DIFFERENTIABILITY OF EFFECTIVE FRONTS IN THE CONTINUOUS SETTING IN TWO DIMENSIONS

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ABSTRACT. We study the effective front associated with first-order front propagations in two dimensions (n=2) in the periodic setting with continuous coefficients. Our main result says that that the boundary of the effective front is differentiable at every irrational point. Equivalently, the stable norm associated with a continuous \mathbb{Z}^2 -periodic Riemannian metric is differentiable at irrational points. This conclusion was obtained decades ago for smooth metrics ([4, 6]). To the best of our knowledge, our result provides the first nontrivial property of the effective fronts in the continuous setting, which is the standard assumption in the PDE theory. Combining with the sufficiency result in [15], our result leads to a realization type conclusion: for continuous coefficients, a polygon could be an effective front if and only if it is centrally symmetric with rational vertices and nonempty interior.

1. Introduction

1.1. **Settings.** We first give a minimalistic introduction to the periodic homogenization of Hamilton-Jacobi equations. For each $\varepsilon > 0$, let $u^{\varepsilon} \in C(\mathbb{R}^n \times [0, \infty))$ be the viscosity solution to

(1.1)
$$\begin{cases} u_t^{\varepsilon} + H\left(\frac{x}{\varepsilon}, Du^{\varepsilon}\right) = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u^{\varepsilon}(x, 0) = g(x) & \text{on } \mathbb{R}^n. \end{cases}$$

Here, the Hamiltonian $H=H(y,p):\mathbb{R}^n\times\mathbb{R}^n\to\mathbb{R}$ is a given continuous function satisfying

$$(1.2) \quad \begin{cases} \text{for } p \in \mathbb{R}^n, \ y \mapsto H(y,p) \text{ is } \mathbb{Z}^n\text{-periodic;} \\ H \text{ is coercive in } p, \text{ that is, } \lim_{|p| \to \infty} \min_{y \in \mathbb{R}^n} H(y,p) = +\infty. \end{cases}$$

For the initial data q, we assume

$$(1.3) g \in \mathrm{BUC}(\mathbb{R}^n) \cap \mathrm{Lip}(\mathbb{R}^n),$$

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where BUC (\mathbb{R}^n) is the set of bounded, uniformly continuous functions on \mathbb{R}^n .

Under assumptions (1.2)–(1.3), u^{ε} converges to u locally uniformly on $\mathbb{R}^n \times [0, \infty)$ as $\varepsilon \to 0$, and u solves the effective equation (see [14, 10, 20])

(1.4)
$$\begin{cases} u_t + \overline{H}(Du) = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u(x, 0) = g(x) & \text{on } \mathbb{R}^n. \end{cases}$$

The effective Hamiltonian $\overline{H} \in C(\mathbb{R}^n)$ depends nonlinearly on H, and is determined by the cell (ergodic) problems as follows.

Definition 1 (Effective Hamiltonian). For each $p \in \mathbb{R}^n$, there exists a unique constant $\overline{H}(p) \in \mathbb{R}$ such that the following cell problem has a continuous \mathbb{Z}^n -periodic viscosity solution

(1.5)
$$H(y, p + Dv) = \overline{H}(p) \qquad in \ \mathbb{T}^n = \mathbb{R}^n / \mathbb{Z}^n.$$

Note that v = v(y, p) is not unique even up to additive constants in general.

If H is convex in p, then so is \overline{H} . In this case, the effective Hamiltonian is also given by an inf-max formula (see, e.g., [20])

(1.6)

$$\overline{H}(p) = \inf_{\phi \in C^{\infty}(\mathbb{T}^n)} \max_{y \in \mathbb{T}^n} H(y, p + D\phi(y)) = \inf_{\phi \in C^1(\mathbb{T}^n)} \max_{y \in \mathbb{T}^n} H(y, p + D\phi(y)).$$

It is clear that \overline{H} is defined in a very implicit way. A central and fundamental goal in the homogenization theory is to understand qualitative and quantitative properties of \overline{H} . To date, not much is known about fine properties of \overline{H} .

For clarity of presentation, throughout this paper, we focus on the case H(y,p)=a(y)|p| for $a\in C(\mathbb{T}^n,(0,\infty))$ that arises from the modeling of first-order front propagations (e.g., crystal growth, flame propagation), which is probably one of the most physically relevant examples in the homogenization theory. In this situation, $\overline{H}(p)$ represents the effective propagation speed. Thanks to the above inf-max formula,

(1.7)

$$\overline{H}(p) = \inf_{\phi \in C^{\infty}(\mathbb{T}^n)} \max_{y \in \mathbb{T}^n} a(y) |p + D\phi(y)| = \inf_{\phi \in C^1(\mathbb{T}^n)} \max_{y \in \mathbb{T}^n} a(y) |p + D\phi(y)|.$$

Clearly, \overline{H} is convex, even, and positively homogeneous of degree 1. We sometime write $\overline{H} = \overline{H}_a$ to emphasize the dependence on the function a. Due to those properties of \overline{H}_a , its 1-sublevel set

$$S_a := \left\{ p \in \mathbb{R}^n : \overline{H}_a(p) \le 1 \right\}$$

belongs to W, which denotes the collection of all convex sets in \mathbb{R}^n that are centrally symmetric with nonempty interior. The convex dual D_a of S_a , determined by

$$D_a = \partial \overline{H}_a(0),$$

the subdifferential of \overline{H}_a at the origin, is called the effective front, which also belongs to W. The following realization problem is of our main interests.

Question 1. For what kind of $W \in \mathcal{W}$ does there exist a function $a \in C(\mathbb{T}^n, (0, \infty))$ such that $S_a = W$ (or $D_a = W$)?

We would like to mention that this kind of problems has been one of the major research focuses in the context of stable norms in geometry/dynamical system and first passage percolation theory in probability. See discussions about connections later.

Definition 2. We say that a point p is rational if $p \in \lambda \mathbb{Z}^n$ for some $\lambda \in \mathbb{R}$. A point p is called irrational if it is not rational. Moreover, for a n-dimensional convex set K, ∂K is said to be differentiable at $p \in \partial K$ if there exists a unique unit vector $q \in \mathbb{R}^n$ such that $q \cdot (p' - p) \leq 0$ for all $p' \in \partial K$.

When n=2 and $a\in C^2(\mathbb{T}^2,(0,\infty))$, the following extra restrictions are known in equivalent forms of stable norms in metric geometry or β -functions in the Aubry-Mather theory. See [3, 4, 7, 19] for instance.

- (i) ∂S_a is C^1 . Equivalently, D_a is strictly convex.
- (ii) ∂S_a is not strictly convex (i.e., it contains line segments) unless a is constant. Equivalently, ∂D_a is not C^1 unless a is constant.
- (iii) ∂S_a does not contain a line segment of irrational slope. Equivalently, ∂D_a is differentiable at every irrational point.

The C^2 regularity was needed in the proofs of (i)–(iii) to ensure that corresponding Hamiltonian systems (or geodesics in the stable norm context) have unique solutions, which implies that two distinct orbits minimizing the associated actions (or minimal geodesics) cannot intersect twice. Together with two dimensional topology, this leads to a beautiful identification of the minimizing orbits with circle maps that provides a nice characterization of structures of these orbits (see [3]). For example, minimal orbits in the same Aubry set are well ordered. Also, there exists a Hölder continuous selection of viscosity solution $v(\cdot, p)$ of (1.5) in two dimensions with respect to a suitable family of parameters ([8]).

However, in the merely continuous situation, uniqueness of solutions to the corresponding ODEs and intersection restrictions of distinct minimizing orbits cease to exist. Two minimizing orbits might intersect multiple (even infinitely many) times. Structures of minimizing orbits could be topologically very bad and contain various pathological behaviors. As one of the consequences, it is known now that ∂S_a might not be C^1 for $a \in C(\mathbb{T}^2, (0, \infty))$, i.e., the above property (i) fails. In particular, it was proved in [15] that for any $\alpha \in (0,1)$, every polygon with rational slopes in \mathcal{W} can be S_a for some a in $C^{1,\alpha}(\mathbb{T}^2, (0,\infty))$ (i.e., realizable), which implies that realizable sets are at least dense in \mathcal{W} . It is then tempting to think that every shape in \mathcal{W} might be realizable in the class of $C(\mathbb{T}^2, (0,\infty))$. In this paper, we give a negative answer to this by showing that the above property (iii) still holds in the continuous setting.

1.2. Main results.

Theorem 1.1. Assume that n=2, and H(y,p)=a(y)|p| for $(y,p) \in \mathbb{T}^2 \times \mathbb{R}^2$ for some $a \in C(\mathbb{T}^2,(0,\infty))$. Then, ∂S_a does not contain a line segment of irrational slope. Equivalently, ∂D_a is differentiable at every irrational point.

Combining with the existence result in [15], Theorem 1.1 implies that when n = 2, for continuous coefficients, a polygon could be an effective front D_a if and only if it is centrally symmetric with rational vertices and nonempty interior. To the best of our knowledge, this might be the first realization result within an interesting class of functions.

The above result can be proved either from the geometric point of view using minimizing geodesics or from the PDE point of view using characteristics of solutions of the cell problem (1.5). In this paper, we choose to use the latter approach. For that, let us consider the closely related mechanical Hamiltonian

$$H(y,p) = \frac{1}{2}|p|^2 + V(y)$$
 for $(y,p) \in \mathbb{T}^n \times \mathbb{R}^n$,

for some $V \in C(\mathbb{T}^n)$. Let $\overline{H}(p)$ be the associated effective Hamiltonian. Note that for $\overline{H}(p) > \max_{\mathbb{T}^n} V$,

$$\frac{1}{2}|p+Dv|^2+V(y)=\overline{H}(p)\quad \Rightarrow\quad \frac{1}{\sqrt{2(\overline{H}(p)-V(y))}}|p+Dv|=1.$$

Accordingly, for $c > \max_{\mathbb{T}^n} V$ and $a = \frac{1}{\sqrt{2(c-V)}}$,

$$F_c = \{ p \in \mathbb{R}^2 : \overline{H}(p) = c \} = \partial S_a.$$

Here is our second main result, which implies directly Theorem 1.1.

