

# 1 First-principles study of MXene properties with 2 varying hydrofluoric acid concentration

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## 10 SUMMARY

11 With varying hydrofluoric acid (HF) concentrations under three etching conditions, we presented a comparative study of the  
12 effects of both the ordered and randomly mixed terminated  $Ti_3C_2T_x$  surfaces with a wide variation of O/OH/F  
13 stoichiometry on the thermodynamic stability and electronic properties. Regardless of the HF concentration, a OH-rich surface  
14 is found to be thermodynamically stable and the electrical conductivity of  $Ti_3C_2T_x$  is substantially affected by the OH  
15 concentration. The charge density difference and electron localization function demonstrated a significant electron localization  
16 at the hydroxyl group on the O/OH/F mixed terminated surface, which could yield a locally induced dipole on the surface that  
17 renders favorable reaction sites on the functionalized surface. In addition, a large tunability in the work function ( $\Delta\Phi \sim 3.5$  eV)  
18 is predicted for  $Ti_3C_2T_x$ . These findings provide a pathway for strategically tuning the electronic and structural properties of  
19  $Ti_3C_2$  MXenes etched with HF.

## 20 INTRODUCTION

21 The active research of two-dimensional (2D) materials conducted since the successful exfoliation of graphene<sup>1</sup> in  
22 2004 has led to the discovery of a new, emerging class of 2D materials consisting of carbides and nitrides of transition metals,  
23 known as MXenes<sup>2</sup>. MXenes are two-dimensional materials with the general formula  $M_{n+1}X_nT_x$ , where M is an early transition  
24 metal (e.g., Ti, V, Cr), X is carbon, nitrogen or carbonitride, and T is the surface termination group comprised of O, OH, F,  
25 and/or Cl<sup>2</sup>. Due to their compelling physical, electronic, and chemical properties, MXenes have attracted immense theoretical  
26 and experimental research interests in a variety of applications, such as Li-ion batteries<sup>3,4</sup>, gas sensors<sup>5</sup>, hydrogen storage  
27<sup>6</sup>, and thermoelectrics<sup>7</sup>. Among those studies, nearly 70 % have been devoted to  $Ti_3C_2T_x$ , the first ever experimentally  
28 synthesized MXene<sup>8</sup>. It is considered the most comprehensively studied MXene to date.

29  $Ti_3C_2T_x$  can be selectively etched from its MAX phase with hydrofluoric acid (HF), where A is an A group element  
30 typically from groups 13 and 14 of the periodic table (A=Al for  $Ti_3C_2T_x$ )<sup>8</sup>. Due to the highly reactive Ti surface after etching,  
31 the exfoliated  $Ti_3C_2T_x$  often consists of randomly distributed surface functional groups (*i.e.*, O, OH, F), which are collectively  
32 expressed as  $T_x$ <sup>9</sup>. However, due to the complexity and computational cost of simulating mixed terminated surfaces, most of  
33 the theoretical studies have considered either the bare  $Ti_3C_2$ <sup>10,11</sup> or the uniformly terminated  $Ti_3C_2T_x$  with a single functional  
34 group<sup>4,7,12-14</sup>. This is often regarded as the first- and second-generation model of MXenes<sup>15</sup>. Early experimental efforts such  
35 as powder X-ray diffraction (XRD)<sup>8</sup>, high-resolution transmission electron microscopy (TEM)<sup>8,9,16</sup>, and X-ray atomic pair  
36 distribution function (PDF)<sup>17</sup> were used to gain insight into the distribution of the functional group compositions. However,  
37 each method was hindered by its insensitivity to hydrogen, which is essential for understanding the surface termination<sup>15</sup>.  
38 Thus, using the atomic pair distribution function supported by a high-quality neutron total scattering method, Wang et al.<sup>15</sup>  
39 obtained the first resolution of the  $Ti_3C_2T_x$  structure synthesized under different conditions and proposed a multilayered  
40 structural model of  $Ti_3C_2T_x$  to be the next-generation model of MXenes.

