



The Width of Magnetic Ejecta Measured near 1 au: Lessons from STEREO-A Measurements in 2021–2022

Noé Lugaz^{1,4} , Bin Zhuang¹ , Camilla Scolini^{1,2} , Nada Al-Haddad¹ , Charles J. Farrugia¹ , Réka M. Winslow¹ , Florian Regnault¹ , Christian Möstl³ , Emma E. Davies³ , and Antoinette B. Galvin¹

¹ Institute for the Study of Earth, Oceans, and Space, University of New Hampshire, Durham, NH, USA; [noe.lugaz@unh.edu](mailto: noe.lugaz@unh.edu)

² Solar-Terrestrial Centre of Excellence/SIDC, Royal Observatory of Belgium, Brussels, Belgium

³ Austrian Space Weather Office, GeoSphere Austria, Graz, Austria

Received 2023 October 12; revised 2023 December 6; accepted 2023 December 19; published 2024 February 20

Abstract

Coronal mass ejections (CMEs) are large-scale eruptions with a typical radial size at 1 au of 0.21 au but their angular width in interplanetary space is still mostly unknown, especially for the magnetic ejecta (ME) part of the CME. We take advantage of STEREO-A angular separation of 20°–60° from the Sun–Earth line from 2020 October to 2022 August, and perform a two-part study to constrain the angular width of MEs in the ecliptic plane: (a) we study all CMEs that are observed remotely to propagate between the Sun–STEREO-A and the Sun–Earth lines and determine how many impact one or both spacecraft in situ, and (b) we investigate all in situ measurements at STEREO-A or at L1 of CMEs during the same time period to quantify how many are measured by the two spacecraft. A key finding is that out of 21 CMEs propagating within 30° of either spacecraft only four impacted both spacecraft and none provided clean magnetic cloud-like signatures at both spacecraft. Combining the two approaches, we conclude that the typical angular width of an ME at 1 au is ∼20°–30°, or 2–3 times less than often assumed and consistent with a 2:1 elliptical cross section of an ellipsoidal ME. We discuss the consequences of this finding for future multi-spacecraft mission designs and for the coherence of CMEs.

Unified Astronomy Thesaurus concepts: Solar coronal mass ejections (310); Ejecta (453); Interplanetary physics (827)

1. Introduction

The properties of coronal mass ejections (CMEs) are known primarily from remote observations, such as coronagraphic and heliospheric imagers, and from in situ measurements from interplanetary probes that directly measure their magnetic field and plasma properties. By combining hundreds of in situ measurements, we have learned that the average radial size of a CME at 1 au is ∼0.30 au (Richardson & Cane 2010), which can be decomposed into ∼0.09 au for the dense sheath region that is present in ∼80% of CMEs (Salman et al. 2021) and ∼0.21 au (Lepping et al. 2005; Kilpua et al. 2017) for the magnetic ejecta (ME; sometimes referred to as magnetic obstacle). In this manuscript, we use the term CME to describe the whole eruption whether it is observed remotely or measured in situ. The term ICME has been used in the past to refer to the *interplanetary counterpart* of a CME. Because CMEs are now routinely measured in situ by Parker Solar Probe (Fox et al. 2016; Raouafi et al. 2023) and Solar Orbiter (Müller et al. 2020) at distances where they are also observed remotely by coronagraphs or heliospheric imagers, distinguishing between ICMEs and CMEs is meaningless. The term ME is meant to refer to the typically low- β , magnetically dominated ejecta portion of a CME and is favored as it has been used for decades (Cane et al. 1998; Lugaz et al. 2012). While some authors

prefer the term *magnetic obstacle*, in our opinion, it gives the impression that CMEs are *obstacles*, whereas in fact they are propagating and expanding structures. Magnetic clouds (MCs), which have well-defined properties as described first in Burlaga et al. (1981) are a special case of MEs, of which they represent 30%–40% of the measurements near 1 au (Richardson & Cane 2010; Jian et al. 2018).

The radial properties of CMEs and MEs are relatively well understood; in addition to their size, we know the average radial speed and radial expansion speed (Gopalswamy et al. 2015; Jian et al. 2018). The same is, however, not true for CME properties in the nonradial direction. Some recent work has focused on the magnitude of nonradial flows within CMEs and MEs (Al-Haddad et al. 2022), but a key aspect is the ME longitudinal extent. The latitudinal extent of CMEs and MEs can be estimated from remote observations from coronagraphs for the bright front and dark cavity. Estimates for the bright front have been 45°–70° (Hundhausen 1993; St. Cyr et al. 2000b; Yashiro et al. 2004), excluding partial and full halo. There has not been any dedicated statistical work on the latitudinal width of the dark cavity, as far as we are aware. Note that, throughout this work, we list the total width rather than the half-width as it is sometimes done for remote observations.

The angular extent (especially the longitudinal width) of MEs is a critical physical characteristic, both to develop more accurate models and theories of CMEs as well as for space weather forecasting to know (a) how far from the Sun–Earth line measurements can be made that are useful for CME forecasts, and (b) how far from the Sun–Earth line a CME must be propagating to impact Earth. Regarding CME models, we note that, under the flux rope paradigm, if the CME leg-to-leg angular extent is comparable to the nonradial cross section, flux

⁴ Corresponding author.

rope models would need to be significantly modified to account for the CME major radius of curvature and the magnetic flux balance within a *thick* flux rope. The angular extent is also a critical quantity to investigate ME coherence, an area of recent focus in CME studies (Owens et al. 2017; Scolini et al. 2023).

To determine the longitudinal width from remote observations, there are a few possible approaches. Under the assumption that CMEs can be represented by a sphere (the so-called ice cream cone model; Xue et al. 2005), the longitudinal and latitudinal widths are equal. A more common assumption is that the ME part of CMEs can be represented as an axially symmetric highly twisted magnetic flux rope with a *croissant* shape (see, for example, the graduated cylindrical shell (GCS) model of Thernisien et al. 2009). In situ studies have revealed that the inclination of MEs (a parameter that affects the latitudinal and longitudinal widths) is more or less random (Janvier et al. 2015), i.e., there is an equipartition of inclination angles. Under the assumption that the number of CMEs measured remotely is high enough, the average longitudinal width should be the same as the average latitudinal width.

For individual events, or small statistics, fitting and reconstruction techniques can be used to determine the CME angular width from remote observations. Howard & DeForest (2012) analyzed the 2008 December 12 CME and found latitudinal widths of 39° and 31° for the bright front and dark cavity, respectively, and longitudinal widths of 56° and 24° . For the same event, Byrne et al. (2010) found an angular width of the bright front (assumed to be spherically symmetric) to increase with distance from 45° to 60° from 0.05 to 0.2 au. Lugaz et al. (2010) developed a stereoscopic method that can return the CME longitudinal width, and it was found for two events that the longitudinal width decreases from 90° to 70° as the CME propagates to 0.5 au. Rodriguez et al. (2011) performed the analysis of 26 CMEs using the GCS model (see Thernisien et al. 2009; Thernisien 2011) during the early years of the STEREO missions (late 2007 to mid-2008). The average reconstructed longitudinal width was found to be 51° and the average latitudinal width was 36° . Zhao et al. (2017) studied 33 CMEs with the GCS model and found an average width of 48° (from leg to leg, i.e., ignoring the thickness of the CME). The same average value was found by Martinić et al. (2022) for 22 CMEs. Lastly, the KINCAT catalog of the Advanced Forecast For Ensuring Communications Through Space project has 122 CMEs with an average angular size (leg to leg) of 44° . We conclude this paragraph by noting two important points. First, in the inner heliosphere, researchers often rely on the self-similar expansion model of Davies et al. (2012) with a width of 60° or 90° , which is higher than the typical width of CMEs as reported from these past studies. Second, there is an overall consensus in the community that the CME angular width remains constant as the CME propagates. This is typically attributed to St. Cyr et al. (2000b; see their Section 8.5), which was a preliminary study of CME lateral expansion in the LASCO/C3 field of view. Whether or not this holds true for the ME and beyond 0.15 au (the end of LASCO C3 field of view) has not been investigated.

Multiple single-spacecraft in situ measurements combined with fitting techniques can be used statistically to derive the typical shape of MEs and CMEs (Janvier et al. 2014; Démoulin et al. 2016); however, this requires one to assume the CME and ME angular extents. In those works, the authors assume extents

of 70° and 60° , respectively, for the CME and ME. Without such an assumption, the analysis of the distribution of impact parameters resulted in a determination of the CME cross section to have an aspect ratio of about 3:1. We note that this corresponds to a longitudinal extent of $\sim 35^\circ$ based on the radial size of 0.21 au for a highly inclined flux rope. Multi-spacecraft measurements provide the only means by which to determine the longitudinal extent of MEs. Cane et al. (1998) investigated multi-spacecraft measurements from Helios and noted that “except for a potential event measured at 53°], there was only one case in which an ejecta was seen at two spacecraft separated by more than 40° . In addition, in a number of cases, two spacecraft were separated by less than 40° , but an ejecta was seen at only one spacecraft.” Kilpua et al. (2011) provided an overview of multi-spacecraft measurements both before the launch of STEREO as well as during the first two years of the missions when the two STEREO spacecraft were separated by 0° – 45° . This includes the well-studied 2007 November event (Farrugia et al. 2011) with three spacecraft measurements of the same ME over $(40^\circ, 8^\circ)$ separations, but also two cases in 2007 December with only one clear spacecraft measurement of the ME (but potentially inconclusive measurements at a second spacecraft) for total separations of $\sim 43^\circ$ – 44° , indicating that two spacecraft separated by 20° – 25° may not observe the same ME.

One of the most in-depth studies of CME longitudinal extent was performed by Good & Forsyth (2016) using a database of more than 400 CMEs measured in the inner heliosphere by MESSENGER, Venus Express, Wind/ACE, and STEREO-A and STEREO-B. Their key finding was that there was a probability of $\sim 80\%$ for an ME to be measured by two spacecraft separated by less than 15° but it was $\sim 14\%$ for separations of 30° – 60° , and no MEs were clearly observed by two spacecraft separated by more than 60° . This study has the significant advantage of relying on a very large number of CMEs, but has two main limitations: (i) the spacecraft in the innermost heliosphere (MESSENGER and Venus Express) only had magnetic field measurements, and no plasma measurements, which means that the identification relied fully on the magnetic field properties, (ii) the spacecraft used in the study were at different heliocentric distances, ranging from 0.31 au (MESSENGER when Mercury is at perihelion) to 1.09 au (STEREO-B at a maximum heliocentric distance). This last point means that the CMEs and MEs do evolve for up to a few days between the measurements at the two spacecraft, which could include deflection and nonradial expansion. In addition, the combination of these two limitations means that ensuring that the same ME is measured by two spacecraft separated by more than 15° – 20° is especially challenging. As such, the number from this study should be taken, in our opinion, as an upper bound.

Since late 2020, STEREO-A has been within 60° from the Sun–Earth line at a heliocentric distance of 0.95–0.98 au and during the ascending phase of solar cycle 25. This provides the opportunity to perform similar studies as those described above but with more CMEs than during 2007–2009 and for two spacecraft at the same heliocentric distance. Lugaz et al. (2022) described an ME measured by STEREO-A and Wind in 2021 February while the two spacecraft were separated by 55° in longitude. The CME was propagating almost exactly in between the Sun–STEREO-A and Sun–Earth lines, giving an angular extent of this event of $\sim 60^\circ$. While the authors

conclusively showed that the same ME was measured at the two spacecraft, there were significant global differences in the measurements, such as (i) the ME was driving a shock at STEREO-A but not at Wind, (ii) the duration of the ME was about twice shorter at Wind than at STEREO-A. Many of these differences were due to the interaction of the ME with a high-speed solar wind stream that impacted STEREO-A before the ME but into which the part of the ME that impacted Wind was engulfed.