Theorem 1.2. Assume that $H(y,p) = \frac{1}{2}|p|^2 + V(y)$ for $(y,p) \in \mathbb{T}^n \times \mathbb{R}^n$ for some $V \in C(\mathbb{T}^n)$. Then, the following properties hold.

(1) For any $n \in \mathbb{N}$, $p_0, p_1 \in \mathbb{R}^n$, if $\overline{H}(p_0) \neq \overline{H}(p_1)$, then for all $\lambda \in (0, 1)$,

$$\overline{H}(\lambda p_1 + (1-\lambda)p_0) < \lambda \overline{H}(p_0) + (1-\lambda)\overline{H}(p_1),$$

that is, \overline{H} is strictly convex along directions that are not tangential to the level set in any dimension.

(2) For n = 2 and $c > \max_{\mathbb{T}^2} V$, F_c does not contain a line segment of irrational slope.

Property (1) in the above theorem was known for smooth V (see [11]). As an immediate corollary, we obtain that

Corollary 1.3. Assume the settings in Theorem 1.2. Then, the effective Lagrangian $\overline{L}(q)$, defined as

$$\overline{L}(q) = \sup_{p \in \mathbb{R}^2} \left\{ p \cdot q - \overline{H}(p) \right\},\,$$

is differentiable at every irrational point q if q is not singular (i.e., $q \notin \partial \overline{H}(p)$ for $p \in \{\overline{H} = \max_{\mathbb{T}^2} V\}$.)

The non-singular assumption added in the above corollary is just for technical convenience so that we do not need to deal with the minimum level set $\{p \in \mathbb{R}^2 : \overline{H}(p) = \max_{\mathbb{T}^2} V\}$, which we believe can be removed. This is left to interested readers as an exercise. The same conclusion is known to be true for smooth Hamiltonians. See [9, 18] for instance.

1.3. Connection with stable norms. For $a \in C(\mathbb{T}^n, (0, \infty))$, we define a corresponding periodic Riemannian metric on \mathbb{R}^n as

$$g = \frac{1}{a(x)} \sum_{i=1}^{n} dx_i^2.$$

Let $d_a(\cdot, \cdot)$ denote the distance function induced by this metric. The *stable* norm associated with q (or a) is defined as

(1.8)
$$||x||_a = \lim_{\lambda \to \infty} \frac{d_a(0, \lambda x)}{\lambda} \quad \text{for } x \in \mathbb{R}^n.$$

Properties of stable norms associated with more general Riemannian metric have been extensively studied in the community of dynamical systems and geometry. A central question there is what kind of metric could be a stable norm (realization problem). See [5] for more background. The limit (1.8) can also be viewed as a homogenization problem of the following static Hamilton-Jacobi equation

$$\begin{cases} a\left(\frac{x}{\varepsilon}\right)|Dw^{\varepsilon}| = 1 & \text{for } x \in \mathbb{R}^n \setminus \{0\}, \\ w^{\varepsilon}(0) = 0. \end{cases}$$

Here, w^{ε} is the maximal viscosity solution to the above. By the optimal control formula, for $x \in \mathbb{R}^n$,

$$w^{\varepsilon}(x) = \varepsilon d_a\left(0, \frac{x}{\varepsilon}\right).$$

As $\varepsilon \to 0$, $w^{\varepsilon} \to w = \|\cdot\|_a$ locally uniformly on \mathbb{R}^n , and $w = \|\cdot\|_a$ is the maximal viscosity solution of

$$\begin{cases} \overline{H}_a(Dw) = 1 & \text{for } x \in \mathbb{R}^n \setminus \{0\}, \\ w(0) = 0. \end{cases}$$

The effective front $D_a = \partial \overline{H}_a(0)$ is then exactly the unit ball of the stable norm $\|\cdot\|_a$

$$D_a = \{ x \in \mathbb{R}^n : ||x||_a \le 1 \}.$$

Our Theorem 1.1 is equivalent to the conclusion that the stable norm $||x||_a = w(x)$ is differentiable at irrational points when n = 2 for merely continuous metric, which was known for smooth metric [4]. See [6] for some extensions to higher dimensions. We would like to mention that the connection between Hamilton-Jacobi equations and geometry is well-known to experts and these

equations are sometimes called metric problems in the PDE literature. The general static Hamilton-Jacobi equation corresponds to a Finsler metric. See [16, 17, 20] and the references therein. However, there are very few works concerning deep properties of effective quantities in the PDE literature where the focus is mainly on the wellposedness issue.

We also want to point out that if an i.i.d metric function a(x) is considered on a lattice, this essentially becomes the first passage percolation (FPP) problem in the probability theory and the corresponding effective front is called "limiting shape" there. Also, the realization question is a major open problem ([2]). Homogenization in PDE, stable norm in geometry and limiting shape if FPP are basically the same type of averaging.

1.4. **Open problems.** We list here several open questions that are of interests.

Question 2. Does there exist a nonconstant $a \in C(\mathbb{T}^2, (0, \infty))$ such that S_a is a strictly convex set (e.g., a disk)?

Question 3. Is any set in W realizable if we look at $a \in L^{\infty}(\mathbb{T}^2, (0, \infty))$ with positive essential lower bound?

Question 4. What can we say about the effective front in higher dimensions $(n \geq 3)$, which is much harder due to the lack of topological restrictions? A counterexample constructed in [6] says that Theorem 1.1 is not always true when $n \geq 3$ even for smooth a(y). In addition, by a clever approach that does not rely on two dimensional topology, a suitable generalization of Theorem 1.1 was proved in [6] for $a \in C^3(\mathbb{T}^n, (0, \infty))$. An interesting question is whether the C^3 assumption there can be relaxed, which is closely related to regularity of solutions to Hamilton-Jacobi equations. See Remark 1. In fact, the so called "generalized coordinates" in [6] are equivalent to viscosity subsolutions $u = p \cdot y + v(y)$ to the cell problem

$$a(y)|Du(y)| = a(y)|p + Dv(y)| = 1$$
 in \mathbb{T}^n .

Organization of the paper. The paper is organized as follows. Some preliminary results are given in Section 2. We then give the proof of Theorem 1.2 in Section 3. One of the key ideas is to rearrange pieces of intersecting minimizing orbits (Remark 3 and Remark 4) in order to simplify the topology and reduce to situations that are kind of similar to the classical scenarios in [4].

2. Preliminaries

In this section, we review some relevant concepts and facts from weak KAM theory [12] (see also [9]) and the Aubry-Mather theory [3] in the continuous setting. See [13] for extensions to more general quasiconvex Hamiltonians. We assume throughout this section

$$H(y,p) = \frac{1}{2}|p|^2 + V(y)$$
 for all $(y,p) \in \mathbb{T}^n \times \mathbb{R}^n$,

where $V \in C(\mathbb{T}^n)$. Then, the cell problem reads

(2.1)
$$\frac{1}{2}|p + Dv(y)|^2 + V(y) = \overline{H}(p).$$

It is important to note that we only have V is merely continuous on \mathbb{T}^n . Except the differentiability property Lemma 2.2, most of the definitions and proofs are straightforward extensions of those in the smooth setting.

Definition 3. Let $\eta \in AC([a,b], \mathbb{R}^n)$ be a given curve. Then, η is called an absolute minimizer of the action $\int (\frac{1}{2}|\dot{\gamma}(t)|^2 - V(\gamma(t)) + c) dt$ if

$$\int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\eta}(t)|^2 - V(\eta(t)) + c \right) dt \le \int_{s_1}^{s_2} \left(\frac{1}{2} |\dot{\gamma}(t)|^2 - V(\gamma(t)) + c \right) dt$$

for any $[t_1, t_2] \subset [a, b]$, and $\gamma \in AC([s_1, s_2], \mathbb{R}^n)$ satisfying that $\gamma(s_i) = \eta(t_i)$ for i = 1, 2. Here $AC(I, \mathbb{R}^n)$ represents the set of absolutely continuous curves defined on the interval I.

The following lemma is quite basic and well-known (see [12, 20]).

Lemma 2.1. Let U be an open subset of \mathbb{R}^n . Assume that for some $c \in \mathbb{R}$, $w \in W^{1,\infty}(U)$ satisfies that

$$\frac{1}{2}|Dw|^2 + V(x) \le c \quad \text{for a.e. } x \in U.$$

Then, for any $\eta \in AC([a,b], U)$,

$$I[\eta, (a, b)] = \int_a^b \left(\frac{1}{2}|\dot{\eta}(t)|^2 - V(\eta(t)) - c\right) dt \ge w(\eta(b)) - w(\eta(a)).$$

Next is a differentiability property, an important and new result in the merely continuous setting.

Lemma 2.2. Let U be an open subset of \mathbb{R}^n . Assume that for some $c \in \mathbb{R}$, $w_i \in W^{1,\infty}(U)$ for i = 1, 2 satisfies that

$$\frac{1}{2}|Dw_i|^2 + V(x) \le c \quad \text{for a.e. } x \in U.$$

Assume further that there exists a curve $\eta \in AC([a,b],U)$ such that, for i=1,2,

$$\int_{a}^{b} \left(\frac{1}{2} |\dot{\eta}(t)|^{2} - V(\eta(t)) + c \right) dt = w_{i}(\eta(b)) - w_{i}(\eta(a)).$$

Then, the following properties hold.

(1) For a.e. $t \in [a, b]$,

(2.2)
$$\frac{1}{2}|\dot{\eta}(t)|^2 + V(\eta(t)) = c.$$

(2) If η is differentiable at $t_0 \in (a,b)$, then w_1 and w_2 are differentiable at $x = \eta(t_0)$ and

$$Dw_1(\eta(t_0)) = Dw_2(\eta(t_0)) = \dot{\eta}(t_0).$$

(3) For all $x \in \eta((a,b))$, $w_1 - w_2$ is differentiable at x and $D(w_1 - w_2)(x) = 0.$

Proof. Without loss of generality, we assume that $0 \in (a,b)$ and $\eta(0) = 0$.