41 Inspired by the work from Wang et al., several theoretical investigations focused on the effect of mixed functional  
42 group terminations (O, OH, F). Caffrey<sup>18</sup> proposed an empirical model to study the change in the structural and  
43 electrochemical properties of the mixed terminated  $Ti_3C_2T_x$  and  $V_2CT_x$  structures compared to that of the uniformly terminated  
44 surfaces. According to the Caffrey study, the empirical model reproduces the lattice parameter, electron density of states, and  
45 work function that are consistent with experimental data. To date, the most comprehensive study on the distribution and  
46 composition of the surface functionalization of 2D MXenes employing the cluster expansion method was conducted by  
47 Ibragimova et al.<sup>19</sup>. In that study, the optimum O:OH:F composition was 50:25:25 under standard hydrogen electrode (SHE)  
48 conditions, with similar distribution patterns that were unaffected by thickness and type of MXene.

49 Nevertheless, the design pathway for tuning the mixed surface termination is still absent from the literature. Using  
50 energy dispersive X-ray spectroscopy (EDX) in the PDF characterization, Wang et al.<sup>15</sup> estimated the average atomic ratio of  
51 O:F in the multilayer  $Ti_3C_2T_x$  sample to be 0.85 and 1.4 when etched with 48 wt. % and 10 wt. % HF, respectively. Based on  
52 those O:F ratios, Wang et al. derived the stoichiometry of  $T_x$ , which is equal to  $O_{0.1}(OH)_{0.8}F_{1.1}$  and  $O_{0.13}(OH)_{1.04}F_{0.83}$ . In addition,  
53 the overall crystallinity and ordering are also affected by the concentration of HF. A higher concentration of HF yields higher  
54 F composition in the surface termination. Intuitively, this is consistent with an increased probability of F being available to  
55 terminate the freshly etched surface of Ti with higher HF concentrations. Hence, inspired by the findings from the Wang et al.  
56

1 and the cluster expansion studies <sup>19</sup>, the thermodynamic stability and electronic properties of  $Ti_3C_2T_x$  with different surface  
2 terminations were investigated.

## 3 RESULTS AND DISCUSSION

### 4 Formation energy and thermodynamic stability

5 As a benchmark for comparison with the O-, OH-, F-mixed terminated surfaces, a single-layer  $Ti_3C_2T_x$  sheet uniformly  
6 terminated with the O, OH, and F functional groups (*i.e.*,  $Ti_3C_2T_x$ , where  $T_x$  is represented by  $O_pOH_qF_r$  and  $p+q+r = 2.0$ ), is  
7 considered. A  $5\times5\times1$  supercell of  $Ti_3C_2$  was constructed, where both surfaces were populated with the functional group  
8 adatoms (Figure 1), corresponding to a total of 25 adsorption sites per surface. For a pristine 2D  $Ti_3C_2$  surface, there are three  
9 possible high-symmetry adsorption sites for the adatoms, namely Top (on top of the outer layer Ti atoms), face centered cubic  
10 (FCC) (on top of the middle layer Ti atoms in the hollow site), and hexagonal close packed (HCP) (on top of the C atoms in  
11 the hollow site) as shown in Figure 1a and 1b. The formation energies are calculated as  $E_{formation/adatom} =$   
12  $(E_{Ti_3C_2T_x} - E_{Ti_3C_2} - nE_{T_x})/n$ , where  $E_{formation/adatom}$ ,  $E_{Ti_3C_2T_x}$ ,  $E_{Ti_3C_2}$  and  $E_{T_x}$  are the formation energy (per adatom), total  
13 energy of  $Ti_3C_2T_x$ , total energy of the bare  $Ti_3C_2$  structure, and the energy of the functional group adatoms in the gas phase.  
14 The number of adsorbed adatoms are denoted by  $n$ .

15 From the reported literature <sup>4,7,10-14</sup>, most of the previous theoretical studies on a homogeneous or mixed-surface  
16 termination focused solely on a fully-terminated surface. In addition, they have found that the FCC site is the most favorable  
17 adsorption site for the full termination of O/OH/F single species. Complementary to these prior studies <sup>4,7,10-13</sup>, a basic  
18 understanding of the change in the thermodynamic stability transitioning from the bare  $Ti_3C_2$  structure to fully terminated  
19  $Ti_3C_2T_x$  structure provides valuable insight into the influential factors of surface termination. By gradually increasing the percent  
20 coverage of adatoms on the bare  $Ti_3C_2$  surface the thermodynamic stability of  $Ti_3C_2T_x$  can be systematically evaluated. The  
21 preferred binding sites for different single-adatom coverages can aid in constructing the mixed terminated  $Ti_3C_2T_x$  surface.  
22 Among all the  $Ti_3C_2T_x$  distributions considered at different percentage coverages, the adatoms preferred the FCC sites,  
23 regardless of the percent coverage on the  $Ti_3C_2$  surface.

