

# Arbitrary Waveforms From a Fine-Resolution Microresonator-Based Spectral Shaper

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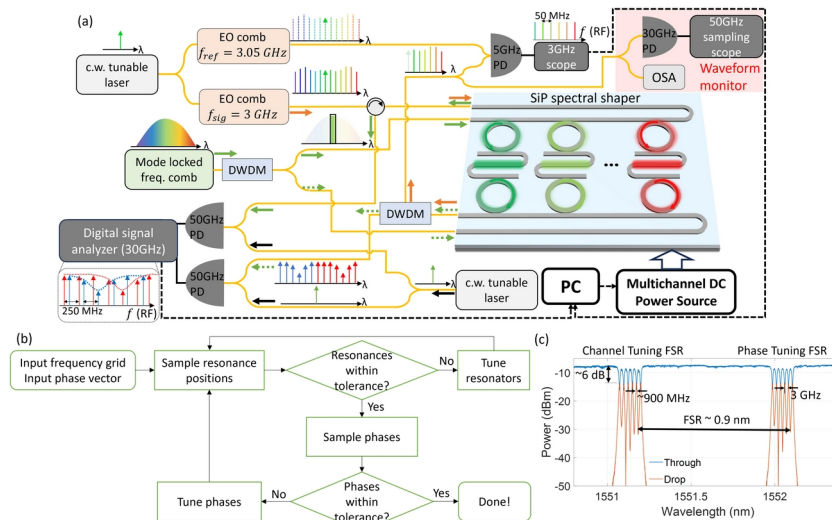
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**Abstract:** We report a scheme for programming microresonator-based spectral shapers and demonstrate it with a six-channel, sub-GHz linewidth, silicon photonic spectral shaper to generate arbitrary waveforms from optical lines of a 3 GHz electro-optic comb.

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Optical pulse shaping is a useful technique involving the manipulation of the relative amplitudes and phases of spectral slices from a broadband optical spectrum in order to shape the waveform in the time domain. Extensive research over the past couple decades has resulted in robust shaper systems utilizing bulk free space optics with components like diffraction gratings, lenses, and liquid crystal on silicon spatial light modulators [1]. These setups are well-suited for femtosecond repetition rate waveform synthesis, but they can be excessively large and incur a significant penalty in terms of optical loss when assembled to achieve fine spectral resolution (order of GHz) necessary for low repetition rate waveforms. There has been a recent desire for fine spectral resolution pulse shaping in, for example, generating picosecond waveforms with a low repetition rate for the control of electronic transitions within molecular ions or atoms [2, 3], or as the driving source for superconducting based waveform synthesizers [4], and in spectral quantum information processing [5]. Photonic integrated circuits (PICs) offer a low size, weight, and power (SWaP) chip-scale platform for pulse shaping. To this end, the arrayed waveguide grating (AWG) device has been extensively researched. However, the AWG requires long optical paths to be operated at a fine-resolution, resulting in a large footprint ( $\sim \text{cm} \times \text{cm}$ ) and significant optical loss and errors, especially for high-index contrast platforms like silicon-on-insulator (SOI) [6]. A comparatively small number of integrated shaper designs have implemented microresonator-based spectral dispersers [7, 8]. Microresonators benefit from a small footprint, resonance frequency tunability via the thermo-optic effect, and are amenable to fine-resolution filtering [9]. Yet, microresonators are challenging to control. Here, we demonstrate a silicon photonic spectral shaper on the SOI platform using fine-resolution microresonator filtering elements. We develop a control scheme to manage the thermal crosstalk within the system and enable programmability of the shaper. We finally demonstrate the synthesis of arbitrary waveforms in the time domain using optical lines from a 3 GHz electro-optic (EO) comb. A concept illustration of our six-channel SOI shaper is shown in Fig. 1(a). The integrated chip, which



**Fig. 1:** (a) Diagram of the experimental setup. Yellow lines indicate optical fiber lines, and dashed black lines indicate USB cables. PC: Personal Computer, OSA: Optical Spectrum Analyzer, DWDM: Dense Wavelength Division Multiplexer, PD: Photodiode, DC: Direct Current. (b) Flowchart of our shaper control scheme. (c) Transition spectrum of the spectral shaper programmed to compress six lines from a 3 GHz EO comb, as measured from a laser.

