

Molecular Ball Joints: Mechanochemical Perturbation of Bullvalene Hardy–Cope Rearrangements in Polymer Networks

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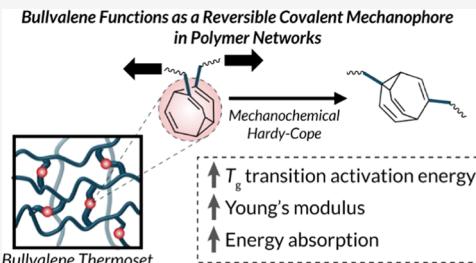
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ABSTRACT: The solution-state fluxional behavior of bullvalene has fascinated physical organic and supramolecular chemists alike. Little effort, however, has been put into investigating bullvalene applications in bulk, partially due to difficulties in characterizing such dynamic systems. To address this knowledge gap, we herein probe whether bullvalene Hardy–Cope rearrangements can be mechanically perturbed in bulk polymer networks. We use dynamic mechanical analysis to demonstrate that the activation barrier to the glass transition process is significantly elevated for bullvalene-containing materials relative to “static” control networks. Furthermore, bullvalene rearrangements can be mechanically perturbed at low temperatures in the glassy region; such behavior facilitates energy dissipation (i.e., increased hysteresis energy) and polymer chain alignment to stiffen the material (i.e., increased Young’s modulus) under load. Computational simulations corroborate our work that showcases bullvalene as a reversible “low-force” covalent mechanophore in the modulation of viscoelastic behavior.



INTRODUCTION

Recent advances in polymer mechanochemistry have led to the development of force-responsive materials.^{1,2} Incorporating force-responsive small molecules (e.g., mechanophores) into polymer networks introduces novel properties such as cargo release^{3–5} or stress visualization.^{6–8} Many related systems are prized, however, for their “sacrificial bonds” that can dissipate exogenous force and enhance mechanical properties. Materials with such bonds can generally be separated into two categories, noncovalent (“low-force” activation) (Figure 1A) and covalent (“high-force” activation) (Figure 1B), depending on how the mechanically active units are structured. Noncovalent systems take advantage of supramolecular interactions and/or mechanical bonds;^{9–13} specific examples include soft materials containing folded proteins,¹⁴ π – π aromatic stacking,^{15–18} hydrogen bonding,^{19–25} metal–organic motifs,^{26–31} rotaxanes,^{32–34} and catenanes.^{19,35–40} Adaptations in these systems are generally reversible with low activation barriers. On the other hand, systems bearing covalent mechanophores use higher energy barrier chemical transformations to adapt to applied mechanical force (e.g., pericyclic reaction). Specific examples include reversible retro-[4 + 2] Diels–Alder reactions for adaptable materials and nonreversible cycloreversions for tear resistance or toughening.^{41–46} While there are no examples to date of low-force materials mediated by covalent bonds (i.e., chemical reactions), we envisioned the potential for “shapeshifting” molecular cages to provide access

to this class of materials with a novel energy absorption motif (Figure 1C).

One of the most studied molecular shapeshifters, bullvalene, was first synthesized by Doering, Roth, and Schröder in the 1960s and since then has fascinated chemists due to its unique molecular architecture.^{47,48} Three alkenes emanating from a central carbon atom, organized symmetrically in a boat conformation, and connected by a strained cyclopropane ring sets up ideal conditions for rapid Hardy–Cope rearrangements.⁴⁹ Unsubstituted bullvalene has more than 1.2 million ($10^{13}/3$) degenerate isomers rapidly interconverting (E_a = ca. 11 kcal/mol) at room temperature (Figure 2A).

Supramolecular interactions can perturb the interconversion of fluxional isomers for applications in molecular recognition (Figure 2B).^{52–55} Such substituted bullvalenes, when unperturbed, exist in an equilibrium distribution of valence isomers. Upon addition of analytes, bullvalene engages in guest–host interactions, shifting to a new equilibrium population. Accessing such multisubstituted bullvalenes, however, is difficult due to lengthy synthetic sequences.^{56,57} Recent progress in concisely accessing substituted bullvalenes

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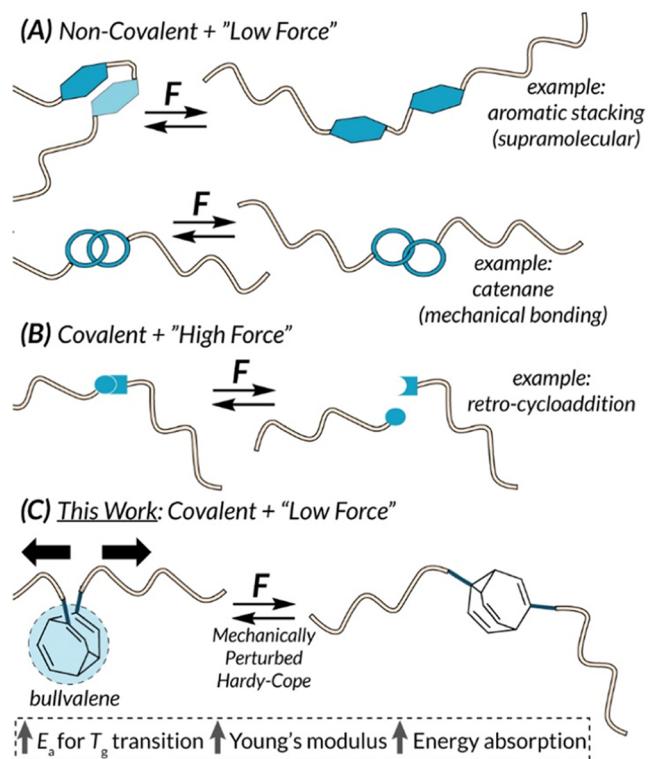


Figure 1. Overview of mechanically responsive soft materials making use of (A) low-force supramolecular interactions or mechanical bonds, (B) high-force covalent bonds, and (C) low-force shape-shifting covalent molecular cages (e.g., bullvalene).

