The January 2022 Hunga eruption cooled the southern hemisphere in 2022 and 2023

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Abstract

The 2022 Hunga volcanic eruption injected a significant quantity of water vapor into the stratosphere while releasing only limited sulfur dioxide. It has been proposed that this excess water vapor could have contributed to global warming, potentially pushing temperatures beyond the 1.5°C threshold of the Paris Climate Accord. However, given the cooling effects of sulfate aerosols and the contrasting impacts of ozone loss (cooling) versus gain (warming), assessing the eruption's net radiative effect is essential. Here, we quantify the Hunga-induced perturbations in stratospheric water vapor, sulfate aerosols, and ozone using satellite observations and radiative transfer simulations. Our analysis shows that these components induce clear-sky instantaneous net radiative energy losses at both the top of the atmosphere and near the tropopause. In 2022, the Southern Hemisphere experienced a radiative forcing of -0.55 ± 0.05 W m^{-2} at the top of the atmosphere and -0.52 ± 0.05 W m^{-2} near the tropopause. By 2023, these values decreased to -0.26 ± 0.04 W m^{-2} and -0.25 ± 0.04 W m^{-2}, respectively. Employing a two-layer energy balance model, we estimate that these losses resulted in cooling of about -0.10 ± 0.02 K in the Southern Hemisphere by the end of 2022 and 2023. Thus, we conclude that the Hunga eruption cooled rather than warmed the Southern Hemisphere during this period.

Introduction

On 15 January 2022, the Hunga Tonga volcano erupted from a shallow (~200 m deep) submarine vent, injecting unprecedented amounts of water vapor[1-2, 4-6] (~ 120-150 Tg) into the stratosphere. In addition, in situ and space-borne observations show that the eruption released a moderate amount of sulfur dioxide (~0.4-0.7 Tg) into the stratosphere, which was subsequently converted to sulfate aerosols over several weeks[2,4,6,9-19]. These two components, sulfate aerosol and water vapor, have opposite climate effects. Sulfate aerosols cool by scattering solar radiation, while water vapor, as a greenhouse gas, leads to stratospheric cooling and surface warming[12]. Its net radiative effect at the top of the atmosphere (TOA) depends on altitude, potentially causing either warming or cooling[12-15]. Given the unprecedented stratospheric water vapor injection, several studies have hypothesized that the Hunga eruption resulted in net climate warming[11-12]. Assessing this possibility is crucial, as volcanic-induced warming could increase the likelihood of temporarily exceeding the 1.5°C threshold set by the Paris Agreement[5] in the coming years. In contrast, past large subaerial (continental arc) eruptions—such as the 1991 Pinatubo, 1883 Krakatau, and 1815 Tambora events—led to prolonged global cooling due to the dominant influence of sulfate aerosols[18-24].

The efficacy of sulfate aerosol-induced cooling is influenced by aerosol properties such as number concentration, mass concentration, particle size distribution, and residence time[3,6,15,17], all of which substantially impact mass extinction efficiency (MEE)—a measure of an aerosol's ability to attenuate radiation per unit mass[25]. MEE depends on particle radius through two primary factors: (1) Extinction efficiency: This parameter exhibits a non-monotonic dependence on particle size, with its first peak approximately 25% greater than its second[25]. (2) Total geometric cross-section: For a fixed aerosol mass, the total geometric cross-section is inversely proportional to particle radius[25]. Consequently, smaller particles have a larger total cross-sectional area per unit mass, which enhances scattering and increases MEE.

Explosive volcanic eruptions can potentially disrupt stratospheric ozone concentrations, particularly when sulfate aerosols interact with a moist stratosphere[25-32]. Stratospheric ozone depletion leads to net cooling at Earth's TOA due to reduced shortwave absorption by ozone[22,26-37]. The climatic effects of shallow submarine eruptions can differ significantly from those of

110 subaerial eruptions due to variations in water vapor, SO_2, and halogen gas emissions, which

111 influence ozone chemistry[41-43]. Thus, accurately determining the net impact of the Hunga

112 eruption on stratospheric composition and radiative forcing is crucial. Most submarine eruptions

113 occur in deep water and release minimal volcanic material into the atmosphere. However, the

114 Hunga eruption—the largest shallow submarine eruption of the satellite era—offers a rare

115 opportunity to assess the climate impacts of such events[38-40]. Thus, this requires assessing the

116 perturbation of three key radiatively active species: water vapor, ozone, and sulfate aerosols.

117

118   Previous studies have either focused on a single atmospheric constituent or relied on

119 simplified scaling analyses to estimate radiative forcing in the first and second years after the

120 eruption[14,15]. This has likely contributed to conflicting conclusions, with some studies suggesting

121 that the Hunga eruption caused warming[5], while others found it led to cooling[14,15]. For instance,

122 Schoeberl[14] estimated a net negative (cooling) radiative forcing by combining a scaling analysis

123 based on past subaerial eruptions (for aerosols) with radiative calculations for water vapor[43-46]. In

124 a separate study, Schoeberl[15] applied this simplified aerosol radiative forcing scaling method to

125 demonstrate a global reduction in downward radiative flux at the tropopause over two years,

126 attributing this decline to changes in water vapor, aerosols, and ozone. In contrast, Jenkins[5]

127 concluded that the Hunga eruption could contribute to global warming. However, this study only

128 considered the impact of stratospheric water vapor (SH_2O) perturbation while neglecting the

129 effects of aerosols, which are critical to the overall radiative balance.

