taylor, A. H. On the nature of temporal variability of the Gulf Stream path from 75 to 55 W. *Earth Interact.* **20**, 1–17 (2016).
17. Bisagni, J. J., Gangopadhyay, A. & Sanchez-Franks, A. Secular change and inter-annual variability of the Gulf Stream position, 1993–2013, 70° – 55° W. *Deep Sea Res. Part I Oceanogr. Res. Papers* **125**, 1–10 (2017).
18. Dong, S. & Baringer, M. O. Slow down of the Gulf Stream during 1993–2016. *Sci. Rep.* **9**, 1 (2019).
19. Taylor, A. H. & Stephens, J. A. The North Atlantic oscillation and the latitude of the Gulf Stream. *Tellus A: Dyn. Meteorol. Oceanogr.* **50**, 134–142 (1998).
20. Joyce, T. M., Deser, C. & Spall, M. A. The relation between decadal variability of subtropical mode water and the North Atlantic Oscillation. *J. Clim.* **13**, 2550–2569 (2000).
21. Chi, L., Wolfe, C. L. & Hameed, S. The distinction between the Gulf Stream and its North Wall. *Geophys. Res. Lett.* **46**(15), 8943–8951 (2019).
22. Chaudhuri, A. H., Gangopadhyay, A. & Bisagni, J. J. Response of the Gulf Stream transport to characteristic high and low phases of the North Atlantic Oscillation. *Ocean Model.* **39**, 220–232 (2011).
23. Rossby, T. & Benway, R. L. Slow variations in mean path of the Gulf Stream east of Cape Hatteras. *Geophys. Res. Lett.* **27**(1), 117–120 (2000).
24. Ezer, T. & Atkinson, L. P. Accelerated flooding along the US East Coast: On the impact of sea-level rise, tides, storms, the Gulf Stream, and the North Atlantic Oscillations. *Earth's Future* **2**(8), 362–382 (2014).
25. Marshall, J. *et al.* North Atlantic climate variability: Phenomena, impacts and mechanisms. *Int. J. Climatol. J. R. Meteorol. Soc.* **21**(15), 1863–1898 (2001).
26. Gill, A. E. *Atmosphere–ocean dynamics* (Academic press, 1982).
27. Gangopadhyay, A., Cornillon, P. & Watts, D. R. A test of the Parsons–Veronis hypothesis on the separation of the Gulf Stream. *J. Phys. Oceanogr.* **22**, 1286–1301 (1992).
28. Dengg, J. The gulf stream separation problem. *The Warmwatersphere of the North Atlantic Ocean* **1996**, 254–290 (1996).
29. Haidvogel, D. B. & Beckmann, A. *Numerical ocean circulation modeling*. (Imperial College Press, 1999).
30. Taylor, A. H. & Gangopadhyay, A. A simple model of interannual displacements of the Gulf Stream. *J. Geophys. Res. Oceans* **106**(C7), 13849–13860 (2001).
31. Qiu, A. & Chen, S. Eddy-mean flow interaction in the decadally modulating Kuroshio Extension system. *Deep Sea Res. Part II Top. Stud. Oceanogr.* **57**, 1098–1110 (2010).
32. Yang, B. *et al.* The onshore intrusion of Kuroshio subsurface water from February to July and a mechanism for the intrusion variation. *Prog. Oceanogr.* **167**, 97–115 (2018).
33. Kelly, K. A. *et al.* Western boundary currents and frontal air-sea interaction: Gulf Stream and Kuroshio Extension. *J. Clim.* **23**, 5644–5667 (2010).
34. Gangopadhyay, A. & Chao, Y. Sensitivity of the gulf stream path to the cyclonic wind stress curl. *Global Atmos. Ocean Syst.* **7**, 151–178 (2000).
35. Hellerman, R. & Rosenstein, M. Normal monthly wind stress over the world ocean with error estimates. *J. Phys. Oceanogr.* **13**(7), 1093–1104 (1983).
36. Barnier, B., Siefridt, L. & Marchesiello, P. Thermal forcing for a global ocean circulation model using a three-year climatology of ECMWF analyses. *J. Mar. Syst.* **6**(4), 363–380 (1995).
37. Gifford, I. H., Gangopadhyay, A., Andres, M., Gawarkiewicz, G., Oliver, H., & Silver, A. Wind stress, wind stress curl, and upwelling velocities in the Northwest Atlantic (80–45°W, 30–45°N) during 1980–2019. <https://zenodo.org/record/8200832> (2023).
38. Talley, L. D. *Descriptive physical oceanography: an introduction* (Academic press, 2011).
39. Seidov, D., Mishonov, A., Reagan, J. & Parsons, R. Resilience of the Gulf Stream path on decadal and longer timescales. *Sci. Rep.* **9**, 1 (2019).