We first prove (2.2). In fact, by standard mollification of w_1 and approximations, we have that

$$w_1(\eta(b)) - w_1(\eta(a)) = \int_a^b p_1(t) \cdot \dot{\eta}(t) dt$$
$$= \int_a^b \frac{1}{2} \left(|p_1(t)|^2 + |\dot{\eta}(t)|^2 - |p_1(t) - \dot{\eta}(t)|^2 \right) dt$$

for some $p_1(t) \in \partial w_1(\eta(t))$ for $t \in (a, b)$. Here,

$$\partial w_1(x) = \operatorname{co}(K(x)),$$

where co(K(x)) is the convex hull of the set

 $K(x) = \{ p \in \mathbb{R}^n : \exists \{x_k\} \to x \text{ s.t. } Dw_1(x_k) \text{ exists, and } \{Dw_1(x_k)\} \to p \}.$ Apparently,

$$\frac{1}{2}|p_1(t)|^2 + V(\eta(t)) \le c \quad \text{for all } t \in [a, b].$$

Accordingly, we must have that

$$\frac{1}{2}|p_1(t)|^2 + V(\eta(t)) = c$$
 and $p_1(t) = \dot{\eta}(t)$ for a.e. $t \in [a, b]$.

Hence (2.2) holds. Note that this also implies that η is Lipschitz continuous.

To prove (2) and (3), it suffices to show that for a sequence $\{\lambda_m\} \subset (0, \infty)$ with $\lim_{m\to\infty} \lambda_m = 0$, if

$$\lim_{m \to \infty} \frac{\eta(\lambda_m t)}{\lambda_m} \quad \text{ exists for all } t \in \mathbb{R},$$

then, for all $t \in \mathbb{R}$, $x \in \mathbb{R}^n$, and i = 1, 2,

$$\lim_{m \to \infty} \frac{\eta(\lambda_m t)}{\lambda_m} = qt \quad \text{and} \quad \lim_{m \to \infty} \frac{w_i(\lambda_m x) - w_i(0)}{\lambda_m} = q \cdot x$$

for some q satisfying $\frac{1}{2}|q|^2 + V(0) = c$. We only need to prove this claim for w_1 as the proof for w_2 is the same. Let

$$\bar{\eta}(t) = \lim_{m \to \infty} \frac{\eta(\lambda_m t)}{\lambda_m}.$$

Owing to Lemma 2.1, we have that, for any $a \le t_1 < t_2 \le b$,

$$w_1(\eta(t_2)) - w_1(\eta(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2}|\dot{\eta}(t)|^2 - V(\eta(t)) + c\right) dt.$$

In fact, η is an absolute minimizer of the action $\int (\frac{1}{2}|\dot{\gamma}(t)|^2 - V(\gamma(t)) + c) dt$ (see Definition 3). Then, $\eta(\lambda_m t)/\lambda_m$ is an absolute minimizer of the action

$$\int \left(\frac{1}{2}|\dot{\gamma}(t)|^2 - V(\lambda_m \gamma(t)) + c\right) dt$$

over $[a/\lambda_m, b/\lambda_m]$. By the stability of minimizing curves, $\bar{\eta}$ is an absolute minimizer of the action

$$\int \left(\frac{1}{2}|\dot{\gamma}(t)|^2 - V(0) + c\right) dt$$

over any finite interval of \mathbb{R} . By the Euler-Lagrange equations, $\bar{\eta}(t)$ must be a line passing through the origin 0, that is,

$$\bar{\eta}(t) = q \cdot t$$

for some $q \in \mathbb{R}^n$. For convenience, denote $M = \sqrt{2(c - V(0))}$. Due to (2.2), $|q| \leq M$.

By passing to a subsequence if necessary, we assume that

$$\lim_{m \to \infty} \frac{w_1(\lambda_m x) - w_1(0)}{\lambda_m} = u(x) \quad \text{ for all } x \in \mathbb{R}^n.$$

Then, $u \in W^{1,\infty}(\mathbb{R}^n)$, and u satisfies that

$$\frac{1}{2}|Du(x)|^2 + V(0) \le c \quad \text{ for a.e. } x \in \mathbb{R}^n.$$

Equivalently, $|Du(x)| \leq M$ for a.e. $x \in \mathbb{R}^n$. Thanks to (2.2),

$$w_1(\eta(\lambda_m t)) - w_1(\eta(0)) = \int_0^{\lambda_m t} \left(\frac{1}{2}|\dot{\eta}(s)|^2 - V(\eta(s)) + c\right) ds$$
$$= \int_0^{\lambda_m t} 2\left(c - V(\eta(s))\right) ds.$$

Dividing both sides by λ_m , and sending $m \to \infty$, we derive that

$$u(qt) = tM^2$$
 for all $t \in \mathbb{R}$.

Since $|Du| \leq M$,

$$u(qt) \le M|q|t \le M^2t.$$

Therefore, we deduce that |q| = M, and

$$u(et) = Mt$$
 for all $t \in \mathbb{R}$

for $e = \frac{q}{|q|} = \frac{q}{M}$. Then, the standard tightness argument in [1] leads to

(2.3)
$$u(x) = q \cdot x \quad \text{for all } x \in \mathbb{R}^n.$$

Let us give a proof (2.3) here for completeness. Due to $|Du| \leq M$,

$$|u(x) - Mt| = |u(x) - u(et)| \le M|x - et|$$
 for all $x \in \mathbb{R}^n$, $t \in \mathbb{R}$.

Taking square of both sides leads to

$$u(x)^{2} - 2Mtu(x) + M^{2}t^{2} \le M^{2}(|x|^{2} - 2x \cdot et + t^{2}),$$

which is reduced to

$$|u(x)|^2 - M^2|x|^2 \le 2Mt(u(x) - q \cdot x).$$

Since for fixed $x \in \mathbb{R}^n$, the above holds for all $t \in \mathbb{R}$, we must have

$$u(x) = q \cdot x$$
.

Remark 1. In the setting of the above lemma, for smooth V, it is well known that η is C^2 and both w_1 and w_2 are differentiable along η . See [13] for differentiability results for Lipschitz continuous coefficients. In our setting, the subtle point is that it is not very clear to us whether w_1 and w_2 are differentiable along η . However, $w_1 - w_2$ is indeed differentiable along η with $D(w_1 - w_2) = 0$, which is enough for our purpose. It remains an interesting question to prove or disprove the differentiability of w_i along η . In particular, if a proper uniform C^1 regularity of $s(x) = w_1 - w_2$ near η (i.e., $|s(x) - s(\eta(0))| = o(d(x, \eta))$) can be established, the C^3 assumption of a(x) in [6] might be relaxed to the mere continuity assumption via suitable adjustment of the methods there.

For t > 0, and $p, x, y \in \mathbb{R}^n$, let

$$G_{t,p}(x,y) = \inf_{\substack{\xi \in AC([0,t],\mathbb{R}^n), \\ \xi(0) = x, \ \xi(t) \in y + \mathbb{Z}^n}} \int_0^t \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) - p \cdot \dot{\xi}(s) + \overline{H}(p)\right) ds,$$

and

$$G_p(x,y) = \liminf_{t \to \infty} G_{t,p}(x,y).$$

By Lemma 2.1, $G_{t,p}(x,y) \ge v(y) - v(x)$. Hence,

$$G_p(x,y) \ge v(y) - v(x)$$

for any viscosity solution v of (2.1). In particular, $G_p(x,x) \geq 0$. In addition, it is easy to see that $G_p(x,y)$ is \mathbb{Z}^n -periodic and Lipschitz continuous in x and y.

Now we define the Aubry set associated with $p \in \mathbb{R}^n$ as

$$\mathcal{A}_p = \{ x \in \mathbb{R}^n : G_p(x, x) = 0 \}.$$

Definition 4. Given $p \in \mathbb{R}^n$, a Lipschitz continuous curve $\xi : \mathbb{R} \to \mathbb{R}^n$ is called a global characteristic associated with a viscosity solution v of the cell problem (2.1) if, for all $t_1 < t_2$,

$$v(\xi(t_2)) - v(\xi(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) - p \cdot \dot{\xi}(s) + \overline{H}(p) \right) ds,$$

or equivalently, for $u(x) = p \cdot x + v(x)$ for $x \in \mathbb{R}^n$,

$$u(\xi(t_2)) - u(\xi(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2}|\dot{\xi}(s)|^2 - V(\xi(s)) + \overline{H}(p)\right) ds.$$

By Lemma 2.1, to show that ξ is a global characteristics of v is sufficient to show that there exist $T_m \to +\infty$ and $T_m' \to -\infty$ as $m \to +\infty$ such that

$$v(\xi(T_m)) - v(\xi(T_m')) = \int_{T_m'}^{T_m} \left(\frac{1}{2}|\dot{\xi}(s)|^2 - V(\xi(s)) - p \cdot \dot{\xi}(s) + \overline{H}(p)\right) ds,$$

In addition, owing to Lemma 2.2, if ξ is a global characteristic associated with v solving (2.1) for a given $p \in \mathbb{R}^n$, then

(2.4)
$$\frac{1}{2}|\dot{\xi}(s)|^2 + V(\xi(s)) = \overline{H}(p) \quad \text{for a.e. } t \in \mathbb{R}.$$

Remark 2. By Lemma 2.1, for $p \in \mathbb{R}^n$ and v is a viscosity solution of (2.1), every global characteristic ξ associated with v is an absolute minimizer of the action $\int (\frac{1}{2}|\dot{\gamma}(t)|^2 - V(\gamma(t)) + \overline{H}(p)) dt$. See Definition 3.

Lemma 2.3. For $p \in \mathbb{R}^n$ and v is a viscosity solution of (2.1), suppose that $\{\xi_m\}$ is a sequence of global characteristics associated with v such that

$$\lim_{m \to \infty} \xi_m = \xi \quad locally \ uniformly \ in \ \mathbb{R}.$$

Then, ξ is also a global characteristic of v.

Proof. Fix $t_1 < t_2$. For $u(x) = p \cdot x + v(x)$ for $x \in \mathbb{R}^n$, and $m \in \mathbb{N}$, we have that

$$u(\xi_m(t_2)) - u(\xi_m(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}_m(s)|^2 - V(\xi(s)) + \overline{H}(p) \right) ds.$$

Sending $m \to \infty$, by the lower semicontinuity of the integral, we have that

$$u(\xi(t_2)) - u(\xi(t_1)) \ge \liminf_{m \to \infty} \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}_m(s)|^2 - V(\xi_m(s)) + \overline{H}(p) \right) ds$$

$$\ge \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) + \overline{H}(p) \right) ds.$$

Combining this with Lemma 2.1, we get the desired result.