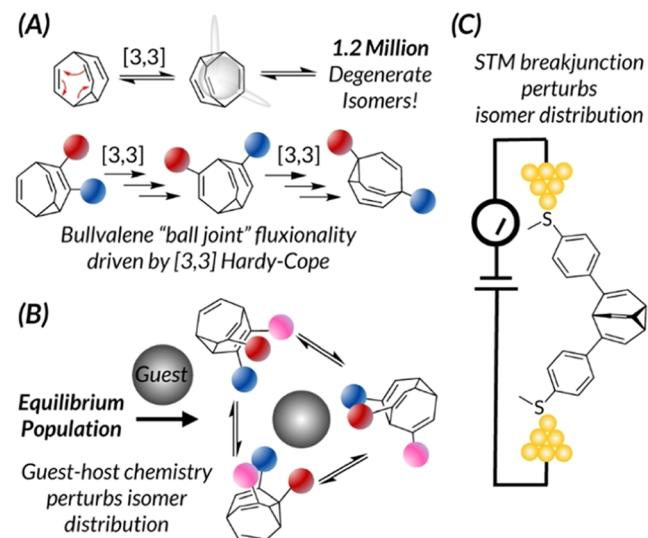


Figure 2. (A) Illustration of rapid bullvalene Hardy-Cope rearrangements. (B) Bullvalene adjusts equilibrium population upon guest-host interactions. (C) Bullvalene employed as a single-molecule circuit junction capable of adapting to the STMBJ movement.

reported by Fallon reignited interest in bullvalene applications.^{58,59} Most applications of bullvalene explored so far are the solution state, however.

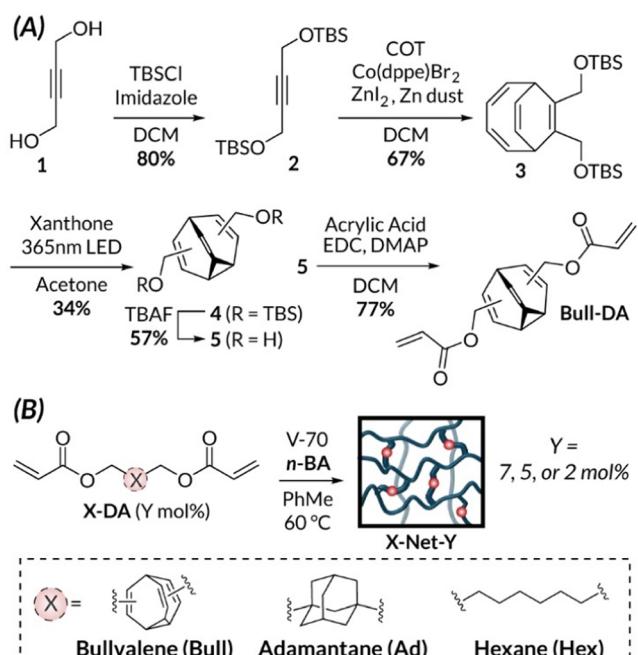
Interestingly, examples of bullvalene rearrangements in the solid state also exist.^{60,61} Specifically, dynamic behaviors were observed in unsubstituted bullvalene single crystals with an activation barrier for Hardy-Cope rearrangements resembling those in isotropic and liquid crystalline solutions. More

recently, Darwish and Fallon demonstrated the mechanical control of bullvalene fluxionality in a scanning tunneling microscopy breakjunction (STMBJ). Perturbed isomerization facilitates dynamic piezoresistance (Figure 2C).⁶² Our previous work on incorporating bullvalene into π -rich thermoplastics showed that stochastic bullvalene isomers could effectively modulate the rigidity of such polymers but left many unanswered questions about resultant thermomechanical behaviors.⁶³

Since bullvalene's valence isomer distribution can adapt to exogenous stimuli in solution and in the bulk (Figure 2),^{52–54,60–62} we imagined that upon applied mechanical force within a bulk material, the equilibrium population of bullvalene could adapt to the applied force vector. We envisioned that studying the thermomechanical properties of these materials would be an ideal way to understand molecular-level sigmatropic rearrangements within polymer networks. If reversible Hardy-Cope rearrangements can be perturbed through mechanical force, we might expect unique thermoset properties by bridging the gap between "low force" and "covalently bonded" force-responsive materials in a strictly nonscissile system (Figure 1C). Herein, we incorporate bullvalene into elastomeric polymer networks to study their thermal behavior and temperature-dependent viscoelastic properties alongside control networks containing "static" cross-linkers. Using dynamic mechanical analysis (DMA), we demonstrate that bullvalene rearrangements not only increase the activation barrier for the glass transition process but also lead to distinct stiffening and energy dissipation behaviors below the glass transition temperature (i.e., glassy region). Furthermore, we computationally explore the impact of external force on discrete bullvalene isomers and identify two plausible "mechanophore" pathways that may contribute to the observed properties. These collective thermal and mechanical properties unique to bullvalene thermosets provide credence for the consequences of molecular fluxionality on bulk materials properties.