40. Rossby, T. & Zhang, H. M. The near-surface velocity and potential vorticity structure of the Gulf Stream. *J. Mar. Res.* **59**(6), 949–975 (2001).
41. Gifford, I. H. The Synchronicity of the Gulf Stream Free Jet and the Wind Induced Cyclonic Vorticity Pool. MS Thesis, University of Massachusetts Dartmouth. 63 pp. https://umassd.primo.exlibrisgroup.com/permalink/01MA_DM_INST/1q2aha6/alm9914351958901301 (2023).
42. Gifford, I. H., Silver, A., Gangopadhyay, A., Andres, M., Gawarkiewicz, G., & Oliver, H. Monthly and Annual contour lines of the zero and the positive maximum of the Wind Stress Curl over Western North Atlantic during 1980–2019 and the Gulf Stream path during 1993–2019. 10.5281/zenodo.8217388 (2023).
43. Albert, A., Echevin, V., Lévy, M. & Aumont, O. Impact of nearshore wind stress curl on coastal circulation and primary productivity in the Peru upwelling system. *J. Geophys. Res. Oceans* **115**(C12), 1. <https://doi.org/10.1029/2010JC006569> (2010).
44. Kumar, S. P. *et al.* Physical control of primary productivity on a seasonal scale in central and eastern Arabian Sea. *J. Earth Syst. Sci.* **109**(4), 433–441. <https://doi.org/10.1007/BF02708331> (2000).
45. Macías, P. J. S., Franks, M. D. & Ohman, M. R. Landry, Modeling the effects of coastal wind- and wind–stress curl-driven upwellings on plankton dynamics in the Southern California current system. *J. Mar. Syst.* **94**, 107–119 (2012).
46. Rosenblatt, M. Remarks on some nonparametric estimates of a density function. *Ann. Math. Stat.* **27**(3), 832–837. <https://doi.org/10.1214/aoms/1177728190> (1956).
47. Parzen, E. On estimation of a probability density function and mode. *Ann. Math. Stat.* **33**(3), 1065–1076. <https://doi.org/10.1214/aoms/1177704472> (1962).
48. Stommel, H. The westward intensification of wind-driven ocean currents. *Eos, Trans. Am. Geophys. Union* **29**(2), 202–206 (1948).
49. Vallis, G.K. *Atmospheric and oceanic fluid dynamics* (Cambridge University Press, 2006).
50. Sverdrup, H. U., Johnson, M. W. & Fleming, R. H. *The Oceans: Their physics, chemistry, and general biology* (Prentice-Hall, 1942).
51. Munk, W. H. On the wind-driven ocean circulation. *J. Atmos. Sci.* **7**, 80–93 (1950).
52. Veronis, G. Wind-driven ocean circulation—Part 1. Linear theory and perturbation analysis. Deep sea research and oceanographic abstracts (Vol. 13, No. 1, pp. 17–29) (Elsevier, 1966).
53. Veronis, G. Wind-driven ocean circulation—Part 2. Numerical solutions of the non-linear problem. Deep sea research and oceanographic abstracts (Vol. 13, No. 1, pp. 31–55) (Elsevier, 1966).
54. Munk, W. H. & Carrier, G. F. The wind-driven circulation in ocean basins of various shapes. *Tellus* **2**(3), 158–167 (1950).
55. Carrier, G. F. & Robinson, A. R. On the theory of the wind-driven ocean circulation. *J. Fluid Mech.* **12**, 49–80 (1962).
56. Veronis, G. On parametric values and types of representation in wind-driven ocean circulation studies. *Tellus* **17**, 77–84 (1965).
57. Verron, J. & Le Provost, C. Response of eddy-resolved general circulation numerical models to asymmetrical wind forcing. *Dyn. Atmos. Oceans* **15**(6), 505–533 (1991).
58. Parsons, A. T. A two-layer model of Gulf Stream separation. *J. Fluid Mech.* **39**, 511–528 (1969).
59. Veronis, G. Model of world ocean circulation: I Wind-driven, two-layer. *J. Mar. Res.* **31**, 228–288 (1973).
60. Silver, A., Gangopadhyay, A., Gawarkiewicz, G., Taylor, A. & Sanchez-Franks, A. Forecasting the Gulf Stream path using buoyancy and wind forcing over the North Atlantic. *J. Geophys. Res. Oceans* **126**(8), 7614 (2021).
61. Fofonoff, N. Steady flow in a frictionless homogeneous ocean (1955).
62. Zhang, R. & Vallis, G. K. The role of bottom vortex stretching on the path of the North Atlantic western boundary current and on the northern recirculation gyre. *J. Phys. Oceanogr.* **37**, 2053–2080 (2007).

