

# Numerical Investigation of CO<sub>2</sub> Convective Transport in Stochastically Generated Heterogenous Media: Implications for Long-Term CO<sub>2</sub> Sequestration in Saline Aquifers

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# **Abstract**

Dissolution trapping is one of the most dominant mechanisms for CO<sub>2</sub> storage in subsurface porous media saturated with brine. The CO<sub>2</sub> dissolution rate and overall fluid flow dynamics in subsurface formations can vary significantly based on permeability variation. Although some numerical simulations have focused on these factors, detailed flow behavior analysis under nonuniform permeability distribution needs further study. For this purpose, we conduct simulations on the flow behavior of CO<sub>2</sub>-dissolved brine in two different heterogeneous media. The spatial permeability variations in the cell enable the analysis of complex subsurface storage phenomena, such as changes in finger morphology and preferential dissolution path. Finally, the amount of CO<sub>2</sub> dissolved was compared between each case, based on which we draw informed conclusions about CO<sub>2</sub> storage sites. The results demonstrated a preferential movement of CO<sub>2</sub>-dissolved regions toward high permeability regions, whereas a poor sweep efficiency was observed due to minimum dissolution in areas with lower permeability. Furthermore, simulation results also reveal uneven CO<sub>2</sub> concentration inside the convective fingers. This study provides fundamental insight into the change in flow behavior at heterogeneous regions, which could be translated into saline aquifer conditions. The proposed workflow in this study could be extended further to analyze complex heterogeneous storage systems at different flow regimes.

**Keywords:** CO<sub>2</sub> Sequestration, Dissolution Trapping, Relative Permeability Hysteresis, Capillary pressure, Diffusion

### Introduction

The atmospheric concentration of CO<sub>2</sub> increased by almost 50% compared to the beginning of the industrial revolution, causing irreversible damage to ecosystems (Aldred et al., 2021). The effects of rising surface temperatures due to increased CO<sub>2</sub> emissions are also evident from frequent natural disasters, which are reported to have caused around \$2.28 trillion in losses from 1980 (Shahriar & Khanal, 2023b). Recent studies have focused on CO<sub>2</sub> capture and storage technologies to achieve environmental sustainability

through CO<sub>2</sub> mitigations. One of the most promising ways to reduce the harmful effects of rising CO<sub>2</sub> concentrations is subsurface sequestration in saline aquifers, a major geological storage option because of the huge storage potential and wide availability (Shahriar & Khanal, 2023b). CO<sub>2</sub> is stored in saline aquifers through four trapping mechanisms: structural trapping, residual trapping, dissolution trapping, and trapping through mineralization. As observed in previous studies, dissolution trapping is the primary trapping mechanism in saline aquifers, where almost two-thirds of the injected CO<sub>2</sub> is stored through dissolution in brine (Khanal & Shahriar, 2022; Khosrokhavar et al., 2014).