24 In this study, an ordered and random distribution is considered when terminating the  $Ti_3C_2$  surface, wherein the  
25 surface adatoms are distributed in an orderly manner such that each O, OH and F termination species is located close to each  
26 other forming a cluster-like region for the ordered distribution. This is in contrast to the randomly distributed adatoms for the  
27 random distribution (Figure 1c). The random model was constructed by assigning the FCC sites adatoms and performing  
28 structural optimization. The details of the model construction can be found in the supplemental information (SI). The surface  
29 distributions presented here are the representative cases with the lowest energy for each stoichiometry.

30 Table 1 shows the formation energy of various surface terminations under three different etching conditions. The  
31 atomic ratio O:F corresponds to the concentration of HF during etching, representing 48 and 10 wt. % for O:F = 0.85 and 1.4,  
32 respectively (as adapted from Wang *et. al* <sup>15</sup>). Employing the same terminology, the O:F ratio that mimics a weak acidic  
33 environment (~ 5 wt. % HF), is derived from the stoichiometry of  $T_x = O_{1.04}(OH)_{0.48}F_{0.48}$ . Due to the limitation of the simulation  
34 cell, the O:F ratios considered in this work are 0.8, 1.3, and 3.2 for 48, 10, and 5 wt. % HF, respectively.

35 In previous work <sup>15</sup>, the O:F ratio was obtained by considering all oxygen atoms present in  $Ti_3C_2T_x$ , including those  
36 associated with the hydroxyl group (OH). In addition, the OH fraction was explicitly set as a constant in the derivation of the  
37  $T_x$  stoichiometry. In this work, to study the effect of varying the OH fraction, the F fraction is fixed, while the O and OH fractions  
38 are varied. This results in two different surface coverages: OH fraction that is either higher (OH-rich) or lower (OH-poor) than  
39 the O fraction (Table 1).

40 From DFT calculations, the  $Ti_3C_2T_x$  stability is more sensitive to the relative stoichiometry of O, OH and F adatoms  
41 on the surface rather than the local site distribution of adatoms (Table 1). Using gas phase  $O_2/F_2$  and isolated OH as the  
42 reference energy, our calculations show a decreasing trend in the formation energy as the OH fraction decreases. As shown  
43 in Table 1, the OH-rich coverages are thermodynamically stable, regardless of the HF concentration. Furthermore, the ordered  
44 and random terminated structures are found to be degenerate in terms of formation energy (Table 1). Both the ordered and  
45 random distribution of O/OH/F adatoms (Figure 1c) exhibit nearly similar thermodynamic stability, where the structures with  
46 random distribution are ~0.01 - 0.04 eV more stable than that of the ordered distribution (Table 1).

47 Based on the thermodynamic stability study of the partially-terminated surfaces, local clustering of O adatoms tends  
48 to be less stable (~0.2-0.3 eV/adatom), especially for the lower coverages (*i.e.*,  $\leq 60\%$ ). This trend is less profound in partial  
49 surface termination with local clustering of OH and F adatoms, especially for F adatoms (Figure S1 in SI). Hence, the difference  
50 in the formation energy between the ordered and random distribution is determined by local O-clustering especially at low  
51 concentrations, as seen in the nearest neighbor analysis (see Figure S5-S10 in SI). The O-O fraction is predominantly the first  
52 nearest neighbor in the ordered distribution instead of the second and third nearest neighbors, implying a signature of local  
53 clustering of O adatoms could be a factor in distinguishing an ordered from a random distribution.