130

131   Here, we conduct a comprehensive analysis of the instantaneous net radiative response to

132 the Hunga eruption in 2022 and 2023 by: (a) Utilizing satellite remote sensing observations to

133 quantify the spatiotemporal distribution of SH_2O, ozone, and sulfate aerosols (including their

134 size distributions). (b) Using these observations as inputs for idealized 1D radiative transfer

135 model simulations. Here, the model is run independently for each grid cell across the near-global

136 domain, producing a three-dimensional output that captures both spatial and vertical variability

137 within the stratosphere.

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Results

We used space-borne remote sensing data from SAGE III[47-48] aboard the International Space Station (ISS) (solar occultation, v5.3; available since June 2017) to quantify the three-dimensional distribution of key radiatively active stratospheric species affected by the Hunga eruption: SH_2O, aerosols, and ozone (see Methods and Supplementary Figure 1). SAGE III/ISS employs a solar occultation technique, capturing up to 31 measurements daily during sunrise and sunset to generate dawn and dusk stratospheric profiles. This method requires approximately one month to cover latitudes from ~60°N to 60°S[47-48], providing detailed monthly coverage of the tropics and mid-latitudes, with seasonal variations influencing the sampling pattern. Additionally, we use multi-wavelength SAGE III/ISS observations to retrieve stratospheric aerosol particle size before and after the Hunga eruption.

For our radiative forcing analysis, we used SAGE III/ISS data due to its exceptional vertical resolution of 0.5 km. This high-resolution dataset on radiatively active species—water vapor, ozone, and aerosols—is critical for accurately estimating atmospheric heating rates and radiative forcings at the TOA and tropospause[49]. However, SAGE III/ISS data have known limitations, including a dry bias in water vapor measurements[50]. Earlier versions of SAGE III showed a ~10% dry bias in the stratosphere, which improved to ~5% in Version 5.2 and 5.3 for the mid-stratosphere[51,52]. Additionally, water vapor data below ~20 km are noisier due to aerosol and cloud interference, contributing to ~10–20% uncertainty in radiative forcing estimates, particularly in the lower stratosphere[50]. Furthermore, the infrequent sampling of aerosol extinction profiles at multiple wavelengths introduces biases in stratospheric aerosol optical depth (SAOD) estimates. To address this, we compare SAGE III-derived SAOD values with those from multiple instruments, including OMPS-NASA[53-56] (Ozone Mapping and Profiler Suite; NASA aerosol retrieval algorithms), OMPS-SASK[53-56] (Ozone Mapping and Profiler Suite; University of Saskatchewan aerosol retrieval algorithms), OSIRIS[56] (Optical Spectrograph and Infrared Imaging System), and GloSSAC[56,57] (Global Space-based Stratospheric Aerosol Climatology) (see Supplementary Figure 2).

We calculated the changes in SH_2O, sulfate aerosols, and ozone during the first and second years following the Hunga eruption ("Hunga-2022" and "Hunga-2023") relative to their

172   variations in the pre-eruption period (7 June 2017 – 9 December 2021; "CLIM") using SAGE-

173   III/ISS observations. Perturbed stratospheric profiles for water vapor and aerosols from the

174   eruption are identified as those with values exceeding two standard deviations above the

175   background climatology. Similarly, ozone perturbations from the eruption are identified using a

176   similar filtering criterion as described by Wilmouth[35]. To compare the SAGE-III based

177   perturbations in SH_2O and ozone following the eruption, we used long-term (2005–2023) data

178   from the Aura Microwave Limb Sounder (MLS[58]). Here, the absolute changes in water vapor

179   and ozone are compared with SAGE-III observations (Supplementary Figure 3a–d for SH_2O and

180   ozone from MLS[58]; 3e–h for SH_2O and ozone from SAGE-III).

181

182   We computed the instantaneous net (shortwave and longwave) atmospheric radiative

183   heating rates and perturbations to the TOA and near the tropopause energy budgets using

184   the LibRadtran[59] radiative transfer model (see Supplementary Figure 1 and Methods). These

185   energy budget perturbations serve as proxies for the overall atmospheric energy balance,

186   accounting for the contributions of stratospheric water vapor, ozone, and sulfate aerosols. In the

187   following sections, we first examine the individual radiative effects of each perturbed species

188   before analyzing their combined net impact.

189

190   **Changes in stratospheric water vapor after Hunga eruption**

191   We show the absolute changes in the zonal-mean latitude-altitude variations of

192   stratospheric H_2O mixing ratio between the (Fig. 1a) "Hunga-2022" and "CLIM" periods and

193   (Fig. 1b) "Hunga-2023" and "CLIM" periods based on the SAGE III/ISS observations.

194

195   We observe that the SH_2O injected by the Hunga eruption spread throughout the middle

196   stratosphere from the southern mid-latitudes to the northern tropics during the first year (Fig. 1a).