Several experimental works have tried to unearth the fundamental aspects of dissolution trapping in saline aquifers using a Hele-Shaw cell, an apparatus built with two transparent plates with a small gap width between them (Faisal et al., 2015; Taheri et al., 2017). Using the Hele-Shaw, it is possible to visualize fluid flow on a 2D scale, including density-driven CO<sub>2</sub> transport, i.e., dissolution trapping. Furthermore, by varying the gap width and packing the Hele-Shaw cell with glass beads or strips with different thicknesses, different permeability and porosity conditions are obtainable, mimicking transport mechanisms for different saline aquifers (W. S. Amarasinghe et al., 2021; De et al., 2022; De Paoli et al., 2020; Singh et al., 2019). (Taheri et al., 2017) performed qualitative and quantitative tests to explore the density-driven natural convection during CO<sub>2</sub> storage under various conditions and different apertures and dipping angles. The objective parameters of their study were the onset time of convection, convective finger morphology, and dissolution flux of CO<sub>2</sub> in water. Dissolution of CO<sub>2</sub> under high pressure (1.5 to 5 bar) using Hele-Shaw cells was investigated in the work of (Outeda et al., 2014), where it was observed that higher pressure conditions increase the CO<sub>2</sub> mass influx, thus promoting faster subsurface CO<sub>2</sub> transport. Another study at highpressure conditions by (Mahmoodpour et al., 2019) at around 36.9 bar reported similar findings, stating that CO<sub>2</sub> sequestration is favorable at high-pressure, high-temperature (HPHT) conditions. Several studies also investigated the effects of heterogeneity in subsurface CO<sub>2</sub> storage using Hele-Shaw cells. Amidst different forms of heterogeneities present in porous media, such as chemical heterogeneity (change in mineralogical properties), petrophysical heterogeneity (permeability and porosity variation), and biological heterogeneity (biofilm growth), petrophysical-heterogeneities have a significant effect in dictating the dissolution pattern during CO<sub>2</sub> subsurface aquifer storage. Therefore, most studies have considered changes in permeability or porosity while conducting experiments or numerical simulations. For example, (W. Amarasinghe et al., 2020) conducted CO<sub>2</sub> dissolution experiments at 100 bar and 50 °C in a Hele-Shaw cell using glass bead size of 50 to 600 µm. They observed piston-like displacement in lower permeability cases, while finger shape patterns are visible in high permeability cases. Furthermore, a preferential fingering of CO<sub>2</sub>-rich highdensity water towards the highly permeable regions was observed. Similar findings were reported in the work of (Vosper et al., 2014) using glass beads with diameters of 0.4–0.6 mm. Meanwhile, a different way of changing permeability through the Hele-Shaw cell was implemented in our previous work (Shahriar & Khanal, 2023b), where, by changing the placement of sidewall shims, we got a complex permeability distribution inside the Hele-Shaw cell. It was observed that there is a preferential finger pathway towards the high permeability regions, translating to an uneven CO<sub>2</sub> dissolution profile in heterogeneous media.

Various studies have also considered numerical simulations to report CO<sub>2</sub> dissolution behavior in subsurface porous media; however, despite the more realistic nature of heterogeneous media during CO<sub>2</sub> sequestration operations, most simulation studies on CO<sub>2</sub> dissolution in saline aquifers use homogeneous models (Alipour et al., 2020; Taheri et al., 2017; Zhang et al., 2020). The effects of heterogeneity were numerically simulated by (Soltanian et al., 2017), where the facies spatial organization was reported to be a dictating parameter for CO<sub>2</sub> convective dissolution. Moreover, (Wang et al., 2023) reported a deteriorated CO<sub>2</sub> sequestration efficiency for reservoirs with a high degree of heterogeneity. Although these studies have considered the effects of heterogeneity, a comparative analysis of the CO<sub>2</sub> dissolution behavior between homogeneous and heterogeneous media (different heterogeneous distributions while keeping the same average permeability) needs further study. To address this issue, we report the CO<sub>2</sub> dissolution between

one homogeneous and two heterogeneous media with different permeability distributions. Moreover, in order to provide a comparative evaluation, we keep the average permeability of the three cases the same, allowing us to highlight only the effects of different heterogeneity formations. The remainder of the paper is organized as follows: First, we describe the governing equations, initial and boundary conditions, and solution method. Subsequently, we validate the numerical simulation model by matching the simulation results at different times to experimental results from a homogeneous model. Following this, we use the validated model to simulate two types of heterogeneous media to understand the dissolution profile. Furthermore, we compare the results to the homogeneous model to further explain the effects of permeability variation in 2D heterogeneous systems and discuss possible future studies. Finally, we present the main conclusions from this study.

# **Numerical Method**

## **Governing Equations**

The fluid inside the Hele-Shaw cell can be considered as continuum and incompressible under the experimental conditions (presented in Table 1); therefore, the Navier-Stokes equation and continuity were used to numerically simulate the fluid dynamics and mass transfer inside the cell (Chevalier et al., 2015; Faisal et al., 2015). The incompressible approximation is valid only because the density change of water with dissolved CO<sub>2</sub> is relatively small compared to the initial reference density(Emami-Meybodi et al., 2015). Therefore, the global formulation of the conservative equation for a porous media with CO<sub>2</sub> (gas phase, g) and water (liquid phase, l) is expressed in Eq. 1:

$$\frac{\partial}{\partial t} \left[ \sum_{\gamma = l, g} \left( \phi \rho_{\gamma} S_{\gamma} \omega_{\gamma}^{i} \right) \right] = - \sum_{\gamma = l, g} \nabla \left( \rho_{\gamma} V_{\gamma} \omega_{\gamma}^{i} \right) - \sum_{\gamma = l, g} \nabla \left( J_{\gamma}^{i} \right) + \sum_{\gamma = l, g} \nabla \left( m_{d\gamma} \omega_{\gamma}^{i} \right)$$

$$\tag{1}$$

Where,  $\gamma$  and *i* represent the phase (liquid or gas) and species (CO<sub>2</sub> or H<sub>2</sub>O), respectively. Furthermore,  $\rho$ , $\phi$ ,S, $\omega$ , and  $m_d$  are the density, porosity, saturation, mass fraction, and mass rate density, respectively. Boussinesq approximation is used for density calculations as the fluid properties are independent of the solute concentration at small CO<sub>2</sub> content in water. Furthermore, thermal effects caused by CO<sub>2</sub> absorption are negligible and can be ignored. A similar approach has been used in the numerical investigation of Zhang et al. (Zhang et al., 2020). The diffusive-convection process of CO<sub>2</sub> dissolution in water can be explained by the nonlinear governing equations as follows (Eq. 2–4):

$$\nabla . v = 0 \tag{2}$$

$$\frac{\partial \rho}{\partial t} + \nabla (\rho v \cdot v) = -\nabla P + \mu \nabla^2 v + F \tag{3}$$

$$\frac{\partial C_A}{\partial t} + v \nabla C_g = D_{CO_2} \nabla^2 C_g \tag{4}$$

Where,  $\rho$ , P,  $C_g$ ,  $\mu$  and  $D_{CO_2}$  represents the density, pressure,  $CO_2$  concentration, viscosity, and diffusion coefficient of  $CO_2$  in water. The liquid velocity is represented by V, and the buoyancy source in the gravity direction, represented by F, is given by:

$$F = \rho_o \left( 1 + \beta_g + C_g \right) g \tag{5}$$

Where,  $\beta_g$  represents the solutal expansion coefficient of CO<sub>2</sub>, and g is the gravitational acceleration.

Physical Model (Graphical Illustration)		Experimental Conditions and Fluid Properties	
Pure $CO_2$ Interface $P_{gas} = P_{atm}$		Dimensions	259 mm (width) and 242 mm (height)
Aqueous Phase (Water Saturated)	$\frac{dC}{dy} = 0$	Temperature	22 °C
		Pressure	1 atm
		Salinity	No salinity
		Aperture <sup>a</sup>	1 mm
		Water dynamic Viscosity	$0.95475 \times 10^{-3} \mathrm{Pa}\mathrm{s}$
		Density Variation <sup>b</sup>	0.29 kg/m <sup>3</sup>
$\frac{dC}{dx} = 0$		Diffusion coefficient of CO <sub>2</sub> in water	$1.4424 \times 10^{-9} \mathrm{m}^2/\mathrm{s}$

Table 1—Physical model and the experimental conditions considered for the simulation model.

<sup>a</sup>Considered only for the homogeneous base case, <sup>b</sup>Density difference between water saturated with CO<sub>2</sub> (higher) and pure water (lower)

## **Initial and Boundary Conditions**

This study's physical model (base case) is built similarly to our previous experimental configuration (Shahriar & Khanal, 2023a). The flow domain is at 22 °C and 1 atm. As shown in Table 1, the bottom and lateral boundaries have no-flow conditions, whereas the top boundary was set as a pure CO<sub>2</sub> interface.

#### **Solution Method**

With Ansys Fluent multiphysics software, the governing equations with initial and boundary conditions were numerically solved using the finite volume methodology. The pressure-correction method was used through the implicit iterative scheme SIMPLE algorithm, a transient algorithm commonly used for solving governing equations for density-driven convective flow (Ferziger et al., 2020; Soboleva, 2018; Zhang et al., 2020). We also selected a small time step of 0.1s, corresponding to the Newton-Raphson convergence criterion equal to  $10^{-6}$  (Faisal et al., 2015). Pressure discretization was performed by the pressure staggering option (PRESTO), whereas quadratic upstream interpolation for convection dynamics (QUICK) was used to discretize CO<sub>2</sub> components and the momentum term.