63. Schoonover, J., Dewar, W. K., Wienders, N. & Deremble, B. Local sensitivities of the Gulf Stream separation. *J. Phys. Oceanogr.* **47**, 353–373 (2017).
64. Rossby, T. On gyre interactions. *Deep Sea Res. Part II Top. Stud. Oceanogr.* **46**, 139–164 (1999).
65. Hameed, S. & Piontkovski, S. The dominant influence of the Icelandic Low on the position of the Gulf Stream northwall. *Geophys. Res. Lett.* **31**, 9 (2004).
66. Sanchez-Franks, A., Hameed, S. & Wilson, R. E. The Icelandic low as a predictor of the Gulf Stream north wall position. *J. Phys. Oceanogr.* **46**, 817–826 (2016).
67. Pickart, R. S. & Smethie, W. M. How does the deep western boundary current cross the Gulf Stream?. *J. Phys. Oceanogr.* **23**(12), 2602–2616 (1993).
68. Bryden, H. L. Wind-driven and buoyancy-driven circulation in the subtropical North Atlantic Ocean. *Proc. Math. Phys. Eng. Sci.* **477**(2256), 1. <https://doi.org/10.1088/rspa.2021.0172> (2021).
69. Rossby, T., Palter, J. & Donohue, K. What can hydrography between the New England Slope, Bermuda and Africa tell us about the strength of the AMOC over the last 90 years?. *Geophys. Res. Lett.* **49**, 9173. <https://doi.org/10.1029/2022GL099173> (2022).
70. Caesar, L., McCarthy, G. D., Thornalley, D. J. R., Cahill, N. & Rahmstorf, S. Current Atlantic meridional overturning circulation weakest in last millennium. *Nat. Geosci.* **14**, 118–120 (2021).
71. Lozier, M. S. Deconstructing the conveyor belt. *Science* **328**, 1507–1511 (2010).
72. Frajka-Williams, E. *et al.* Atlantic meridional overturning circulation: Observed transport and variability. *Front. Mar. Sci.* **1**, 260 (2019).
73. Lozier, M. S. Overturning in the North Atlantic. *Annu. Rev. Mar. Sci.* **4**, 291–315 (2012).
74. Mielke, E. & Frajka-Williams, J. Baehr, observed and simulated variability of the AMOC at 26 N and 41 N. *Geophys. Res. Lett.* **40**, 1159–1164 (2013).
75. Enriquez, A. G. & Friehe, C. A. Effects of wind stress and wind stress curl variability on coastal upwelling. *J. Phys. Oceanogr.* **25**, 1651–1671 (1995).
76. Wang, Q. *et al.* Wind stress curl and coastal upwelling in the area of Monterey Bay observed during AOSN-II. *J. Phys. Oceanogr.* **41**(5), 857–877 (2011).
77. Risien, C. M. & Chelton, D. B. A global climatology of surface wind and wind stress fields from eight years of QuikSCAT scatterometer data. *J. Phys. Oceanogr.* **38**(11), 2379–2413 (2008).
78. Oliver, H. *et al.* Diatom hotspots driven by western boundary current instability. *Geophys. Res. Lett.* **48**(11), 1943 (2021).
79. Ma, J. & Smith, W. O. Primary productivity in the Mid-Atlantic Bight: Is the shelf break a location of enhanced productivity?. *Front. Mar. Sci.* **9**, 824303 (2022).
80. Minobe, S., Miyashita, M., Kuwano-Yoshida, A., Tokinaga, H. & Xie, S. P. Atmospheric response to the Gulf Stream: Seasonal variations. *J. Clim.* **23**(13), 3699–3719 (2010).
81. Minobe, S., Kuwano-Yoshida, A., Komori, N., Xie, S. P. & Small, R. J. Influence of the Gulf Stream on the troposphere. *Nature* **452**(7184), 206–209 (2008).
82. Shi, Q. & Bourassa, M. A. Coupling ocean currents and waves with wind stress over the gulf stream. *Remote Sens.* **11**(12), 1476. <https://doi.org/10.3390/rs11121476> (2019).
83. Kobayashi, S. *et al.* The JRA-55 reanalysis: General specifications and basic characteristics. *J. Meteorol. Soc. Japan. Ser. II* **93**(1), 5–48 (2015).
84. Gelaro, R. *et al.* The modern-era retrospective analysis for research and applications, version 2 (MERRA-2). *J. Clim.* **30**(14), 5419–5454 (2017).
85. Large, W. G. & Pond, S. Open ocean momentum flux measurements in moderate to strong winds. *J. Phys. Oceanogr.* **11**(3), 324–336 (1981).
86. Trenberth, K. E., Large, W. G. & Olson, J. G. The effective drag coefficient for evaluating wind stress over the oceans. *J. Clim.* **2**(12), 1507–1516 (1989).
87. Wen, C., Kumar, A. & Xue, Y. Uncertainties in reanalysis surface wind stress and their relationship with observing systems. *Clim. Dyn.* **52**, 3061–3078 (2019).