197   Around the eruption latitude (20°S–0°S latitude), we observed a 2 km thick layer, extending

198   from 27 to 29 km, with a 90% enhancement in SH_2O mixing ratio relative to the background

199   climatology (Fig. 1a and Supplementary Figure 4a). Further south, from 20°S to 50°S, the

200   increase in stratospheric water vapor mixing ratio was more modest, indicating lower

201   concentrations of water vapor due to its latitudinal transport by stratospheric wind patterns. The

202   near-global SH_2O content map, integrated from the lower to the upper stratosphere (up to an
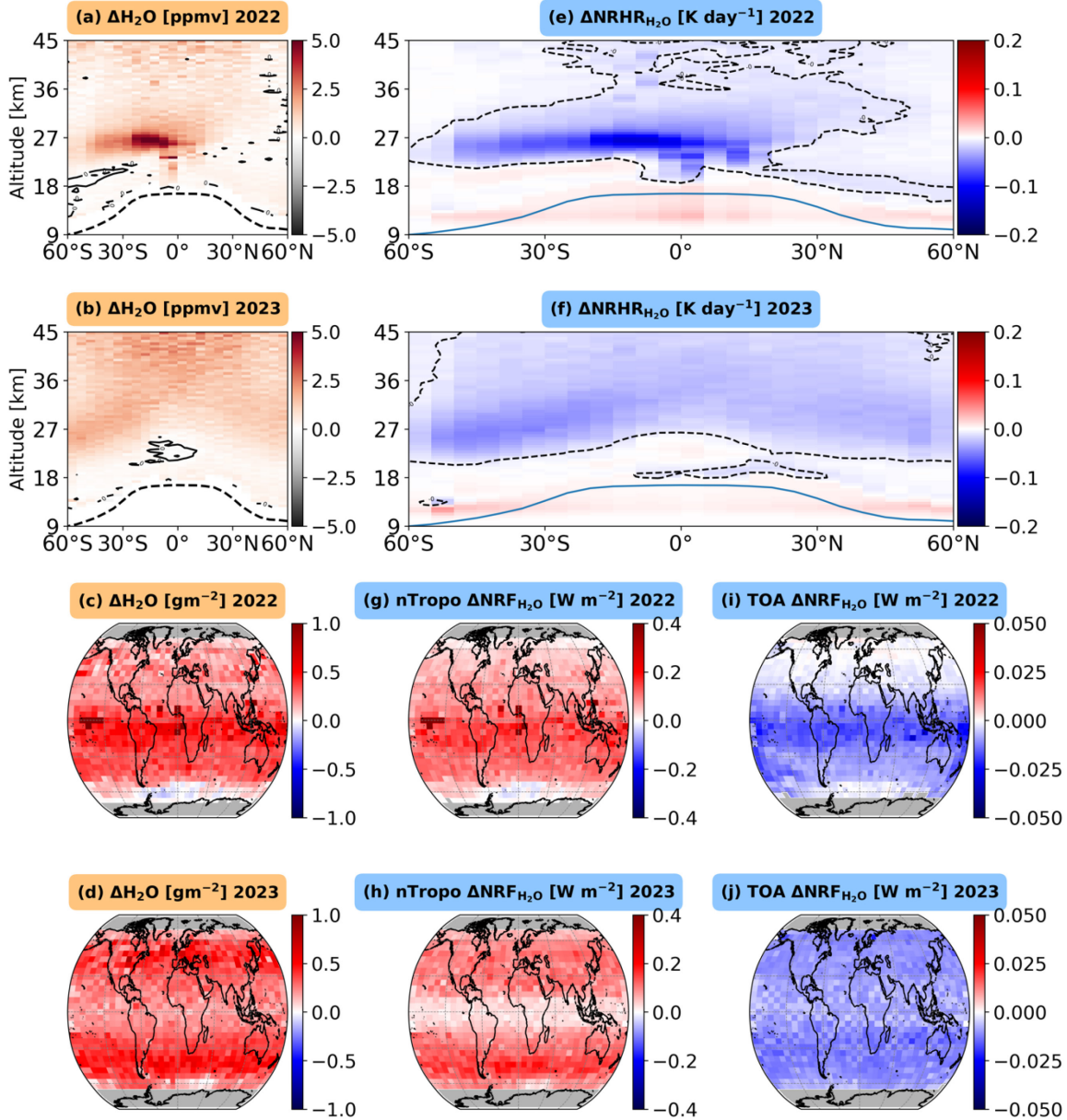
Figure 1: Observed enhancement of stratospheric water vapor and the associated simulated instantaneous radiative forcing during 2022 and 2023 following the Hunga eruption. Panels (a) and (b) show the zonal- and annual-mean latitude–altitude variation in the change of the stratospheric water vapor (H_2O) mixing ratio [ppmv] for 2022 and 2023, respectively, relative to the reference climatology (CLIM) period from 2017 to 2021. Contours denote near-zero changes, and the dashed black line represents the tropopause height. Panels (c) and (d) show the

213   *corresponding vertically integrated changes from the lower stratosphere to the upper*

214   *stratosphere. Panels (e) and (f) show the latitude–altitude distribution of the difference in net*

215   *(longwave + shortwave) radiative heating rates (ΔNRHR; K day⁻¹) of water vapor within the*

216   *stratosphere between each post-eruption year and the CLIM period. Panels (g) and (h) depict*

217   *changes in near-tropopause net radiative forcing (ΔNRF_{H_2O}; W m⁻²) for 2022 and 2023,*

218   *highlighting perturbations primarily observed near the lowest levels of the lower stratosphere*

219   *(above the tropopause height). Panels (i) and (j) show the changes at the top-of-atmosphere*

