



A New Asteroseismic Kepler Benchmark Constrains the Onset of Weakened Magnetic Braking in Mature Sun-like Stars

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Abstract

Stellar spin down is a critical yet poorly understood component of stellar evolution. In particular, results from the Kepler Mission imply that mature age, solar-type stars have inefficient magnetic braking, resulting in a stalled spin-down rate. However, a large number of precise asteroseismic ages are needed for mature (≥ 3 Gyr) stars in order to probe the regime where traditional and stalled spin-down models differ. In this paper, we present a new asteroseismic benchmark star for gyrochronology discovered using reprocessed Kepler short cadence data. KIC 11029516 (Papayu) is a bright ($Kp = 9.6$ mag) solar-type star with a well-measured rotation period (21.1 ± 0.8 days) from spot modulation using 4 yr of Kepler long-cadence data. We combine asteroseismology and spectroscopy to obtain $T_{\text{eff}} = 5888 \pm 100$ K, $[\text{Fe}/\text{H}] = 0.30 \pm 0.06$ dex, $M = 1.24 \pm 0.05 M_{\odot}$, $R = 1.34 \pm 0.02 R_{\odot}$, and age of 4.0 ± 0.4 Gyr, making Papayu one of the most similar stars to the Sun in terms of temperature and radius with an asteroseismic age and a rotation period measured from spot modulation. We find that Papayu sits at the transition of where traditional and weakened spin-down models diverge. A comparison with stars of similar zero-age main-sequence temperatures supports previous findings that weakened spin-down models are required to explain the ages and rotation periods of old solar-type stars.

Unified Astronomy Thesaurus concepts: Asteroseismology (73); Stellar ages (1581); Stellar evolution (1599); Stellar evolutionary tracks (1600); Stellar rotation (1629)

1. Introduction

The search for stars with solar-like oscillations has been revolutionized in the past decades with capabilities of space telescopes such as the NASA Kepler mission (Gilliland et al. 2010a). Solar-like oscillations are convection-driven pressure waves that propagate through the interiors of stars, probing their internal structures and compositions (Chaplin & Miglio 2013; García & Ballot 2019). Fundamental stellar properties such as mass, mean density, surface gravity, and radius can be derived through global asteroseismic parameters (Huber et al. 2013; Chaplin et al. 2014; Serenelli et al. 2017). Additionally, modeling of individual oscillation frequencies yields precise stellar age measurements (Silva Aguirre et al. 2015, 2017; Metcalfe et al. 2014; Creevey et al. 2017).

Independent from asteroseismology, stellar rotation periods can also be used to predict ages of stars through

gyrochronology (Barnes 2007). Gyrochronology is built upon the characteristic property of cool main-sequence dwarfs to lose angular momentum as they age (Skumanich 1972). This spin down is triggered by the magnetic braking generated through the interaction of the highly ionized stellar wind with the stellar magnetic field (Weber & Davis 1967). However, using gyrochronology as an age estimator requires a well-calibrated relationship between the rate of stellar rotation and age, which depends predominantly on the depth of convection zone. For example, hot stars ($T_{\text{eff}} > 6250$ K) beyond the Kraft break (Kraft 1967) experience minimal spin down as they have thin surface convective envelopes that are unable to sustain magnetic braking (Durney & Latour 1978). On the other hand, stars cooler than the Kraft break ($T_{\text{eff}} < 6250$ K) have deeper convection zones, resulting in greater spin down. Physical changes from stellar evolution, such as increasing radii, further complicate the angular momentum loss or spin-down process.

Stellar spin down has a direct impact on stellar activity. Stellar dynamo theory predicts that the magnetic field is modulated through meridional circulation, convection, and differential rotation (Charbonneau 2010). Thus, since the rate



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of stellar rotation has significant impact on stellar activity and spot modulation (Wright et al. 2011, 2018; Santos et al. 2021; McQuillan et al. 2014), stellar activity proxies are also related to the Rossby number, which is defined as the ratio between the rotation period and the convective overturn timescale (Lehtinen et al. 2020; Noyes et al. 1984). The Rossby number is representative of the ratio between the inertial and coriolis forces within a rotating star, and appears to be a crucial metric for the effectiveness of magnetic braking.

Recent results have highlighted interesting tensions between ages from asteroseismology and those predicted by gyrochronology. Specifically, van Saders et al. (2016) found that mature-aged solar-type stars rotate faster than expected from standard spin-down models, which assumes stars spin down at a rate inversely proportional to the square root of their age (Skumanich 1972). Additionally, limits imposed on the Rossby number, such as allowing magnetic braking to weaken after a critical Rossby number of 2.16 (hereafter Ro_{crit} ; van Saders et al. 2016), were suggested to yield empirical agreement with expected ages. Metcalfe et al. (2016) further examined the correlation between magnetism, stellar activity, and weakened magnetic braking to suggest that the stalling of such braking is caused by a decrease in the dipole component of the global magnetic field as it is disrupted by differential stellar rotation (Réville et al. 2015). More recent evidence in support of stalled spin down include the detection of a pileup of slowly rotating stars (David et al. 2022) and the confirmation of rotation periods using asteroseismology (Hall et al. 2021).

Despite a mounting body of evidence in its favor, a challenge for interpreting stalled spin-down models is the lack of older stars and nonsolar metallicities in our gyrochronology and age calibration sample. Nonsolar metallicities are relevant to the stalled spin-down hypothesis as metallicity impacts stellar convection, which is critical for spin down (Tayar et al. 2017). Thus, our stellar rotation samples are limited and biased toward young stars in open clusters (Mamajek & Hillenbrand 2008; Meibom et al. 2015; Douglas et al. 2017; Curtis et al. 2020), and the rate at which older stars spin down is not well understood (Metcalfe & van Saders 2017). In particular, the region where magnetic braking shuts down is not well calibrated due to a lack of solar-type stars.

While asteroseismic ages have been measured for hundreds of stars, only a small fraction have robustly measured rotation periods. For example, only 10 out of the 66 solar-like oscillators from the Kepler “Legacy” sample have rotation periods measured via spot modulation (Silva Aguirre et al. 2017; McQuillan et al. 2014; Santos et al. 2021). The reprocessing of Kepler short-cadence data (Thompson et al. 2016) now allows a comprehensive search for additional solar-like oscillators in the Kepler data. Here we present the discovery of a new asteroseismic gyrochronology benchmark star that is similar to the Sun in terms of rotation period and age, and thus provides an excellent opportunity to test rotational spin-down models.

2. Observations

2.1. Kepler Photometry

2.1.1. Solar-like Oscillations

For the first four quarters of the Kepler Mission, the Kepler Asteroseismic Science Consortium (KASC) surveyed ≈ 2000 targets, yielding ≈ 500 main-sequence and subgiant detections

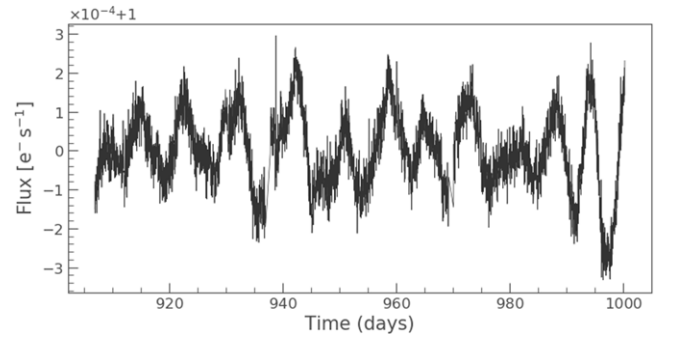


Figure 1. Long-cadence lightcurve of the Quarter 10 of Kepler data for Papayu. The lightcurve is generated using the `lightkurve` package (Barentsen et al. 2020).

and ≈ 1500 nondetections of asteroseismic solar-like oscillators (Chaplin et al. 2011a). The nondetections can be partially explained by increased stellar activity (Chaplin et al. 2011a, 2011b) or the nonoptimal apertures in past Kepler data releases. Reprocessed Kepler data (DR25; Thompson et al. 2016) has already allowed new detections (Mathur et al. 2022).