RESULTS AND DISCUSSION

Network Synthesis. The requisite cross-linked elastomers for thermal and mechanical testing were synthesized from *n*-butyl acrylate and an appropriate diacrylate cross-linker under thermal free radical conditions (V-70, 60 °C) that we have confirmed bullvalene tolerates (Figure S1). We selected cross-linker loadings between 2 and 7 mol % in the preparation for thermal and mechanical experiments (i.e., X-Net-Y). The cross-linker for target bullvalene networks (Bull-Net-Y), bullvalene diacrylate (Bull-DA), is derived from dimethanol bullvalene 5, the product of Co-catalyzed formal [6 + 2] cycloaddition of silyl-protected 2-butyne-1,4-diol 2 and cyclooctatetraene (COT) after deprotection.^{51,58,64} Photochemical di- π -methane rearrangement^{65–67} to 4 and subsequent esterification with acrylic acid affords Bull-DA (Scheme 1A). In the selection of control static cross-linkers, we first chose adamantane diacrylate (Ad-DA). Bullvalene and adamantane are both C₁₀ hydrocarbon cages with similar molar masses. Conveniently, substituents on adamantane are also situated in a nonplanar fashion, akin to bullvalene's geometry. As Bull-DA and Ad-DA should have similar steric profiles, we can probe the effect of cross-linker topology and fluxionality on bulk materials properties through the synthesis of Ad-Net-Y. These two X-DA cross-linkers were used to access homogeneous X-Net-Y materials (Y = 7, 5, and 2 mol %) as

Scheme 1. Synthesis of X-Net-Y Polymer Networks^a

well as “mixed” Bull_{2.5}Ad_{2.5}-Net-5 and Bull_{1.25}Ad_{3.75}-Net-5 materials to further interrogate the role of bullvalene within cross-linked materials. In addition, we controlled for cross-linker rigidity by utilizing commercially available 1,6-hexanediol diacrylate (Hex-DA) as a second control cross-linker toward Hex-Net-Y (Scheme 1B).

Comparison of Physical Properties. The physical properties of the resulting networks (ca. 80% monomer conversion; Table S1) Bull-Net-7, Ad-Net-7, and Hex-Net-7 were characterized using a variety of analytical techniques. Thermal gravimetric analysis (TGA; Figure S2) confirms that all materials have similar decomposition temperatures (T_d = 364–377 °C), and helium pycnometry demonstrates that all materials have the same densities (ρ = 1.06–1.09 g/cm³; Table S2) at 25 °C. To further interrogate network structure, we performed small-angle X-ray scattering (SAXS) on Bull-Net-7 and Ad-Net-7 (Figure S3). Both bulk samples exhibit relatively

featureless SAXS curves indicative of no long-range order; a lack of sharp features indicates a lack of crystalline and/or highly ordered domains. Tensile plateau moduli (E'_0) measured using DMA frequency sweep experiments at 20 °C confirm that all materials have E'_0 (0.261, 0.269, 0.405 MPa for Bull-Net-7, Ad-Net-7, and Hex-Net-7, respectively) on the same order of magnitude (Figure S4). These data suggest that Bull-Net-7 and Ad-Net-7 have similar entanglement molecular weights, while Hex-Net-7 has shorter distances between cross-links. The glass transition temperatures of these materials were also characterized through modulated differential scanning calorimetry (MDSC). Analysis of the total heat flow shows that Bull-Net-7 and Ad-Net-7 have similar $T_{g,DSC}$ (−34 and −37 °C, respectively) (Figure 3A), while Hex-Net-7 has a lower $T_{g,DSC}$ (−42 °C), likely due to reduced cross-linker rigidity and/or increased cross-linker density.⁶⁸ For the less densely cross-linked materials (X-Net-5 and X-Net-2), adamantine, bullvalene, and mixed networks’ $T_{g,DSC}$ remained similar (Figures 3B,C and S5). We also performed swelling experiments in several organic solvents (methylene chloride, acetone, ethyl acetate, dimethylformamide) to compare cross-linking densities across all X-Net-7 materials (Figure S6). While Bull-Net-7 consistently has the highest equilibrium swelling ratios, they are generally within 10–15% of that of Ad-Net-7. Cross-linker fluxionality (i.e., flexibility) may play a role in facilitating increased swelling for Bull-Net-7. On the other hand, Hex-Net-7 always has a much lower swelling ratio than Bull-Net-7 in all solvents; we attribute these inconsistencies in swelling ratios potentially to differences in cross-linking steric and density. Due to the general similarities in monomer conversion, T_d , ρ , SAXS, $T_{g,DSC}$, E'_0 , and swelling between Bull-Net-7 and Ad-Net-7, we assume that network structures are comparable for the purposes of subsequent thermomechanical analyses. On the other hand, the comparative differences noted above for Hex-Net-7 suggest that its network structure is not comparable and, therefore, serves as a negative control for subsequent studies.