220   *(TOA) net radiative forcing (ΔNRF_{H_2O}; W m⁻²) for 2022 and 2023, respectively.*

221

222   In the second year (2023) following the eruption, the injected SH_2O has spread from the

223   southern to the northern hemisphere, covering the middle to the upper stratosphere. The

224   substantial dispersion of SH_2O during the second year has led to a decrease in its mixing ratio

225   and concentration (Fig. 1b, d; see relative changes in SH_2O in Supplementary Figure 4a, b). This

226   accumulated SH_2O has reduced by around 10% relative to the first year, as observed from

227   SAGE-III (Fig. 1d; Table 1). The widespread vertical distribution of SH_2O in 2023, in contrast to

228   2022, is attributed to the stratospheric Brewer-Dobson circulation in combination with some

229   perturbations in stratospheric circulation patterns due to stratospheric

230   temperature perturbations following the Hunga eruption[33,38].

231

232   We also quantify the impact of the 2022 Hunga eruption on the seasonal timeline of

233   changes in SH_2O concentration from 2022 to 2023, compared to its seasonal 5-year climatology

234   value between 2017 and 2021 (Fig. 2a). We observe a distinct increase in SH_2O concentration

235   from January 2022 onwards from SAGE-III observations. The peak in absolute changes in SH_2O

236   concentration value is approximately 0.5 g m⁻² (30% higher relative to the background value) in

237   the SH during spring season (MAM) of 2022 (Fig. 2a). The increase in SH_2O gradually decreases

238   in the summer (JJA) and autumn (SON) of 2022 and other seasons of 2023. Note that the

239   decrease in SH_2O concentration in the SH from 2022 to 2023 is minimal with a decline of around

240   10-15% (Fig. 2a). The NH and near-global domain also experience a substantial increase in the

241   magnitude of SH_2O in 2022 and 2023. The strong vertical transport of tropical middle

242   stratosphere SH_2O into the upper stratosphere and towards mid- and high-latitudes regions in

243   2023 in the SH and NH are attributed to Brewer-Dobson circulation[33,38, 46], as evident from the

244   vertical and spatial map of SH_2O in Fig. 1a-d.

245

246   **Net radiative heating rate and radiative forcing of the Hunga-perturbed SH_2O: 2022 versus**

247   **2023**

248   We used idealized radiative transfer model simulations (see Methods) to assess the

249   stratospheric radiative heating rates (Fig. 1e, f) and radiative forcing of the perturbed SH_2O near

250   the tropopause (Fig. 1g, h) and at the TOA (Fig. 1i, j). In 2022, in the middle stratosphere (24-26

251   km) around the eruption latitude, the additional SH_2O caused a cooling of -0.20 K day^{-1} due to

252   enhanced emission of longwave radiation (Supplementary Figure 5a). This cooling was partially

253   offset by heating of up to +0.05 K day^{-1} from the additional absorption of SW radiation by the

254   SH_2O (Supplementary Figure 5g). These effects combined to produce a net radiative cooling rate

255   of up to -0.15 K day^{-1} (dark blue color in Fig. 1e). During 2023, the vertical spreading of the

256   SH_2O layer in both hemispheres exhibited a net cooling effect below -0.1 K day^{-1} (Fig. 1f).

257

258   During 2022 and 2023, the SH_2O perturbation in the middle and upper stratosphere

259   following the eruption substantially affected the clear-sky net instantaneous radiative forcing

260   near the tropopause (Fig. 1g, h) and at the TOA (Fig. 1i, j) relative to background climatology.

261   Enhanced SH_2O, a potent greenhouse gas, substantially increases downward longwave radiation

262   emission from the water vapor layer, contributing to a higher net radiative forcing within Earth's

263   atmosphere while having minimal impact on shortwave radiation (Supplementary Figure 5g-l).

264   We find that the Hunga eruption's impact on SH_2O in 2022 increased the net radiative flux by

265   +0.12 $\pm$ 0.01 Wm^{-2} near the tropopause in the SH and +0.10 $\pm$ 0.01 Wm^{-2} in the near-global

266   domain (Table 1). In contrast, at Earth's TOA, the enhanced SH_2O in 2022 decreased the net

267   radiative flux by ~-0.010 $\pm$ 0.001 Wm^{-2} in the SH and ~-0.008 $\pm$ 0.001 Wm^{-2} in the near-global

268   domain (Table 1). This is due to the enhanced longwave radiation emission from water vapor in

269   the upper stratosphere.

270

271   During the summer of 2022, we quantified a peak positive instantaneous net radiative

272   flux of +0.15 $\pm$0.02 Wm^{-2} near the tropopause and a negative instantaneous net radiative flux of

273   -0.014 $\pm$0.001 Wm^{-2} at TOA in the SH, driven by enhanced SH_2O (Fig. 2e, i). In the SH, the

274 increased positive net instantaneous radiative flux of $+0.15 \pm 0.02$ Wm^{-2} near the tropopause

275 contributed to a surface warming of $+0.015$ K ± 0.004 K by the end of 2022, as estimated using

276 the FaIR[7-8] simple two-layer energy balance climate emulator (see Methods; light cyan color in

277 Supplementary Figure 6e). The increased net instantaneous radiative forcing and its associated

278 uncertainty near the tropopause and at the TOA are calculated solely based on changes in SH$_2$O,

279 assuming a constant background temperature profile for both pre- and post-eruption periods.