To search for new solar-like oscillators, we used reprocessed data for all Kepler stars observed in short cadence. We used stellar properties from Gaia (Berger et al. 2018) to recalculate detection probabilities using the same formalism as described in Chaplin et al. (2011a, 2011b). Out of the 4785 stars with short-cadence data, we selected only the targets with high (>0.9) Gaia-derived asteroseismic detection probabilities, premeasured rotation periods, and a single quarter (30 days) of short-cadence data (Gilliland et al. 2010b; Berger et al. 2018; McQuillan et al. 2014). This narrowed our search down to 256 stars. We focus on one benchmark star in our sample, KIC 11029516 (hereafter Papayu¹⁶). Papayu was observed in long cadence for 17 quarters, and short cadence for one quarter via the NASA Kepler guest investigator program # GO20022 (PI Guzik). Additional new detections, including some with measured rotation periods, will be presented in M. Sayeed et al. (2024, in preparation). We focus on KIC 11029516 as one of the highest signal-to-noise ratio (S/N) examples that allows individual frequency modeling.

2.1.2. Rotation Period

Figure 1 shows a long-cadence lightcurve for one of the quarters of Papayu’s long-cadence data. Santos et al. (2021) reports a rotation period of P_{rot} of 22.1 ± 1.9 days, McQuillan et al. (2014) estimates 21.1 ± 0.8 days, and a clear and consistent rotational modulation signal can be seen via visual inspection (Figure 1) with a rotation period of ≈ 20 days. Slight differences in P_{rot} estimates from Santos et al. (2021) and McQuillan et al. (2014) arise from the use of all 17 quarters and the removal of photometric pollution from nearby stars in Santos et al. (2021) compared with the use of AutoACF (McQuillan et al. 2013) in McQuillan et al. (2014). We adopt a rotation period of 21.1 ± 0.8 days for the remainder of this study.

¹⁶ Papayu, or પપયુ, translates to Papaya in the primary author’s mother’s mother tongue, Gujarati.

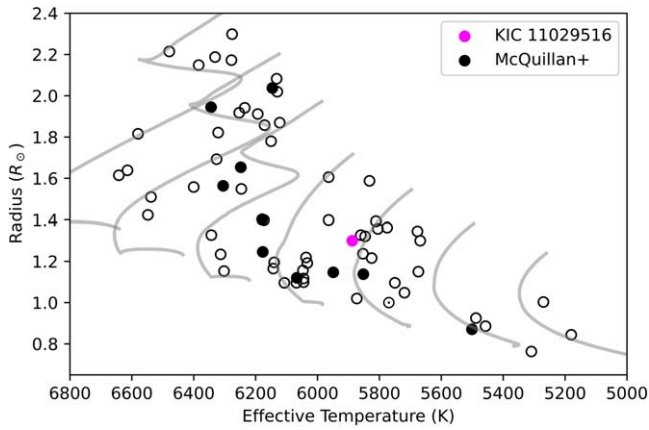


Figure 2. Stellar radius vs. effective temperature of the Kepler Legacy sample (Silva Aguirre et al. 2017). Filled circles show stars with rotation periods determined by McQuillan et al. (2014). KIC 11029516 (Papayu) is shown in pink.

2.2. High-resolution Spectroscopy

We obtained a high-resolution spectrum of Papayu on 2019 August 14 using the HIRES spectrograph (Vogt et al. 1994) at the Keck I 10 m telescope on Maunakea observatory, Hawai‘i. The spectrum was obtained and reduced as part of the California Planet Search queue (Howard et al. 2010). We obtained a 1.5 minute integration using the C5 decker, resulting in an S/N per pixel of 100 at ~ 600 nm with a spectral resolving power of $R \sim 60,000$.

To measure atmospheric parameters, we applied Specmatch-synth (Petigura 2015), which fits a synthetic grid of model atmospheres and has been extensively validated through the California Kepler Survey (Petigura et al. 2017; Johnson et al. 2017). The Specmatch-synth resulted in $T_{\text{eff}} = 5888 \pm 100$ K, $\log g = 4.14 \pm 0.1$ dex, $[\text{Fe}/\text{H}] = 0.31 \pm 0.06$ dex, and $v \sin i = 3.02 \pm 1.0$ km s $^{-1}$. Additionally, the Specmatch-empirical parameters were $T_{\text{eff}} = 5759.00 \pm 110.00$ K and $[\text{Fe}/\text{H}] = 0.31 \pm 0.09$ dex (Yee et al. 2017). Thus, these Specmatch-synth parameters are consistent with the latest Kepler-Gaia stellar properties catalog (Berger et al. 2020) and within $\approx 1\sigma$ to Specmatch-empirical.

Figure 2 shows an asteroseismic H-R diagram of stellar radius versus effective temperature utilizing MIST isochrones (Choi et al. 2016). There are only 10 stars out of the Kepler LEGACY sample of 66 stars with well-determined asteroseismic properties and rotation periods from spot modulation that can be used to calibrate the relationship between gyrochronology and stellar age. Figure 2 shows Papayu is one of the most similar stars to the Sun in terms of radius and temperature with a measured rotation period from spot modulation.

2.3. Bolometric Flux and Parallax

We calculated a bolometric flux of $f_{\text{bol}} = 3.68 \pm 0.11 \times 10^{-9}$ erg s $^{-1}$ cm $^{-2}$ by using Two Micron All Sky Survey (2MASS) K-band photometry (Skrutskie et al. 2006) and applying a bolometric correction from MIST isochrones (Choi et al. 2016), as implemented in *isoclassify* (Huber et al. 2016). We used a bolometric correction error floor of 0.03 mag, consistent with the expected systematic offsets for bolometric fluxes (Zinn et al. 2019; Tayar et al. 2022). Interstellar extinction from 3D dust maps (Green et al. 2015) was found to be negligible, consistent with the short distance (≈ 130 pc). Finally, we

combined our f_{bol} value with the Gaia DR3 parallax (Lindgren et al. 2021) to calculate luminosity of $1.91 \pm 0.05 L_{\odot}$.

As an independent check we performed a spectral energy distribution (SED) fit following the procedures described in Stassun & Torres (2016), Stassun et al. (2017, 2018). We used $B_T V_T$ magnitudes from Tycho-2, JHK_S magnitudes from 2MASS, W1–W4 magnitudes from Wide-field Infrared Survey Explorer, $G_{\text{BP}} G_{\text{RP}}$ magnitudes from Gaia, the near-ultraviolet magnitude from GALEX, and the U magnitude from the Kepler-INT Survey. We performed a fit using PHOENIX stellar atmosphere models (Husser et al. 2013), with the free parameters being the effective temperature (T_{eff}) and the extinction A_V , which we limited to maximum line-of-sight value from the Galactic dust maps of Schlegel et al. (1998). We also adopted the metallicity ($[\text{Fe}/\text{H}]$) from the spectroscopic analysis above. Integrating the (unreddened) best-fit model SED gives the bolometric flux at Earth, $F_{\text{bol}} = 3.873 \pm 0.090 \times 10^{-9}$ erg s $^{-1}$ cm $^{-2}$. Combining F_{bol} with the Gaia parallax gives the luminosity $L_{\text{bol}} = 1.998 \pm 0.047 L_{\odot}$, in good agreement with the value derived using bolometric corrections.

3. Asteroseismic Data Analysis

3.1. Data Preparation

We used the single quarter of available short-cadence data for our analysis. Pre-search Data Conditioning SAP flux was preferred over Simple Aperture Photometry (SAP) flux, as it has been detrended for any systematic effects (Smith et al. 2012; Stumpe et al. 2012). We ignored data points with nonzero quality flags and applied a 3σ clipping of the lightcurve with a box width of 50 data points to remove any outliers. Furthermore, to focus on the high-frequency asteroseismic oscillations, we removed data artifacts due to Kepler observing quarters and normalized the lightcurve using a box filter of 0.05 days to remove long periodic variations caused by instrumental effects (thermal variations, drifts) and/or stellar variability (stellar rotation, stellar activity). The power spectrum generated after making these corrections is shown in Figure 3, which shows a clear detection of solar-like oscillations around 2000 μHz .

3.2. Background Modeling

In order to separate the asteroseismic oscillations from the granulation signal and background noise, we used the Bayesian modeling software package DIAMONDS (Corsaro & De Ridder 2014). We used a composite stellar background model that consisted of constant photon noise, stellar activity variations, a sampling response function, and two Harvey models (Harvey 1985). Eight total parameters were involved in the fit, including the white noise threshold (in ppm 2 μHz^{-1}), the Harvey model parameters with a characteristic frequency (μHz) and amplitude (ppm), along with the frequency of maximum power (ν_{max}), the height (in ppm 2 μHz^{-1}), and the FWHM of the power excess.