Definition 5. Given $p \in \mathbb{R}^n$, a Lipschitz continuous curve $\xi : \mathbb{R} \to \mathbb{R}^n$ is called a universal global characteristic associated with p if it is a global characteristic for every viscosity solution v of the cell problem (2.1).

We denote by

 $\mathcal{U}_p = the \ collection \ of \ all \ universal \ characteristics \ associated \ with \ p.$

Then for every $\xi \in \mathcal{U}_p$,

(2.5)
$$\frac{1}{2}|\dot{\xi}(s)|^2 + V(\xi(s)) = \overline{H}(p) \quad \text{for a.e. } t \in \mathbb{R}.$$

Also, owing to Lemma 2.3, the set \mathcal{U}_p is closed for limit of orbits.

Lemma 2.4. Given $p \in \mathbb{R}^n$, every viscosity solution v of (2.1) has a global characteristic.

Proof. Let $u(x) = p \cdot x + v(x)$ for $x \in \mathbb{R}^n$. Then, $w(x,t) = u(x) - \overline{H}(p)t$ for $(x,t) \in \mathbb{R}^n \times [0,\infty)$ solves the following Cauchy problem

$$\begin{cases} w_t + \frac{1}{2}|Dw|^2 + V(x) = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ w(x, 0) = u(x) & \text{on } \mathbb{R}^n. \end{cases}$$

By the optimal control formula, for each $k \in \mathbb{N}$,

$$w(0,k) = \inf_{\substack{\gamma \in \mathrm{AC}([0,k],\mathbb{R}^n),\\ \gamma(k) = 0}} \left\{ \int_0^k \left(\frac{1}{2} |\dot{\gamma}(s)|^2 - V(\gamma(s)) \right) \, ds + u(\gamma(0)) \right\}.$$

There exists $\xi_k \in AC([-k,0],\mathbb{R}^n)$ with $\xi_k(0) = 0$ such that

$$w(0,k) = \int_{-k}^{0} \left(\frac{1}{2} |\dot{\xi}_k(s)|^2 - V(\xi_k(s)) \right) ds + u(\xi_k(-k)).$$

This is equivalent to

$$u(\xi_k(0)) - u(\xi_k(-k)) = \int_{-k}^0 \left(\frac{1}{2}|\dot{\xi}_k(s)|^2 - V(\xi_k(s)) + \overline{H}(p)\right) ds$$

Owing to Lemma 2.2,

$$\frac{1}{2}|\dot{\xi}_k(s)|^2 + V(\xi_k(s)) = \overline{H}(p)$$
 for a.e. $t \in [-k, 0]$.

By sending $k \to \infty$ and passing to a subsequence if necessary, we have that $\xi_k \to \xi$ locally uniformly on $(-\infty, 0]$, and, similar to Lemma 2.3, $\xi: (-\infty, 0] \to \mathbb{R}^n$ with $\xi(0) = 0$ is a backward characteristic associated with v. More precisely, for any $t_1 < t_2 \le 0$,

$$u(\xi(t_2)) - u(\xi(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) + \overline{H}(p) \right) ds.$$

Again, due to Lemma 2.2, for a.e. $t \leq 0$,

$$\frac{1}{2}|\dot{\xi}(t)|^2 + V(\xi(t)) = \overline{H}(p).$$

Next, we create a global characteristic from this backward characteristic ξ . For $m \in \mathbb{N}$, let $\gamma_m : (-\infty, m] \to \mathbb{R}^n$ be such that

$$\gamma_m(t) = \xi(t-m) + k_m$$
 for all $t \le m$,

where $k_m \in \mathbb{Z}^n$ is chosen so that $\gamma_m(0) = \xi(-m) + k_m \in [0,1]^n$. Again, by passing to a subsequence if necessary, $\gamma_m \to \gamma$ locally uniformly on \mathbb{R} . We obtain that γ is a global characteristic associated with v.

Lemma 2.5. For any $p \in \mathbb{R}^n$,

$$\mathcal{A}_{p}\neq\emptyset$$
.

Proof. Let v be a viscosity solution of (2.1) and $\xi : \mathbb{R} \to \mathbb{R}^n$ be a global characteristic associated with v. By projecting ξ to \mathbb{T}^n and a suitable translation in time, we may find sequences $\{t_m\} \to +\infty$ and $\{x_m\} \subset [0,1]^n$ such that

$$\lim_{m \to +\infty} (t_{m+1} - t_m) = +\infty,$$

$$\xi(t_m) = x_m + k_m \quad \text{for some } k_m \in \mathbb{Z}^n,$$

and

$$\lim_{m \to \infty} x_m = x_0 \in [0, 1]^n,$$

and

$$v(x_{m+1}) - v(x_m) = \int_{t_m}^{t_{m+1}} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) - p \cdot \dot{\xi}(s) + \overline{H}(p) \right) ds.$$

Then,

$$G_{t_{m+1}-t_m,p}(x_m, x_{m+1}) = v(x_{m+1}) - v(x_m).$$

This implies

$$G_p(x_0, x_0) \le \liminf_{m \to \infty} G_{t_{m+1} - t_m, p}(x_m, y_m) = 0.$$

Thus, $x_0 \in \mathcal{A}_p$.

Lemma 2.6. For any $x \in \mathcal{A}_p$, there exists $\xi \in \mathcal{U}_p$ such that $\xi(0) = x$. In particular, this implies that $\mathcal{U}_p \neq \emptyset$.

Proof. According to the definition of \mathcal{A}_p , there exist $\{t_m\} \to \infty$ and a sequence of curves $\gamma_m : [0, t_m] \to \mathbb{R}^n$ such that $\gamma_m(0) = x$, and $\gamma_m(t_m) = x + k_m$ for some $k_m \in \mathbb{Z}^n$, and

$$\lim_{m \to \infty} \int_0^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p \cdot \dot{\gamma}_m(s) + \overline{H}(p) \right) ds = 0.$$

Given a viscosity solution v of (2.1), owing to Lemma 2.1, for any fixed L > 0,

$$v(\gamma_m(L)) - v(x)) \le \int_0^L \left(\frac{1}{2}|\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p \cdot \dot{\gamma}_m(s) + \overline{H}(p)\right) ds$$

and

$$v(\gamma_m(t_m)) - v(\gamma_m(L))) \le \int_L^{t_m} \left(\frac{1}{2}|\dot{\gamma}_m|^2 - V(\gamma_m) - p \cdot \dot{\gamma}_m + \overline{H}(p)\right) ds.$$

Together with $0 = v(\gamma_m(L)) - v(x) + v(\gamma_m(t_m)) - v(\gamma_m(L))$, we have that

$$\lim_{m \to \infty} \left(\int_0^L \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p \cdot \dot{\gamma}_m(s) + \overline{H}(p) \right) ds - (v(\gamma_m(L)) - v(x)) \right) = 0.$$

Similarly,

$$\lim_{m \to \infty} \left(\int_{t_m - L}^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p \cdot \dot{\gamma}_m(s) + \overline{H}(p) \right) ds - \left(v(\gamma_m(t_m)) - v(\gamma(t_m - L)) \right) \right) = 0.$$

Define $\xi_m: [-t_m/2, t_m/2] \to \mathbb{R}^n$ as

$$\xi_m(t) = \begin{cases} \gamma_m(t) & \text{for } t \in \left[0, \frac{t_m}{2}\right], \\ \gamma_m(t_m + t) - k_m & \text{for } t \in \left[-\frac{t_m}{2}, 0\right] \end{cases}$$

Clearly, for any fixed L > 0, $\{\|\xi_m\|_{H^1((-L,L))}\}$ is uniformly bounded. Up to a subsequence if necessary, we may assume that

$$\lim_{m \to \infty} \xi_m(t) = \xi(t) \quad \text{locally uniformly in } \mathbb{R}.$$

for $\xi \in AC(\mathbb{R}, \mathbb{R}^n)$.

Then, using Lemma 2.1 and the lower semicontinuity of the integral (similar to the proof of Lemma 2.3), we see that, for any L > 0,

$$\int_{-L}^{L} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) - p \cdot \dot{\xi}(s) + \overline{H}(p) \right) ds = v(\xi(L)) - v(\xi(-L)).$$

Hence ξ is a universal global characteristic associated with p.

For smooth V, two different orbits in the same Aubry set cannot intersect and two different absolute minimizers of the same action cannot intersect twice. However, both situations could happen with merely continuous V. Consequently, the structure of orbits on \mathcal{U}_p might be very complicated. Below we provide two procedures to join different pieces of two global characteristics, which will be used later to select nice minimizing orbits and then simplify the topology of interacting curves.

Remark 3. Given $p \in \mathbb{R}^n$ and v is a viscosity solution of (2.1), suppose that ξ_1 and ξ_2 are two global characteristics associated with v satisfying that for some $t_1, t_2 \in \mathbb{R}$,

$$\xi_1(t_1) = \xi_2(t_2).$$

Define

$$\xi_3(t) = \begin{cases} \xi_1(t) & \text{for } t \le t_1 \\ \xi_2(t - t_1 + t_2) & \text{for } t \ge t_1. \end{cases}$$

See Figure 2.1.

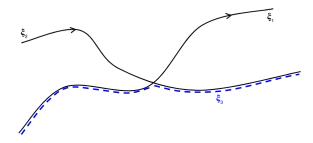


FIGURE 2.1. Formation of the curve ξ_3

Then, it is easy to see that ξ_3 is also a global characteristic associated with v. In fact, for $a \le t_1 \le b$,

$$v(b) - v(t_1) = \int_{t_1}^{b} \left(\frac{1}{2} |\dot{\xi}_3(s)|^2 - V(\xi_3(s)) - p \cdot \dot{\xi}_3(s) + \overline{H}(p) \right) ds$$

and

$$v(t_1) - v(a) = \int_a^{t_1} \left(\frac{1}{2} |\dot{\xi}_3(s)|^2 - V(\xi_3(s)) - p \cdot \dot{\xi}_3(s) + \overline{H}(p) \right) ds.$$

Thus

$$v(b) - v(a) = v(b) - v(t_1) + v(t_1) - v(a)$$

$$= \int_a^b \left(\frac{1}{2}|\dot{\xi}_3(s)|^2 - V(\xi_3(s)) - p \cdot \dot{\xi}_3(s) + \overline{H}(p)\right) ds.$$

Remark 4 (Crossing of two universal global characteristics). Suppose that $p, p' \in F_c$, and ξ and $\tilde{\xi}$ are orbits in \mathcal{U}_p and $\mathcal{U}_{p'}$, respectively. Assume that there exist $t_1, t_2, t'_1, t'_2 \in \mathbb{R}$ such that, for i = 1, 2,

$$P_i = \xi(t_i) = \tilde{\xi}(t_i').$$

Now we present how to construct new orbits on \mathcal{U}_p and $\mathcal{U}_{p'}$ by joining different pieces of ξ and $\tilde{\xi}$. Without loss of generality, we assume that $t_1 < t_2$. There are two cases.