280 Notably, these calculations do not explicitly account for the influence of clouds and temperature

281 adjustments, which may lead to an underestimation of uncertainties.

282   In 2023, Hunga-perturbed SH$_2$O increased the net radiative flux near the tropopause

283 by $+0.09 \pm 0.01$ Wm^{-2} in the SH and $+0.08 \pm 0.01$ Wm^{-2} in the NH. These changes in radiative

284 flux led to a mean surface warming of $+0.022 \pm 0.003$ K in the SH and $+0.017 \pm 0.002$ K in the

285 NH (see blue and cyan colors in Supplementary Figure 6e). At Earth's TOA, Hunga-perturbed

286 SH$_2$O decreased the net instantaneous radiative flux by -0.01 ± 0.001 Wm^{-2} in the SH and -0.009

287 ± 0.001 Wm^{-2} in the near-global domain.

288   In aggregate, the SH$_2$O in the first year (2022) and second year (2023) following the

289 eruption contributed to increased radiative forcing near the tropopause and decreased radiative

290 forcing at Earth's TOA, primarily in the SH, but without a strong longitudinal pattern (Fig. 1c, g,

291 i). In 2023, the widespread dispersion of SH$_2$O led to tropospheric warming near the tropopause

292 in both hemispheres (Fig. 1d, h).

293

294   The spread of stratospheric water vapor is driven by the Brewer–Dobson circulation, a

295 large-scale atmospheric circulation pattern that transports tropical tropospheric air into the

296 stratosphere, followed by its poleward movement and descent at higher latitudes[6,46].

297 Furthermore, the perturbed water vapor levels may influence this background circulation. A

298 more detailed investigation of these interactions would require the application of advanced