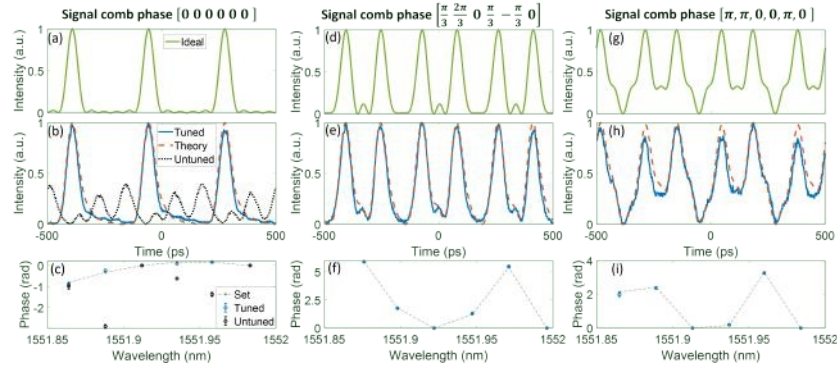


Fig. 2: (a,d,g) Ideal target waveforms in the time domain (green traces) and signal comb spectral phase (top) for the compressed, Talbot, and forked state, respectively. (b,e,h) Measured waveforms (blue trace) and the ideal waveforms (dashed orange traces) for the compressed, Talbot, and forked state, respectively. An additional waveform is shown in (b) (black dotted trace) with an untuned phase and phase shifters driven to  $\sim \pi$  phase. (c,f,i) Target (gray dashed) and measured (blue scattered with error bars) signal comb spectral phase for the compressed, Talbot, and forked states, respectively. An additional measurement is shown in (c) (black scattered with error bars) for the untuned phase state, with phase shifters driven to induce  $\sim \pi$  phase.

channel, a resonator first demultiplexes a spectral slice of an input optical signal and routes it to a thermo-optic phase shifter before being multiplexed back onto a common output waveguide via a second identical resonator. A six-channel system (twelve resonators and six phase shifters) is built on the SOI platform fabricated by Photonics. Resonators are designed using multimode waveguides for the benefit of a reduced sidewall roughness and propagation loss. A single shaper channel exhibits a  $\sim 900$  MHz linewidth,  $\sim 115$  GHz FSR, and  $\sim 6$  dB insertion loss. Doped silicon slabs placed adjacent to the waveguides are used for thermo-optic tuning of resonators and phase shifters. Loopback waveguides are formed from both the input and output of the shaper. Our measurement setup is also shown in Fig. 1(a). Using the loopback waveguide bidirectionally and employing spectral filters, we simultaneously perform multi-heterodyne spectroscopy (MHS) and dual-comb spectroscopy (DCS) measurements on two adjacent free spectral ranges (FSRs) of the shaper (Fig. 1(c)) to rapidly ( $\sim \mu$ s) query the resonant frequencies of resonators on both sides of the shaper and the phases applied to each line, respectively. For the power from one of the EO combs used for DCS and routing it to the waveform monitor section in Fig. 1(b), we can view the temporal profile of six comb lines after propagation through the shaper in real time. A Python routine (block diagram in Fig. 1(b)) is written to actively tune channels onto a particular frequency grid and distinct phases with feedback from MHS and DCS signals.

As a demonstration of our method, we synthesize three waveforms from six optical lines of the 3 GHz EO comb. The ideal waveforms, the measured and ideally measured (given our measurement system response) waveforms, and the measured phase on each shaper channel are shown in Fig. 2(a-i). The first demonstration is to combine the six lines from the EO comb to form transform-limited pulses. The MHS scheme is run to align channels onto a 3 GHz grid without tuning the phases. However, in our scheme phase shifters are driven at the start to a value slightly higher than  $\pi$  phase. The temporal waveform is shown in Fig. 2(b) (black dotted trace) for these six channels. With phases tuned, the pulses are compressed (blue trace in Fig. 2(b)). The spectral transmission response of the system for this state is shown in Fig. 2(c) and the untuned and tuned phases are shown in Fig. 2(c). Next, we target two additional waveforms: a doubled repetition-rate waveform following the temporal Talbot effect and a forked waveform. The temporal and spectral results for these states are shown in Fig. 2(d-i). These waveforms are highly sensitive to phase deviations and therefore phase tuners needed to be iteratively adjusted in fine steps to achieve the agreement shown in the figures. This requirement could be eliminated with a direct measurement of the reference EO comb phase.

In conclusion, we experimentally demonstrated a scheme for programming microresonator-based spectral shapers. A set of three waveforms are synthesized using a six-channel sub-GHz resolution SOI shaper. This work shows promise towards realizing fully integrated fine-resolution pulse shaping systems for the next generation of pulse shaping applications.

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## References

1. A. M. Weiner, Opt. Commun. 284(15), 3669-3692 (2011).
2. Y. Ma et al., Opt. Express 28(12), 17171-17187 (2020).
3. P. R. Stollenwerk et al., Phys. Rev. Lett. 125, 113201 (2020).
4. D. Lee et al., APL Photonics 8, 086115 (2023).
5. H.-H. Lu, et al., Optica, 5(11), 1455-1460 (2018).
6. M. Gehl et al., Opt. Express, 25(6), 6320-6334 (2017).
7. A. Agarwal et al., J. Light. Technol., 24(1), 77-87 (2006).
8. J. Wang et al., Nat. Commun., 6, 5957 (2015).
9. L. M. Cohen et al., arXiv Preprint, 2309.17222, (2023).