54 Based on the partial termination study, we have found an empirical relationship to predict the relative stability  
55 between the structures. The O adatom has a stronger effect on the overall stability of  $Ti_3C_2T_x$ , with an energy difference of  
56 ~0.6 eV/adatom between the bare and full termination surfaces, followed by OH with a ~0.2 eV/adatom difference. In contrast,  
57 F adatoms do not significantly vary the thermodynamic stability of the structure, where nearly a constant of -5.3 eV/adatom  
58 formation energy is observed when transitioning from the bare to full F-terminated surface. This could explain the subtle  
59 difference in formation energy (~0.2 - ~0.3 eV) between the structures etched with different wt. % HF, as the thermodynamic  
60 stability is insensitive to the variation of F concentration, but mostly determined by the relative content of O and OH adatoms.  
61

1 Due to the fact that MXene surface is highly tunable with a variation possibility on the M, X, and  $T_x$  positions, therefore multiple  
2 factors can be readily tuned in MXenes surface chemistry. Our partial termination study (shown in Figure S1) also  
3 demonstrated the importance of local clustering of O/OH/F species, which motivates the construction of local ordered  
4 distribution in surface termination. In Figure S1, it is found that for partial F termination, the ordered and random surface  
5 termination are energetically degenerate. For OH termination, a subtle energy difference of less than 0.05 eV/adatom between  
6 the ordered and distribution is observed when OH percent coverage is increased, indicating competitive stability between the  
7 two types of distributions. While the O termination exhibits a slightly different trend where a local clustering of O at lower  
8 percent coverage is less favorable, whereas at higher percent coverage (framed in red), the energy difference between  
9 ordered and random surface termination is less than ~0.10 eV/adatom. This observation highlights the importance of studying  
10 the different types of surface distributions of termination groups, besides the variation in mixed stoichiometry of surface  
11 termination.

### 12 **Electronic properties**

#### 13 **Electron density of states (eDOS)**

14 To identify the general features of the electronic structure among different O/OH/F terminations on the  $Ti_3C_2T_x$   
15 structure, the electron density of states (eDOS) is calculated in an energy range near the Fermi level,  $E_F$ . Overall, the metallic  
16 nature of  $Ti_3C_2T_x$  is evident due to the presence of finite electronic states at Fermi level ( $E_F$ ) regardless of the terminating  
17 adatoms and their distribution. From Figure 2, the Ti 3d orbital dominates the electronic states throughout the energy range  
18 from -0.5 to 0.5 eV within the vicinity of  $E_F$  in all cases. This indicates that the 3d orbital would contribute the most to the  
19 tunneling current in STM characterization. The eDOS at  $E_F$ , i.e.,  $N(E_F)$  of all structures is within the range of ~30 to ~50  
20 states/eV (see Figure S3-S4). Wang et. al<sup>20</sup> studied the surface properties of  $Ti_3C_2T_x$  and observed the tunable conductivity  
21 with different surface termination species. With uniform termination species, OH termination has the highest electronic states  
22 at  $E_F$ . Our calculation shows a consistent observation where the  $N(E_F)$  of a uniformly O-terminated surface is the highest  
23 compared to that of an O/F-terminated surface.

24 Figure 2 shows the partial density of states (pDOS) of  $Ti_3C_2T_x$  with ordered surface distribution etched with 48 wt. %  
25 HF. Interestingly,  $Ti_3C_2T_x$  OH-rich surfaces consist of a larger magnitude of eDOS at  $E_F$ , ( $N(E_F)$  ~50 states/eV) than that of  
26 OH-poor surfaces ( $N(E_F)$  ~30 states/eV). Similar results were observed for randomly terminated surfaces. While there are no  
27 clear empirical relationships between the O/OH/F concentrations and  $N(E_F)$ , it is noted that in general, the surface with excess  
28 OH concentration has an upper limit of  $N(E_F)$  saturated at ~50 states/eV, compared to the lower limit of  $N(E_F)$  around ~30  
29 states/eV with excess O concentration. Thus, lower electronic conductivity is expected for surfaces with lower OH  
30 concentration. Of note, the  $N(E_F)$  of the uniformly OH-terminated  $Ti_3C_2$  surface is nearly twice (~90 states/eV) as much as the  
31 uniformly F/O-terminated surfaces (see Figure S2 in SI). However, when the mixed terminated O/OH/F surface is present,  
32 the  $N(E_F)$  is substantially suppressed to fall within the range closer to that of the F/O-terminated surface (i.e., ~40 states/eV).  
33 Thus, we believe the increase of the electron states at the Fermi energy level  $E_F$  with increasing OH content could be explained  
34 by the  $N(E_F)$  of the uniformly O/OH/F-terminated surface. Upon full termination of uniform OH species, we observe that the  
35  $N(E_F)$  is almost twice as large compared to that of the full O/F termination. Thus, the OH-rich surface should resemble a  
36 surface closer to that of the uniformly OH-terminated surface, resulting in an increase in  $N(E_F)$ . The change in  $N(E_F)$  indicates  
37 the tunable electronic conductivity of  $Ti_3C_2T_x$  via varying the OH/O adatom concentrations on the surface. Higher electronic  
38 conductivity, in principle, could potentially be achieved with higher OH concentration.