Additionally, in order to induce the convection phenomenon or the fingering pattern during the  $CO_2$  dissolution, we chose to include small sinusoidal perturbations), in the initial condition (at time  $t \le 0.1$  sec) instead of using numerical rounding errors to trigger the instabilities, i.e., (Farajzadeh et al., 2011; Zhang et al., 2020).

$$c = 1 + 0.01\sin(2\pi x/\lambda) \tag{6}$$

Where, c represents the dimensionless concentration,  $\lambda = \frac{1}{12}$ , and x specifies the horizontal direction, following the work of Farajzadeh et al. (Farajzadeh et al., 2011). The perturbations were located at the top interface where the saturation of dissolved  $CO_2$  is 1, similar to the work of Faisal et al. (Faisal et al., 2015). In real-life scenarios, the convection phenomenon is due to thermodynamic instabilities or other forms of heterogeneities (Farajzadeh et al., 2008; Kneafsey & Pruess, 2010).

# Validation of Simulation with Experimental Results

The operating conditions for the base case simulation were set to be the same as our previous experimental work (Shahriar & Khanal, 2023a). We validate the simulation results by qualitatively comparing the CO<sub>2</sub> finger morphology at different timescales, as shown in Fig. 1. For similar conditions and times, the simulation results are in good agreement with the experimental finger morphology. Furthermore, the overall CO<sub>2</sub> dissolution behavior is observed to be similar, where at the initial stage, diffusion is the primary

mass transfer mechanism, and the diffusive layer moves downward vertically without any instability or deformation. However, throughout time, the density difference between the CO<sub>2</sub>-dissolved water and the freshwater leads to a density-driven convective flow, which is evident through the formation of small fingers.

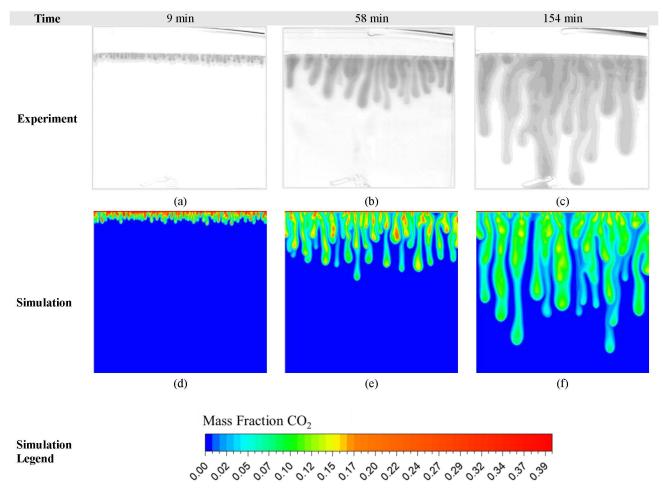


Figure 1—Validation of simulation results (d-f) with experimental results (a-c) at similar conditions for different timescales in the same order.

Furthermore, Fig. 1 shows that CO<sub>2</sub>-dissolved fingers merge and grow in width throughout time, a phenomenon accurately portrayed in the simulations. Additionally, due to cell-scale convection, i.e., CO<sub>2</sub> dissolved water at the top getting replaced by the freshwater due to convective flow, nascent fingers form at the top of the gas-water interface, which is better visible through the simulation results. The simulation results also reveal uneven CO<sub>2</sub> concentration inside the convective fingers, which has also been reported in previous studies (Shahriar & Khanal, 2023b; Zhang et al., 2020). Another factor worth noting is the slight discrepancy in the CO<sub>2</sub> finger morphology and vertical travel for the experimental and simulation results. This can be attributed to the fact that the Hele-Shaw cell aperture is not perfectly homogeneous for experimental studies due to the precision of the plexiglass used. Furthermore, it should be noted that several studies also report some minor dissimilarities in convective finger morphology while performing repeatability tests (Faisal et al., 2015; Shahriar & Khanal, 2023b). Nevertheless, similar CO<sub>2</sub> dissolution behavior and comparable finger travel between the results ensure the validity of the simulation.