In the present work, we take advantage of this return of STEREO-A toward the Sun–Earth line during a time of increased solar activity. Our key goal is to put constraints on the longitudinal extent of MEs through a multistep approach. We first identify all CMEs that erupted and propagated $\sim 30^\circ$ from the Sun–Earth or Sun–STEREO-A line. We do so from a combination of analyzing remote coronagraphic observations and performing GCS fittings, as described in Section 2.1. We then use the existing database and analyze in situ measurements of CMEs and MEs to determine how many of these CMEs are measured in situ by none, one, or two spacecraft (see Section 2.3). We also analyze in this section specific examples of CMEs measured in situ by one or two spacecraft. In Section 3, we look at all in situ measurements of CMEs at STEREO-A and Wind and determine how many of them are measured by the other spacecraft. We discuss our results in Section 4, in particular the consequences for CME coherence and design of future multi-spacecraft missions. We conclude in Section 5.

2. CMEs Propagating toward L1 or STEREO-A from 2020 November to 2022 August

2.1. Remote Observations

We focus on the time period from 2020 November to 2022 August as this corresponds to STEREO-A being 60° – 20° from the Sun–Earth line (always on the east side). During this time period (2020 November–2022 August), the monthly sunspot number was 46.5 ± 27.5 , corresponding to the first half of the rising phase of solar cycle 25. We start our analysis by identifying all partial and full halo CMEs observed by LASCO (Brueckner et al. 1995) and STEREO/SECCHI/COR2 (Howard et al. 2008). To do so, we use the SOHO/LASCO catalog (Yashiro et al. 2004) and the COR2 CACTUS automatic catalog (Robbrecht & Berghmans 2004). For the LASCO catalog, we start from all CMEs with an angular width larger than 120° . For the CACTUS catalog, we start from all events marked as partial or full halo (I, II, III, or IV). This results in 179 CMEs.

We visually inspected LASCO and COR2 movies for each of these events as well as on-disk signatures from Solar Dynamics Observatory/Atmospheric Imaging Assembly (Lemen et al. 2012) to eliminate back-sided events and partial halos from LASCO perspective which are clearly western limb events (i.e., which are propagating far from the Sun–STEREO-A line). This resulted in a list of 51 events that could potentially impact both spacecraft from visual inspection. For these 51 events, two members of the investigation performed independent GCS (using a Python version; von Forstner 2021) visual fits to the COR2 and C2 images. We primarily focus on the longitudinal direction of the CMEs with respect to the Sun–Earth and Sun–STEREO-A lines and we use the average and minimum (away from the Sun–Earth line) direction from the

two visual fits to categorize the CMEs into four groups. Cat I (21 events) are those CMEs that could impact both spacecraft as they are found to propagate within 30° of both Sun–spacecraft lines for at least one reconstruction. Cat II (14 events) are those CMEs propagating within 30° from the Sun–Earth line but more than 30° from the Sun–STEREO-A line. Cat III (four events) are those CMEs propagating within 30° from the Sun–STEREO-A line but more than 30° from the Sun–Earth line. Cat IV (12 events) are those events propagating more than 30° from both Sun–spacecraft lines. The number of CMEs in each category is not representative of their true occurrence rate, as there is a bias toward Earth-directed CMEs by using the more extensive LASCO CME catalog as compared to the automatic CACTUS COR2 catalog. Tables 1 and 2 list the information about each Cat I CME and all other CMEs, respectively. Some examples of the stereoscopic remote observations for the CMEs discussed below are shown in Figure 6 in the Appendix. We also estimate the CME speed around $10 R_\odot$ based on two consecutive GCS reconstructions, where we simply do a first-order differentiation of the height versus time). This is given in the text in the following section, while we use the LASCO/CDAW linear speed in Table 1.

Figure 1 shows a visual sketch of the projected direction of propagation of the 21 Cat I CMEs onto the ecliptic. One of the two GCS fittings for CME 16 was 7° east of the Sun–Earth line, corresponding to an angular separation of 26° with the Sun–STEREO-A line at that time, even though the average longitudinal direction was 4° west of the Sun–Earth line. The 2021 February CME studied in Lugaz et al. (2022) corresponds to CME 2, whereas the 2021 November 2 CME studied in Regnault et al. (2024) is CME 9.

2.2. In Situ Data and Data Analysis

We use in situ measurements from Wind Magnetic Field Investigation (MFI; Lepping et al. 1995) and 3DP (Lin et al. 1995) as well as measurements from IMPACT (Luhmann et al. 2008) and PLASTIC (Galvin et al. 2008) on board STEREO-A to investigate the impact of the MEs at the two spacecraft and their overall properties. In the following sections, we discuss the CME arrival times and properties of the two spacecraft. Determining the fast magnetosonic shock arrival, when one is present, is relatively straightforward, with the rapid and simultaneous increase in the magnetic field, density, velocity, and temperature. Determining the start and end times of the ME is not as straightforward. The start of the ME is often marked by a decrease in the proton β , a decrease in temperature below the expected temperature of Lopez (1987), an increase in magnetic field strength, a decrease in the magnetic field variability, and often a discontinuity in the magnetic field vector. As discussed in Zurbuchen & Richardson (2006), there is no single property that is always measured within MEs. Here, we focus in general on the low β region, since it indicates a magnetically dominated period. Choosing a start time that corresponds to a discontinuity in the magnetic field vector often allows to match the start time at the two spacecraft (see, for example, CME 5 in the top panels of Figure 2). When there is a clear change in the velocity, β , or magnetic field, we use this as the end time (see, for example, CME 12 in the bottom panels of Figure 2). Other times, the end time of the ME does not correspond to a clear discontinuity and is harder to determine.

Table 1
Overview of the 21 Cat I CMEs Studied Here

CME No.	$\Delta\phi$ (deg)	Date	Time (UT)	V (km s $^{-1}$)	ϕ_{L1} (deg)	ϕ_{STA} (deg)	λ_{L1} (deg)	λ_{STA} (deg)	Impact
1	57	2020 Dec 1	07:12	735	33	24	17	12	None
2	55.6	2021 Feb 20	11:25	555	31	25	-10	-15	Both
3	54.7	2020 Mar 20	01:25	337	29	26	3	1	STA
4	53	2020 Apr 17	17:48	353	17	36	10	11	None
5	51	2021 May 22	08:48	381	33	18	-6	-1	Both
6	42.4	2021 Aug 23	06:48	440	21	21	4	7	Wind
7	42.1	2021 Aug 26	18:48	644	20	23	4	7	STA
8	37.4	2021 Oct 28	08:48	245	2	36	-34	-36	STA
9	37.2	2021 Nov 2	02:48	1473	6	31	14	11	Wind
10	36.2	2021 Nov 24	14:26	390	4	32	-17	-21	Wind
12	34.7	2022 Jan 29	23:36	530	12	23	12	16	Both
13	34	2022 Feb 19	17:48	274	28	7	-17	-19	None
14	33.4	2022 Mar 16	13:25	491	19	15	13	13	STA
15	33.1	2022 Mar 25	05:48	495	27	6	-4	-3	None
16	32.9	2022 Mar 28	12:00	702	-4	37	16	16	Wind
17	30	2022 May 10	14:36	415	21	9	-22	-20	None
18	26.2	2022 Jun 26	19:48	200	3	23	5	8	STA
19	21.5	2022 Aug 13	18:07	606	4	18	-27	-26	None
20	21.1	2022 Aug 15	15:24	450	8	13	-32	-30	None
21	21	2022 Aug 16	02:24	495	10	11	-37	-35	Both

Notes. The columns show, from left to right, the CME ID number, the longitudinal separation between the Sun–Earth and Sun–STEREO-A lines, the date and time of the CME first appearance in LASCO images, the LASCO first-order speed, the GCS longitude separation with L1 (ϕ_{L1}) and with STEREO-A (ϕ_{STA}), the GCS latitude separation with L1 (λ_{L1}) and with STEREO-A (λ_{STA}) (a negative sign indicates that the CME propagates south of the Sun–spacecraft line) and whether there was an impact in situ.

2.3. Statistics

To identify the in situ counterparts to these CMEs, we use the HELIO4CAST catalog (Möstl et al. 2020) for CMEs at STEREO-A and Wind as well as the catalog of CMEs at ACE by Richardson & Cane (2010). We visually inspected all in situ time periods at Wind and STEREO-A around the expected time of the arrival of the CMEs (typically \sim 2–5 days after the initial appearance in coronagraph images to ensure that we do not miss any potential events).

Overall, from the 21 Cat I CMEs, propagating within 30° of both spacecraft, only four CMEs are measured at both spacecraft (three events discussed below plus the one from Lugaz et al. (2022), namely, CMEs 2, 5, 12, and 21), seven are not measured in situ at all (CMEs 1, 4, 13, 15, 17, 19, and 20), six at STEREO-A only (CMEs 3, 7, 8, 11, 14, and 18), and four at Wind only (CMEs 6, 9, 10, and 16). As such 19% (four out of 21) of CMEs propagating in between both spacecraft impacted both, a number that we consider smaller than expected. There are other ways to look at these data: (i) 14 of these 21 CMEs (67%) impacted at least one spacecraft, and (ii) there were 18 in situ measurements of MEs from 42 time periods (21 at STEREO-A and 21 at Wind) corresponding to potential CME impacts (43%).

We consider whether or not CMEs that impact a spacecraft have a smaller angular separation between their direction of propagation (from the GCS model) and the Sun–spacecraft line than those that do not impact a particular spacecraft. We find that there is statistically no difference in the angular separation between the CME direction from GCS and the Sun–spacecraft line for CMEs that impact one spacecraft (18 CME–spacecraft pairs with an average separation of $18.2^\circ \pm 10.6^\circ$) versus CMEs that do not impact a spacecraft (24 CME–spacecraft pairs with an average separation $19.4^\circ \pm 10.9^\circ$). This shows that the study

is not just looking at CMEs that are propagating closer to one spacecraft versus the other.

Hereafter, we discuss specific examples corresponding to different scenarios. CME 2 is described in detail in Lugaz et al. (2022), where the authors concluded that the ME impacted both STEREO-A and Wind but with very different in situ signatures (absence of shock at Wind, ME duration differs by a factor of 2 between STEREO-A and Wind, etc.). Below, we discuss in detail all CME events associated with multi-spacecraft in situ ME measurements to showcase how different the in situ measurements look at the two spacecraft. We also discuss one case without any impact, two cases with impact at only one spacecraft, and one complex case.