Thermomechanical Analysis. With an enhanced understanding of X-Net-Y physical and thermal behaviors, we next turned to thermomechanical analysis. First, DMA temperature-sweep experiments were performed in the linear viscoelastic region under tension (1 Hz); thermomechanical $T_{g,DMA}$ of Bull-Net-7, Ad-Net-7, and Hex-Net-7 were measured to be 4.6, −9.1, and −18 °C, respectively (Figure 4A), at the apex of each tan (δ) peak. Bull-Net-7 has the broadest tan (δ) peak,

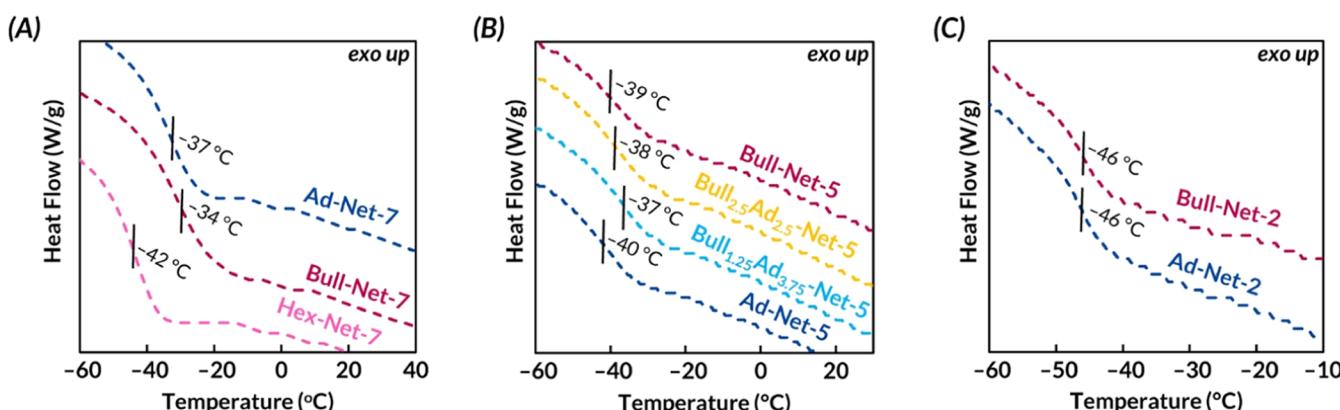


Figure 3. MDSC X-Net-Y total heat flow curves for (A) 7 mol % cross-linker loading (Y = 7), (B) 5 mol % cross-linker loading (Y = 5), including mixed materials derived from Bull-DA and Ad-DA, and (C) 2% cross-linker loading (Y = 2).

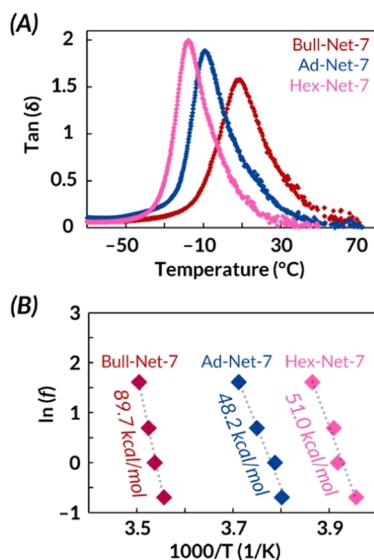


Figure 4. (A) Temperature-sweep experiments (1 Hz) of X-Net-7 materials. (B) Arrhenius plots of $T_{\text{g,DMA}}$ of X-Net-7 materials.

which suggests a more complex damping mechanism (*vide infra*) relative to **Ad-Net-7** and **Hex-Net-7**. It should be noted that while the span of $T_{\text{g,DMA}}$ values is larger than that for $T_{\text{g,DSC}}$, the onset of $T_{\text{g,DMA}}$ is quite similar across all materials. To quantify differences in these thermomechanical glass transition processes, the respective activation energies were calculated through multifrequency temperature-sweep experiments and fitting to the Arrhenius equation (Figures 4B and S7–S9).⁶⁹ Here, E_a for the **Bull-Net-7** $T_{\text{g,DMA}}$ transition (89.7 kcal/mol) is significantly higher than those for **Ad-Net-7** (48.2 kcal/mol) and **Hex-Net-7** (51.0 kcal/mol). This broader $\tan(\delta)$ peak and elevated E_a for **Bull-Net-7** could be caused by an increase in fluxional rearrangements as the material approaches the glass transition temperature. As we describe in more detail below (see the **Computational Mechanistic Investigations** section), at lower temperatures under tension, the Hardy–Cope fluxionality may exist as a combination of mechanically guided thermal rearrangements (i.e., slowly interconverting bullvalenes “align” to the force vector) and/or mechanically activated rearrangements (i.e., discrete bullvalene isomers function as force-driven mechanophores).⁷⁰ These are not exhaustive options but rather serve as two plausible pathways that are conceptually related to STMBJ-driven bullvalene isomerization reported by Darwish and Fallon.⁶²

In the case of a broader $\tan(\delta)$ peak, the exact rearrangement pathway traversed by a particular bullvalene cross-linker likely varies across the continuum of the material due to local network inhomogeneities and/or cooperativity from neighboring bullvalene moieties. In other words, **Bull-Net-7** exhibits a wide range of molecular motion stemming from an ensemble of bullvalene rearrangements; the corresponding $T_{\text{g,DMA}}$ transition, therefore, occurs over a broader range of temperatures compared to **Ad-Net-7** and **Hex-Net-7**. In the case of **Bull-Net-7**’s elevated E_a for the $T_{\text{g,DMA}}$ transition, long-range segmental motion is restricted below $T_{\text{g,DMA}}$ and bullvalene has access to only a few immediate isomers that do not cause structural changes to the polymer network. As the temperature continues to increase while approaching $T_{\text{g,DMA}}$, molecular mobility dramatically increases. Consequently, there are more degrees of freedom for bullvalene to access additional

isomers while adapting to the external force vector through a series of Hardy–Cope rearrangements. In doing so, energy is dissipated within the bullvalene cage (i.e., additional energy is required for sigmatropic rearrangements) and the observed E_a for **Bull-Net-7** is elevated relative to those of the control networks.