Case 1. $t'_1 < t'_2$. Define

$$\xi_2(t) = \begin{cases} \xi(t) & \text{for } t \le t_1, \\ \tilde{\xi}(t+t_1'-t_1) & \text{for } t_1 \le t \le t_1+t_2'-t_1', \\ \xi(t+t_2-t_1-t_2'+t_1') & \text{for } t_1+t_2'-t_1' \le t. \end{cases}$$

See Figure 2.2.

Case 2. $t'_1 > t'_2$. Define

$$\xi_3(t) = \begin{cases} \xi(t) & \text{for } t \le t_1, \\ \tilde{\xi}(t_1 + t_1' - t) & \text{for } t_1 \le t \le t_1 + t_1' - t_2', \\ \xi(t + t_2 - t_1 - t_1' + t_2') & \text{for } t_1 + t_1' - t_2' \le t. \end{cases}$$

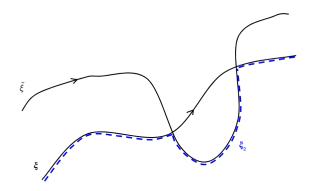


Figure 2.2. Combining two universal global characteristics

Since both ξ and $\tilde{\xi}$ are absolute action minimizing curves connecting P_1 and P_2 , we have that

$$\int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) + c \right) ds = \left| \int_{t_1'}^{t_2'} \left(\frac{1}{2} |\dot{\tilde{\xi}}(s)|^2 - V(\tilde{\xi}(s)) + c \right) ds \right|.$$

Accordingly, we have the following conclusion.

Corollary 2.7. The above ξ_2 or ξ_3 belongs to \mathcal{U}_p .

Definition 6. We say that ξ_2 or ξ_3 constructed in Remark 4 the adjustment of ξ with respect to $\tilde{\xi}$ between t_1 and t_2 .

Lemma 2.8. Let $p \in \mathbb{R}^n$, and v be a viscosity solution of (2.1). Let ξ : $\mathbb{R} \to \mathbb{R}^n$ be a global characteristic associated with v. Assume that there exists $\{t_m\} \to \pm \infty$ such that $\xi(t_m)/t_m$ converges as $m \to \infty$. Then,

$$\lim_{m \to \infty} \frac{\xi(t_m)}{t_m} \in \partial \overline{H}(p).$$

Proof. Suppose that

$$\lim_{m \to \infty} \frac{\xi(t_m)}{t_m} = q$$

for some $q \in \mathbb{R}^n$. It suffices to show that

$$(2.6) \overline{H}(p') > \overline{H}(p) + q \cdot (p' - p) \text{for all } p' \in \mathbb{R}^n.$$

Let v' be a viscosity solution of (2.1) with p = p'. By Lemma 2.1,

$$p' \cdot \xi(t_m) + v'(\xi(t_m)) - (p' \cdot \xi(0) + v'(\xi(0))) \le \int_0^{t_m} \left(\frac{1}{2} |\dot{\xi}|^2 - V(\xi) + \overline{H}(p')\right) ds.$$

Meanwhile,

$$p \cdot \xi(t_m) + v(\xi(t_m)) - (p \cdot \xi(0) + v(\xi(0))) = \int_0^{t_m} \left(\frac{1}{2}|\dot{\xi}|^2 - V(\xi) + \overline{H}(p)\right) ds.$$

Taking the difference of the above two equations, dividing both sides by t_m , and sending $m \to \infty$, we derive (2.6).

Definition 7. For $p \in \mathbb{R}^n$ and v is a viscosity solution of (2.1), a global characteristic $\xi : \mathbb{R} \to \mathbb{R}^n$ associated with v is called periodic if there exist T > 0 and $q \in \mathbb{Z}^n$ such that

$$\xi(t+T) - \xi(t) = q$$
 for all $t \in \mathbb{R}$.

In this case, q/T is called the rotation vector of ξ .

Owing to Lemma 2.8, the rotation vector

$$\frac{q}{T} \in \partial \overline{H}(p).$$

Also, it is clear that every periodic global characteristic associated with some v must be a universal global characteristic.

Corollary 2.9. For $p \in \mathbb{R}^n$ and v is a viscosity solution of (2.1), let ξ : $\mathbb{R} \to \mathbb{R}^n$ be a global characteristic associated with v. If there exist $t_1 < t_2$ such that

$$\xi(t_2) - \xi(t_1) = q \in \mathbb{Z}^n,$$

then

$$\frac{q}{t_2 - t_1} \in \partial \overline{H}(p).$$

Proof. Let $T = t_2 - t_1$. Define $\tilde{\xi} : \mathbb{R} \to \mathbb{R}^n$ as

$$\tilde{\xi}(t) = \xi(t - kT + t_1) + kq$$
 if $t \in [kT, (k+1)T]$,

for all $k \in \mathbb{Z}$. Then $\tilde{\xi}(0) = \xi(t_1)$ and $\tilde{\xi}(t+T) = \tilde{\xi}(t) + q$ for all $t \in \mathbb{R}$. Since

$$0 = v(\xi(t_2)) - v(\xi(t_1)) = \int_{t_1}^{t_2} \left(\frac{1}{2} |\dot{\xi}(t)|^2 - V(\xi(t)) - p \cdot \dot{\xi}(t) + \overline{H}(p) \right) dt,$$

for any $m \in \mathbb{N}$, we have that

$$\begin{split} I(\tilde{\xi}, [-mT, mT]) &= \sum_{k=-m}^{m-1} I(\tilde{\xi}, [kT, (k+1)T]) \\ &= \sum_{k=-m}^{m-1} I(\tilde{\xi}, [0, T]) = \sum_{k=-m}^{m-1} I(\xi, [t_1, t_2]) \\ &= 0 = v(\tilde{\xi}(mT)) - v(\tilde{\xi}(-mT)). \end{split}$$

Here.

$$I(\eta, [a, b]) = \int_a^b \left(\frac{1}{2}|\dot{\eta}(t)|^2 - V(\eta(t)) - p \cdot \dot{\eta}(t) + \overline{H}(p)\right) dt.$$

Then, $\tilde{\xi}$ is a periodic global characteristic associated with v. In fact, $\tilde{\xi} \in \mathcal{U}_p$. The conclusion follows from Lemma 2.8.

Lemma 2.10. When n = 2, for every $q \in \mathbb{Z}^2$ and $c > \max_{\mathbb{T}^n} V$, there exists $p_q \in F_c$ such that \mathcal{U}_{p_q} has a periodic orbit ξ such that, for some T > 0,

$$\xi(t+T) - \xi(t) = q$$
 for all $t \in \mathbb{R}$.

Proof. This result is well-known for smooth V. Similar to the proof of the stability of global characteristics in Lemma 2.3, the merely continuous version can be established by approximating V with smooth periodic functions under the maximum norm.

Lemma 2.11. Suppose that there exist $p_0, p_1 \in \mathbb{R}^n$ and $\lambda \in (0, 1)$ such that, for $p_{\lambda} = \lambda p_0 + (1 - \lambda)p_1$,

$$\overline{H}(p_{\lambda}) = \lambda \overline{H}(p_0) + (1 - \lambda) \overline{H}(p_1).$$

Then,

$$\overline{H}(p_{\lambda}) = \overline{H}(p_0) = \overline{H}(p_1), \qquad \mathcal{A}_{p_{\lambda}} \subset \mathcal{A}_{p_0} \cap \mathcal{A}_{p_1},$$

and

$$\mathcal{U}_{p_{\lambda}} \cap \mathcal{U}_{p_0} \cap \mathcal{U}_{p_1} \neq \emptyset.$$

Proof. We divide the proof into two steps.

Step 1. For $x \in \mathcal{A}_{p_{\lambda}}$, there exist $\{t_m\} \to \infty$ and a sequence of curves $\gamma_m : [0, t_m] \to \mathbb{R}^n$ such that $\gamma_m(0) = x, \gamma_m(t_m) \in x + \mathbb{Z}^n$, and

$$\lim_{m \to \infty} \int_0^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p_\lambda \cdot \dot{\gamma}_m(s) + \overline{H}(p_\lambda) \right) ds = 0.$$

Let

$$A_m = \lambda \int_0^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p_0 \cdot \dot{\gamma}_m(s) + \overline{H}(p_0) \right) ds,$$

and

$$B_m = (1 - \lambda) \int_0^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p_1 \cdot \dot{\gamma}_m(s) + \overline{H}(p_1) \right) ds.$$

Then

$$\int_0^{t_m} \left(\frac{1}{2} |\dot{\gamma}_m(s)|^2 - V(\gamma_m(s)) - p_\lambda \cdot \dot{\gamma}_m(s) + \overline{H}(p_\lambda) \right) ds = A_m + B_m.$$

By Lemma 2.1, $A_m, B_m \ge 0$. Hence,

$$\lim_{m \to \infty} A_m = \lim_{m \to \infty} B_m = 0.$$

Meanwhile, by the definition of $G_p(x,x)$, it is obvious that

$$0 \le \lambda G_{p_0}(x, x) \le \lim_{m \to \infty} A_m$$
 and $0 \le (1 - \lambda)G_{p_1}(x, x) \le \lim_{m \to \infty} B_m$.