2.4. Multi-spacecraft ME Measurements

2.4.1. CME 5: 2021 May 22 Eruption

This CME erupted on 2021 May 22 with a first LASCO image at 08:48 UT, at a time when the L1–Sun–STEREO-A angle was 51° . It is associated with an active region that was located on N17E37 at the time of the eruption. The CME is listed as a full halo in the LASCO catalog, but it looks more like a partial eastern halo, which is how it is listed in the CACTUS C2 catalog (with an angular width of 112°). GCS reconstruction yields a speed of \sim 450 km s $^{-1}$. In the STEREO-A/COR2 field of view, it appears as a full halo with a dominant western direction. The two separate GCS reconstructions give directions of N8E41 and N8E25. We take an average longitudinal direction of $33^\circ \pm 8^\circ$, and we consider that a direction slightly east of the bisector of the L1–Sun–STEREO-A angle is most likely in light of the relative appearance in LASCO/C2 and SECCHI/COR2 images. (For simplicity, in this manuscript, we list the direction east of the Sun–Earth line as positive and west of the Sun–Earth line as negative.) This

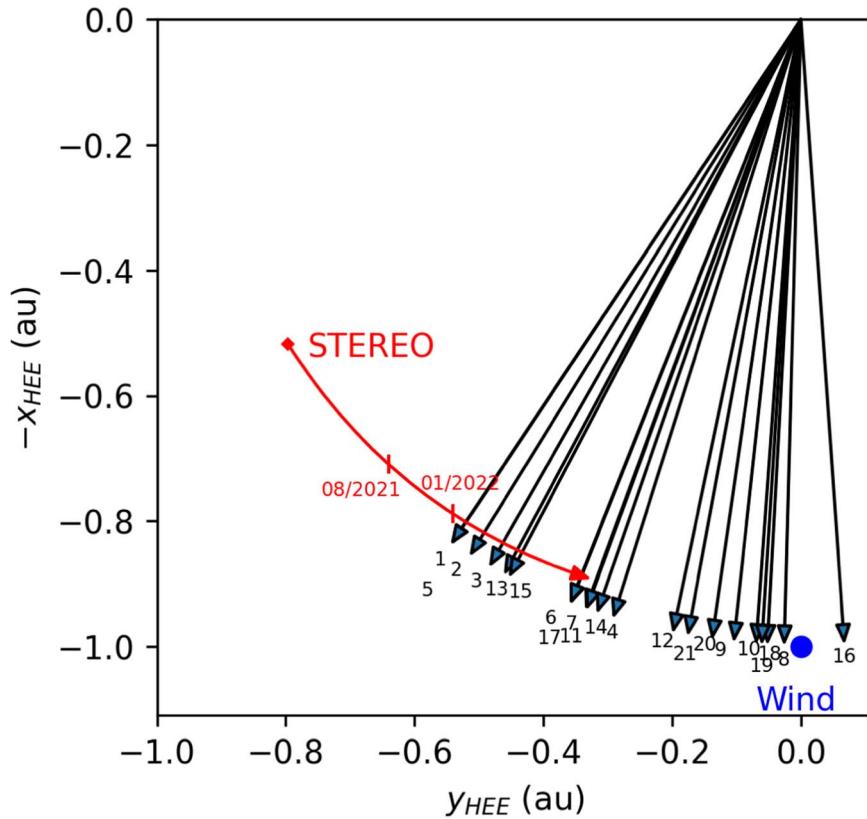


Figure 1. Longitudinal direction from the GCS of the 21 CMEs from 2020 November to August that could impact both STEREO-A and Wind. The Sun is at the origin (0,0) and this is a top-down view on the ecliptic with each CME average direction shown with an arrow. The plot is in the Heliocentric Earth Ecliptic coordinate system. The CMEs are numbered chronologically from 2020 November to August. The position of STEREO-A at the time of CME 1 is shown with the diamond and the end of the arrow corresponds to CME 21. In 2021 August, STEREO-A was 42° from the Sun–Earth line (CME 6) and it was (34.5°) in 2022 January (CME 11). Further information about these CMEs can be found in Table 1.

CME can be expected to impact 1 au at the end of May 25 or early on May 26. Note that STEREO-A heliocentric distance was 0.96 au at that time, which would result in a transit time of $\sim 3-4$ hr less for a CME speed of $400-450 \text{ km s}^{-1}$ near 1 au.

An ME can be identified in both STEREO-A and Wind in situ measurements starting on May 26 as shown in the top row of Figure 2. At Wind, there is a fast-forward (FF) shock at 11:44 UT followed by a relatively complex ME whose starting time can be chosen as 21:30 UT when the proton temperature goes below the expected proton temperature of Lopez (1987) at the same time as there is a clear discontinuity in the R and N components of the magnetic field vector. The ME is characterized by enhanced magnetic field strength, which peaks in its front half, multiple magnetic field vector rotations, and a slowly increasing speed profile. In the database of Richardson & Cane (2010), it is reported as an MC-0 event, i.e., an event that “lacks most of the typical features of a magnetic cloud.” There is no specific feature that can define the end boundary, but the magnetic field strength goes back to typical about 30 hr after the beginning of the ME. The end boundary of this CME is not shown in Figure 2, as it is unclear.

At STEREO-A, there is a clearer ME preceded by a wave (not a shock) at 19:34 UT on May 25 (shown at time -4.43 hr). There is a similar discontinuity in the R and N components of the magnetic field at 9:43 UT on May 26 as the one at Wind. The ME is characterized by an enhanced magnetic field strength with a *classical* peak in the center, a relatively smooth rotation of the magnetic field vector, and a decreasing speed

profile but a temperature that closely follows the expected temperature. This would probably be categorized as an MC-1 as it lacks the clear low-temperature signatures that are most common for true MCs.

The similarities between the ME measurements at Wind and STEREO-A are limited to the front half of the ME: the ME start boundary as described above and a period of ~ 8 hr at STEREO-A with $B_N < 0$ and ~ 6 hr at Wind, followed by a sharp (at STEREO-A) rotation of B_N and B_T , which corresponds to the peak of the magnetic field strength at STEREO-A. This time period appears similar at both spacecraft. The period after the sharp rotation in the magnetic field component lasts ~ 20 hr at STEREO-A but only 2 hr at Wind. Instead of a slowly changing magnetic field vector at Wind, there is a more complex behavior of the magnetic field components. Overall, we conclude from the combination of remote and in situ measurements that the CME on 2021 May 22 impacted both STEREO-A and Wind, with a clearer impact at STEREO-A, which is consistent with the visual inspection of the images. The comparison of the magnetic field and plasma measurements at the two spacecraft indicates that, while it is the same ME at both spacecraft, the ME is barely, if at all, coherent at this angular separation. This indicates that the ME angular size is $>51^\circ$ (the Wind-STEREO-A separation) but that the coherence length is smaller than this. We also note that, contrary to the event investigated in Lugaz et al. (2022), the B_R and B_T components of the magnetic field do not seem to indicate a crossing through a different leg of a flux rope.

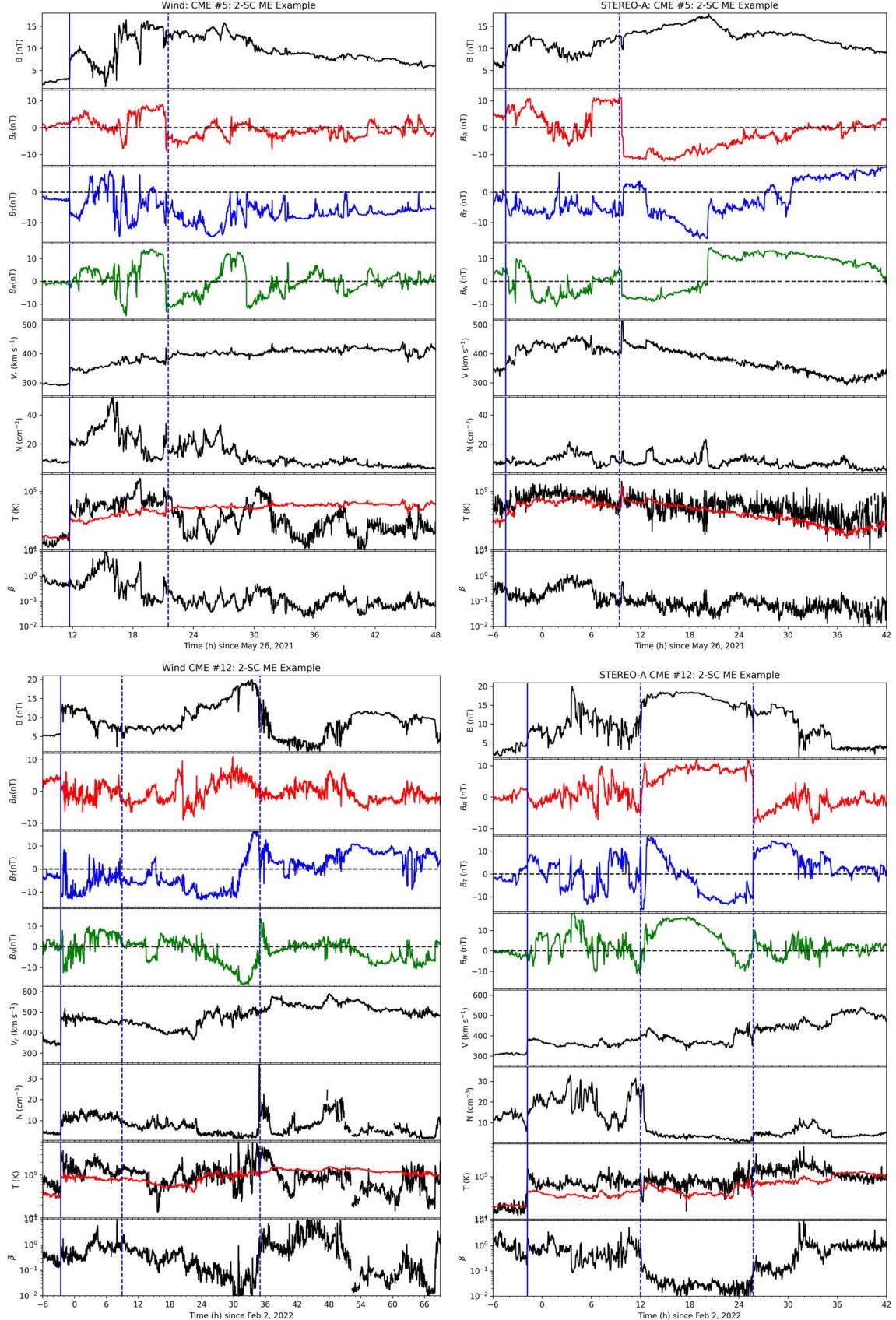


Figure 2. Two probable multi-spacecraft measurements of CMEs by Wind (left) and STEREO-A (right). The top row shows CME 5 and the bottom shows CME 12. Each plot shows from top to bottom, the magnetic field strength, the magnetic field components in the RTN coordinate system, the proton velocity, density, and temperature (with the expected temperature in red) and proton β . The plots of Wind and STEREO-A have the same y-scale for each quantity (but different x-scale). FF shocks are marked with a solid blue vertical line and the ME boundaries with blue dashed vertical lines (the back boundary for CME 5 is not clear and is not shown).

2.4.2. CME 12: 2022 January 29 Eruption

This CME erupted on 2022 January 29 with a first LASCO image at 23:36 UT, at a time when the L1–Sun–STEREO-A angle was $(34^\circ.7)$. The CME is listed as a full halo in the LASCO catalog with a speed of $\sim 500 \text{ km s}^{-1}$ and it has a clear eastern asymmetry indicating a propagation east of the Sun–Earth line. In STEREO-A/COR2 field of view, it appears as a full halo with a dominant western direction. The two separate GCS reconstructions give directions of N8E10 and N12E13. We take an average longitudinal direction of $11^\circ \pm 2^\circ$, which is consistent with the slightly more symmetric halo view from L1 than from STEREO-A. This CME can be expected to impact 1 au on or around February 2.

An FF shock occurs at Wind at 21:26 UT (time -2.57 hr) on February 1 as shown in the bottom left row of Figure 2. The next 36 hr are clearly associated with a CME, although the boundaries of the ME are not straightforward to determine. The database of Richardson & Cane (2010) lists a start time of 16:00 UT on February 2, but the duration of 17.7 hr between the shock and the beginning of the ME is about twice longer than typical for CMEs at 1 au. A start time of the ME around 09:00 UT is also possible based on a discontinuity in B_T . While Richardson & Cane (2010) list this event as an MC-2, we note that the magnetic field strength has a relatively complex profile, the temperature inside the ME is not particularly low, but there is a clear rotation of the magnetic field vector starting at 16:00 UT (or 09:00 UT if considering a slightly more complex period) and the speed shows the typical monotonically decreasing profile until 23:00 UT. The ME speed is about 450 km s^{-1} . The HELIO4CAST database uses a start time of 20:46 UT on February 2 corresponding to a clear discontinuity in B_R . Both databases list an end time around 10:40 UT on February 3 (34.67 hr). About 13 hr after the end of the ME, there is a second ME with somewhat more typical characteristics.