Figure 3 shows the fitted background model and its individual components calculated via DIAMONDS. Initial guesses were estimated through visual inspection of the power spectrum between 270 and 8333 μHz . The goodness and convergence of fit was assessed through DIAMONDS via the posterior probability distributions (Appendix). The final values for the background fit, calculated as the median and 1σ

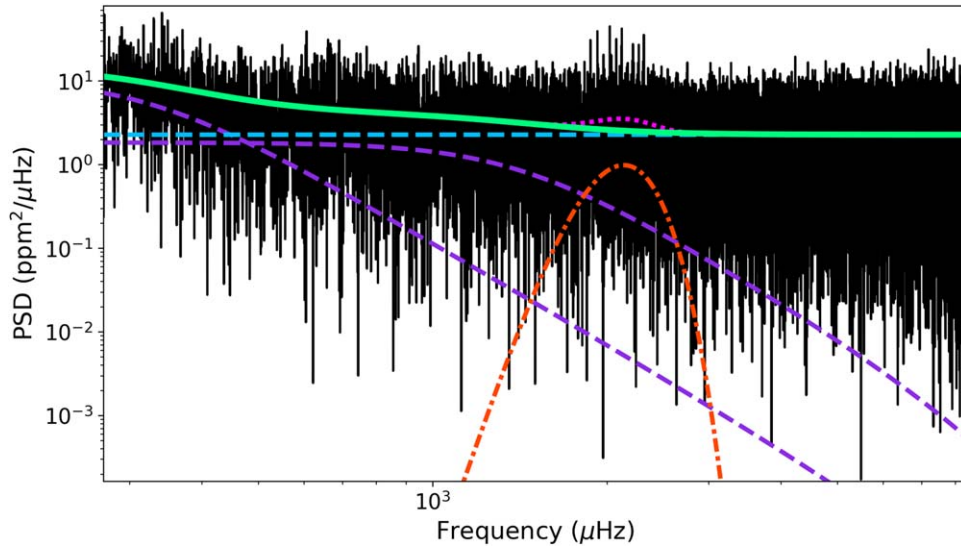


Figure 3. Power spectrum of the single quarter of short-cadence data for Papayu after outlier correction and high-pass filtering. The solid green line is the modeled granulation background. The blue dashed line is a constant white noise, the purple dashed curves are the Harvey models, the dashed-dotted orange curve is the height of the oscillation, and the dotted pink curve is the frequency of maximum power and the FWHM of the oscillation.

Table 1
Model Parameters of the Background Fit

Parameter	Median Value
w ($\text{ppm}^2 \mu\text{Hz}^{-1}$)	$2.28^{+0.02}_{-0.01}$
σ_{long} (ppm)	$62.17^{+3.27}_{-2.70}$
b_{long} (μHz)	$322.14^{+19.01}_{-25.33}$
σ_1 (ppm)	$53.12^{+1.47}_{-1.17}$
b_1 (μHz)	$1,378.86^{+170.40}_{-201.89}$
H_{osc} ($\text{ppm}^2 \mu\text{Hz}^{-1}$)	$1.04^{+0.10}_{-0.11}$
ν_{max} (μHz)	$2,141.18^{+29.81}_{-28.74}$
σ_{env} (μHz)	$241.58^{+44.79}_{-49.56}$

Note. White noise (w), first Harvey model amplitude (σ_{long}), first Harvey model turnoff (b_{long}), second Harvey model amplitude (σ_1), second Harvey model turnoff (b_1), height of oscillation envelope (H_{osc}), frequency of maximum oscillation (ν_{max}), and FWHM of oscillation envelope (σ_{env}).

confidence limits of the posterior distributions, are listed in Table 1.

3.3. Individual Frequencies

After correcting the power spectrum for the background, we extracted individual frequencies using the Peakbagging module of DIAMONDS (Corsaro & De Ridder 2014). The module operates by generating a set of multiple Lorentzian profiles.

Similar to the Backgrounds module, the PeakBagging module requires initial guesses of each peak’s linewidth or Γ_i (μHz), amplitude or A_i (ppm), and central frequency or $\nu_{0,i}$ (μHz). We identified 15 possible peaks based on manual cross-identification. Thus, 45 initial guesses (Γ_i , A_i , and $\nu_{0,i}$ for $0 < i \leq 15$) with symmetric upper and lower limits were provided to DIAMONDS in a PeakBagging run.

We performed mode identification of extracted frequencies using an échelle diagram (Grec et al. 1983). An échelle diagram

is a compact visualization of a power spectrum, as it breaks the power spectrum up into $\Delta\nu$ -sized sections (radial orders) and stacks them atop each other. We expect an échelle diagram to show vertical ridges of quadrupole ($l = 2$), radial ($l = 0$), and dipole ($l = 1$) modes.

The initial guesses for linewidths and amplitudes were based on visual inspection and then tuned by fitting only a few sets of peaks at a time. Cumulative addition of the peaks allowed tuning of the linewidth and amplitude guesses at each step. Figure 4 shows the fit to the background-corrected power spectrum, with Lorentzian profiles derived from these 45 parameters. The final values and posteriors of all 45 parameters are shown in Table 2 and Figure 11.

Figure 5 shows an échelle diagram of the extracted frequencies. We successfully extracted five radial modes, six dipole modes, and four quadrupole modes. The value of large frequency separation, $\Delta\nu$, was measured to be $97.31 \pm 0.14 \mu\text{Hz}$ by taking the mean separation between consecutive radial modes. Similarly, using the mean of the small frequency separations for each radial order, we calculated a $\delta\nu_{02}$ value of $6.12 \pm 0.15 \mu\text{Hz}$.

4. Asteroseismic Modeling

4.1. Scaling Relations and the C-D Diagram

We calculated a preliminary asteroseismic age using an asteroseismic H-R diagram (hereafter C-D diagram; Christensen-Dalsgaard 1984) and asteroseismic scaling relationships (Bellinger 2019). Physically, the large frequency separation scales with the square root of the mean stellar density, while the small separation scales with the sound speed gradient in the core (Ulrich 1986). Thus the small separation is sensitive to stellar age on the main sequence.

Figure 6 shows a C-D diagram of $\delta\nu_{02}$ against $\Delta\nu$ for Papayu’s metallicity ($[\text{Fe}/\text{H}] = 0.30$), with stellar-evolution models from ASTEC and ADIPLS (Christensen-Dalsgaard 2008) as described in White et al. (2011). In the C-D diagram, stars evolve from the top right to the bottom left. The range of stellar masses in the C-D diagram is from 1.0 to 1.6 M_{\odot} , with the latter end slightly above the Kraft break (Kraft 1967). Using Papayu’s values of $\Delta\nu$ and $\delta\nu_{02}$ from

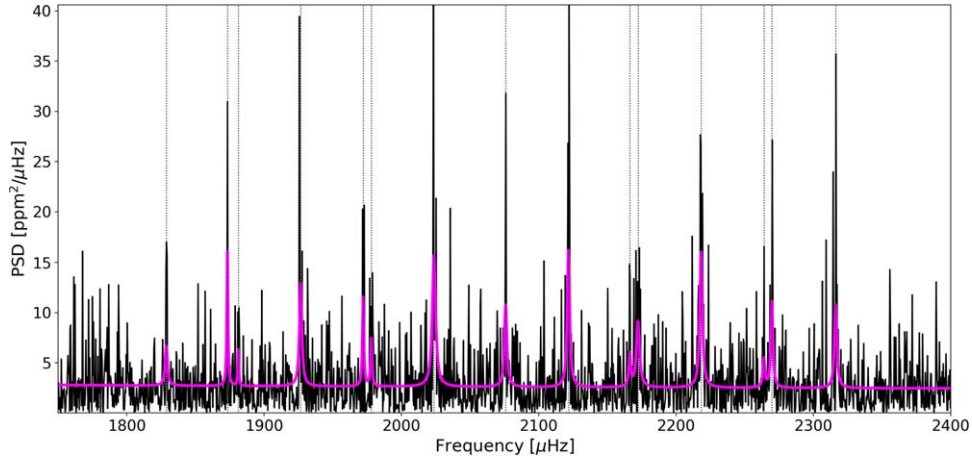


Figure 4. Identified oscillation frequencies for Papayu. The background-subtracted power spectrum is in black; the Lorentzian profiles that fit each of the 15 oscillation frequencies are in pink. Posteriors and final values for the parameters used in this fit are given in the [Appendix](#).