# Methodology and Results of Heterogeneous Cases

## **Methodology for Creating Heterogeneous Regions**

This study considers two heterogeneous cases 1 and 2, with different permeability variations but the same mean permeability. The mean permeability (k) of the heterogeneous cases is set to be  $8.33 \times 10^{-8}$  m<sup>2</sup>, similar to the homogeneous case presented in **Section 3**, corresponding to a Hele-Shaw cell aperture (b) of 0.001 m, calculated as follows.

$$k = \frac{b^2}{12} \tag{7}$$

Sixteen different regions are created within the domain, assigning different permeability values corresponding to Hele-Shaw aperture of 0.0013, 0.0012, 0.0011, 0.0009, 0.0008, and 0.0007 m to the regions, as shown in the permeability profile in Fig. 2 and 3. In heterogeneity case 1, the permeability profile is set so that a high permeability region follows a low permeability region, and so on. In case 2, the low permeability regions are set together, allowing better connectivity between the high permeability regions. It should be noted that while the permeability distribution is different between the two cases, the mean permeability is the same, thus allowing us to specifically highlight only the effects of permeability distribution in heterogeneous media. A similar methodology is adopted in the work of (Mahyapour et al., 2022), where they considered heterogeneous media with permeability of 200 mD and 400 mD.

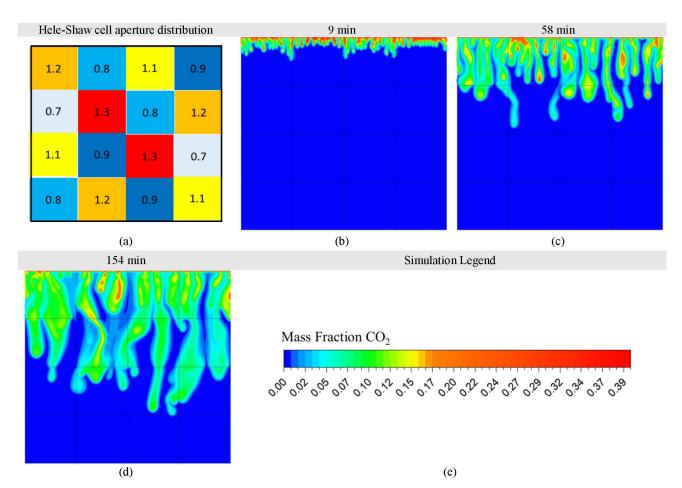


Figure 2—(a) Hele-Shaw cell aperture distribution (expressed in millimeters) and (b-d) the evolution of CO<sub>2</sub> dissolution through time for heterogeneous Case 1. The mass fraction of CO<sub>2</sub> is shown in (e).

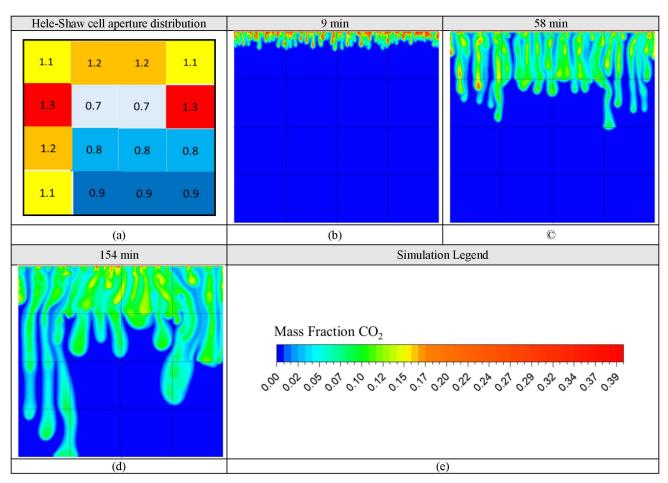


Figure 3—(a) Hele-Shaw cell aperture distribution (expressed in millimeters) and (b-d) the evolution of CO<sub>2</sub> dissolution through time for heterogeneous Case 2. The mass fraction of CO<sub>2</sub> is shown in (e)