At STEREO-A, there is a shock-like feature at 22:11 UT on February 1 (time -1.82 hr). There is a clear ME starting between 12:00 and 12:50 UT on February 2 corresponding to a large drop in the proton density and a clear discontinuity in the magnetic field. We consider 01:45 UT on February 3 (25.75 hr) as the end time of the ME, corresponding to an increase in the proton β and a clear discontinuity in the magnetic field. The ME has a speed of about 360 km s^{-1} .

Based solely on the in situ measurements, there are few, if any, signs that STEREO-A and Wind measured the same event. However, the beginning of the sheath at the two spacecraft is within 1.5 hr from each other and there is just one CME that propagated in between the Sun–STEREO-A and Sun–L1 lines at a time and with a speed that matches the start time of the event. We consider that it is likely the same event, but this is primarily based on the remote observations and the fact that STEREO-A and Wind were separated by less than 35° at that time. This would indicate that the ME angular size is 35° but the coherence length is significantly less than 35° .

2.4.3. CME 21: 2022 August 15 Eruption

There were at least five front-sided full or partial halo CMEs from 2022 August 13–17 with the ones on August 13 and 15 potentially propagating in between the Sun–Earth and Sun–STEREO-A lines (CMEs 20 and 21 in our list). At the time, STEREO-A was separated by 21° – 21.5° from the Sun–Earth

line. We focus here on the August 15 CME, which has a first appearance in LASCO images at 15:24 UT with an apparent angular width of 125° . It originates from an active region, which was around S20W01 at that time. As STEREO-A and SOHO were only separated by 21° , the GCS reconstruction has more uncertainty, but it appears visually that the CME propagates in between the two spacecraft, being slightly eastward from the SOHO/LASCO point-of-view and slightly westward from the STEREO-A/COR point of view. The two GCS reconstructions give the directions of S25E07 and S25E09, so that we take an average longitudinal direction of $8^\circ \pm 1^\circ$. Based on the speed measured in coronagraphs, we can expect an arrival at 1 au on or around August 19.

Wind measures an FF shock at 16:51 UT on August 19 (see the top plot of Figure 3) followed by an ME that starts around 03:00 UT on August 20, although the list of Richardson & Cane (2010) has a later start time. The end time is at around 18:30 UT on August 20. The ME has a relatively weak magnetic field strength ($\sim 6 \text{ nT}$), no clear rotation of the magnetic field vector, a temperature similar to the expected temperature, and not very low proton β . However, the magnetic field variability is small and there is a clear decreasing speed profile. The ME has an average speed of 550 km s^{-1} with an expansion speed of $\sim 95 \text{ km s}^{-1}$. There is no ME listed in the HELIO4CAST catalog but an MC-1 is listed in the list of Richardson & Cane (2010).

STEREO-A measures an FF shock at 20:57 UT on August 19 followed by an ME, which starts around 01:00 UT on August 20 (HELIO4CAST lists a 05:16 UT start time) and lasts until 19:00 UT on August 20. The magnetic field strengths in the sheath and ME are clearly stronger at STEREO-A than at Wind (ME field strength of $\sim 12 \text{ nT}$). There is a clearer portion of the ME with low proton β and a lower-than-expected proton temperature. The ME has a speed of $\sim 550 \text{ km s}^{-1}$ with almost no expansion. Similar to the previous event, if it were not for the near-simultaneous measurements, the close proximity of L1 and STEREO-A ($\sim 21^\circ$), and the fact that there is a CME propagating in between the two spacecraft, it is possible that these two events would not be linked. While this might be related to the origin of the CME south of the ecliptic, we note that the STEREO-A measurements differ significantly from those from Wind.

2.5. Single-spacecraft In Situ ME Measurements and No Impact

Every other 17 CME that is found remotely to propagate in between the Sun–Earth and Sun–STEREO-A lines impacts at most one spacecraft, either Wind or STEREO-A. We give four examples, one for a detection at STEREO-A only, and one without impact at any spacecraft near 1 au, as well as two events that are harder to interpret.

2.5.1. CME 8 2021 October 28 Eruption: A Complex Event

This CME erupted on 2021 October 28 with a first LASCO image at 8:48 UT, at a time when the L1–Sun–STEREO-A angle was $(37^\circ.4)$. GCS reconstructions give a direction of S30E2 for the CME in the corona, so it is directed almost exactly along the Sun–Earth line (but south of the ecliptic) and $\sim 35^\circ$ from the Sun–STEREO-A line.

There is a potential ME event measured at STEREO-A lasting 14 hr and starting around 12:20 UT on October 31 (see

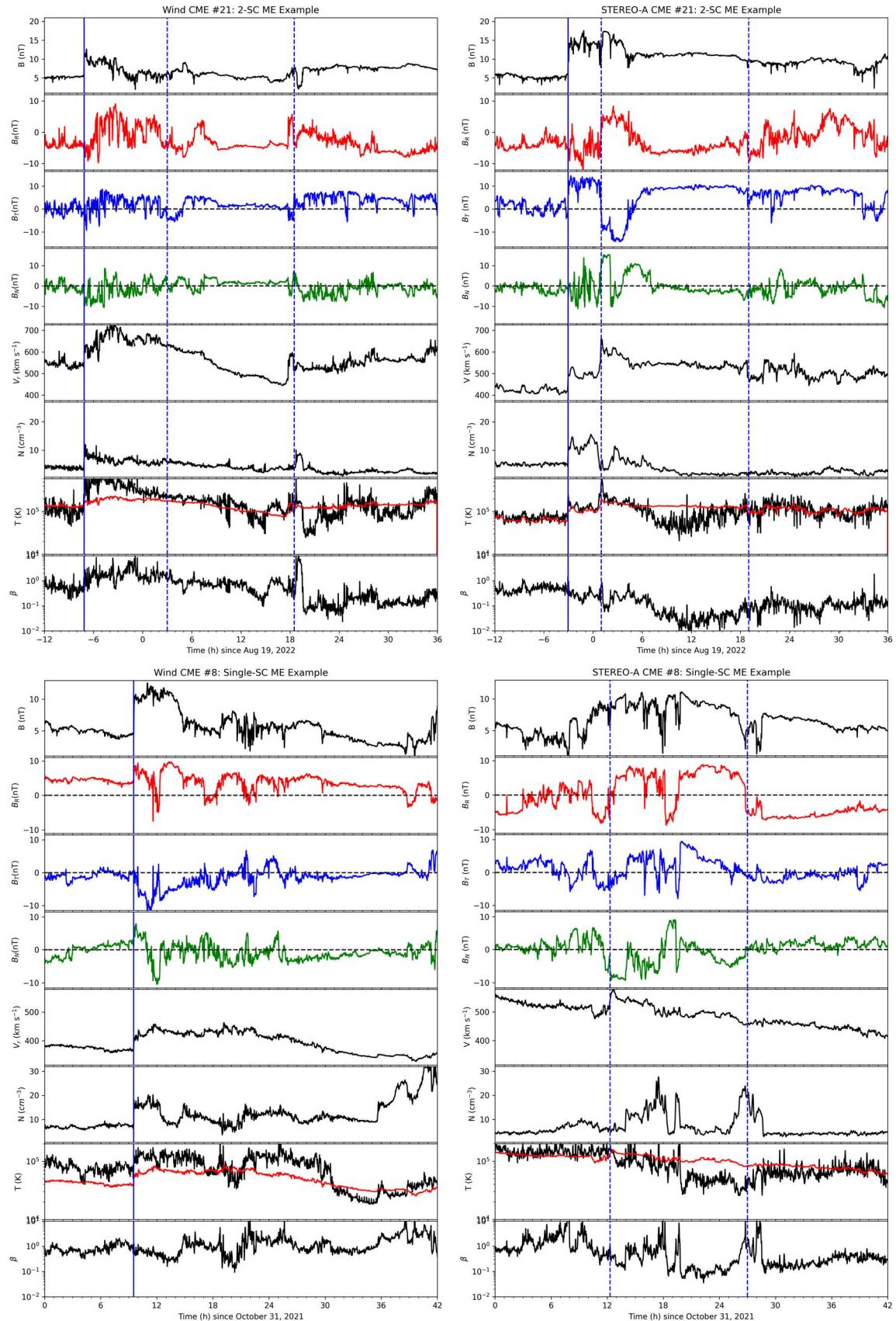


Figure 3. Wind (left) and STEREO-A (right) measurements of a two-spacecraft ME measurement (CME 21, top) and a complex single-spacecraft ME measurement (CME 8, bottom, ME at STEREO-A, shock at Wind). Same format as Figure 2 but both x -scale and y -scale are the same between STEREO-A and Wind.

the bottom right plot of Figure 3). We consider this as a sheath-less ME with complex magnetic field vector rotations with B_N mostly negative and a magnetic field magnitude around 10 nT. There is a period of colder-than-expected temperatures and low proton β for ~ 8 hr in the second half of the event. At a similar time at 9:32 UT on October 31, Wind measures a clear FF shock followed by a weak magnetic field and period of typical proton β . About 20 hr after the shock, there is a short period (~ 4 hr) of lower-than-expected temperature. We consider that this event is a shock-less and sheath-less ME measured at STEREO-A and a *driverless* shock measured at Wind. Whether or not it is associated with the October 28 CME is beyond the scope of this work as well as confirmation that the shock measured at STEREO-A is the same shock as measured at Wind.

2.5.2. CME 7: 2021 August 26 Eruption: Single-spacecraft ME at STEREO-A

This CME erupted on 2021 August 26 with a first LASCO image at 18:48 UT, at a time when the L1–Sun–STEREO-A angle was 42° and associated with an active region around N20E14. The CME is listed as a full halo in the LASCO catalog with a speed of ~ 750 km s $^{-1}$. In STEREO-A/COR2 field of view, it appears as a partial halo with a dominant western direction. GCS reconstructions gave a direction of N11E20.

An FF shock occurs at STEREO-A at 10:08 UT on August 30 (see the top right plot in Figure 4). An ME starts around 16:00 UT on August 30 with a rotation of the magnetic field vector and a decrease in proton density. The magnetic field and density profiles are complex but the speed is monotonically decreasing for more than 36 hr with an average speed of ~ 350 km s $^{-1}$ and an expansion speed ~ 50 km s $^{-1}$. The magnetic field strength has an average ~ 15 nT with multiple complete rotations of the magnetic field vector. HELIO4CAST lists two MEs separated by about 10 hr but the period in between has a similar speed profile and we consider that the in situ measurements are consistent with a multi-MC event as described in Lugaz et al. (2017). At the same time at Wind, there are no clear signatures of any ME. There is a period of the proton temperature being slightly lower than the expected temperature, which corresponds to a period with slightly less variability in the magnetic field vector. However, the magnitude of the magnetic field is ~ 5 nT and the B_N component which has three clear periods of strong negative B_N at STEREO-A is close to 0 throughout this time period at Wind.

The direction of the CME is clearly in between STEREO-A and Wind, the clear signature of a complex event at STEREO-A and the lack of transient signatures at Wind is confusing for a spacecraft separation of 42° and a propagation direction close to the angle bisector. It is possible that the interaction of multiple CMEs may have played a role in the deflection of the CMEs toward the east (STEREO-A) although confirming this is beyond the scope of this study. We note (see Appendix Figure 7) that the CME has a clear component in the ecliptic and a propagation clearly out of the ecliptic cannot be invoked to explain the in situ measurements. The measurements would indicate that the ME angular size is less than 42° but, if no deflection is invoked, the ME angular size would be less than 25° .