Table 2
Best-fit Values for Papayu’s Oscillation Frequencies

Peak #	n_i	l_i	$\nu_{0,i}$ (μHz)	A_i (ppm)	Γ_i (μHz)
1	17	1	$1,828.73^{+0.23}_{-0.25}$	$5.13^{+0.66}_{-0.64}$	$2.02^{+0.53}_{-0.59}$
2	17	2	$1,873.23^{+0.10}_{-0.11}$	$6.73^{+0.72}_{-0.78}$	$1.03^{+0.40}_{-0.33}$
3	18	0	$1,881.22^{+0.29}_{-0.33}$	$3.07^{+0.41}_{-0.42}$	$0.80^{+0.17}_{-0.17}$
4	18	1	$1,926.60^{+0.28}_{-0.28}$	$8.46^{+0.77}_{-0.73}$	$2.12^{+0.37}_{-0.42}$
5	18	2	$1,972.06^{+0.21}_{-0.14}$	$5.93^{+0.69}_{-0.60}$	$1.21^{+0.28}_{-0.30}$
6	19	0	$1,978.00^{+0.26}_{-0.27}$	$4.42^{+0.40}_{-0.38}$	$1.25^{+0.24}_{-0.24}$
7	19	1	$2,023.67^{+0.20}_{-0.24}$	$11.16^{+0.73}_{-0.91}$	$2.88^{+0.38}_{-0.40}$
8	20	0	$2,075.85^{+0.15}_{-0.18}$	$7.73^{+0.56}_{-0.60}$	$2.21^{+0.51}_{-0.47}$
9	20	1	$2,121.61^{+0.14}_{-0.15}$	$9.75^{+0.62}_{-0.54}$	$2.10^{+0.42}_{-0.48}$
10	20	2	$2,166.26^{+0.19}_{-0.17}$	$4.13^{+0.43}_{-0.45}$	$1.62^{+0.25}_{-0.24}$
11	21	0	$2,172.25^{+0.38}_{-0.31}$	$7.73^{+0.57}_{-0.52}$	$2.73^{+0.44}_{-0.48}$
12	21	1	$2,218.20^{+0.16}_{-0.16}$	$10.17^{+0.58}_{-0.54}$	$2.29^{+0.29}_{-0.36}$
13	21	2	$2,263.91^{+0.26}_{-0.27}$	$4.04^{+0.58}_{-0.66}$	$1.73^{+0.25}_{-0.31}$
14	22	0	$2,269.76^{+0.16}_{-0.17}$	$7.13^{+0.53}_{-0.43}$	$1.76^{+0.22}_{-0.25}$
15	22	1	$2,316.22^{+0.20}_{-0.19}$	$7.71^{+0.49}_{-0.51}$	$2.14^{+0.25}_{-0.17}$

Note. Column names denote the peak number, i (increasing from left to right in Figure 4), radial order (n_i), spherical degree (l_i), central peak frequency ($\nu_{0,i}$), peak amplitude (A_i), and peak linewidth (Γ_i).

Section 3 we are able to constrain Papayu’s age from the C-D diagram to ≈ 4.0 Gyr. Papayu’s age can also be measured by substituting its stellar properties from Table 3 into the age-scaling relationship in Bellinger (2019). From this we obtain an age estimate of 3.64 ± 0.31 Gyr.

4.2. Frequency Modeling

To derive a more robust age, we modeled Papayu using the stellar-evolution code MESA (Paxton et al. 2011, 2013, 2015, 2018, 2019) using effective temperature, metallicity, $\log g$,

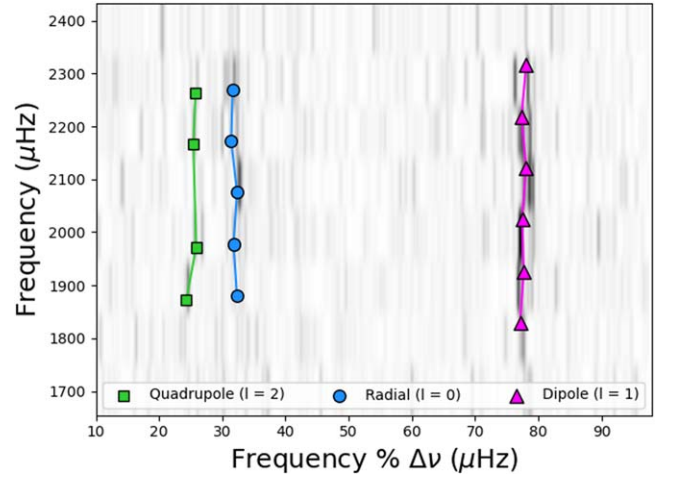


Figure 5. Échelle diagram using frequencies extracted by the PeakBagging module of the DIAMONDS package. The colors correspond to the different degrees of spherical harmonics as denoted in the legend, with the radial mode denoted in blue, the dipole mode denoted in pink, and the quadrupole mode denoted in green. The échelle diagram is created via *echelle* (Hey & Ball 2020).

$\Delta\nu$ (μHz), $\delta\nu_{02}$ (Table 3), and all the frequencies obtained using the DIAMONDS module (Table 2). All observed frequencies and parameters were added to a MESA inlist. This inlist uses the *astero* and *GYRE* modules (Townsend & Teitler 2013) to generate a pre-main-sequence model and perform a simplex search. The simplex search executes multiple runs in order to obtain the stellar model with the best fit to the observed parameters and frequencies. While the *astero* module obtains the best fits for each parameter, the linked *GYRE* stellar oscillation module computes modeled frequencies using a parallelized multiple Magnusson shooting scheme.

Specifically, the input physics was selected for variable mixing length parameter (α), helium abundance (Y), metallicity ($[\text{Fe}/\text{H}]$), and mass. Minimum and maximum thresholds for these varying parameters and initial values were provided. The target value and error in $\log g$ was computed through asteroseismic scaling relations, similar to the initial guess for mass. The initial guesses for effective temperature and metallicity were provided through spectroscopic observations

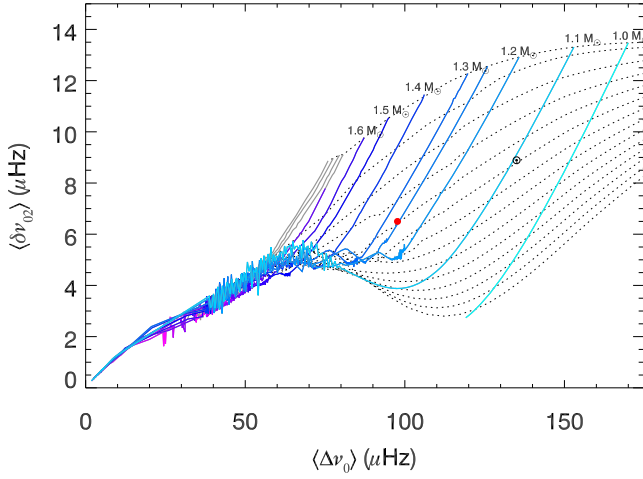


Figure 6. A C-D diagram, with Papayu’s position shown in red and the Sun’s position shown in black. The solid lines are stellar-evolution models tailored to Papayu’s metallicity of 0.30 dex; therefore the Sun does not cross the 1 M_{\odot} model. Stars on these model tracks evolve from the top right to the bottom left. The dashed lines are isochrones from zero to 12 Gyr (top to bottom) in steps of 2 Gyr.