#### Results and Discussions

As shown in Fig. 2, the permeability distribution is set to be in a certain manner so that a high permeability region follows a low permeability one, all while keeping the average permeability to be  $8.33 \times 10^{-8}$  (Hele-Shaw cell aperture of 1 mm). Although the initial stages of the  $CO_2$  dissolution are similar to the homogeneous case, the tip of the  $CO_2$  dissolved finger-shaped regions exhibits a preferential flow towards the high permeable regions, clearly observable on the second row at 58 minutes (Fig. 2c). Similar preferential movement is also observed in our previous study (Shahriar & Khanal, 2023b). Furthermore, this vertical finger travel is also high in high permeability regions. Compared to the homogeneous case (Fig. 1), heterogeneous case 1 has lower vertical travel at the end of the simulation (154 minutes). Furthermore, at the end of the simulation period, it is found that more  $CO_2$  is dissolved in the homogeneous case by comparing the average  $CO_2$  concentration for the homogeneous case (5.797 × 10<sup>-4</sup> mol/L) and heterogeneous case 1 (5.787× 10<sup>-4</sup> mol/L). (Taheri et al., 2018) also observed higher cumulative  $CO_2$  dissolved in homogeneous conditions than in heterogeneous cases. This shows that reservoirs with no or low permeability heterogeneity are preferable for higher  $CO_2$  dissolution.

Although the heterogeneous case 1 has permeability variation, the overall CO<sub>2</sub> dissolved region seems to move downwards uniformly. This can be attributed to the effects of low connectivity between the high permeability zones, preventing the formation of any particular preferential flow region. However, the effect of connectivity in creating preferential flow regions is observable in heterogeneous case 2 (Fig. 3), where the high permeability regions are stacked vertically, and the low permeability regions are placed together in the bottom right corner of the cell, as observed in the aperture distribution, shown in Fig. 3a.

As shown in Fig. 3, the CO<sub>2</sub> dissolved region follows a preferential path and an uneven CO<sub>2</sub> sweep in the low permeability regions. This channeling effect can be attributed to the good connectivity between the high permeability zones. Another interesting phenomenon is observed in the heterogeneous case 2; although faster vertical travel is exhibited in the high permeability zones (CO<sub>2</sub> dissolved region reaches the bottom of the cell in about 154 minutes), the overall CO<sub>2</sub> dissolution in heterogeneity case 2 is the lowest with average CO<sub>2</sub> concentration of  $5.769 \times 10^{-4}$  mol/L, compared to the homogeneous case ( $5.797 \times 10^{-4}$  mol/L) and heterogeneous case 1 ( $5.787 \times 10^{-4}$  mol/L). The lower CO<sub>2</sub> dissolution in this scenario can be attributed to the poor CO<sub>2</sub> dissolution in the low permeability regions. Since all the cases considered in the study have similar average permeability, it is safe to assume that the permeability distribution in different regions can lead to faster or slower CO<sub>2</sub> storage, also dictated by the connectivity of the permeability zones.

Although this study considers the effect of permeability distribution for an average permeability corresponding to a cell aperture of 1 mm, i.e., fluid flow in the Darcy regime based on the numerical considerations, future studies can consider both Hele-Shaw and 3D flow regimes for CO<sub>2</sub> dissolution in heterogeneous media. Another factor worth noting is that although this study considers fluid flow in the Darcy regime, where the effect of gap-induced dispersion is negligible, future studies can incorporate 3D models to report any change in dissolution behavior (Shahriar & Khanal, 2023b). Further insights into the effect of permeability distribution are obtainable through multiple variations of the heterogeneity cases and will be considered in future studies.

# **Conclusions**

This study explored the  $CO_2$  dissolution behavior due to permeability distribution using two heterogeneous cases with the same average permeability and compared them with the homogeneous case. Nonuniform finger movement is observed due to spatial permeability heterogeneity, where alongside faster vertical travel, the preferential finger movement is towards the high permeability zones. Furthermore, the simulation results highlighted the uneven  $CO_2$  concentration inside the convective fingers. The average  $CO_2$  concentration after 154 minutes follows the order of homogeneous case (5.797 × 10<sup>-4</sup> mol/L) > heterogeneous case 1 (5.787× 10<sup>-4</sup> mol/L) > heterogeneous case 2 (5.769 × 10<sup>-4</sup> mol/L), illustrating the importance of considering connectivity channels and permeability distribution during  $CO_2$  dissolution. The preferential flow path is more prominent in heterogeneous cases with high connectivity. The proposed methodology in this study can also be extended to analyze more complex heterogeneous systems for different fluid flow regimes. The findings of this study will add to the existing knowledge on the effect of permeability distribution in  $CO_2$  dissolution in saline aquifers.

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