Accordingly, $G_{p_0}(x,x) = G_{p_1}(x,x) = 0$. Then $x \in \mathcal{A}_{p_0} \cap \mathcal{A}_{p_1}$. This implies that

$$\mathcal{A}_{p_{\lambda}} \subset \mathcal{A}_{p_0} \cap \mathcal{A}_{p_1}$$
.

In addition, as in the proof of Lemma 2.6, a suitable reparametrization of $\{\gamma_m\}$ gives a sequence of curves that converges to a common orbit in $\mathcal{U}_{p_\lambda} \cap \mathcal{U}_{p_0} \cap \mathcal{U}_{p_1}$.

Step 2. Choose an orbit $\xi \in \mathcal{U}_{p_0} \cap \mathcal{U}_{p_1}$. Then, by (2.5),

$$\frac{1}{2}|\dot{\xi}(t)|^2 + V(\xi(t)) = \overline{H}(p_0) \text{ for a.e } t \in \mathbb{R},$$

and

$$\frac{1}{2}|\dot{\xi}(t)|^2 + V(\xi(t)) = \overline{H}(p_1) \quad \text{for a.e } t \in \mathbb{R}.$$

Therefore, $\overline{H}(p_0) = \overline{H}(p_1)$.

3. Proof of Theorem 1.2

We are in the setting of Theorem 1.2 in this section. Part (1) follows immediately from Lemma 2.11. We now prove part (2). Throughout this section, we assume that $c > \max_{\mathbb{T}^2} V$. Hence any orbit on \mathcal{U}_p for $p \in F_c$ does not intersect with itself since $u = p \cdot x + v$ is strictly increasing along any orbit.

We argue by contradiction. Suppose that F_c contains a line segment of an irrational slope. Assume that p_0 and p_1 are two points in the interior of the line segment. According to Lemma 2.11,

$$\mathcal{A}_{p_0} = \mathcal{A}_{p_1},$$

and

$$p_0 - p_1$$
 is an irrational vector.

Then, the outward unit normal vector \vec{n} is also irrational, and

$$\partial \overline{H}(p_0) = \partial \overline{H}(p_1) = \{\lambda \vec{n} : \lambda \in [\alpha, \beta]\}$$

for two positive numbers $0 < \alpha < \beta$. Without loss of generality, we assume that

$$(3.1) \vec{n} \cdot (1,0) > 0.$$

Let v_0 and v_1 be viscosity solutions to (2.1) corresponding to $p = p_0$ and $p = p_1$, respectively. Write

$$\mathcal{U}=\mathcal{U}_{p_0}\cap\mathcal{U}_{p_1},$$

and

$$S = \bigcup_{\xi \in \mathcal{U}} \xi(\mathbb{R}) \subset \mathbb{R}^2.$$

Owing to Lemma 2.11, $\mathcal{U} \neq \emptyset$. Also, by Lemma 2.2, $u_0 - u_1$ is differentiable at $x \in S$, and

$$(3.2) D(u_0 - u_1)(x) = 0 \text{for } x \in S.$$

Here,
$$u_0(x) = p_0 \cdot x + v_0(x)$$
, and $u_1(x) = p_1 \cdot x + v_1(x)$ for $x \in \mathbb{R}^2$.

Thanks to Lemma 2.10, we can choose $p' \in F_c$ such that $\mathcal{U}_{p'}$ contains a periodic orbit η such that for some T > 0

(3.3)
$$\eta(t+T) - \eta(t) = (0,1) = e_2$$
 for all $t \in \mathbb{R}$.

For $e_1 = (1, 0)$, denote

(3.4)
$$\Lambda = \max\{|e_1 \cdot (x - y)| : x, y \in \eta(\mathbb{R})\}.$$

Choose a positive integer $J > \Lambda + 1$ and for $k \in \mathbb{Z}$, denote

$$\eta_k = \eta + k(J, 0).$$

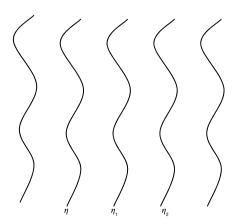


FIGURE 3.1. Family of $\{\eta_k\}_{k\in\mathbb{Z}}$

Clearly, these curves are mutually disjoint. See Figure 3.1.

Throughout this section, ξ represents a given orbit on \mathcal{U} . Then, owing to Lemma 2.8,

(3.5)
$$\lim_{t \to \infty} \frac{\xi(t)}{|\xi(t)|} = \lim_{t \to -\infty} \frac{-\xi(t)}{|\xi(t)|} = \vec{n}.$$

Then, ξ intersects each η_k .

For $k \in \mathbb{Z}$, write

$$t_{k,+} = \max\{t \in \mathbb{R} : \xi(t) \in \eta_k(\mathbb{R})\}, \qquad t_{k,-} = \min\{t \in \mathbb{R} : \xi(t) \in \eta_k(\mathbb{R})\}.$$

Owing to (3.5), both $t_{k,+}$ and $t_{k,-}$ are finite. See Figure 3.2.

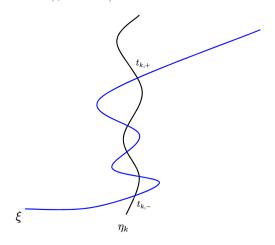


FIGURE 3.2. Intersections of η_k and ξ

Lemma 3.1. For $k \in \mathbb{Z}$, assume that

$$\xi(t_{k,+}) = \eta_k(\theta_+)$$
 and $\xi(t_{k,-}) = \eta_k(\theta_-)$

for $\theta_-, \theta_+ \in \mathbb{R}$. Then,

(1)

$$\xi(\mathbb{R}) \cap \eta_k(\mathbb{R}) \subset \{\eta_k(t) : \min\{\theta_+, \theta_-\} \le t \le \max\{\theta_+, \theta_-\}\};$$

(2)

$$|\theta_+ - \theta_-| < T$$
.

Here, T is the constant defined in (3.3).

Proof. We first prove (1). Were the conclusion of (1) not true, there would exist $t_0 \in (t_{k,-}, t_{k,+})$ such that

$$\xi(t_0) = \eta_k(\theta)$$

for some $\theta < \min\{\theta_+, \theta_-\}$ or $\theta > \max\{\theta_+, \theta_-\}$. Without loss of generality, we assume that $\theta > \theta_+ > \theta_-$. Then,

$$\underbrace{\int_{\theta_{-}}^{\theta_{+}} \left(\frac{1}{2} |\dot{\eta}_{k}(s)|^{2} - V(\eta_{k}(s)) + c\right) \, ds}_{A} = \underbrace{\int_{t_{k,-}}^{t_{k,+}} \left(\frac{1}{2} |\dot{\xi}(s)|^{2} - V(\xi(s)) + c\right) \, ds}_{B}$$

as both η_k and ξ are absolute minimizers of the action connecting $\eta_k(\theta_-)$ and $\eta_k(\theta_+)$. Also,

$$\underbrace{\int_{\theta_{-}}^{\theta} \left(\frac{1}{2} |\dot{\eta}_{k}(s)|^{2} - V(\eta_{k}(s)) + c\right) ds}_{C} = \underbrace{\int_{t_{k,-}}^{t_{0}} \left(\frac{1}{2} |\dot{\xi}(s)|^{2} - V(\xi(s)) + c\right) ds}_{D}$$

since both η_k and ξ are absolute minimizers of the action connecting $\eta_k(\theta_-)$ and $\eta_k(\theta)$. On the other hand, it is obvious that

$$B > D$$
 and $A < C$.

This is a contradiction.

Next we prove (2). Again we argue by contradiction. Without loss of generality, assume that $\theta_+ - \theta_- \ge T$. See Figure 3.3.

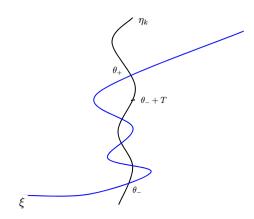


FIGURE 3.3. The situation where $\theta_+ - \theta_- \ge T$

Let ξ_1 be the adjustment of ξ with respect to η_k between $t_{k,-}$ and $t_{k,+}$ (see Definition 6). Then, owing to (1), $\xi_1(\mathbb{R}) \cap \eta_k(\mathbb{R}) = \eta_k([\theta_-, \theta_+])$. In particular, ξ_1 would contain two points $A = \eta_k(\theta_-)$ and $B = \eta_k(\theta_- + T) = A + (0, 1)$. Owing to Corollary 2.9, \vec{n} , the normal vector of F_c at p_0 is parallel to (0, 1), which contradicts the assumption that \vec{n} is irrational.

Remark 5. (1) in the the above lemma actually says that the intersection parameters are monotonic. Precisely speaking, for $t_1 < t_2 < t_3$, if

$$\xi(t_i) = \eta_k(\theta_i)$$
 for $i = 1, 2, 3,$

then we have either $\theta_1 < \theta_2 < \theta_3$ or $\theta_1 > \theta_2 > \theta_3$.

Write

$$L_A = \int_0^T \left(\frac{1}{2}|\dot{\eta}(s)|^2 - V(\eta(s)) + c\right) ds,$$

which is exactly the action of one cycle of η .

Owing to (3.1) and (3.5), ξ will intersect η_k before it intersects η_{k+1} . The following lemma says that ξ will not intersect η_k again after it intersects η_{k+1} for large enough J.

Lemma 3.2. Assume that

$$J > \max \left\{ \Lambda + 1, \ \frac{L_A}{\sqrt{c - \max_{\mathbb{T}^2} V}} + \Lambda \right\}.$$

Then,

$$t_{k,+} < t_{k+1,-}$$
 for all $k \in \mathbb{Z}$.

Recall that $\Lambda = \max\{|e_1 \cdot (x - y)| : x, y \in \eta(\mathbb{R})\}.$

Proof. If the conclusion of the lemma were false, then we would have

$$t_{k+1,-} \in (t_{k,-}, t_{k,+}).$$

See Figure 3.4.