299 climate modeling, which is beyond the scope of this study.
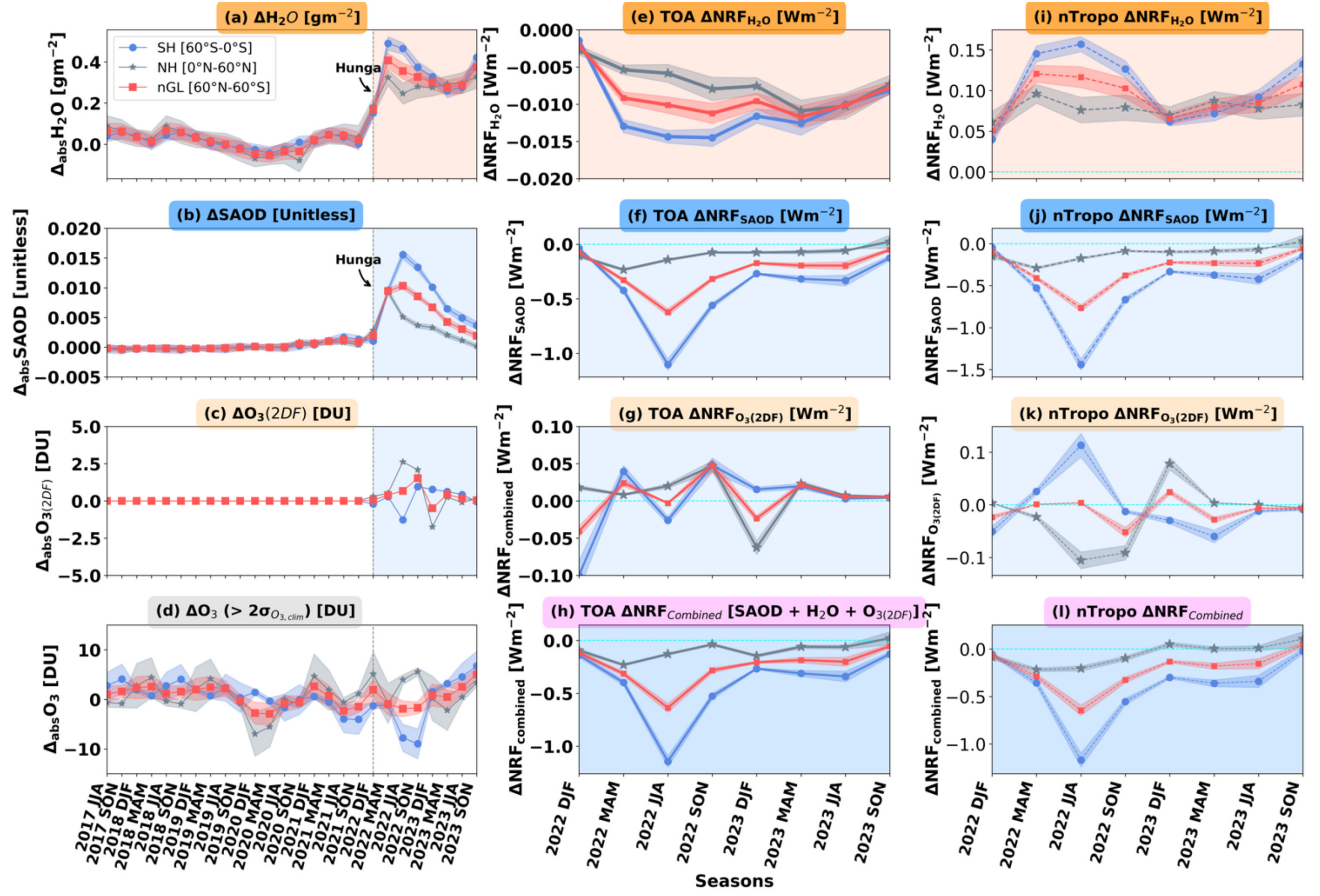
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**Figure 2: Evolution of the absolute changes and corresponding top-of-atmosphere (TOA) and
near-tropopause instantaneous net radiative forcing (ΔNRF) of three radiatively important
stratospheric species perturbed by the Hunga eruption.** *Panels (**a**)–(**d**) show the evolution of the
seasonal mean absolute changes (relative to the reference climatology [CLIM] period) from
2017 until 2023 for: (**a**) stratospheric water vapor (SH$_2$O) concentration, (**b**) stratospheric
aerosol optical depth (SAOD), (c) ozone concentration (with 2D-filtered condition), and (**d**)
ozone concentration (unfiltered, showing values exceeding 2σ above background climatology,
UNF). These values are presented for the Southern Hemisphere (SH; blue line), the Northern
Hemisphere (NH; grey line), and near-global (60°S–60°N; nGL; red line). Panels (**e**)–(**h**) show
the clear-sky net (longwave plus shortwave) radiative forcing (ΔNRF) at TOA due to post-
eruption perturbations in: (**e**) SH$_2$O, (**f**) SAOD, (**g**) ozone (2D-filtered; 2DF), and (**h**) the
combined effect of SH$_2$O + SAOD + ozone (2D-filtered). Panels (**i**)–(**l**) show the corresponding
dashed blue, grey, and red lines indicate the clear-sky ΔNRF near the tropopause. Dashed cyan
lines indicate the horizontal bar near zero in (**e-h**) and (**i-l**). Shaded regions represent the
interannual standard deviation (1σ) of the background values (except for 2D-filtered ozone,*
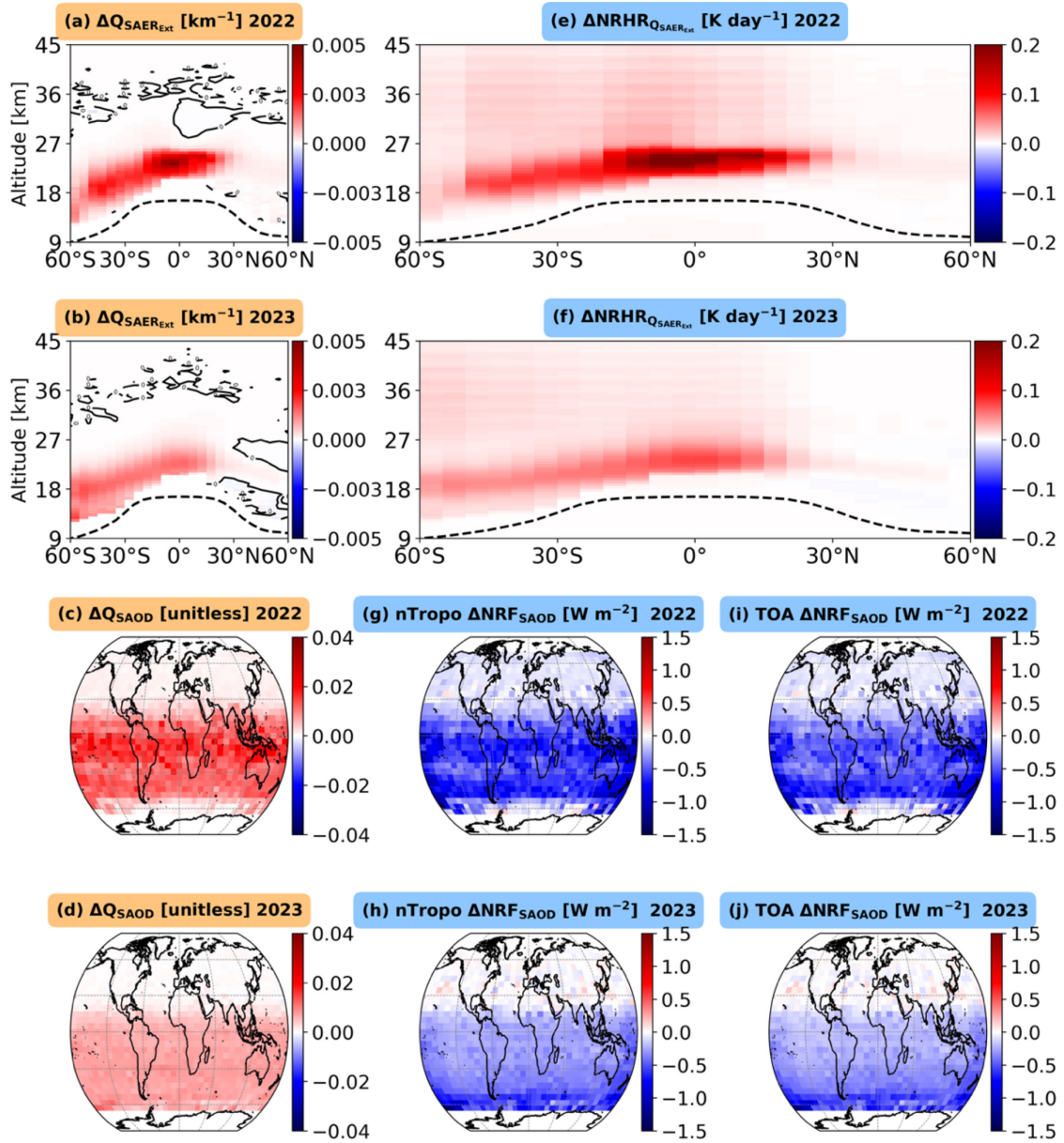
318   *which is extracted above 2σ; see Methods). The perturbations in SH₂O, SAOD, and ozone (2D-*
319   *filtered) following the eruption exceed their interannual variability, resulting in the combined*
320   *instantaneous radiative forcing shown in panel (**l**).*