#### 41 **Charge density difference and electron localization function (ELF)**

42 To further study the charge distribution and the electronic properties of different surface terminations, charge density  
43 difference and electron localization function (ELF) calculations were performed. A charge density difference plot was  
44 constructed by taking the difference in electron density between the  $Ti_3C_2T_x$ , pristine  $Ti_3C_2$ , and the adatoms to identify the  
45 local charge accumulation at the terminated sites. The ELF is often employed to describe and visualize the chemical bonding  
46 in molecules or solids<sup>21</sup> and is bound between 0 and 1, where ELF = 1 indicates perfect localization and ELF = 0.5 corresponds  
47 to uniform electron gas.

48 Figure 3 shows the charge density difference of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$  with an ordered and random surface  
49 distribution. In both cases, charge depletion and accumulation tend to happen at the surface adatoms. Large inhomogeneity  
50 of charge depletion/accumulation (ranging from ~-0.1 e/Å<sup>3</sup> to ~+0.1 e/Å<sup>3</sup>) was observed, as expected, particularly for OH  
51 adatom, due to the high polarity of the hydroxyl group. Overall, the dominant charge depletion region is localized at the  
52 hydrogen of the hydroxyl group, whereas the charge accumulation is distributed around the electronegative O and F sites of  
53 adatoms. Interestingly, from the charge density difference analysis, the electron accumulation/depletion regions are found to  
54 be similar among different surface distributions for the termination O/OH/F adatoms (Figure 3). Therefore, the difference in  
55 the local distribution of the functional adatoms (i.e., ordered vs. random) does not contribute significantly to the charge density  
56 difference. The plot of ELF in Figure 4 gives further insight into the local bonding nature of  $Ti_3C_2T_x$ . Similar to the charge  
57 density difference distribution, the ELF is very similar to both the ordered and random termination. In contrast to this intriguing  
58 similarity, the spatial distribution of electrons in the ELF plot is highly localized at the hydroxyl group (ELF ≈ 1), particularly at  
59 the H atoms (Figure 4), but not at the O or F adatoms. This is an intriguing observation since the electron localization at the  
60 hydroxyl group, regardless of whether the termination is ordered or random, could give rise to the local imbalance surface  
61 charge distribution relative to O and F adatoms on  $Ti_3C_2T_x$  surfaces as favorable reaction sites, as also indicated by the charge  
62 density difference analysis (Figure 3). As a result, this will generate unique locally induced dipoles on the  $Ti_3C_2T_x$  surfaces

1 (Figure 3), which might yield important applications in surface chemistry (e.g., gas adsorption and catalysis) as favorable  
2 reaction sites on the functionalized  $Ti_3C_2T_x$  surface.

3 **Work function**

4 The effects of the surface distribution of different termination groups on the work function for  $Ti_3C_2T_x$  were also  
5 studied. The work function,  $\Phi$ , was calculated by computing the planar average of the electrostatic potential along the surface  
6 normal direction, as:  $\Phi = eV_{\text{vacuum}} - E_{\text{Fermi}}$ , where  $V_{\text{vacuum}}$  is the vacuum potential and  $E_{\text{Fermi}}$  is the Fermi level. Regardless of the  
7 HF concentration, there was an increasing trend in the work function with increasing O concentration.