Table 3
Stellar Parameters for Papayu

Star Information	
KIC ID	11029516
Tycho ID	TYC 3547-1118-1
2MASS ID	J19270298+4835118
Coordinates (ICRS)	19h 27 m 02 s, +48°35′11″
Kepler magnitude (mag)	9.6
Observables from Kepler	
ν_{\max} (μHz)	2121.7 ± 0.1
$\Delta\nu$ (μHz)	97.1 ± 0.1
$\delta\nu_{02}$ (μHz)	6.1 ± 0.2
Rotation period (days)	21.1 ± 0.8
Photometry and Gaia	
f_{bol} (10^{-9} erg s $^{-1}$ cm $^{-2}$)	3.68 ± 0.11
Luminosity (L_{\odot})	1.91 ± 0.05
Spectroscopy	
T_{eff} (K)	5888 ± 100
[Fe/H] (dex)	0.30 ± 0.06
Asteroseismic Modeling	
Mass (M_{\odot})	1.23 ± 0.03
Radius (R_{\odot})	1.34 ± 0.01
$\log g$ (cgs)	4.28 ± 0.01
Density (gcc)	0.728 ± 0.005
Age (Gyr)	4.0 ± 0.4

as described in Section 2.2. Initial values for the helium abundance and mixing length were 0.27 ± 0.03 and 1.8 ± 0.2 respectively. A cubic surface correction scheme, utilizing the default correction factor from Ball & Gizon (2017), is applied to adjust for the systematic disparity between modeled and observed mode frequencies due to inaccuracies in modeling near-surface stellar layers. The temperature structure of the atmosphere was defined using the Eddington gray relation, with opacity varying with optical depth. Elemental diffusion is allowed, and a radiative turbulence coefficient of unity is

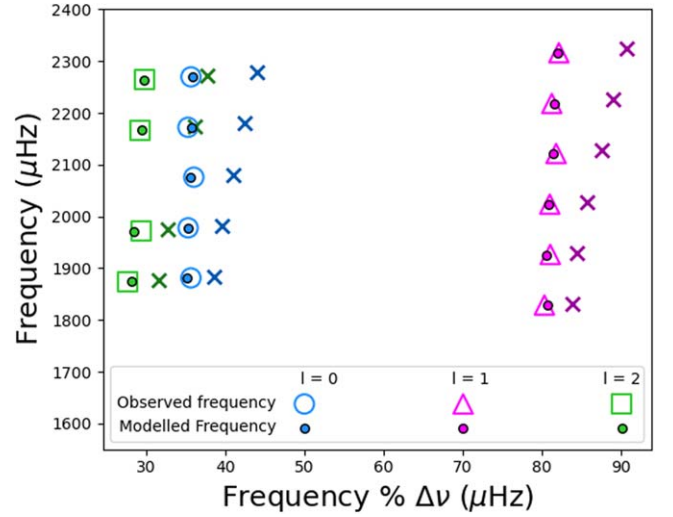


Figure 7. Échelle diagram comparing observed frequencies (unfilled shapes) to frequencies obtained from our best-fit MESA module (filled dots and crosses). The colors of the symbols represent the various spherical degrees, with the radial mode in blue, dipole mode in pink, and quadruple mode in green. Original MESA derived frequencies are denoted by colored crosses, and the surface-corrected MESA frequencies are shown in filled dots.

selected in accordance with Morel & Thévenin (2002). For varying helium abundance and mixing length, we performed runs varying overshooting, keeping it fixed to the default value of 0.015 and keeping it fixed to 0, but found no significant differences in MESA outputs. Papayu is on the verge of developing a convective core at $1.24 \pm 0.05 M_{\odot}$, so it appears overshooting is not a significant concern yet (Claret & Torres 2018). Additionally, no value for ν_{\max} was provided during our runs, and the χ^2 seismic fraction was increased to 0.8, compared to the default value of 0.667. Partial evaluation controls, chi-squared based time-step controls and stopping conditions were slightly altered to fine-tune our measurements, and all other parameters were unchanged.

MESA outputs multiple runs, each with frequencies and stellar parameters. Frequency and stellar parameter outputs were plotted with respect to run number to ensure compliance with input physics. Observed and modeled frequencies for the best-fitting model are shown in Figure 7, demonstrating good agreement. Stellar parameter values along with their corresponding statistical uncertainties for the best-fit MESA model are as follows: mass = $1.23 \pm 0.01 M_{\odot}$, radius = $1.34 \pm 0.01 R_{\odot}$, $\log g = 4.276 \pm 0.001$ cgs, density = 0.728 ± 0.004 gcc, and age = 4.0 ± 0.2 Gyr. To test for systematic errors, multiple coauthors modeled observed oscillation frequencies using other stellar-evolution codes and modeling methods, including BeSPP, YREC, BASTA, MESA, and AMP (Demarque et al. 2008; Christensen-Dalsgaard 2008; Metcalfe et al. 2009; Paxton et al. 2011, 2013, 2015, 2018, 2019; Ball & Gizon 2017; Serenelli et al. 2017; Aguirre Børsen-Koch et al. 2022). Model inputs include individual frequencies (Table 2), spectroscopic (Keck/HIRES) T_{eff} and [Fe/H], along with $\Delta\nu$, ν_{\max} , and Gaia-derived luminosity. Modeling efforts yielded consistent results, both with MESA derived parameters and outputs from various pipelines. We adopted the best-fit values from MESA modeling described above, with uncertainties estimated by adding in quadrature the formal uncertainty from MESA modeling and the standard deviation of the values over all methods. The final adopted values are listed in Table 3. Our final age uncertainty is

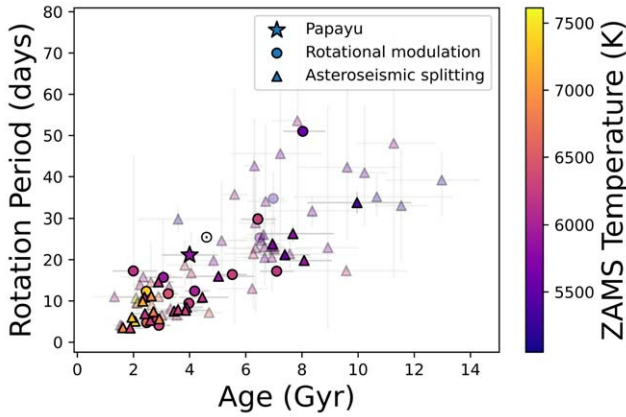


Figure 8. Rotation period vs. asteroseismic age for known asteroseismic gyrochronology benchmark stars color-coded by ZAMS temperature. Rotation periods are obtained through stellar modulation (circles) or asteroseismic frequency splittings (triangles). Fainter points show rotation period measurements with uncertainties exceeding 10%.

10%, and the value agrees within 4% to alternative estimates described in Section 4.1. Furthermore, Papayu’s rotation period and age align with the green M67 isochrone depicted in Figure 4 of Barnes et al. (2016). This suggests the gyrochrone age of a star exhibiting Papayu’s rotation period corresponds to the age we computed in this section.

4.3. Comparison to Age from Chromospheric Activity

The Keck/HIRES spectrum (Section 2.2) allows a measurement of the chromospheric activity from Ca H&K lines. We measured $\log R'_{\text{HK}} = -5.061$ using the method of Isaacson & Fischer (2010). After correcting for the metallicity, we calculated $\log R'_{\text{HK}}[T_{\text{eff}}] = -4.968$ on the scale of Lorenzo-Oliveira et al. (2018). The corresponding age inferred from chromospheric activity is $3.9^{+1.1}_{-0.9}$ Gyr, with an error bar estimated from the uncertainty in the age-activity calibration. The estimated age from chromospheric activity provides a good match to the asteroseismic age of 4.0 ± 0.4 Gyr.