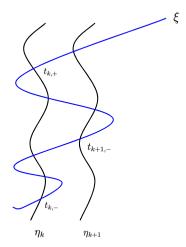


FIGURE 3.4. The situation where $t_{k+1,-} \in (t_{k,-}, t_{k,+})$

Note first that

$$2(J-\Lambda) \le |\xi(t_{k+1,-}) - \xi(t_{k,-})| + |\xi(t_{k+1,-}) - \xi(t_{k,+})| \le \int_{t_k}^{t_{k,+}} |\dot{\xi}(s)| \, ds.$$

Meanwhile,

$$\int_{t_{k,-}}^{t_{k,+}} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) + c \right) ds \ge \sqrt{2} \int_{t_{k,-}}^{t_{k,+}} \sqrt{c - V(\xi(s))} |\dot{\xi}(s)| ds$$

$$\ge \sqrt{c - \max_{\mathbb{T}^2} V} \int_{t_{k,-}}^{t_{k,+}} |\dot{\xi}(s)| ds.$$

Also,

$$\int_{t_{k,-}}^{t_{k,+}} \left(\frac{1}{2} |\dot{\xi}(s)|^2 - V(\xi(s)) + c \right) ds$$

$$= \left| \int_{\theta_{k,-}}^{\theta_{k,+}} \left(\frac{1}{2} |\dot{\eta}_k(s)|^2 - V(\eta_k(s)) + c \right) ds \right| \le L_A.$$

Here, $\xi(t_{k,+}) = \eta_k(\theta_{k,+})$, and $\xi(t_{k,-}) = \eta_k(\theta_{k,-})$. The inequality in the above is due to (2) in Lemma 3.1. The absolute sign $|\cdot|$ is added because $\theta_{k,-}$ might be bigger than $\theta_{k,+}$. Therefore,

$$2(J - \Lambda) \le \frac{L_A}{\sqrt{c - \max_{\mathbb{T}^2} V}},$$

which contradicts the choice of J.

For two orbits ξ_1 and ξ_2 on \mathcal{U} , we define the distance

$$d(\xi_1, \xi_2) = \sum_{k \in \mathbb{Z}} \frac{\arctan(d(\xi_1, \xi_2, k))}{|k|^2 + 1}.$$

Here.

$$d(\xi_1, \xi_2, k) = \min\{|\theta_1 - \theta_2| : \theta_1 \in A_{1,k}, \ \theta_2 \in A_{2,k}\},\$$

where, for i = 1, 2,

$$A_{i,k} = \{ \theta \in \mathbb{R} : \eta_k(\theta) \in \xi_i(\mathbb{R}) \}.$$

Clearly, the distance function $d(\cdot, \cdot, k)$ (hence $d(\cdot, \cdot)$) is lower semicontinuous with respect to orbits on \mathcal{U} , i.e., if $\xi_{i,n} \to \xi_i$ locally uniformly for i = 1, 2, then

$$\liminf_{n \to \infty} d(\xi_{1,n}, \xi_{2,n}) \ge d(\xi_1, \xi_2).$$

Define

$$(3.6) \quad \mathcal{I} = \{\theta \in \mathbb{R} : \text{there exists an orbit } \xi \in \mathcal{U} \text{ such that } \eta(\theta) \in \xi(\mathbb{R})\}.$$

Due to the T-periodicity of η and the fact that the set \mathcal{U} is closed under limits of orbits, \mathcal{I} is a T-periodic closed set.

To finish the proof, the next step is to show that for $u_0(x) = p_0 \cdot x + v_0(x)$, and $u_1(x) = p_1 \cdot x + v_1(x)$ for $x \in \mathbb{R}^2$,

(3.7)
$$u_0 - u_1$$
 is constant on \mathcal{I} ,

which will lead to a contradiction since $\lim_{\theta\to\infty} |u_0(\eta(\theta)) - u_1(\eta(\theta))| = \infty$. Suppose that

$$\mathbb{R}\backslash\mathcal{I}=\bigcup_{i=1}^{\infty}(a_i,b_i),$$

where $\{(a_i, b_i)\}_{i \geq 1}$ are disjoint open intervals. Obviously, $(a_i, b_i) \subseteq (a_i, a_i + T)$ for each $i \in \mathbb{N}$ since $\xi \in \mathcal{U} \Rightarrow \xi + (0, 1) \in \mathcal{U}$.

Lemma 3.3. For all $j \in \mathbb{N}$,

(3.8)
$$u_0(\eta(a_j)) - u_1(\eta(a_j)) = u_0(\eta(b_j)) - u_1(\eta(b_j)).$$

Proof. We only need to prove the claim for j=1. Owing to the lower semicontinuity of the distance function, we may choose ξ_1 and ξ_2 in \mathcal{U} such that

- (i) $\xi_1(0) = \eta(a_1)$ and $\xi_2(0) = \eta(b_1)$.
- (ii) $d(\xi_1, \xi_2)$ attains the minimum value among all curves in \mathcal{U} satisfying (i).

To simplify the associated topology between curves, we adjust ξ_1 and ξ_2 with respect to each η_k between $t_{i,k,-}$ and $t_{i,k,+}$ for i=1,2 respectively (see Definition 6). Here

$$t_{i,k,+} = \max\{t \in \mathbb{R} : \xi_i(t) \in \eta_k(\mathbb{R})\}, \qquad t_{i,k,-} = \min\{t \in \mathbb{R} : \xi_i(t) \in \eta_k(\mathbb{R})\}.$$

By Lemma 3.2, for i = 1, 2, the time intervals $(t_{i,k,-}, t_{i,k,+})$ are mutually disjoint, i.e.,

$$\dots < t_{i,-1,-} \le t_{i,-1,+} < t_{i,0,-} \le t_{i,0,+} < t_{i,1,-} \le t_{i,1,+} < t_{i,2,-} \le t_{i,2,+} < \dots$$

Hence the adjustments are well defined.

In addition, thanks to (1) of Lemma 3.1, the distance between two adjusted orbits is not greater than $d(\xi_1, \xi_2)$. Thus two adjusted orbits also satisfy (i)-(ii) above. By abuse of notations, we still use ξ_1 and ξ_2 to represent corresponding adjusted orbits. Accordingly, we may assume that for each $k \in \mathbb{Z}$, and i = 1, 2,

(3.9)
$$\xi_i(\mathbb{R}) \cap \eta_k(\mathbb{R}) = \xi_i([t_{i,k,-}, t_{i,k,+}]) = \eta_k([\theta_{i,k,-}, \theta_{i,k,+}]).$$

Here $\xi_i(t_{i,k,+}) = \eta_k(\theta_{i,k,+})$ and $\xi(t_{i,k,-}) = \eta_k(\theta_{i,k,-})$. It could happen that $\theta_{i,k,+} < \theta_{i,k,-}$. In terms of topology, the above basically plays the role like that ξ_i and η_k only intersect once for smooth V. We consider two cases.

Case 1. $\xi_1(\mathbb{R}) \cap \xi_2(\mathbb{R}) \neq \emptyset$. Assume that $\xi_1(t_1) = \xi_2(t_2)$ for some $t_1, t_2 \in \mathbb{R}$. See Figure 3.5.

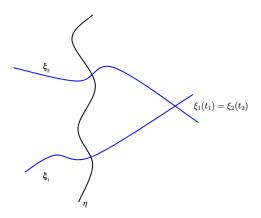


FIGURE 3.5. The situation where $\xi_1(\mathbb{R}) \cap \xi_2(\mathbb{R}) \neq \emptyset$

Then,

$$u_0(\xi_1(t_1)) - u_0(\xi_1(0)) = u_1(\xi_1(t_1)) - u_1(\xi_1(0))$$
$$= \int_0^{t_1} \left(\frac{1}{2}|\dot{\xi}_1(s)|^2 - V(\xi_1(s)) + c\right) ds,$$

and

$$u_0(\xi_2(t_2)) - u_0(\xi_2(0)) = u_1(\xi_2(t_2)) - u_1(\xi_2(0))$$
$$= \int_0^{t_2} \left(\frac{1}{2}|\dot{\xi}_2(s)|^2 - V(\xi_2(s)) + c\right) ds.$$

Taking the difference of the two equations above leads to the claim.

Case 2. $\xi_1(\mathbb{R}) \cap \xi_2(\mathbb{R}) = \emptyset$. For each $k \in \mathbb{Z}$, let

$$d_1(k) = \max\{\theta : \eta_k(\theta) \in \xi_1(\mathbb{R}) \cap \eta_k(\mathbb{R})\},\$$

and

$$d_2(k) = \min\{\theta : \eta_k(\theta) \in \xi_2(\mathbb{R}) \cap \eta_k(\mathbb{R})\}.$$

Owing to two dimensional topology and (3.9), we have that

$$d_1(0) = a_1$$
 and $d_2(0) = b_1$,

and

$$d_1(k) < d_2(k)$$
 for all $k \in \mathbb{Z}$.

See Figure 3.6.

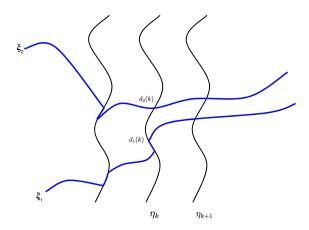


FIGURE 3.6. Positions of $d_1(k), d_2(k)$

We claim that

(3.10)
$$\sum_{k \in \mathbb{Z}} (d_2(k) - d_1(k)) \le T,$$

which will be proved in Lemma 3.4 below. In particular, (3.10) implies

$$\lim_{k \to \infty} (d_2(k) - d_1(k)) = 0.$$

Similar to Case 1 above, since $\eta_k(d_i(k)) \in \xi_i(\mathbb{R})$ for i = 1, 2,

$$u_0(\eta_k(d_1(k))) - u_0(\xi_1(0)) = u_1(\eta_k(d_1(k))) - u_1(\xi_1(0)),$$

and

$$u_0(\eta_k(d_2(k))) - u_0(\xi_2(0)) = u_1(\eta_k(d_2(k))) - u_1(\xi_2(0)).$$

Taking the difference of the two equations and sending $k \to \infty$, we obtain the claim.

Proof of Theorem 1.2. To finish the proof of the main result, we only need to prove (3.7). We proceed by using (3.8).

Write

$$g(t) = u_0(\eta(t)) - u_1(\eta(t)).$$

Then, g(t) is Lipschitz continuous and by (3.2),

$$g'(t) = 0$$
 for $t \in \mathcal{I}$.