2.5.3. CME 9 2021 November 2 Eruption: Single-spacecraft ME at Wind

This CME erupted on 2021 November 2 with a first LASCO image at 02:48 UT, at a time when the L1–Sun–STEREO-A angle was $(37^\circ, 2)$. GCS reconstructions give a direction of S19E6, so directed very close to the Sun–Earth line (but south of the ecliptic) and $\sim 30^\circ$ from the Sun–STEREO-A line. This is a relatively fast CME with a speed of more than 1400 km s $^{-1}$.

There is a clear ME at Wind with an FF shock at 19:40 UT on November 2 (see the bottom left plot in Figure 4) about 41 hr after the first image of the CME on LASCO. In situ speeds above 750 km s $^{-1}$ are measured, both in the sheath and in the front half of the ME, which starts around 6:30 UT on November 3. There is a clear rotation of B_T from negative to positive with a positive B_N . The ME is characterized by a low temperature, low proton β , and a relatively clear expansion. STEREO-A measures near the same time a complex event that starts with a discontinuity at 4:15 UT on November 3, about 8.5 hr after the shock at Wind. There are more complex magnetic field signatures as well as high density and a speed of 450–500 km s $^{-1}$. B_T is primarily positive. Confirming whether or not the shock is driven by the same CME at Wind and STEREO-A is beyond the scope of this study. In any case, there is only an ME at Wind.

2.5.4. CME 15: 2022 March 25 Eruption: No Impact

This CME erupted on 2022 March 25 with a first LASCO image at 05:48 UT, at a time when the L1–Sun–STEREO-A angle was 33° and associated with an active region around S20E36. The CME is listed as a full halo in the LASCO catalog and it has a clear eastern asymmetry indicating a propagation east of the Sun–Earth line. It has a speed of ~ 600 km s $^{-1}$. In the STEREO-A/COR2 field of view, it appears as a full halo with a very slight asymmetry to the west. GCS reconstructions gave a direction of S10E27 with uncertainties of $\pm 4^\circ$. The CME is directed almost exactly toward STEREO-A but because the two spacecraft separation was small, an impact at both spacecraft can be expected.

There is no CME listed either in HELIO4CAST or the database of Richardson & Cane (2010). At Wind, there are plasma and magnetic field measurements that are consistent with very weak ejecta, starting around 6:00 UT on March 28 and lasting about 1 day (see the top left plot of Figure 5). There is a low-density, lower-than-expected temperature, and decreasing speed profile. However, the magnetic field is very weak, below 5 nT and the proton β is not low indicating that this is not a magnetically dominated time period. At STEREO-A, there are clear indications of the passage of a stream interaction region (SIR), followed by a high-speed solar wind stream. There is no period with a lower-than-expected temperature or low proton β . It is not possible to prove that the signatures late on March 27 and early on March 28 are not associated with an ME, but, if they do, they are very unusual.

Interaction with the high-speed solar wind stream may have caused a deflection of the CME, which resulted in the lack of any detection at Wind as well as a change in the morphology of the CME, which makes the detection at STEREO-A more complex. Since the CME was initially propagating within about 5° of the Sun–STEREO-A line and 25° from the Sun–Earth line, this nondetection can only be understood by invoking the

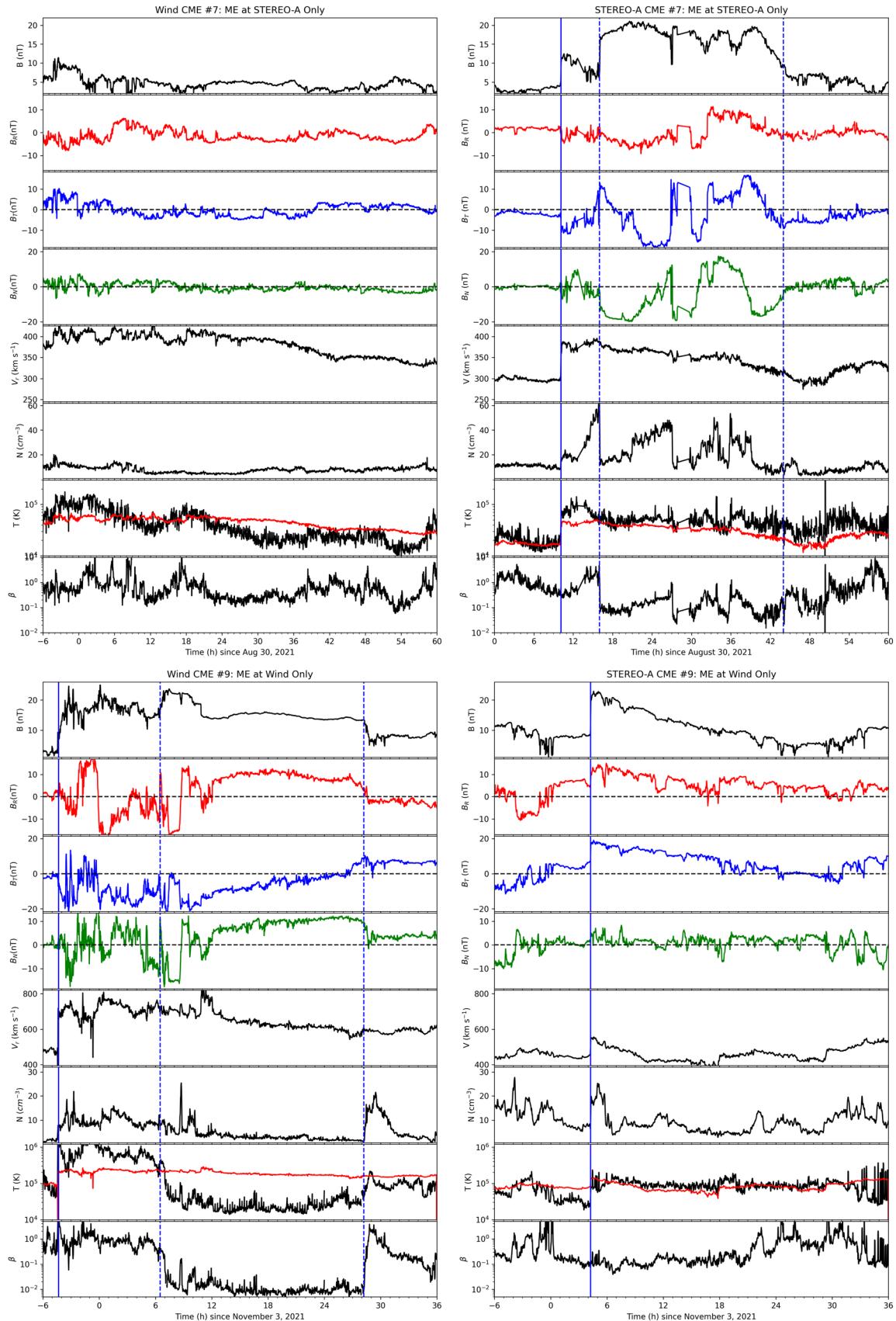


Figure 4. Wind (left) and STEREO-A (right) measurements of a clear single-spacecraft ME measurement: CME 7 at STEREO-A (top) and CME 9 at Wind Only. Same format as Figure 2.

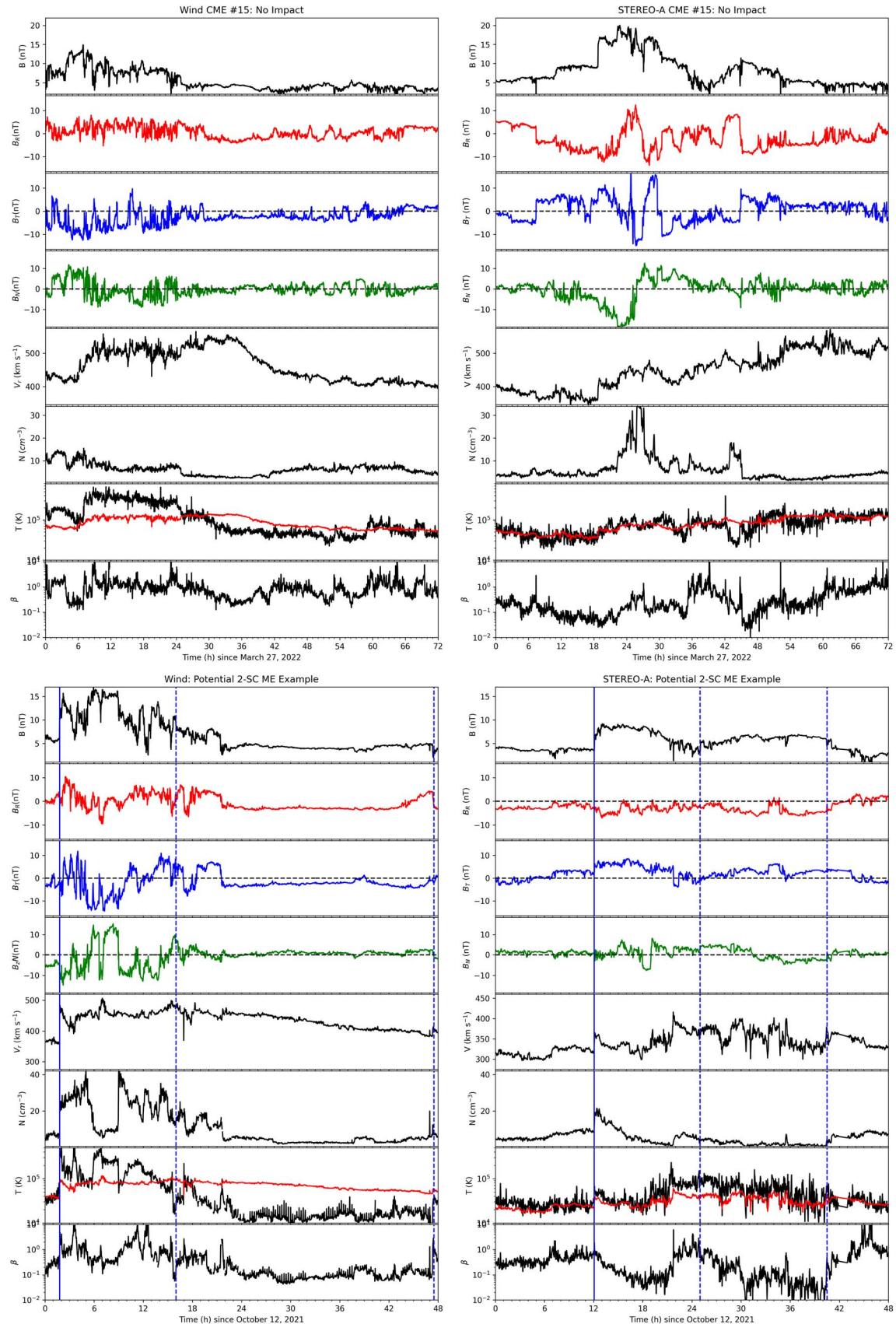


Figure 5. Wind (left) and STEREO-A (right) measurements of a nondetection of an ME for CME 15 and a potential additional two-spacecraft measurement for a CME on 2021 October 12. Same format as Figure 2.

interaction with the SIR and high-speed stream. However, in light of the relatively low separations between STEREO-A and L1, this is still a surprising finding.

2.6. Effect of the CME Propagation Away from the Ecliptic

Table 1 shows in the eighth and ninth columns the latitudinal angles between the Sun–spacecraft line and the CME direction of propagation from the GCS reconstruction for Wind and STEREO-A. While both Wind and STEREO-A are in the ecliptic, they have different latitudes with respect to the solar equator, varying between -6° and 6° . The average latitudinal separation between the two spacecraft for these 21 CMEs is 2.5° .