8 Among all mixed terminated structures, the minimum work function corresponds to the structures with the most OH  
9 content for each HF concentration. As expected, as the acidic environment (or HF concentration) increases, the work functions  
10 of  $Ti_3C_2T_x$  become gradually higher due to the increase in terminated F content. All the calculated work functions are within  
11 the upper (6.04 eV with full O termination) and lower limit (1.81 eV with full OH termination) of the uniformly terminated (O,  
12 OH, F) surfaces. With fixed F percentage coverage on surface termination, a high work function of  $Ti_3C_2T_x$  is tunable by  
13 decreasing OH termination and increasing O termination. As shown in Figure 5, a linear dependence of the work function with  
14 respect to O/OH content is found with the varying stoichiometry. This observed trend is consistent with the reported literature  
15 18,22.

16 Despite having different types of surface distributions, no noticeable difference is seen in the work function between  
17 the ordered and randomly terminated surfaces, analogous to their relative thermodynamic stability (Table 1). Hence, the work  
18 function is, in fact, a function of stoichiometry and is insensitive to the surface termination distribution. Thus, it is possible to  
19 experimentally estimate the OH/O composition of  $Ti_3C_2T_x$  by measuring the work function. Likewise, the work function can be  
20 tuned to the desired value by etching with the appropriate HF concentration. The large tunability in work function for  $Ti_3C_2T_x$   
21 (i.e.,  $\Delta\Phi \sim 3.5$  eV) due to different wt. % of HF (Figure 5) suggests that  $Ti_3C_2T_x$  can be strategically modified to suit future  
22 electronic device application needs.

23 **CONCLUSION**

24 A systematic analysis was carried out to determine the effects of varying HF etching concentrations on the  
25 thermodynamic stability, electronic properties, and work function of  $Ti_3C_2T_x$  based on DFT calculations. Regardless of the HF  
26 concentration, surfaces with a larger fraction of OH termination were found to be thermodynamically stable. In addition, the  
27 ordered and randomly terminated  $Ti_3C_2T_x$  structures are degenerate in their thermodynamic stability, but subtle differences  
28 arise from the local clustering of the O functional group, which is found to be less favorable for all cases. Overall, a mixed  
29 termination of  $Ti_3C_2T_x$  does not alter the metallic nature of  $Ti_3C_2$  but the electronic conductivity varies significantly with OH  
30 termination concentration and could be potentially enhanced by increasing its content. In terms of bonding analysis, the charge  
31 density difference and ELF demonstrated a significant electron localization at the hydroxyl group, which could yield a locally  
32 induced dipole on the surface that could render unique surface chemistry. From DFT predictions, a large tunability in work  
33 function for  $Ti_3C_2T_x$  (i.e.,  $\Delta\Phi \sim 3.5$  eV) due to different HF wt.% is found. The work function was found to be a function of the  
34 O/OH/F stoichiometry of  $Ti_3C_2T_x$  and is independent of the distribution of surface terminations group. For all mixed terminated  
35 structures, the minimum work function corresponds to the structures with the most OH content for each HF concentration. As  
36 the acidic environment (or HF concentration) increases, the work functions of  $Ti_3C_2T_x$  becomes higher due to the increase in  
37 terminated F content. For fixed F content on surface termination, high work function of  $Ti_3C_2T_x$  is tunable by decreasing in OH  
38 termination and increasing in O termination. The authors believe this work provides valuable insight into the fundamental  
39 characterization of  $Ti_3C_2T_x$  and a useful design pathway to tune the material properties experimentally through the variation of  
40 HF concentration during MAX phase etching to produce the  $Ti_3C_2T_x$  MXene.

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46 **AUTHOR CONTRIBUTIONS**

47 Y.Z.C. and M.H. conducted the calculations and simulations. K.C.L designed the workflow and study. P.W. and K.C.L. secured  
48 funding and provided mentoring. Y.Z.C. wrote the manuscript. M.H., K.C.L, and P.W. revised the manuscript.

49 **DECLARATION OF INTERESTS**

50 The authors declare no competing interests.

51 **MAIN FIGURE TITLES AND LEGENDS**

52 **Figure 1.  $Ti_3C_2T_x$  crystal structure**

1 (A) Top view of the  $Ti_3C_2T_x$  crystal structure, where green: Top Ti, blue: Middle Ti, brown: C.  
 2 (B) Side view of the crystal structure of  $Ti_3C_2T_x$ , where green: Top Ti, blue: Middle Ti, brown: C.  
 3 (C) The demonstration of an ordered and random distribution of surface functional group adatoms, where red: O, green: F,  
 4 blue: OH.