321

322   **Changes in stratospheric aerosols after the Hunga eruption**

323

324   The Hunga eruption injected not only a substantial amount of water vapor but also a
325   moderate quantity of sulfur dioxide (SO_2) into the stratosphere, estimated at approximately 0.4–
326   0.7 Tg[14,15,60]. The total SO_2 transported from the troposphere to the stratosphere is estimated to
327   be around 1 Tg[60]. However, compared to previous volcanic eruptions such as Raikoke in 2019
328   (1.5 Tg[61,62]) and the Pinatubo eruption in 1991 (~17 Tg[63]), the amount of SO_2 injected into the
329   stratosphere by the Hunga eruption was considerably smaller.

330

331   To quantify the impact of the Hunga eruption on the stratosphere, we identified perturbed
332   aerosol extinction in 2022 and 2023 by detecting deviations exceeding above two standard
333   deviations (σ) from the background climatology. This method is based on the expectation that,
334   under normal conditions, stratospheric aerosol extinction remains relatively stable; therefore, any
335   significant deviation likely reflects the additional aerosol load injected by the eruption, which in
336   turn influences radiative forcing and stratospheric dynamics[62]. This approach allowed us to
337   assess the three-dimensional distribution of perturbed stratospheric aerosol extinction and the
338   associated SAOD. We observed an enhanced aerosol extinction of solar radiation over 50°S–
339   20°N in the stratosphere at 22–24 km altitude in 2022 following the eruption (Fig. 3a), with a
340   maximum relative increase in aerosol extinction reaching 800% (Supplementary Fig. 4a). The
341   near-global SAOD map, which represents integrated aerosol extinction from the lower to the
342   upper stratosphere (Fig. 3c), reveals a substantial absolute increase of approximately +0.014 ±
343   0.001 in the SH in 2022 (Fig. 3c; Fig. 2b; Table 1). Similar to SH₂O, stratospheric aerosols are
344   predominantly concentrated in the SH, exhibiting no strong longitudinal pattern (Fig. 3c). During
345   the second-year post-eruption (2023), we observe a decrease in 521 nm sulfate aerosol extinction
346   from 0.004 km⁻¹ to 0.002 km⁻¹ between 50°S and 20°N, at altitudes ranging from 10 km in the
347   southern midlatitudes to 25 km in the tropics (Fig. 3b).

348

349   Unlike water vapor, the spread of SAOD is predominantly confined to the SH (Fig. 3c,d).

350 In 2023, perturbed SAOD decreased by around 50% compared to the first-year aerosol

351 perturbation (Fig. 2b; Fig. 3c, g). With minimal aerosol perturbation in the NH, the cooling

352 impact of SAOD in the SH contributes to a hemispheric asymmetry under idealized conditions.

353 The reduction in SAOD observed in 2023 could be attributed to the gravitational settling of

354 stratospheric sulfate aerosols[64,65].

355

356

357   *Figure 3: Observed enhancement of stratospheric aerosol and the associated simulated*

358   *instantaneous radiative effects during 2022 and 2023 following the Hunga eruption.* Panels

(a) and (b) show the zonal- and annual-mean latitude–altitude variation in the change of the aerosol extinction coefficient at 521 nm ($Q_{SAERext}$; km^{-1}) for 2022 (15 Jan 2022 to 14 Jan 2023) and 2023 (15 Jan 2023 to 31 Dec 2023), respectively, relative to the reference climatology (CLIM) period from 07 June 2017 to 09 Dec 2021. Contours denote near-zero changes, and the dashed black line represents the tropopause height. Panels (c) and (d) display the vertically integrated aerosol extinction coefficient (SAOD) for the perturbed lower to upper stratosphere. Panels (e) and (f) show the latitude–altitude distribution of the difference in net (longwave + shortwave) radiative heating rates (ΔNRHR; K day^{-1}) of stratospheric aerosol between each post-eruption year and the CLIM period. Panels (g) and (h) depict the changes in near-tropopause net radiative flux (ΔNRF$_{SAOD}$; W m^{-2}) for 2022 and 2023, respectively, highlighting perturbations primarily observed in the lowest levels of the lower stratosphere (above the tropopause height). Panels (i) and (j) show the corresponding changes in the top-of-atmosphere (TOA) net radiative flux (ΔNRF$_{SAOD}$; W m^{-2}) for 2022 and 2023, respectively.