5. Discussion

Papayu’s robust rotation period and asteroseismic age allows us to place the star in context with other known asteroseismic gyrochronology benchmarks. Figure 8 shows rotation period as a function of stellar age for all asteroseismic gyrochronology benchmark stars in the LEGACY sample along with Metcalfe et al. (2016) and Hall et al. (2021). The stars in Figure 8 provide a comprehensive overview of known asteroseismic gyrochronology rotators. Almost all the stars in this sample have asteroseismic ages and fundamental stellar properties through Silva Aguirre et al. (2017) and Silva Aguirre et al. (2015). There are a handful of stars that did not have ages or stellar properties measured in both LEGACY surveys but were found in Metcalfe et al. (2016). The rotation periods of stars are obtained either through rotational modulation measurements (García et al. 2014; McQuillan et al. 2014) or through asteroseismic splitting measurements (Davies et al. 2015; Hall et al. 2021). Rotation periods measured via spot modulation were preferred over asteroseismic splittings, with García et al. (2014) being given preference over McQuillan et al. (2014). When periods were

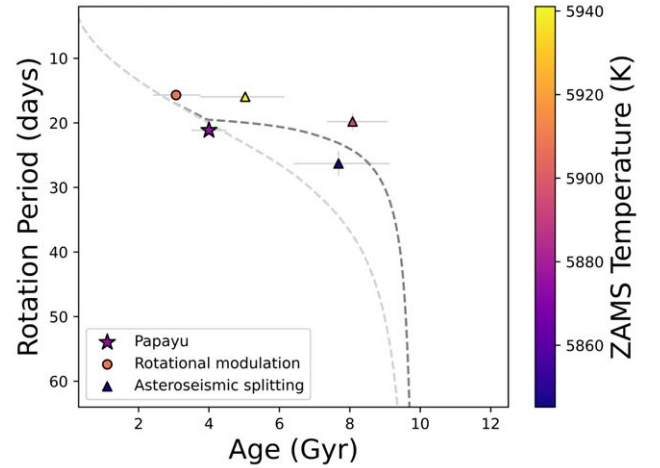


Figure 9. Rotation period vs. asteroseismic age for a subset of asteroseismic gyrochronology benchmark stars within 100 K of Papayu’s ZAMS temperature and with rotation period measurement errors less than 10%. These stars correspond to the brighter colored ones in Figure 8 and are similarly colored by their ZAMS temperatures, which fall in a range ≤ 5974 or ≥ 5774 K, a symmetric interval around Papayu’s ZAMS temperature of 5874 K. The dark gray line corresponds to kiauhoku’s Ro_{crit} model, and the light gray line corresponds to kiauhoku’s fast-launch model for a mass and metallicity that reproduced the mean ZAMS temperature of the sample.

obtained through asteroseismic splittings, Davies et al. (2015) was given preference over Hall et al. (2021).

The stars in Figure 8 are colored by their zero-age main-sequence (ZAMS) temperatures. The ZAMS temperature is a function of mass and metallicity and selects stars with similar structures and convection zone depths. Convective overturn timescales, which scale with convection zone depth, enter in the Rossby number and are therefore an important parameter in braking laws. The ZAMS temperatures were obtained by interpolating model grids for a given stellar mass and metallicity using the *kiauhoku* package (Claytor et al. 2020).

Figure 8 shows that Papayu is similar to the Sun in terms of rotation and age for stars with known spot modulation measurements. Figure 9 takes a closer look into this smaller subset of Figure 8, displaying the period–age space as a function of ZAMS temperatures for a subsample of asteroseismic gyrochronology benchmark stars with well-constrained rotation periods (with errors $\leq 10\%$) and ZAMS temperatures within 100 K of Papayu’s. We also eliminated stars from Hall et al. (2021), which have rotation period uncertainties larger than 10%.

In order to ascertain whether Papayu experiences weakened magnetic braking, we calculated a spin-down track for the average star in this subsample. To construct these tracks, we first averaged the ZAMS temperature among all the stars in the subset. We then calculated a model grid using weakened magnetic braking for the spin down for a fixed solar mass ($1 M_{\odot}$) and varied the metallicity (0.2 dex) until we reached a model with ZAMS temperature that matches the observed ZAMS temperature of the sample. The Ro_{crit} model grid was computed via *kiauhoku* (Claytor et al. 2020), which uses the modified magnetic braking law based on van Saders & Pinsonneault (2013), taking into account stalled stellar spin down based on a critical Rossby number of 2.16 (van Saders et al. 2016). We also calculated standard spin-down models using the fast-launch conditions of van Saders & Pinsonneault (2013) anchored to the rapidly rotating envelope in the Pleiades

and M37 open clusters, and do not take into account the effect of a critical Rossby number for magnetic braking. It is worth noting that assumptions about the initial rotation period do not significantly shift the rotational model tracks (van Saders & Pinsonneault 2013; van Saders et al. 2016; Metcalfe et al. 2020), and that metal-poor stars spin down more effectively than metal-rich stars in the early phases of spin down (Amard & Matt 2020). However, in later spin-down phases this reverses, and metal-poor stars spin down less effectively than metal-rich stars. The *kiauhoku* model grids take this effect into account (Claytor et al. 2020).

Figure 9 shows the Papayu sits at the transition region between both Skumanich-like and Ro_{crit} model tracks, thus, it provides an anchor point where both spin-down model tracks diverge. We note that the models do provide a good match for the Sun since they trace the average mass and metallicity for a star in this sample. We observe that the Ro_{crit} or stalled spin-down model provides a significantly better match to the sample, confirming previous results by implying a preference for the weakened magnetic braking law.

6. Conclusions

We have presented the discovery of a new asteroseismic gyrochronology benchmark star in the Kepler field (KIC 11029516 or Papayu). Our main conclusions are as follows:

1. Papayu was discovered using Kepler short-cadence data post-DR25 reprocessing, and is a bright ($V=9.63$ mag) solar-type star, with a spectroscopic T_{eff} of 5888 ± 100 K and $[\text{Fe}/\text{H}]$ of 0.30 ± 0.06 dex. We used DIAMONDS frequency modeling in conjunction with MESA modeling to obtain its mass of $1.23 \pm 0.03 M_{\odot}$, radius of $1.34 \pm 0.01 R_{\odot}$ and most importantly, asteroseismic age of 4.0 ± 0.4 Gyr.
2. Among oscillating Kepler stars, Papayu is one of the most similar stars to the Sun in terms of asteroseismic age and rotation period. Seventeen quarters of long-cadence data make it possible to use rotational modulation to obtain its rotation period of 21.1 ± 0.8 days (McQuillan et al. 2014).
3. Papayu is an asteroseismic gyrochronology benchmark anchoring the onset of weakened magnetic braking. It is at the transition where both spin-down models diverge, and a comparison with stars of similar ZAMS temperatures supports previous findings that weakened spin-down models are required to explain the ages and rotation periods of old solar-type stars. Its higher-than-average metallicity should allow tests of the composition dependence of the spin-down laws.

The discovery of Papayu shows that several exciting asteroseismic gyrochronology benchmarks are still waiting to be discovered in the Kepler data set. Recent analyzes of the Kepler asteroseismic survey sample have already increased the Kepler yield (Mathur et al. 2022), and upcoming efforts will provide the first homogeneous analysis of the entire Kepler asteroseismic short-cadence data set (M. Sayeed et al. 2024, in preparation). Furthermore, results from K2 and TESS (Metcalfe et al. 2020; Hatt et al. 2023) will continue to contribute to our understanding of the connections between stellar age, activity, and rotation. Lastly, as discussed in Saunders et al. (2024), we emphasize there is uncertainty on \approx Gyr scales regarding the exact point at which solar-like stars begin weakened braking. This is primarily driven by the lack of calibrating sources around the regime where weakened magnetic braking models diverge from the standard.

The process of precisely anchoring the onset of weakened magnetic braking starts with gathering various Papayu-like evolutionary statistics, with this study marking the beginning of this endeavor.

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Software: lightkurve (Barentsen et al. 2020), echelle (Hey & Ball 2020), DIAMONDS (Corsaro & De Ridder 2014), *kiauhoku* (Claytor et al. 2020), MESA (Paxton et al. 2011, 2013, 2015, 2018, 2019), Astropy (Astropy Collaboration et al. 2018), Matplotlib (Hunter 2007), NumPy (Oliphant 2006; Van Der Walt et al. 2011), and Pandas (pandas development team 2020)

Appendix

Figures 10 and 11 show posterior probability distributions for the DIAMONDS parameters used in fitting Papayu's stellar background and oscillation frequencies.