5 **Figure 2. Electron density of states of  $Ti_3C_2T_x$  etched with 48 wt. % HF.**

6 The Fermi level,  $E_F$  is set to zero.

7 **Figure 3. Charge density difference of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$**

8 (A) Ordered termination surface of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$ .

9 (B) Random termination surfaces of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$ .

10 The cyan and yellow regions represent charge depletion and charge accumulation, respectively. Inset shows the 2D slice of  
 11 the charge density difference viewed from (100) plane with the isovalue ranging from -0.03 to 0.03 e/  $\text{\AA}^3$ . The region framed  
 12 in red in the inset indicates a local-induced dipole on the surface.

13 **Figure 4. Plot of electron localization function (ELF) of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$**

14 (A) Ordered termination surface of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$ .

15 (B) Random termination surfaces of  $Ti_3C_2O_{0.24}(OH)_{1.28}F_{0.48}$ .

16 The right panel shows a 2D slice of ELF viewed from the (100) plane. Red and blue regions in the right panel represent  
 17 localized and delocalized charge with an isovalue of 1 and 0 respectively.

18 **Figure 5. Work function of the  $Ti_3C_2T_x$  etched with different HF concentrations.**

19 The reference lines correspond to the calculated work function of the uniformly terminated surface. Color code: Red: O-  
 20 termination (6.03 eV), Green: F-termination (4.95 eV), Blue: OH-termination (1.81 eV).

21  
 22  
 23 **Table 1. Formation energy of various mixed terminated  $Ti_3C_2T_x$  structures with two different local site distribution of  
 24 O/OH/F adatoms (i.e., ordered vs. random).\***

System	$T_x$ stoichiometry	Formation energy/adatom (eV/adatom)	
		Ordered	Random
O:F= 0.8	$O_{0.24}OH_{0.64}F_{1.12}$	-5.46	-5.47
48 wt. % HF	$O_{0.40}OH_{0.48}F_{1.12}$	-5.37	-5.39
	$O_{0.64}OH_{0.24}F_{1.12}$	-5.23	-5.25
	$O_{0.80}OH_{0.08}F_{1.12}$	-5.12	-5.14
O:F=1.3	$O_{0.24}OH_{0.88}F_{0.88}$	-5.54	-5.55
10 wt. % HF	$O_{0.48}OH_{0.64}F_{0.88}$	-5.41	-5.43
	$O_{0.64}OH_{0.48}F_{0.88}$	-5.32	-5.34
	$O_{0.88}OH_{0.24}F_{0.88}$	-5.17	-5.19
O:F=3.2	$O_{0.24}OH_{1.28}F_{0.48}$	-5.65	-5.66

5 wt. % HF	$O_{0.48}OH_{1.04}F_{0.48}$	-5.53	-5.57
	$O_{0.72}OH_{0.80}F_{0.48}$	-5.40	-5.45
	$O_{1.04}OH_{0.48}F_{0.48}$	-5.20	-5.22

1 The O:F ratio corresponds to different HF concentrations according to Ref. <sup>15</sup>.

2 \*Thermal energy corrections are included in the calculation of the energy values in Table 1. As comparison, another reference  
3 energy model considering the experimental conditions following the Ref. <sup>19,22</sup> (including thermal energy corrections) is included  
4 in Table S1.

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9 **STAR METHODS**

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11 **RESOURCE AVAILABILITY**

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13 **Lead contact**

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15 Further information and requests for resources should be directed to the lead contact, Kah Chun Lau ([kahchun.lau@csun.edu](mailto:kahchun.lau@csun.edu)).

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17 **Materials availability**

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19 This study did not yield new unique reagents.

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21 **Data and Code Availability**

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- The simulation data that support the findings of this study are available from the lead contact upon reasonable  
25 request.
- This paper does not report original code.
- Any additional information required to analyze the data reported in this paper is available from the lead contact upon  
26 reasonable request.