Importantly, despite **Ad-Net-7** acting as a “similar” network and **Hex-Net-7** acting as a “dissimilar” network relative to **Bull-Net-7**, **Bull-Net-7** retains the broadest $\tan(\delta)$ peak and highest E_a for $T_{\text{g,DMA}}$ transition across all three materials. In fact, these trends hold true for bullvalene-containing X-Net-5 materials as well (Table 1 and Figures S10–S16), suggesting

Table 1. Summary of X-Net-Y Temperature-Sweep Experiments

sample	$T_{\text{g,DMA}} ({}^{\circ}\text{C})^a$	E_a of T_{g} transition (kcal/mol) ^b
Bull-Net-7	4.6	89
Ad-Net-7	-9.1	48
Hex-Net-7	-18	51
Bull-Net-5	-11	56
Bull ₂ Ad ₂ -Net-5	-14	59
Bull _{1.25} Ad _{3.75} -Net-5	-15	44
Ad-Net-5	-13	34
Bull-Net-2	-25	41
Ad-Net-2	-28	39

^aReported from the apex of $\tan(\delta)$ peak (1 Hz). ^bCalculated from an Arrhenius fit of multifrequency (5, 2, 1, 0.5 Hz) temperature-sweep experiments.

that cross-linker identity rather than network structure is a key driver in these phenomena. At lower cross-linker loadings (i.e., **Bull-Net-2** and **Ad-Net-2**), the minimal difference is observed in temperature-sweep experiments as longer distances between cross-links dilute the impacts of bullvalene described above.

Ad-Net-7 and **Bull-Net-7** were then subjected to frequency sweep experiments using DMA. In these experiments, we studied thermomechanical behavior across a range of temperatures and frequencies (1–100 Hz). We then generated master curves for each of the two networks spanning many decades of frequency using the principle of time–temperature superposition (Figures 5, S17, and S18). Master curves were generated by using appropriate shift factors with $T_{\text{ref}} = T_{\text{g,DMA}} - 20$ $^{\circ}\text{C}$. While the master curves are generally comparable between **Ad-Net-7** and **Bull-Net-7**, further supporting our claim that these two thermosets have similar network structures, there is an interesting difference that is worth pointing out. The onset of the glassy region is more gradual near the edge of the glass transition region (higher frequency/lower temperature) in the case of **Bull-Net-7** compared to **Ad-Net-7**. In other words, E' and E'' level off more rapidly than what is observed for **Bull-Net-7** (Figure 5A). As a result, the corresponding $\tan(\delta)$ curve (Figure 5B) for **Bull-Net-7** is broader than that for **Ad-Net-7**, akin to what was observed in our temperature-sweep experiments (Figure 4). In addition, this broadening observed in the **Bull-Net-7** $\tan(\delta)$ curve is unsymmetrical; the curve tails significantly into the higher-frequency/lower-temperature region because of the more gradual onset of the glassy region as reflected in E' and E'' values. At higher temperatures well above $T_{\text{g,DMA}}$, it is known that bullvalene remains fluxional. When bullvalene is dynamic, we surmise that it can rapidly sample an ensemble of valence isomers and molecular motion is rapid. The rapid bullvalene

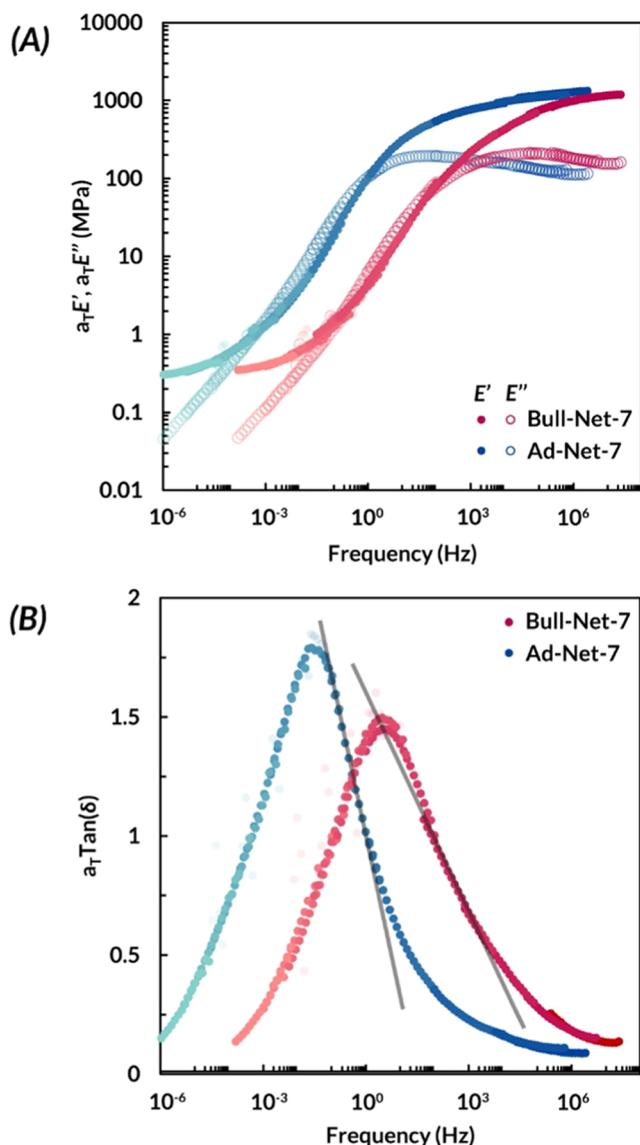


Figure 5. Master curves for Ad-Net-7 and Bull-Net-7 detailing (A) storage (E') and loss (E'') moduli and (B) $\tan(\delta)$ (gray lines indicate differences in the high-frequency region).