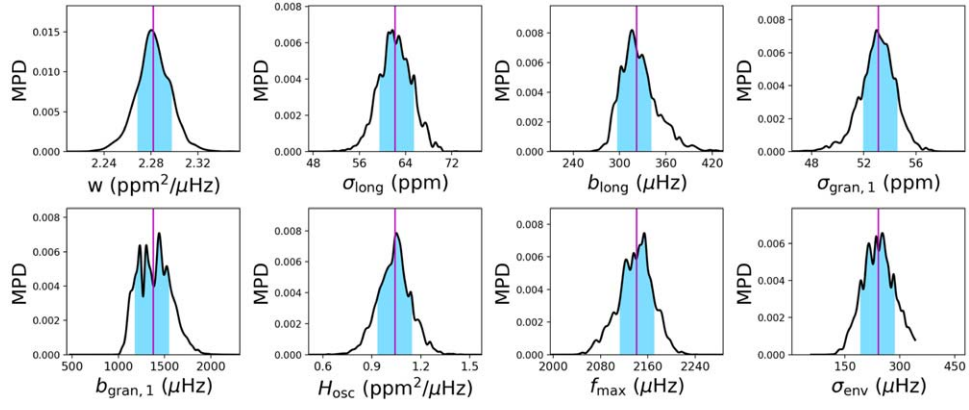


Figure 10. Posteriors for the granulation background fit using DIAMONDS (Corsaro & De Ridder 2014). Marginal probability distribution (MPD) is plotted against the various fit parameters, with the median value indicated by the vertical line.

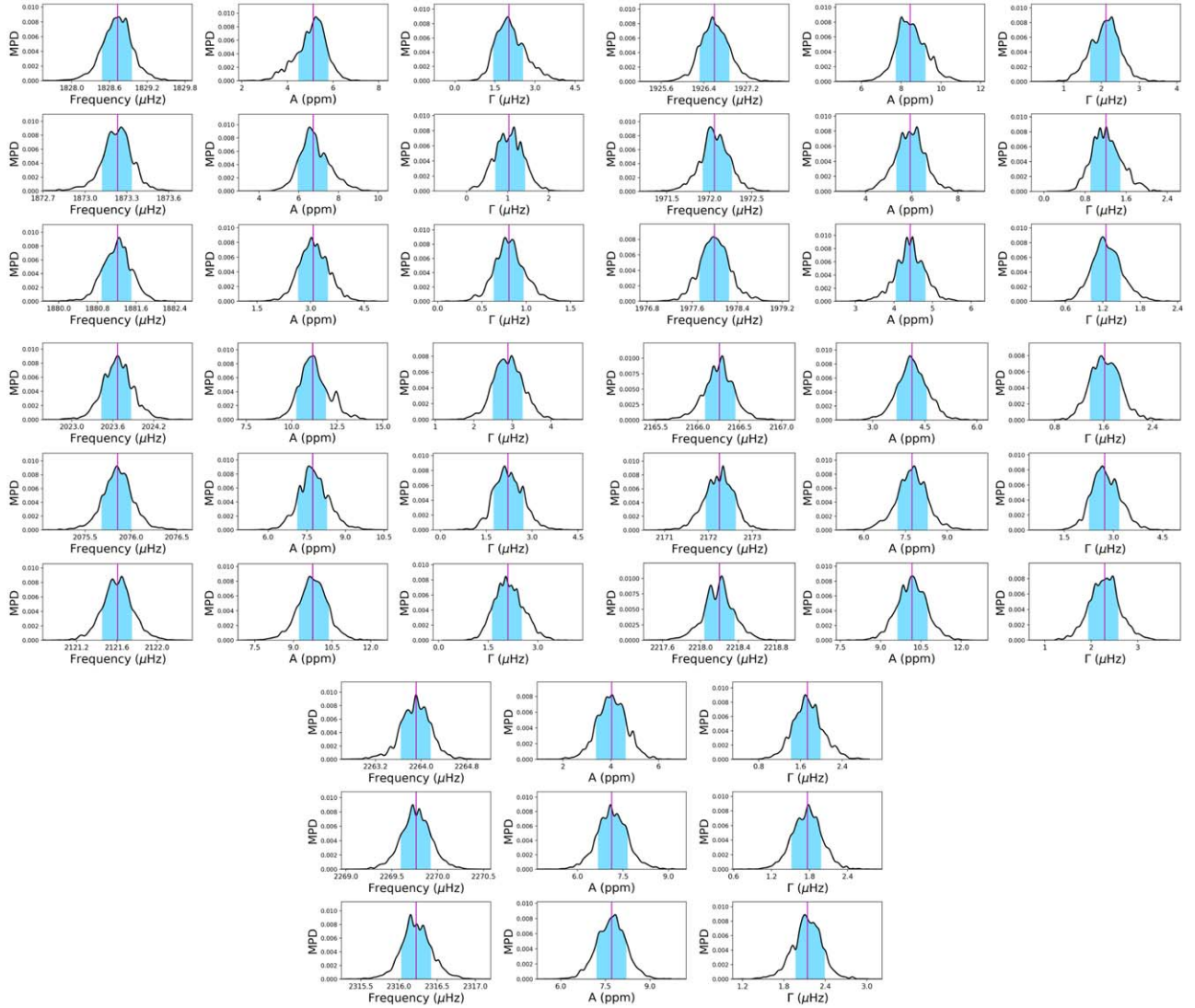
















Figure 11. Posteriors for the PeakBagging frequency modeling fit using DIAMONDS (Corsaro & De Ridder 2014). MPD is plotted against the various fit parameters, with the median value indicated by the vertical line.