27 **METHOD DETAILS**

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29 To model the 2D  $Ti_3C_2T_x$  structures, first-principles density functional theory (DFT) was employed as implemented  
30 in Vienna Ab initio Simulation Package (VASP) <sup>23,24</sup> for the spin-polarized electronic structure calculations. The projector-  
31 augmented-wave (PAW) pseudopotential method, together with the generalized gradient approximation (GGA) and the  
32 Perdew–Burke–Ernzerhof (PBE) exchange-correlation functional were used to represent the exchange-correlation effects in  
33 the DFT simulation <sup>25</sup>. The kinetic energy cutoff for the plane-wave basis was set at 500 eV. For the structural optimization  
34 and electronic calculations, the Brillouin zone was sampled using a  $\Gamma$ -centered grid of (3x3x1) where the convergence in the  
35 total energy (< 0.02 eV) is achieved with respect to the number of k points, while the electron density of states (DOS) was  
36 sampled using a grid of (9x9x1). The electronic self-consistent cycle energy convergence was set at  $1 \times 10^{-5}$  eV and the  
37 residual force on each atom is less than  $1 \times 10^{-4}$  eV/Å. The van der Waals interaction was included in all calculations using  
38 the semi-empirical approach of Grimme (DFT-D3) <sup>26</sup>. For the monolayer structure calculations, a vacuum slab of at least 15 Å  
39 along the z-axis was implemented to avoid interactions between mirror images. To further confirm the thermal stability of these  
40  $Ti_3C_2T_x$  (Table 1) at room temperature, ab initio molecular dynamics (AIMD) was performed on two selected structures for 10  
41 ps in the NVT ensemble employing Nosé–Hoover thermostats, and the results are shown in Figure S11. Our result shows no  
42 significant structure deformation or dissociation of termination groups upon heating at  $T = 300$  K, confirming the thermal  
43 stability of our structures at room temperature. All atomic structures and charge densities are visualized using VESTA<sup>27</sup>.

44

45 To compute the formation energy of various mixed terminated  $Ti_3C_2T_x$  structures with different local site distribution  
46 (i.e. ordered vs. random) of O/OH/F adatoms (Table 1), the reference energies of the adatoms are computed by considering  
47 the total energy and thermal energy corrections of a single molecule for OH and one-half of the total energy of  $O_2$  and  $F_2$  for  
48 O and F, respectively. In gas-phase calculations that employ DFT-GGA, the overestimation of the binding energy often arises  
49 due to the limitations in describing the exchange energy of first-row elements, such as molecular oxygen <sup>28,29</sup>. To overcome  
50 these effects, one could add an empirical correction to the DFT energy <sup>30</sup> or use a more accurate, albeit computationally  
51 expensive, hybrid or meta-GGA functional. In this study, we were interested in the qualitative difference between the adatom  
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1 adsorptions, thus, an empirical correction was not included. Since the calculated formation energies depend on the employed  
2 thermodynamic equations, to further justify the predicted formation energy of mixed terminated  $Ti_3C_2T_x$  structures highlighted  
3 in Table 1, another set of reference energies of the O/OH/F adatoms by considering the splitting of  $H_2O$  and HF proposed  
4 from Ref. 19,22 is adopted (Table S1). Despite using different reference energies of O/OH/F adatoms, some common trends  
5 can be found in Table 1 and S1. As indicated in Table 1 and S1, both the ordered and random distribution of O/OH/F adatoms  
6 on mixed terminated  $Ti_3C_2T_x$  surfaces exhibit nearly similar thermodynamic stability, regardless of different HF wt%.  
7 Interestingly in both cases (*i.e.*, Table 1 and S1), OH-rich surfaces are generally found to be thermodynamically stable in  
8 different HF concentrations, and therefore one can conclude that the presence of OH-rich surfaces at various mixed terminated  
9  $Ti_3C_2T_x$  is highly plausible. For conciseness, we only considered the thermodynamic stability referring to Table 1 in our  
10 discussion. The surface distributions presented here (Table 1 & S1) are the representative cases with the lowest energy  
11 among 2-3 different models per distribution (*i.e.*, ordered and random) for each stoichiometry.  
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#### 14 **SUPPLEMENTAL INFORMATION**

15 Document S1. Supplemental Information, methodology for constructing a random distribution surface, Figures S1–S11, Table  
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