fluxionality alongside freedom in molecular and chain mobility renders Bull-Net-7 similar to Ad-Net-7 at lower frequency/higher temperature. Our previous work studying bullvalene thermoplastics supports operative Hardy–Cope rearrangements throughout the polymer backbone as evidenced by enhanced thermal stability relative to unsaturated analogues.⁶³ As the temperature decreases and approaches $T_{g,DMA}$, macromolecular motion is limited and the bullvalene rearrangements slow down; some bullvalenes may even be kinetically trapped as static valence isomers. As energy is applied through oscillatory tension in frequency sweep experiments, bullvalene may rearrange to applied mechanical force (*vide infra*) by absorbing energy and guiding these rearrangements along the way. As a result, we hypothesize that the broadening of the Bull-Net-7 $\tan(\delta)$ curve and asymmetry at higher frequencies/lower temperatures is the consequence of less thermally active bullvalenes. The result is an ensemble of damping pathways, potentially involving disparate mechan-

ically guided and/or mechanically activated Hardy–Cope rearrangements across a broad range of temperatures.

Tensile Experiments. To glean insight into potential temperature-dependent phenomena in Bull-Net-X, solution-state variable temperature NMR (VT-NMR) experiments (Figure 6A) were conducted on Bull-DA, which reveal slow chemical exchange at 25 °C consistent with E_a = ca. 11 kcal/mol; broad bullvalene vinyl proton (B_V), alkyl proton (B_A), and methylene linker proton (B_L) resonances around 5.9, 2.3, and 4.5 ppm, respectively, support this claim. When Bull-DA is cooled to −40 °C, B_V , B_A and B_L resonances sharpen as the rate of chemical exchange via thermal Hardy–Cope rearrangements decreases further and approaches a static structure. Meanwhile, proton resonances not directly connected to the bullvalene cage (A , A' , A'') actually broaden as bond rotation (E_a < 3 kcal/mol) slows on the NMR time scale. The differing dynamics between the chemical (i.e., bullvalene Hardy–Cope rearrangements) and physical (i.e., bond rotations) allow us to tame bullvalene fluxionality at lower temperatures, rendering bullvalene nearly static and confining it to a smaller quantity of valence isomers. In fact, these observations align well with the unique low-temperature behavior observed in Bull-Net-Y compared to Ad-Net-Y (*vide supra*). Thus, in order to probe mechanochemical perturbation and/or activation of Hardy–Cope rearrangements in bulk materials outside the linear viscoelastic region under tension, we reasoned that experiments needed to be conducted at lower temperatures where thermal bullvalene fluxionality becomes suppressed.

Examination of X-Net-7 bulk energy dissipation behavior was carried out using DMA via cyclic loading experiments. At 20 °C, Bull-Net-7, as well as control network Ad-Net-7, display minimal hysteresis (Figures S19–S22). Bull-Net-7 and Ad-Net-7 also have similar Young's moduli (E = ca. 0.0022 MPa; Table S3). When the same cyclic loading experiments were carried out at −40 °C (Figures 6B and S23–S28), all materials displayed strain–stress curves known for a glassy polymer with little immediate recovery upon unloading.^{71,72} While Ad-Net-7 has noticeable energy dissipation characteristics (E_H = 2110 J/m²), Bull-Net-7 has nearly 3 times the hysteresis energy of Ad-Net-7 (E_H = 6240 J/m²) with enhanced stiffening (E = 8.5 MPa; Table S4).

It should be noted that the mechanism of deformation in materials below glass transition is different from conventional viscoelastic behavior above T_g . Instead of long-range disentanglement through chain sliding, bond cleavage and local molecular motion generally dominate.^{68,71,72} Repeated cyclic loading of all X-Net-7 materials at −40 °C does not show a loss in mechanical performance, thus confirming that minimal covalent bond cleavage occurs at a low temperature in X-Net-7 materials.

To confirm that this behavior (Figure 6) is not due to variations in material density below $T_{g,DMA}$, we also performed thermal expansion experiments between 20 and −40 °C (Figure S29). As no significant differences were observed between Bull-Net-7 and Ad-Net-7, corroboration with the pycnometry data (Table S2) implies comparatively low temperature densities. The observed differences in hysteresis energies may be attributed to differences in the relative stiffness and yield stress of these materials. For example, Bull-Net-7 has a much higher Young's modulus (E = 8.5 MPa) relative to Ad-Net-7 (E = 3.8 MPa). We hypothesize that unlike the sterically demanding adamantane, similarly sized bullvalene rearranges upon applied mechanical force (i.e.,