We first note that there is no clear relationship between the lack or not of impact in situ and the latitudinal direction of the CME. For example, CME 15 did not impact any spacecraft even though it was propagating within 5° of the ecliptic. On the other hand, CME 21 impacted both STEREO-A and Wind even though it was -35° from the ecliptic. While strong latitudinal deflections have been reported in the low-to-middle corona (Kilpua et al. 2009), further latitudinal deflection in the upper corona has also been reported (Byrne et al. 2010). Such deflections typically occur toward the ecliptic, so the number in Table 1, obtained from the GCS reconstruction using COR2 and C3 observations around $10 R_\odot$ should be taken as an upper bound.

Second, for sixteen of the 21 events, the difference in latitude between the Sun–spacecraft line and the CME direction of propagation is smaller than the difference in longitude. Some of the exceptions (CMEs 8, 11, and 21) impacted at least one spacecraft in situ, so the impact of the latitudinal difference cannot be clearly determined. In fact, the current study focuses on the longitudinal size of the ME because STEREO-A and Wind are separated in longitude. An investigation of the ME latitudinal extent would be best performed using a much larger number of events, since it can only rely on single-spacecraft crossings. We discuss the potential effect of ME inclination in the discussion Section (4.1).

Third, it is in theory possible that the tilt between the ME and the ecliptic (where both STEREO-A and Wind are) results in additional misses. As shown in Table 1, there are not any events where the ME propagates latitudinal in between STEREO-A and Wind. However, for MEs with a significant inclination, such effects could still occur, i.e., the CME passes locally above one spacecraft and below the other. The current study is only able to determine the extent of the ME in the ecliptic, which is a combination of the cross-section shape, the ME leg-to-leg extent, and the tilt of the ME with respect to the ecliptic. This is further discussed in the discussion Section (4.2).

3. In Situ Measurements of CMEs from 2020 November to August

3.1. Methodology

We perform a complementary study, where we start from in situ measurements at Wind and STEREO-A to ensure that there were not an additional two spacecraft measurements of MEs during this time period which may not have been clearly associated with coronagraphic sources, either because they are associated with stealth CMEs (from both STEREO-A and LASCO points of view) or because interplanetary deflection,

even if it is limited to a few degrees, resulted in a CME from beyond 30° from the Sun–spacecraft line to impact a given spacecraft.

We start from the list of CMEs at L1 from Richardson & Cane (2010) and the HELIO4CAST database for both CMEs at L1 and STEREO-A. We remove all events with ME duration shorter than 12 hr, as such events have been in the past described as small transients (Cartwright & Moldwin 2010; Yu et al. 2016) and may not always be of coronal origin. We also remove events with weak magnetic field, on average below ~ 6 nT as these are significantly weaker than a typical ME near 1 au. We visually inspected all in situ measurements and ended up with 23 CMEs measured in situ by Wind and 17 by STEREO-A. We note that here MEs measured simultaneously at two spacecraft are double-counted.

We found no additional clear two spacecraft measurements outside of the four highlighted earlier. There is, however, one event, the 2021 October 12 ME, for which there might be multi-spacecraft measurements, which we briefly describe below.

3.2. 2021 October 12 ME

There is an MC-1 in the database of Richardson & Cane (2010) at Wind, with a shock at 01:48 UT on October 12 and an ME from 16 UT on October 12 to the end of October 13 (see the bottom panel of Figure 5). There is a clear FF shock and a dense, hot, and magnetized sheath region followed by a colder-than-expected and relatively low β ME. The magnetic field strength in the ME is on average 6 nT, $B_N \sim 0$ nT and there is only rotation in the first ~ 5 hr followed by steady magnetic field vectors. HELIO4CAST lists a very short ejecta (~ 4 hr) in the sheath region. Due to the steady nature of the magnetic field and the cold temperature, we follow the boundaries of Richardson & Cane (2010). HELIO4CAST lists a CME at STEREO-A starting with an FF shock at 12:04 UT and an ME lasting 15.5 hr starting at 01 UT on October 13. While the magnetic field is also weak, there is a clearer rotation of the magnetic field vector with B_N going from positive to negative. STEREO-A and Wind were separated by 38° in longitude. Based on the in situ measurements, it is impossible, in our opinion, to conclude whether or not this is the same event as there are not many similarities but the two events start within less than 12 hr from each other. The event at Wind is likely to be associated with the halo CME that started on 2021 October 9 around 7:12 UT, for which the GCS reconstruction shows a direction of propagation 3° west of the Sun–Earth line (41° from the Sun–STEREO-A line), a Cat II event.

3.3. Statistics

As such, there are 35 or 36 individual CMEs measured either at Wind, at STEREO-A, or at both spacecraft, four of which have MEs clearly measured by both spacecraft and one is unclear. We do not consider events for which the shock is measured by one spacecraft and the ME by another as it is nearly impossible to confirm that they are associated with the same eruption. Overall, this indicates that 11%–14% of MEs measured in situ during this time period when STEREO-A and L1 were separated by 20° – 60° were measured by both spacecraft. This number is lower than the number of multi-spacecraft measurements based on remote observations ($\sim 20\%$) because it includes (a) impacts from CMEs that were

propagating west of the Sun–Earth line and impacted Wind, and (b) CMEs that were propagating east of the Sun–STEREO-A line and impacted STEREO-A.

4. Discussion

4.1. Additional Effects that May Influence the Results

Under the flux rope paradigm for MEs, the angular extent in the ecliptic is a combination of the ME leg-to-leg width and the ME nonradial cross-section width. First, based on this paradigm, the likelihood of measuring the same ME at two spacecraft is expected to depend on the ME orientation: it should be high for MEs with low inclination and low for MEs with high inclination. Assuming a circular cross section, the ME typical radial size of 0.21 au corresponds to a nonradial cross-section angular width of 12° . If the ME has an elliptical cross section with a 3:1 ratio of the nonradial to radial widths, as found in Démoulin et al. (2013) based on the distribution of impact parameters, the ME would have an angular width of $\sim 37^\circ$. Under the model of Owens et al. (2006), the ME angular width remains constant and should be $\sim 40^\circ$ – 80° . The leg-to-leg width is less well constrained but is typically assumed to be at least 60° .

We first note that most of the MEs plotted in this article are complex, often with multiple magnetic field vector rotations, and therefore determining the ME inclination is difficult. In addition, we do not rely on fitting techniques developed for single-spacecraft measurements, as such techniques can return results with large uncertainties (Al-Haddad et al. 2013) and they rely on the highly twisted axisymmetric flux rope paradigm (Al-Haddad et al. 2019), which we argue may not necessarily hold true in light of the presented measurements. Based on our limited sample, we do not find any indication that the ME magnetic field orientation has an impact on the ME angular width (or the likelihood of measuring the same ME at two different locations). Low-inclined MEs should be easily measured by two spacecraft for separations of 45° – 60° if the expectation about the leg-to-leg width is correct, yet we do not see any indication that more low-inclined MEs were measured at large separations. This puts into question the flux rope paradigm.

CME–CME interaction has been reported to sometimes result in longitudinal deflection (Lugaz et al. 2012) and in compression of the ME in the radial direction (e.g., see Lugaz et al. 2005), and it is likely to also result in changes in the CME angular width. Deflections in particular may result in CMEs having a different direction in the heliosphere than that obtained from GCS. Additionally, CMEs often interact with SIRs as they propagate (Farrugia et al. 2011; Prise et al. 2015; Winslow et al. 2016; Heinemann et al. 2019). This can result in significant differences in the ME properties for events measured near simultaneously by two spacecraft as is presented here (Winslow et al. 2021; Lugaz et al. 2022). Deflection of CMEs in the heliosphere due to their interaction with the structured solar wind has also been hypothesized (Wang et al. 2004) and may influence the presented results.

Overall, the current study focuses on the impact of MEs at the two spacecraft to deduce the angular extent in the ecliptic plane, not the differences in the properties of the MEs at the two spacecraft. For this reason, we did not perform an in-depth study of the interplanetary conditions into which the CMEs

propagate. Therefore, CME–CME and CME–SIR interactions may influence the presented results.

4.2. CME Angular Size in the Ecliptic

11%–14% of MEs measured in situ during the time period from 2020 November to August when STEREO-A and L1 were separated by 20° – 60° in longitude were measured in situ by both spacecraft. This is a slightly lower number than that obtained by Good & Forsyth (2016), who found seven out of 39 events (18%) for two spacecraft separations of 30° – 45° . However, they found a much higher number (19 out of 39, so 48%) for separations of 15° – 30° . Their studies had two main limitations: the two spacecraft were at different distances and one of the two spacecraft typically did not have plasma measurements. We therefore consider that the number obtained by Good & Forsyth (2016) represents an upper bound for the proportion of MEs that may be measured in situ by two spacecraft near 1 au.

One way to analyze the results of the remote-sensing investigation is as follows: for an average separation between the central direction of a CME as obtained from the GCS reconstruction and the Sun–spacecraft line of $(18^\circ.9)$ (for our 21 events), only two-thirds of the CMEs impact a spacecraft. This would indicate that the median angular size in the ecliptic of MEs is close to 40° . However, this does not take into consideration the fact that there were two spacecraft potentially taking in situ measurements. For the 21 Cat I CMEs, the average STEREO-A–Wind separation is only $(37^\circ.6)$; however, there are only four multi-spacecraft ME measurements. This indicates a typical angular size in the ecliptic significantly lower than 30° . Out of these 21 events, there were nine CMEs that are found to propagate almost exactly in between the two spacecraft (within $\pm 5^\circ$ of the bisector) with an average two spacecraft separation of (40.2 ± 13.9) . Two of these CMEs impacted both spacecraft, six impacted one spacecraft and one impacted none (for the largest separation of 57°). This indicates that the ME angular size in the ecliptic might be as small as 20° .

Overall, we consider that this analysis indicates that the ME angular size in the ecliptic is likely to be 20° – 30° , which corresponds to a size of ~ 0.35 – 0.5 au, or 2 ± 0.4 times larger than the radial size of MEs. As we do not see any influence of the ME orientation, we consider that a model where the leg-to-leg and nonradial cross-section width are more or less equal is the most likely interpretation of the data. This would be closer to the geometry of an ice cream cone model or spheromak rather than a flux rope model for the ME. Such models, in recent years, are often considered less realistic than flux rope models of CMEs (see, for example, Luhmann et al. 2020). The ME is however not exactly spherical but corresponds to two scoops of ice cream (aspect ratio of $\sim 2:1$).

The few multi-spacecraft measurements of MEs that we found indicate the magnetic field is not coherent between the two spacecraft, even when two spacecraft are able to measure the ME simultaneously. This is because the magnetic field and plasma signatures differ significantly between the two spacecraft (see Figures 2 and 3). In fact in most of these four cases, without the knowledge that one and only one CME was propagating in between the two spacecraft, it is unlikely that one could be sure that the same event is measured in situ by STEREO-A and Wind. This is particularly true for events 5 and 12. None of these events present MC-like signatures at both

spacecraft (CMEs 5 and 12 are MC-like at STEREO-A only) and performing an analysis of the correlation between the magnetic field measurements as done in Lugaz et al. (2018) would be meaningless since the measurements are too different. Overall, this indicates that past estimates of the ME coherence of $\sim 20^\circ$ (Lugaz et al. 2018; Scolini et al. 2023) may be an upper bound for the true coherence length. We note that, based on the analysis of Owens et al. (2017), an ME with an angular width of 20° should be coherent if Alfvén waves are the carrier of information across the ME.

4.3. Consequences for Future Mission Concepts

We first mention that missions at the Sun–Earth Lagrangian L4 or L5 points may provide critical information from remote observations of CMEs but will not provide in situ multi-spacecraft ME observations, except in extremely rare cases. This is because we did not measure any CMEs for separations beyond (55.6°) as presented in Lugaz et al. (2022). In addition, as shown in Bailey et al. (2020), measurements of the magnetic field (especially B_N) of corotating streams are poorly correlated between L1 and L5.