See (3.6) for the definition of \mathcal{I} . Moreover, due to (3.8), for each $i \in \mathbb{N}$,

$$g(a_i) = g(b_i).$$

Define a new function

$$h(t) = \begin{cases} g(t) & \text{for } t \in \mathcal{I} \\ g(a_i) & \text{for } t \in (a_i, b_i), \text{ and } i \in \mathbb{N}. \end{cases}$$

Clearly, h(t) is Lipschitz continuous and h'(t) = 0 for a.e. $t \in \mathbb{R}$. Hence $h \equiv c$ for some constant $c \in \mathbb{R}$. This leads to

$$q(t) = c$$
 for $t \in \mathcal{I}$.

This is absurd as for $m \in \mathbb{N}$,

$$|g(mT) - g(0)| \ge m|(p_0 - p_1) \cdot (0, 1)| - \max_{\mathbb{T}^2} |v_0| - \max_{\mathbb{T}^2} |v_1|,$$

which leads to $\lim_{t\to\infty} |g(t)| = +\infty$.

Finally, we give a proof of (3.10).

Lemma 3.4. In Case 2 in the proof of Lemma 3.3,

$$\sum_{k \in \mathbb{Z}} (d_2(k) - d_1(k)) \le T.$$

Proof. We first show that, for all $k \in \mathbb{Z}$, the open interval

$$(3.11) (d_1(k), d_2(k)) \subset \mathbb{R} \setminus \mathcal{I},$$

that is, it is one of those open intervals $\{(a_j,b_j)\}_{j\geq 1}$. It suffices to show this for k>0 as the proof for k<0 is similar. We argue by contradiction. If this were not true, then there would exist $k\in\mathbb{N}$ and $\tilde{\xi}\in\mathcal{U}$ such that

(3.12)
$$\tilde{\xi}(0) \in \{ \eta_k(\theta) : \theta \in (d_1(k), d_2(k)) \}.$$

Since $\tilde{\xi}$ cannot pass the portion of η on (a_0, b_0) , we deduce that, if trace backward along $\tilde{\xi}$, it must intersect ξ_1 or ξ_2 before it intersects η . See Figure 3.7.

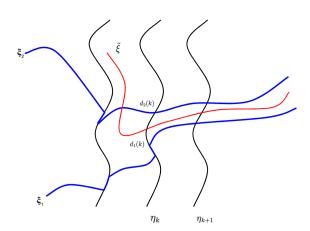


FIGURE 3.7. Relative position of $\tilde{\xi}$

Let

$$t_{-} = \max\{t \le 0 : \tilde{\xi}(t) \in \xi_{1}(\mathbb{R}) \cup \xi_{2}(\mathbb{R})\},\$$

$$t_{+} = \inf\{t \ge 0 : \tilde{\xi}(t) \in \xi_{1}(\mathbb{R}) \cup \xi_{2}(\mathbb{R})\}.$$

Then $t_{-} < 0$ and $t_{+} > 0$. Note that t_{+} could be $+\infty$.

Now we will use the gluing property of Remark 3 to construct a new orbit in \mathcal{U} . By two dimensional topology and (3.9), it is easy to see that, for each $k \in \mathbb{Z}$,

(3.13)
$$\tilde{\xi}((t_-, t_+)) \cap (\eta_k(\mathbb{R})) \subset \{\eta_k(\theta) : \theta \in (d_1(k), d_2(k))\}.$$

Without loss of generality, we assume that $\tilde{\xi}(t_-) \in \xi_2(\mathbb{R})$ and $\tilde{\xi}(t_+) \in \xi_j(\mathbb{R})$ for j = 1 or j = 2 if $t_+ < +\infty$. Suppose that

$$\tilde{\xi}(t_{-}) = \xi_{2}(\bar{t}_{-})$$
 and $\tilde{\xi}(t_{+}) = \xi_{j}(\bar{t}_{+})$.

for $0 < \bar{t}_- < \bar{t}_+$. Here $\bar{t}_+ = +\infty$ if $t_+ = +\infty$. Let

$$\xi_{3}(t) = \begin{cases} \xi_{2}(t) & \text{for } t \leq \bar{t}_{-}, \\ \tilde{\xi}(t+t_{-}-\bar{t}_{-}) & \text{for } \bar{t}_{-} \leq t \leq t_{+}+\bar{t}_{-}-t_{-}, \\ \xi_{j}(t+\bar{t}_{+}+t_{-}-\bar{t}_{-}-t_{+}) & \text{for } t \geq t_{+}+\bar{t}_{-}-t_{-}. \end{cases}$$

Then, by (3.12) and (3.13), we have that

$$\xi_3(0) = \eta(b_1)$$
 and $d(\xi_3, \xi_1) < d(\xi_2, \xi_1)$.

This contradicts the choice of ξ_1 and ξ_2 . Hence our claim (3.11) holds.

Next we show that for $k \neq l$, $(d_1(k), d_2(k))$ is not a T-translation of $(d_1(l), d_2(l))$. In fact, if

$$(d_1(k), d_2(k)) = (d_1(l), d_2(l)) + jT$$

for some $j \in \mathbb{Z}\setminus\{0\}$, then both $\eta(d_1(k))$ and $\eta(d_1(l)) + ((k-l)J, j)$ are on ξ_1 . Owing to Corollary 2.9, the outward normal vector \vec{n} is rational, which contradicts our assumption. Accordingly, after we translate all $(d_1(k), d_2(k))$ into $(a_1, a_1 + T)$, they are disjoint. Therefore, (3.10) holds true.

Remark 6. Our main result can be extended to more general metric Hamiltonians, e.g., $H(y,p) = \sqrt{p \cdot A(y) \cdot p}$, where A(y) is a 2×2 positive definite matrix that is periodic and continuous in y. The proof is similar to that of Theorem 1.1 by employing the Lagrangian associated with $H^2(y,p) = p \cdot A(y) \cdot p$ and a generalized version of the tightness argument used in Lemma 2.2. Here is the generalized version that appeared in the study of L^{∞} - variational problems: suppose that $\hat{H} = \hat{H}(p) \in C^1(\mathbb{R}^n)$ is strictly convex and superlinear. Assume that \hat{u} is a Lipschitz continuous function satisfying that $\hat{H}(D\hat{u}) \leq 0$ for a.e. $x \in \mathbb{R}^n$. If there exists a nonzero vector $q \in \mathbb{R}^n$ such that

$$\hat{u}(qt) - \hat{u}(0) = t\hat{L}(q)$$
 for all $t \in \mathbb{R}$.

Here \hat{L} is the Lagrangian associated with \hat{H} . Then, $\hat{u}(x) = \hat{u}(0) + D\hat{L}(q) \cdot x$ for all $x \in \mathbb{R}^n$ and $\hat{H}(D\hat{L}(q)) = 0$. See page 1123 in [21] for the proof. The details are left to interested readers.

Data available statement: This paper has no associated data.

References

- [1] G. Aronsson, Extension of functions satisfying Lipschitz conditions, Ark. Mat. 6 (1967), 551–561.
- [2] A. Auffinger, M. Damron, Jack Hanson, 50 Years of First-Passage Percolation, University Lecture series, Volume 68, American Mathematical Society.
- [3] V. Bangert, Mather sets for twist maps and geodesics on tori, Dynam. Report. Ser. Dynam. Systems Appl. 1, Wiley, Chichester, 1988, 1–56.
- [4] V. Bangert, Geodesic rays, Busemann functions and monotone twist maps, Calc. Var. Partial Differential Equations 2 (1994), no. 1, 49–63.
- [5] D. Burago, *Periodic metrics*, Adv. Soviet Math. 9, (1992), 205–210.
- [6] D. Burago, S. Ivanov, B. Kleiner, On the structure of the stable norm of periodic metrics, Mathematical Research Letters, 4(6) (1997), 791–808.
- [7] M. J. D. Carneiro, On minimizing measures of the action of autonomous Lagrangians, Nonlinearity 8 (1995), no. 6, 1077–1085.
- [8] C. Q. Cheng, J. Xue, Order property and modulus of continuity of weak KAM solutions, Calculus of variation and PDEs, 57(2018), 65.
- [9] W. E. Aubry-Mather theory and periodic solutions of the forced Burgers equation, Comm. Pure Appl. Math., vol. 52, no. 7, pp. 811–828, 1999.
- [10] L. C. Evans, Periodic homogenisation of certain fully nonlinear partial differential equations, Proc. Roy. Soc. Edinburgh Sect. A 120 (1992), no. 3-4, 245–265.
- [11] L. C. Evans, D. Gomes, Effective Hamiltonians and Averaging for Hamiltonian Dynamics. I, Archive for rational mechanics and analysis 157 (1), 1–33, 2001.
- [12] A. Fathi, The weak KAM theorem in Lagrangian dynamics, Cambridge University Press (2004).
- [13] A. Fathi, A. Siconolfi, PDE aspects of Aubry-Mather theory for quasiconvex Hamiltonians, Calculus of Variations and Partial Differential Equations volume 22, 185–228 (2005).
- [14] P.-L. Lions, G. Papanicolaou and S. R. S. Varadhan, *Homogenization of Hamilton–Jacobi equations*, unpublished work (1987).
- [15] W. Jing, H. V. Tran, and Y. Yu, Effective fronts of polygon shapes in two dimensions, arXiv:2112.10747 [math.AP].
- [16] Y. Li, L. Nirenberg, The distance function to the boundary, Finsler geometry and the singular set of viscosity solutions of some Hamilton-Jacobi equations January 2005, Communications on Pure and Applied Mathematics 58(1), DOI:10.1002/cpa.20051.
- [17] P. L. Lions, Generalized solutions of Hamilton-Jacobi equations, Research Notes in Mathematics; 69, Pitman (Advanced Publishing Program), Boston, Mass.- London, 1982.
- [18] J. N. Mather, Differentiability of the minimal average action as a function of the rotation number, Boletim da Sociedade Brasileira de Matemática-Bulletin/Brazilian Mathematical Society, 21, pages 59–70 (1990).
- [19] J. N. Mather, Action minimizing invariant measures for positive definite Lagrangian systems, Math. Z. 207 (1991), no. 2, 169–207.
- [20] H. V. Tran, Hamilton-Jacobi equations: Theory and Applications, American Mathematical Society, Graduate Studies in Mathematics, Volume 213, 2021.
- [21] Y. Yu, L^{∞} variational problems and weak KAM theory, Comm. Pure Appl. Math, Vol. 60 (2007), no. 8, 1111–1147.

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