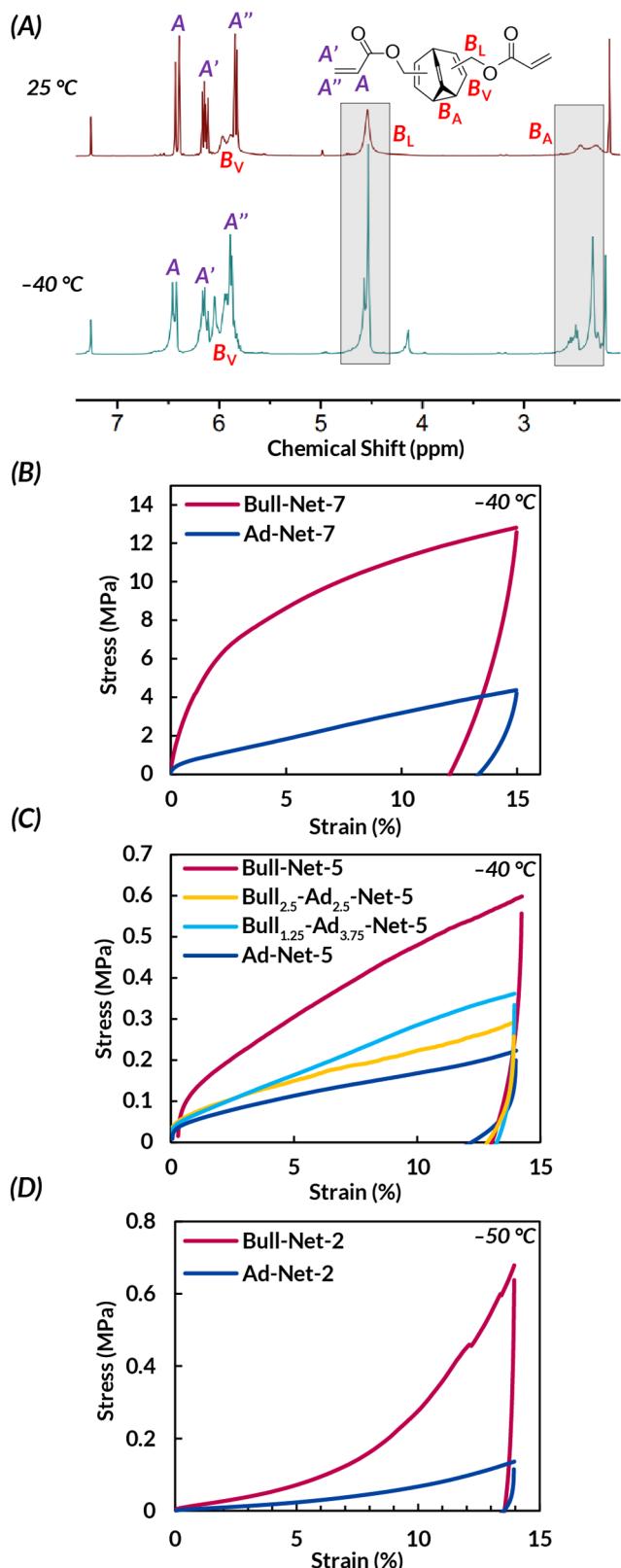


Figure 6. (A) VT ^1H NMR spectrum of Bull-DA. Cyclic loading of X-Net-Y materials: (B) $Y = 7$, (C) $Y = 5$, and (D) $Y = 2$.

“molecular ball joint”) to further realign the already packed polymer strands together and make the material stiffer (Figure 7). In other words, as the rate of bullvalene Hardy–Cope rearrangements decreases at a lower experimental temperature,

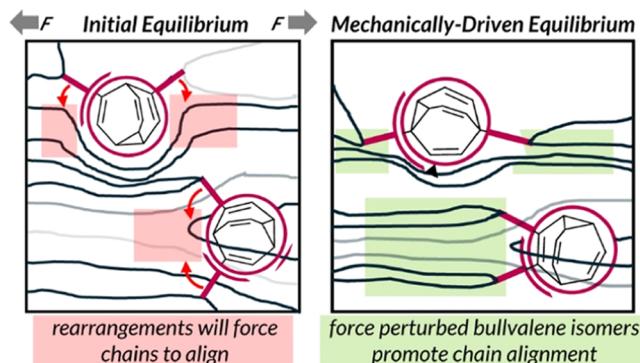


Figure 7. (A) Bullvalene adapts to applied force (B) through chain alignment governed by mechanically guided and/or mechanically activated Hardy–Cope rearrangements.

there is an opportunity to guide the valence isomer trajectory and/or activate certain isomers with external mechanical force. Additionally, without free segmental motion, the initial input mechanical energy can be first “absorbed” in Bull-Net-7 by bullvalene rearrangements and the subsequent realignment of polymer chains. The result is that Bull-Net-7 has a much higher yield stress (ca. 6 MPa at ca. 2% strain) compared to Ad-Net-7 (yield stress <1 MPa at <1% strain) in the shear yielding process.^{71,72} In other words, at low temperatures, Bull-Net-7 has increased stiffness and yield stress relative to control networks; these properties contribute to the high hysteresis energy for Bull-Net-7 driven by the sigmatropic rearrangement. These phenomena are not limited to X-Net-7; analogous Ad-Net-Y and Bull-Net-Y with a lower cross-linker loading ($Y = 2$ or 5 mol %) show similar trends in tensile properties (Figure 6C,D and Table 2). We also tested the same mixed

Table 2. Summary of X-Net-Y Tensile Results

sample	Young's modulus, E (MPa)	hysteresis energy, E_H (kJ/m ²)
Bull-Net-7 ^a	8.50	6.24
Ad-Net-7 ^a	3.80	2.11
Hex-Net-7 ^a	0.110	0.805
Bull-Net-5 ^a	1.00	0.226
Bull _{2.5} -Ad _{2.5} -Net-5 ^a	0.640	0.198
Bull _{1.25} -Ad _{3.75} -Net-5 ^a	0.680	0.141
Ad-Net-5 ^a	0.600	0.0979
Bull-Net-2 ^b	0.810	0.124
Ad-Net-2 ^b	0.029	0.0301

^aData collected at -40 °C. ^bData collected at -50 °C.

systems previously used to study $T_{g,\text{DMA}}$ (Table 1) and found that Bull_{2.5}-Ad_{2.5}-Net-5 and Bull_{1.25}-Ad_{3.75}-Net-5 have intermediary E and E_H relative to Bull-Net-5 and Ad-Net-5; these data further support our claim that bullvalene incorporation enhances mechanical properties. Furthermore, these processes were all demonstrated to be fully reversible; when X-Net-7 materials are allowed to warm to room temperature between successive loading at -40 °C, minimal changes in the resulting stress-strain curves are observed (Figures S30–S32).