Numerous multi-spacecraft in situ mission concepts have been recently proposed to investigate CMEs, shocks, and solar energetic particles (SEPs; e.g., see white papers submitted to the 2023 Decadal Survey (Allen et al. 2022; Akhavan-Tafti et al. 2023; Lugaz et al. 2023; Nykyri et al. 2023), some of them based on earlier concepts (St. Cyr et al. 2000a; Szabo 2005). The inter-spacecraft separations have been proposed to vary between 60° and 90° for the HELIX mission concept and less than 1° for the SWIFT concept of Akhavan-Tafti et al. (2023). We first note that the separations needed to investigate ME magnetic field and coherence may differ significantly from the separations needed to investigate SEP spread due to field line meandering or cross-field diffusion, and from the separations needed to investigate the global structure of shocks. Based on the current study, separations of more than 20° are not appropriate to investigate ME magnetic field as multi-spacecraft measurements are rare at these separations. Also, when they occur, multi-spacecraft measurements at separations of more than 20° do not provide useful information to create more complex models of the ME magnetic field. In fact, longitudinal separations well below 20° (for example $\sim 10^\circ$) appear to be more appropriate.

As our study covers 22 months during the ascending phase of the solar cycle, when there were about 20 CMEs measured in situ at L1, it may also provide an estimate for the number of multi-spacecraft measurements that may occur during the cruise phase of a mission to the Sun–Earth L4 or L5 points, which is typically designed to take 2–3 yr. It is likely that no more than 25% of CMEs will be measured in situ at L1 and the cruising spacecraft. Such numbers should be used to further develop mission concepts and quantifying level 1 requirements.

5. Summary and Conclusions

We have investigated the angular width of MEs by studying all CMEs that were observed remotely to propagate between the Sun–Earth and Sun–STEREO-A lines from 2020 November to August, when the STEREO-A–Sun–Earth angle was between 20° and 60° . We identified 21 such events. The key result is that only four of these 21 CMEs were measured by both STEREO-A and Wind, for separations of (55.6°) , 51° ,

(34.7°) , and 21° . 10 others were measured by only one spacecraft and the last seven by none. In addition, for the four CMEs measured simultaneously by Wind and STEREO-A, the plasma and magnetic field measurements at the two spacecraft differ significantly, which highlights that the coherence length is smaller than the two spacecraft separation.

We complemented the study by looking at all MEs measured in situ by Wind or STEREO-A during the same time period and identifying any additional two spacecraft measurements. We found one potential additional two spacecraft measurement possibly associated with a CME propagating 3° of the Sun–Earth line (and 41° of the Sun–STEREO-A line).

We draw the following conclusions from this study: (i) ME angular width in the ecliptic is likely to be only 20° – 30° for most events, extending to $\sim 50^\circ$ – 55° for the largest events; (ii) there is little indication that low-inclined CMEs have a wider width than highly inclined ones; (iii) therefore, CME models need to account for an elliptical cross section with a relatively small eccentricity (2:1 ratio of major to minor axes) and similar angular sizes from leg to leg as compared to the maximum size of the cross section. These results also highlight that the small number of multi-spacecraft measurements of CMEs during the first two years of the mission (2007–2008) may have been a normal occurrence, not only due to the solar minimum conditions. Kilpua et al. (2011) found four clear multi-spacecraft events during this time period, but with events like the 2007 May 23 ME measured in situ by only one of the two STEREO spacecraft while they were separated by 9° . STEREO-A has now passed in front of the Sun–Earth line in 2023 August and the time period from 2020 August to 2024 July shall provide numerous multi-spacecraft measurements of CMEs for two spacecraft separations of $\pm 20^\circ$ during or close to solar maximum conditions. The time period from mid-2024 to May 2026 would provide a complementary study to the one presented here, for the early part of the descending phase of the solar cycle, if STEREO-A continues to take data in the next 3 yr.

Acknowledgments

The data sets used in this study are from STEREO-A/PLASTIC and IMPACT and Wind MFI and SWE as retrieved from NASA CDAWeb servers.

The HELIO4CAST website is <https://helioforecast.space/icmecat> with the associated doi: 10.6084/m9.figshare.6356420.

All UNH authors acknowledge grant 80NSSC20K0431. N.L. acknowledges additional support from 80NSSC20K0700. B.Z. and C.S. acknowledge additional support from 80NSSC23K1057. B.Z. acknowledges additional support from AGS-2301382. N.A. and F.R. acknowledge grants 80NSSC22K0349, 80NSSC21K0463, and AGS1954983. C.J.F. acknowledges support from grants 80NSSC21K0463, and Wind's grant 80NSSC19K1293. E.D. and C.M. acknowledge funding by the European Union (ERC, HELIO4CAST, 101042188). Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Research Council Executive Agency. Neither the European Union nor the granting authority can be held responsible for them.

Appendix

In Table 2, we list the properties of the Cat II, III, and IV CMEs. In Figures 6 and 7, we show remote images of the CMEs analyzed in this study.

Table 2
List of CAT II, Cat III, and Cat IV CMEs

$\Delta\theta$ (deg)	Date	Time (UT)	V (km s $^{-1}$)	ϕ_{L1} (deg)	Impact
Cat II					
57.5	2020 Dec 7	16:24	1407	13	None
53.1	2021 Apr 22	05:48	355	5	None
52	2021 May 9	12:00	266	4	Wind
38.6	2021 Oct 9	07:12	712	3	Unclear
34	2022 Feb 6	14:00	334	3	Wind
33.8	2022 Mar 7	00:12	241	26	None
33.7	2022 Mar 10	18:48	742	8	Wind
32.7	2022 Apr 3	16:48	502	18	None
32.9	2022 Mar 28	20:24	905	12	None
32.8	2022 Mar 30	18:00	641	24	None
32.2	2022 Apr 11	05:48	940	0	Wind
24.4	2022 Jul 15	15:48	320	19	Wind
21	2022 Aug 17	14:36	901	21	None
21	2022 Aug 14	12:48	481	18	None
Cat III					
52	2021 May 9	15:23	603	61	...
45.2	2021 Jul 24	00:36	587	57	...
47	2021 Oct 31	08:12	366	47	...
27.5	2022 Jun 13	03:12	1150	48	...
Cat IV					
57.8	2020 Nov 29	13:25	2077	96	...
50.5	2021 May 28	23:12	971	56	...
38.7	2021 Oct 7	02:00	299	90	...
34.4	2022 Feb 11	00:48	448	143	...
31	2022 May 1	08:36	700	106	...
30.5	2022 May 8	04:24	263	Back	...
28.9	2022 May 28	15:12	568	Back	...
27.4	2022 Jun 14	01:25	662	Back	...
26.2	2022 Jun 26	03:36	980	Back	...
23.6	2022 Jul 23	19:00	1026	Back	...
22.9	2022 Jul 30	02:24	609	Back	...
20.3	2022 Aug 26	10:36	553	62	...

Notes. Potential impact for CAT III and CAT IV CMEs was not investigated. Back refers to back-sided events as seen from L1.

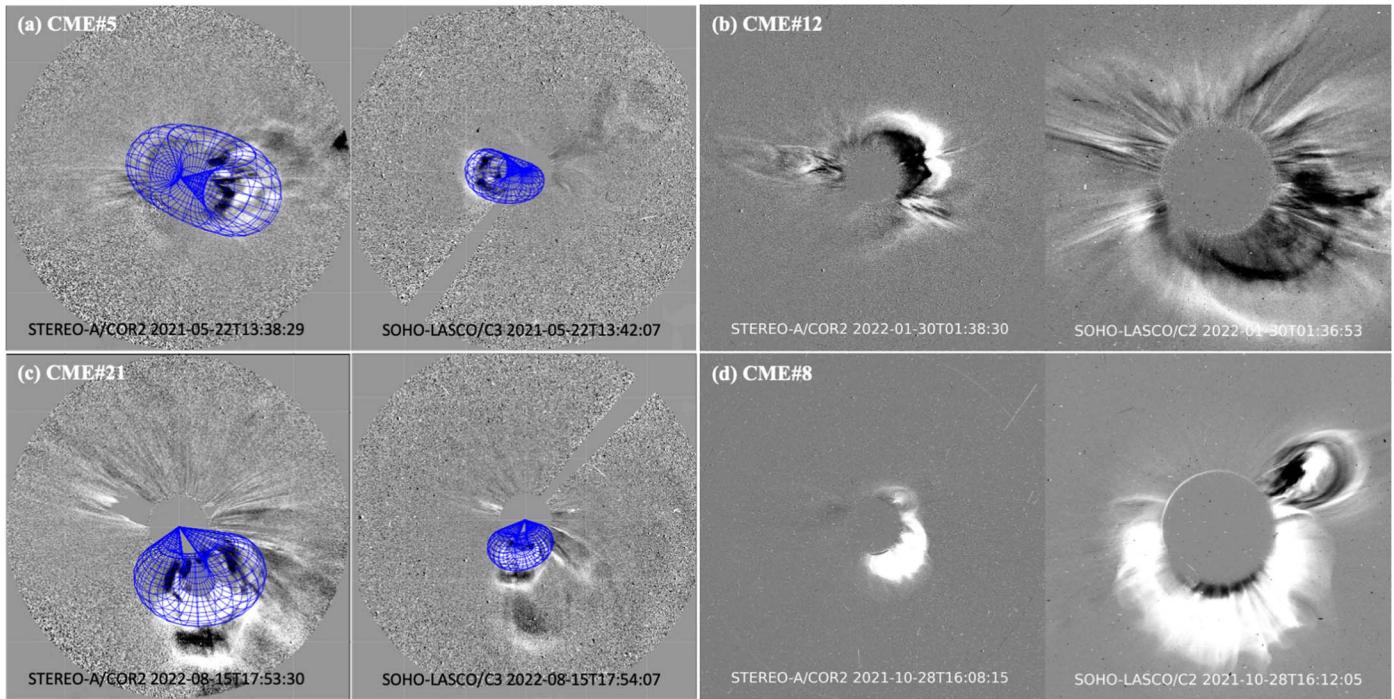


Figure 6. Remote-sensing observations from STEREO-A/COR2 (left) and SOHO/LASCO (right) of the CMEs discussed in Figures 2 and 3, from top left to bottom right, CMEs 5, 12, and 21 (all impacting both spacecraft) and CME 8 (impacting STEREO-A only). When the CME front or direction is not clear, the GCS reconstruction is also shown in blue overlay.