Acknowledgements

We are greatly indebted to the large number of scientists, researchers and fishers who have gone to sea to observe, document and understand the Gulf Stream path and its variability, who inspired us to investigate this wind-GS relationship. We thank the organizations of Japan Meteorological Agency and NASA for the JRA-55 and MERRA-2 reanalysis data products. Our thanks to the AVISO group for processing the altimetry products used to extract the GS path. Special thanks to Professor Chris Wolfe of Stony Brook University for sharing the wind data, to Dr. Chris Piecuch of WHOI for his guidance on using reanalysis products and to Dr. Andre Schmidt for technical and computing help at SMAST. We wish to thank the Editor and Reviewers for their helpful comments to a previous version which improved the presentation of this manuscript. AG, IG and AS are grateful for financial supports from NSF (OCE-1815242 and OCE-2123283), SMAST and UMass Dartmouth. GG was supported by NSF under grant OCE-2122726. MA was supported by NSF (OCE-2123111). HO and AS were also supported by WHOI internal grants.

Author contributions

Conceptualization: AG, IG, MA, GG. Methodology: AG, IG, MA, GG, HO, AS. Investigation: IG, AG, MA, GG, OH. Visualization: IG, AG, AS, MA. Funding Acquisition: AG, MA, GG. Project Administration: AG, GG, MA. Writing—original draft: IG, AG, MA, GG, OH. Writing—review and editing: AG, IG, MA, GG, OH, AS.

Funding

National Science Foundation grant OCE-1851242 and OCE-2123283 (IG, AG, AS). National Science Foundation grant DUE-2030552 (ACCOMPLISH) and S-STEM 2130252 (IG). National Science Foundation grant OCE-2123111 (MA). National Science Foundation grant OCE-2122726 (GG). University of Massachusetts Dartmouth internal grants (IG). Woods Hole Oceanographic Institution internal grants (HO and AS).

Competing interests

The authors declare no competing interests.

Additional information

Supplementary Information The online version contains supplementary material available at <https://doi.org/10.1038/s41598-024-68461-0>.

Correspondence and requests for materials should be addressed to A.G.

Reprints and permissions information is available at www.nature.com/reprints.

Publisher's note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit <http://creativecommons.org/licenses/by/4.0/>.

This is a U.S. Government work and not under copyright protection in the US; foreign copyright protection may apply 2024