Computational Mechanistic Investigations. To explain the unique material behaviors described above, we probed how an external mechanical force impacts bullvalene stability and associated Hardy–Cope rearrangement pathways. To do so, we turned to constrained geometries simulate external force

(CoGEF) simulations (see the Supporting Information for details).^{73–75} We note that these calculations do not consider competitive thermal Hardy–Cope rearrangements and rather only consider the impact of exogenous force on individual static bullvalene valence isomers. To initiate our investigations, we first minimized derivatives of the 12 possible geometric isomers of **Bull-DA** (i.e., dimethylbullvalene analogues **B1–B12**; Figure S33) and calculated the energy profiles for thermal Hardy–Cope interconversions (Figures S35–S41) using unrestricted DFT at the B3LYP-D3/def2svp CPCM(DCM) level of theory. Next, for each bullvalene isomer, we simulated the impact of applying an external force on these isomers to “pull” them apart by increasing the relative distance between the two alkyl substituents (increments of 0.05 Å over 50 steps) on the bullvalene core until a C–C bond homolysis or a mechanochemical event occurred (Figures S41–S43). For isomers **B1–B4** and **B6–B10**, biradical formation occurs via the methyl group homolytic bond extrusion or bullvalene backbone scission (Figures S41–S43 and Table S5).⁷⁵ However, these processes are prohibitively high in energy (>75 kcal/mol between the ground state geometry of each isomer prior to biradical formation) and interestingly do not involve the formation of a radical on the potentially capricious cyclopropane ring in bullvalene. Taken together, based on the reversibility observed in cyclic loading–unloading experiments on **Bull-Net-7** (Figures S30–S32) and these calculations, this high-force (>ca. 4 nN) deleterious mechanochemical bullvalene decomposition is unlikely under our experimental conditions. On the other hand, increasing the distance between substituents in isomers **B11** and **B12** drives the Hardy–Cope rearrangement via the collapse to another isomer (i.e., **B11** → **B4** and **B12** → **B3**). Notably, these events are much lower energy processes (ca. 20–35 kcal/mol) and can proceed with lower calculated forces (<ca. 4 nN) (Figure S43 and Table S5). Thus, in addition to potentially “guiding” the trajectory of thermal Hardy–Cope rearrangements ($\Delta E^\ddagger < \text{ca. } 15 \text{ kcal/mol}$; Figures S34–S40) under mechanical tension, we also present another plausible mechanism that treats two discrete bullvalene valence isomers as mechanophores capable of undergoing mechanically activated Hardy–Cope rearrangements. While it is difficult to probe specific valence isomer interconversions within each **Bull-Net-Y** material, one potential force-responsive pathway could involve two mechanochemical Hardy–Cope rearrangements (large geometrical changes, **B12** → **B3** and **B11** → **B4**) joined by mechanically guided thermal Hardy–Cope rearrangement (minimal geometrical change, **B3** → **B11**) (Figure 8).

CONCLUSIONS

In summary, we report the utility of low-barrier Hardy–Cope rearrangements under strain in polymer networks by incorporating fluxional bullvalenes into *n*-butyl acrylate-based thermoset elastomers. DMA temperature-sweep experiments quantitatively demonstrate that bullvalene can dissipate energy more effectively than static controls, thereby increasing the activation energy for the glass transition process. While thermomechanical testing supports similar behavior between bullvalene (**Bull-Net-Y**) and adamantane (**Ad-Net-Y**) cross-linked materials at room temperature, frequency sweep experiments suggest that rearrangements can be activated mechanochemically as the rate of bullvalene Hardy–Cope rearrangements decrease at a lower temperature. Cyclic loading experiments at temperatures below $T_{g,\text{DMA}}$ demonstrate that

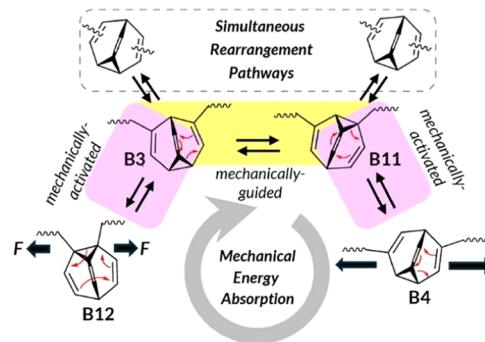


Figure 8. Representative proposed mechanism of action where mechanically activated (pink) and mechanically guided (yellow) Hardy–Cope rearrangements facilitate mechanical energy absorption in **Bull-Net-Y** materials.

Bull-Net-Y is much stiffer with a high hysteresis energy. Comparisons with **Ad-Net-7** suggest that bullvalene’s ability to rearrange upon applied mechanical force facilitates chain alignment and stiffens the material upon loading. Finally, CoGEF simulations are employed to demonstrate the stability of the bullvalene cage under an external mechanical force while also revealing two “mechanically competent” valence isomers capable of engaging in mechanically activated Hardy–Cope rearrangements. Overall, having a simple unimolecular “ball joint” moiety that can dissipate energy through reversible sigmatropic rearrangements shows great potential for enhancing the durability of existing thermosets and opens new opportunities for impact- and vibration-resistant materials.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/jacs.4c04401>.

Experimental and synthetic procedures, solution-state and bulk characterization data, and thermal and mechanical testing data (PDF)

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Notes

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