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References

- Aguirre Børsen-Koch, V., Rørsted, J. L., Justesen, A. B., et al. 2022, *MNRAS*, **509**, 4344
- Amard, L., & Matt, S. P. 2020, *ApJ*, **889**, 108
- Astropy Collaboration, Price-Whelan, A. M., Sipőcz, B. M., et al. 2018, *AJ*, **156**, 123
- Ball, W. H., & Gizon, L. 2017, *A&A*, **600**, A128
- Barentsen, G., Hedges, C., Vinicius, Z., et al. 2020, KeplerGO/lightcurve: Lightcurve v1.9.0, Zenodo, doi:10.5281/zenodo.3687308
- Barnes, S. A. 2007, *ApJ*, **669**, 1167
- Barnes, S. A., Weingrill, J., Fritzewski, D., Strassmeier, K. G., & Platais, I. 2016, *ApJ*, **823**, 16
- Bellinger, E. P. 2019, *MNRAS*, **486**, 4612
- Berger, T. A., Huber, D., Gaidos, E., & van Saders, J. L. 2018, *ApJ*, **866**, 99
- Berger, T. A., Huber, D., van Saders, J. L., et al. 2020, *AJ*, **159**, 280
- Chaplin, W. J., Basu, S., Huber, D., et al. 2014, *ApJS*, **210**, 1
- Chaplin, W. J., Kjeldsen, H., Bedding, T. R., et al. 2011b, *ApJ*, **732**, 54
- Chaplin, W. J., Kjeldsen, H., Christensen-Dalsgaard, J., et al. 2011a, *Sci*, **332**, 213
- Chaplin, W. J., & Miglio, A. 2013, *ARA&A*, **51**, 353
- Charbonneau, P. 2010, *LRSP*, **7**, 3
- Choi, J., Dotter, A., Conroy, C., et al. 2016, *ApJ*, **823**, 102
- Christensen-Dalsgaard, J. 1984, in Proc. Workshop on Space Research Prospects in Stellar Activity and Variability, ed. A. Manganey & F. Praderie (Paris: Obs. Paris-Meudon), 11
- Christensen-Dalsgaard, J. 2008, *Ap&SS*, **316**, 113
- Claret, A., & Torres, G. 2018, *ApJ*, **859**, 100
- Clayton, Z. R., van Saders, J. L., Santos, A. G., et al. 2020, *ApJ*, **888**, 43
- Corsaro, E., & De Ridder, J. 2014, *A&A*, **571**, A71
- Creevey, O. L., Metcalfe, T. S., Schultheis, M., et al. 2017, *A&A*, **601**, A67
- Curtis, J. L., Agüeros, M. A., Matt, S. P., et al. 2020, *ApJ*, **904**, 140
- David, T. J., Angus, R., Curtis, J. L., et al. 2022, *ApJ*, **933**, 114
- Davies, G. R., Chaplin, W. J., Farr, W. M., et al. 2015, *MNRAS*, **446**, 2959
- Demarque, P., Guenther, D. B., Li, L. H., Mazumdar, A., & Straka, C. W. 2008, *Ap&SS*, **316**, 31
- Douglas, S. T., Agüeros, M. A., Covey, K. R., & Kraus, A. 2017, *ApJ*, **842**, 83
- Durney, B. R., & Latour, J. 1978, *GApFD*, **9**, 241
- García, R. A., & Ballot, R. A. 2019, *LRSP*, **16**, 4
- García, R. A., Ceillier, T., Salabert, D., et al. 2014, *A&A*, **572**, A34
- Gilliland, R. L., Brown, T. M., Christensen-Dalsgaard, J., et al. 2010a, *PASP*, **122**, 131
- Gilliland, R. L., Jenkins, J. M., Borucki, W. J., et al. 2010b, *ApJL*, **713**, L160
- Grec, G., Fossat, E., & Pomerantz, M. A. 1983, *SoPh*, **82**, 55
- Green, G. M., Schlafly, E. F., Finkbeiner, D. P., et al. 2015, *ApJ*, **810**, 25
- Hall, O. J., Davies, G. R., van Saders, J., et al. 2021, *NatAs*, **5**, 707
- Harvey, J. 1985, in Proc. ESA Workshop 235, Future Missions in Solar, Heliospheric & Space Plasma Physics, ed. E. Rolfe & B. Battrick (Noordwijk: ESA), 199
- Hatt, E., Nielsen, M. B., Chaplin, W. J., et al. 2023, *A&A*, **669**, A67
- Hey, D., & Ball, W. 2020, Echelle: Dynamic Echelle Diagrams for Asteroseismology, v1.4, Zenodo, doi:10.5281/zenodo.3629933
- Howard, A. W., Johnson, J. A., Marcy, G. W., et al. 2010, *ApJ*, **721**, 1467
- Huber, D., Chaplin, W. J., Christensen-Dalsgaard, J., et al. 2013, *ApJ*, **767**, 127
- Huber, D., Bryson, S. T., Haas, M. R., et al. 2016, *ApJ*, **224**, 2
- Hunter, J. D. 2007, *CSE*, **9**, 90
- Husser, T. O., Wende-von Berg, S., Dreizler, S., et al. 2013, *A&A*, **553**, A6
- Isaacson, H., & Fischer, D. 2010, *ApJ*, **725**, 875
- Johnson, J. A., Petigura, E. A., Fulton, B. J., et al. 2017, *AJ*, **154**, 108
- Kraft, R. P. 1967, *ApJ*, **150**, 551
- Lehtinen, J. J., Spada, F., Käpylä, M. J., Olsper, N., & Käpylä, P. J. 2020, *NatAs*, **4**, 658
- Lindegren, L., Klioner, S. A., Hernández, J., et al. 2021, *A&A*, **649**, A2
- Lorenzo-Oliveira, D., Freitas, F. C., Meléndez, J., et al. 2018, *A&A*, **619**, A73
- Mamajek, E. E., & Hillenbrand, L. A. 2008, *ApJ*, **687**, 1264
- Mathur, S., García, R. A., Breton, S., et al. 2022, *A&A*, **657**, A31
- McQuillan, A., Aigrain, S., & Mazeh, T. 2013, *MNRAS*, **432**, 1203
- McQuillan, A., Mazeh, T., & Aigrain, S. 2014, *ApJS*, **211**, 24
- Meibom, S., Barnes, S. A., Platais, I., et al. 2015, *Natur*, **517**, 589
- Metcalfe, T. S., Creevey, O. L., & Christensen-Dalsgaard, J. 2009, *ApJ*, **699**, 373
- Metcalfe, T. S., Creevey, O. L., Doğan, G., et al. 2014, *ApJS*, **214**, 27
- Metcalfe, T. S., Egeland, R., & van Saders, J. 2016, *ApJ*, **826**, L2
- Metcalfe, T. S., & van Saders, J. 2017, *SoPh*, **292**, 126
- Metcalfe, T. S., van Saders, J. L., Basu, S., et al. 2020, *ApJ*, **900**, 154
- Morel, P., & Thévenin, F. 2002, *A&A*, **390**, 611
- Noyes, R. W., Hartmann, L. W., Baliunas, S. L., Duncan, D. K., & Vaughan, A. H. 1984, *ApJ*, **279**, 763
- Olliphant, T. E. 2006, A Guide to NumPy (Trelgol Publishing USA), <https://web.mit.edu/dvp/Public/numpybook.pdf>
- pandas development team T. 2020, pandas-dev/pandas: Pandas, v2.2.2, Zenodo, doi:10.5281/zenodo.3509134
- Paxton, B., Bildsten, L., Dotter, A., et al. 2011, *ApJS*, **192**, 3
- Paxton, B., Cantiello, M., Arras, P., et al. 2013, *ApJS*, **208**, 4
- Paxton, B., Marchant, P., Schwab, J., et al. 2015, *ApJS*, **220**, 15
- Paxton, B., Schwab, J., Bauer, E. B., et al. 2018, *ApJS*, **234**, 34
- Paxton, B., Smolec, R., Schwab, J., et al. 2019, *ApJS*, **243**, 10
- Petigura, E. A. 2015, PhD thesis, Univ. of California, Berkeley
- Petigura, E. A., Howard, A. W., Marcy, G. W., et al. 2017, *AJ*, **154**, 107
- Réville, V., Brun, A. S., Matt, S. P., Strugarek, A., & Pinto, R. F. 2015, *ApJ*, **798**, 116
- Santos, A. R. G., Breton, S. N., Mathur, S., & García, R. A. 2021, *ApJS*, **255**, 17
- Saunders, N., van Saders, J. L., Lyttle, A. J., et al. 2024, *ApJ*, **962**, 138
- Schlegel, D. J., Finkbeiner, D. P., & Davis, M. 1998, *ApJ*, **500**, 525
- Serenelli, A., Johnson, J., Huber, D., et al. 2017, *ApJS*, **233**, 23
- Silva Aguirre, V., Davies, G. R., Basu, S., et al. 2015, *MNRAS*, **452**, 2127
- Silva Aguirre, V., Lund, M. N., Antia, H. M., et al. 2017, *ApJ*, **835**, 173
- Skrutskie, M. F., Cutri, R. M., Stiening, R., et al. 2006, *AJ*, **131**, 1163
- Skumanich, A. 1972, *ApJ*, **171**, 565
- Smith, J. C., Stumpe, M. C., Cleve, J. E. V., et al. 2012, *PASP*, **124**, 1000
- Stassun, K. G., Collins, K. A., & Gaudi, B. S. 2017, *AJ*, **153**, 136
- Stassun, K. G., Corsaro, E., Pepper, J. A., & Gaudi, B. S. 2018, *AJ*, **155**, 22
- Stassun, K. G., & Torres, G. 2016, *AJ*, **152**, 180
- Stumpe, M. C., Smith, J. C., Cleve, J. E. V., et al. 2012, *PASP*, **124**, 985
- Tayar, J., Clayton, Z. R., Huber, D., & van Saders, J. 2022, *ApJ*, **927**, 31
- Tayar, J., Somers, G., Pinsonneault, M. H., et al. 2017, *ApJ*, **840**, 17
- Thompson, S. E., Caldwell, D. A., Jenkins, J. M., et al. 2016, Kepler Data Release 25 Notes (KSCI-19065-002)
- Townsend, R. H. D., & Teitler, S. A. 2013, *MNRAS*, **435**, 3406
- Ulrich, R. K. 1986, *ApJL*, **306**, L37
- Van Der Walt, S., Colbert, S. C., & Varoquaux, G. 2011, *CSE*, **13**, 22
- van Saders, J. L., Ceillier, T., Metcalfe, T. S., et al. 2016, *Natur*, **529**, 181
- van Saders, J. L., & Pinsonneault, M. H. 2013, *ApJ*, **776**, 67
- Vogt, S. S., Allen, S. L., Bigelow, B. C., et al. 1994, *Proc. SPIE*, **2198**, 362
- Weber, E. J., & Davis, L. J. 1967, *ApJ*, **148**, 217
- White, T. R., Bedding, T. R., Stello, D., et al. 2011, *ApJ*, **743**, 161
- Wright, N. J., Drake, J. J., Mamajek, E. E., & Henry, G. W. 2011, *ApJ*, **743**, 48
- Wright, N. J., Newton, E. R., Williams, P. K. G., Drake, J. J., & Yadav, R. K. 2018, *MNRAS*, **479**, 2351
- Yee, S. W., Petigura, E. A., & von Braun, K. 2017, *ApJ*, **836**, 77
- Zinn, J. C., Pinsonneault, M. H., Huber, D., et al. 2019, *ApJ*, **885**, 166