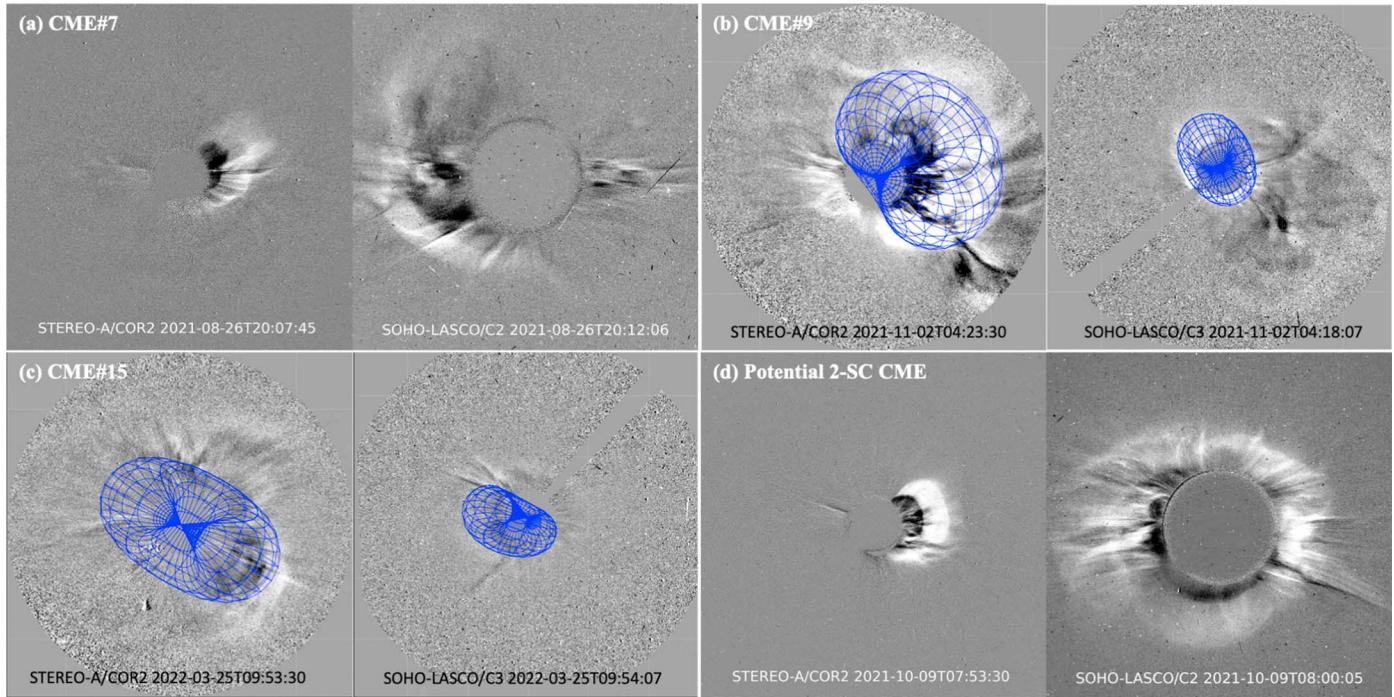


Figure 7. Same as Figure 6 for the CMEs discussed in Figures 4 and 5, from top left to bottom right, CME 7 (impacting STEREO-A only), CME 9 (impacting Wind only), CME 15 (no impact), and the additional CME that may have two spacecraft measurements.

ORCID iDs

Noé Lugaz  <https://orcid.org/0000-0002-1890-6156>
 Bin Zhuang  <https://orcid.org/0000-0002-5996-0693>
 Camilla Scolini  <https://orcid.org/0000-0002-5681-0526>
 Nada Al-Haddad  <https://orcid.org/0000-0002-0973-2027>

Charles J. Farrugia  <https://orcid.org/0000-0001-8780-0673>
 Réka M. Winslow  <https://orcid.org/0000-0002-9276-9487>
 Florian Regnault  <https://orcid.org/0000-0002-4017-8415>
 Christian Möstl  <https://orcid.org/0000-0001-6868-4152>

Emma E. Davies [ID](https://orcid.org/0000-0001-9992-8471) <https://orcid.org/0000-0001-9992-8471>
 Antoinette B. Galvin [ID](https://orcid.org/0000-0003-3752-5700) <https://orcid.org/0000-0003-3752-5700>

References

Akhavan-Tafti, M., Johnson, L., Sood, R., et al. 2023, *FrASS*, **10**, 1185603
 Al-Haddad, N., Galvin, A. B., Lugaz, N., Farrugia, C. J., & Yu, W. 2022, *ApJ*, **927**, 68
 Al-Haddad, N., Nieves-Chinchilla, T., Möstl, C., et al. 2013, *SoPh*, **284**, 129
 Al-Haddad, N., Poedts, S., Roussev, I., et al. 2019, *ApJ*, **870**, 100
 Allen, R. C., Smith, E. J., Anderson, B. J., et al. 2022, *FrASS*, **9**, 1002273
 Bailey, R. L., Möstl, C., Reiss, M. A., et al. 2020, *SpWea*, **18**, e02424
 Brueckner, G. E., Howard, R. A., Koomen, M. J., et al. 1995, *SoPh*, **162**, 357
 Burlaga, L., Sittler, E., Mariani, F., & Schwenn, R. 1981, *JGR*, **86**, 6673
 Byrne, J. P., Maloney, S. A., McAteer, R. T. J., Refojo, J. M., & Gallagher, P. T. 2010, *NatCo*, **1**, 74
 Cane, H. V., Richardson, I. G., & St Cyr, O. C. 1998, *GeoRL*, **25**, 2517
 Cartwright, M. L., & Moldwin, M. B. 2010, *JGRA*, **115**, 8102
 Davies, J. A., Harrison, R. A., Perry, C. H., et al. 2012, *ApJ*, **750**, 23
 Démoulin, P., Dasso, S., & Janvier, M. 2013, *A&A*, **550**, A3
 Démoulin, P., Janvier, M., Masías-Meza, J. J., & Dasso, S. 2016, *A&A*, **595**, A19
 Farrugia, C. J., Berdichevsky, D. B., Möstl, C., et al. 2011, *JASTP*, **73**, 1254
 Fox, N. J., Velli, M. C., Bale, S. D., et al. 2016, *SSRv*, **204**, 7
 Galvin, A. B., Kistler, L. M., Popecki, M. A., et al. 2008, *SSRv*, **136**, 437
 Good, S. W., & Forsyth, R. J. 2016, *SoPh*, **291**, 239
 Gopalswamy, N., Yashiro, S., Xie, H., Akiyama, S., & Mäkelä, P. 2015, *JGRA*, **120**, 9221
 Heinemann, S. G., Temmer, M., Farrugia, C. J., et al. 2019, *SoPh*, **294**, 121
 Howard, R. A., Moses, J. D., Vourlidas, A., et al. 2008, *SSRv*, **136**, 67
 Howard, T. A., & DeForest, C. E. 2012, *ApJ*, **746**, 64
 Hundhausen, A. J. 1993, *JGR*, **98**, 13177
 Janvier, M., Dasso, S., Démoulin, P., Masías-Meza, J. J., & Lugaz, N. 2015, *JGRA*, **120**, 3328
 Janvier, M., Démoulin, P., & Dasso, S. 2014, *A&A*, **565**, A99
 Jian, L. K., Russell, C. T., Luhmann, J. G., & Galvin, A. B. 2018, *ApJ*, **855**, 114
 Kilpua, E., Koskinen, H. E. J., & Pulkkinen, T. I. 2017, *LRSP*, **14**, 5
 Kilpua, E. K. J., Lee, C. O., Luhmann, J. G., & Li, Y. 2011, *AnGeo*, **29**, 1455
 Kilpua, E. K. J., Pomoell, J., Vourlidas, A., et al. 2009, *AnGeo*, **27**, 4491
 Lemen, J. R., Title, A. M., Akin, D. J., et al. 2012, *SoPh*, **275**, 17
 Lepping, R. P., Acuna, M. H., Burlaga, L. F., & Farrell, W. M. 1995, *SSRv*, **71**, 207
 Lepping, R. P., Wu, C.-C., & Berdichevsky, D. B. 2005, *AnGeo*, **23**, 2687
 Lin, R. P., Anderson, K. A., Ashford, S., et al. 1995, *SSRv*, **71**, 125
 Lopez, R. E. 1987, *JGR*, **92**, 11189
 Lugaz, N., Farrugia, C. J., Davies, J. A., et al. 2012, *ApJ*, **759**, 68
 Lugaz, N., Farrugia, C. J., Winslow, R. M., et al. 2018, *ApJL*, **864**, L7
 Lugaz, N., Hernandez-Charpak, J. N., Roussev, I. I., et al. 2010, *ApJ*, **715**, 493
 Lugaz, N., Lee, C. O., Jian, L. K., et al. 2023, *BAAS*, **55**, 249
 Lugaz, N., Manchester, W. B., & Gombosi, T. I. 2005, *ApJ*, **634**, 651
 Lugaz, N., Salman, T. M., Zhuang, B., et al. 2022, *ApJ*, **929**, 149
 Lugaz, N., Temmer, M., Wang, Y., & Farrugia, C. J. 2017, *SoPh*, **292**, 64
 Luhmann, J. G., Curtis, D. W., Schroeder, P., et al. 2008, *SSRv*, **136**, 117
 Luhmann, J. G., Gopalswamy, N., Jian, L. K., & Lugaz, N. 2020, *SoPh*, **295**, 61
 Martinić, K., Dumbović, M., Temmer, M., Veronig, A., & Vršnak, B. 2022, *A&A*, **661**, A155
 Möstl, C., Weiss, A. J., Bailey, R. L., et al. 2020, *ApJ*, **903**, 92
 Müller, D., St. Cyr, O. C., Zouganelis, I., et al. 2020, *A&A*, **642**, A1
 Nykyri, K., Ma, X., Burkholder, B., et al. 2023, *FrASS*, **10**, 1179344
 Owens, M. J., Lockwood, M., & Barnard, L. A. 2017, *NatSR*, **7**, 4152
 Owens, M. J., Merkin, V. G., & Riley, P. 2006, *JGRA*, **111**, 3104
 Prise, A. J., Harra, L. K., Matthews, S. A., Arridge, C. S., & Achilleos, N. 2015, *JGRA*, **120**, 1566
 Raouafi, N. E., Matteini, L., Squire, J., et al. 2023, *SSRv*, **219**, 8
 Regnault, F., Al-Haddad, N., Lugaz, N., et al. 2024, *ApJ*, in press ([arXiv.2311.14046](https://arxiv.org/abs/2311.14046))
 Richardson, I. G., & Cane, H. V. 2010, *SoPh*, **264**, 189
 Robbrecht, E., & Berghmans, D. 2004, *A&A*, **425**, 1097
 Rodriguez, L., Mierla, M., Zhukov, A. N., West, M., & Kilpua, E. 2011, *SoPh*, **270**, 561
 Salman, T. M., Lugaz, N., Winslow, R. M., et al. 2021, *ApJ*, **921**, 57
 Scolini, C., Winslow, R. M., Lugaz, N., & Poedts, S. 2023, *ApJ*, **944**, 46
 St. Cyr, O. C., Mesarch, M. A., Maldonado, H. M., et al. 2000a, *JASTP*, **62**, 1251
 St. Cyr, O. C., Plunkett, S. P., Michels, D. J., et al. 2000b, *JGR*, **105**, 18169
 Szabo, A. 2005, *AdSpR*, **35**, 61
 Thernisien, A. 2011, *ApJS*, **194**, 33
 Thernisien, A., Vourlidas, A., & Howard, R. A. 2009, *SoPh*, **256**, 111
 von Forstner, J. L. F. 2021, *johan12345/gcs_python*: Release 0.2.2, v0.2.2, Zenodo, doi:[10.5281/zenodo.5084818](https://doi.org/10.5281/zenodo.5084818)
 Wang, Y., Shen, C., Wang, S., & Ye, P. 2004, *SoPh*, **222**, 329
 Winslow, R. M., Lugaz, N., Schwadron, N. A., et al. 2016, *JGRA*, **121**, 6092
 Winslow, R. M., Lugaz, N., Scolini, C., & Galvin, A. B. 2021, *ApJ*, **916**, 94
 Xue, X. H., Wang, C. B., & Dou, X. K. 2005, *JGRA*, **110**, 8103
 Yashiro, S., Gopalswamy, N., Michalek, G., et al. 2004, *JGRA*, **109**, A07105
 Yu, W., Farrugia, C. J., Galvin, A. B., et al. 2016, *JGRA*, **121**, 5005
 Zhao, X. H., Feng, X. S., Feng, H. Q., & Li, Z. 2017, *ApJ*, **849**, 79
 Zurbuchen, T. H., & Richardson, I. G. 2006, *SSRv*, **123**, 31