During 2022, following the eruption, we retrieved a mean effective radius of sulfate aerosols of approximately 0.3 μm (with a median of 0.27 μm) in the mid-stratosphere (Table 1 and Supplementary Fig. 7a,e), which is consistent with the findings of Knepp[66]. In 2023, the mean effective radius of sulfate aerosols in the SH was 0.29 μm (Table 1 and Supplementary Fig. 7i,m), nearly identical to that observed in 2022. Boichu[17] and Khaykin[67] showed that during the first two months after the eruption, sulfate aerosol particles were larger; over time, these particles decreased in size, transitioning to finer particles.

We compare the annual mean percentage change in near-global ΔSAOD for 2022, derived from SAGE III, with OMPS-NASA[56], OMPS-SASK[56], OSIRIS[56], and GloSSAC[56] to assess biases (Supplementary Fig. 2; Table 1). Among datasets, OMPS-NASA has the largest ΔSAOD discrepancy (-66.5%), followed by OMPS-SASK (-37.8%), OSIRIS (-12.6%), and SAGE-III/ISS (-15.5%), while GloSSAC shows a slight increase (+1.8%). Despite inter-instrument biases, GloSSAC estimates align well with our SAOD approach based on >2σ perturbations. These differences highlight the need for multi-instrument comparisons to refine SAOD estimates and radiative forcing assessments post-eruption. OMPS-NASA shows a >50%

389 high bias due to its fixed aerosol size assumption, while OMPS-SASK mitigates this by

390 accounting for size variations[54].

391

392 **Net radiative heating rate and radiative forcing of the Hunga-perturbed stratospheric**

393 **sulfate aerosol: 2022 versus 2023**

394   We present the near-global distribution of changes in net radiative flux near the

395 tropopause (Fig. 3g) and at the TOA (Fig. 3i) resulting from the perturbation of sulfate-

396 dominated stratospheric aerosols within the lower-to-middle stratosphere. During 2022, when

397 averaged over the SH, our results indicate that the mean net radiative flux at Earth's TOA

398 decreased by 0.53 ± 0.04 W m^{-2}, leading to a surface cooling of 0.07 ± 0.01 K, as estimated from

399 the FaIR[7,8] model (Supplementary Fig. 6b). This reduction is attributed to an increase in SAOD

400 (i.e., ΔSAOD) of approximately 0.014 ± 0.001 (Fig. 2b; Fig. 3c; Table 1). Similarly, the mean

401 net radiative flux near the tropopause in the SH decreased by 0.66 ± 0.05 W m^{-2} (Fig. 2b; Fig.

402 3g; Table 1), resulting in a surface cooling of 0.09 ± 0.01 K (Supplementary Fig. 6f). On a near-

403 global scale, the mean radiative flux changes at the TOA and near the tropopause are estimated

404 to be -0.34 ± 0.03 W m^{-2} and -0.41 ± 0.04 W m^{-2}, respectively, associated with a ΔSAOD of

405 around 0.008 ± 0.001 (Fig. 2b,f, j; Fig. 5a; Table 1).

406

407   During 2023, the increase in SAOD associated with the Hunga eruption amounted

408 to 0.007 ± 0.001 in the SH, leading to a reduction in net radiative flux at the TOA and near the

409 tropopause by 0.26 ± 0.03 W m^{-2} and 0.31 ± 0.04 W m^{-2}, respectively. Consequently, surface

410 temperatures cooled by around 0.11 ± 0.01 K by the end of the year. The system's thermal

411 inertia helped sustain this cooling, similar to the radiative forcing changes observed in 2022.

412 On a near-global scale, an increase in SAOD of 0.005 ± 0.001 resulted in a reduction in net

413 radiative flux of 0.15 ± 0.03 W m^{-2} at the TOA and 0.19 ± 0.04 W m^{-2} near the tropopause (Fig.

414 2b,f, j; Fig. 5a,b; Table 1).

415

416   Using the above results at TOA, we estimate a radiative forcing efficiency (near-global

417 mean radiative forcing at TOA per unit SAOD) of approximately 25–40 Wm^{-2} during 2022 and

418 2023. This estimate is relatively higher than those reported by Marshall[68] and Schmidt[69] for other

419 subaerial volcanic eruptions.

420

421   We retrieved sulfate aerosol effective radii of approximately 0.3 μm (with a median

422   radius of 0.27 μm) using Mie theory (see Methods), which indicates that these aerosols scatter

423   incoming solar radiation most efficiently per unit mass[65]. This enhanced scattering efficiency

424   may have contributed to the higher radiative forcing efficiency observed for sulfate aerosols,

425   highlighting the need for further chemistry-based sensitivity analyses of radiative efficiency in

426   future studies.

427

428   Our estimated effective (median) radius of sulfate aerosols has a mean value of 0.3 μm

429   (0.27 μm). These particle sizes result in a high sulfate MEE (Table 1; Supplementary Fig. 7d, h,

430   l, p and Supplementary Fig. 8), thereby enhancing sulfate aerosol forcing. The calculated

431   MEE and its associated standard deviation are 4.21 ± 0.13 m² g⁻¹ (4.24 ± 0.13 m² g⁻¹) in

432   the SH, 4.22 ± 0.17 m² g⁻¹ (4.17 ± 0.17 m² g⁻¹) in the NH, and 4.21 ± 0.25 m² g⁻¹ (4.21 ± 0.25

433   m² g⁻¹) for the near-global domain during 2022 (2023), respectively (Table 1).
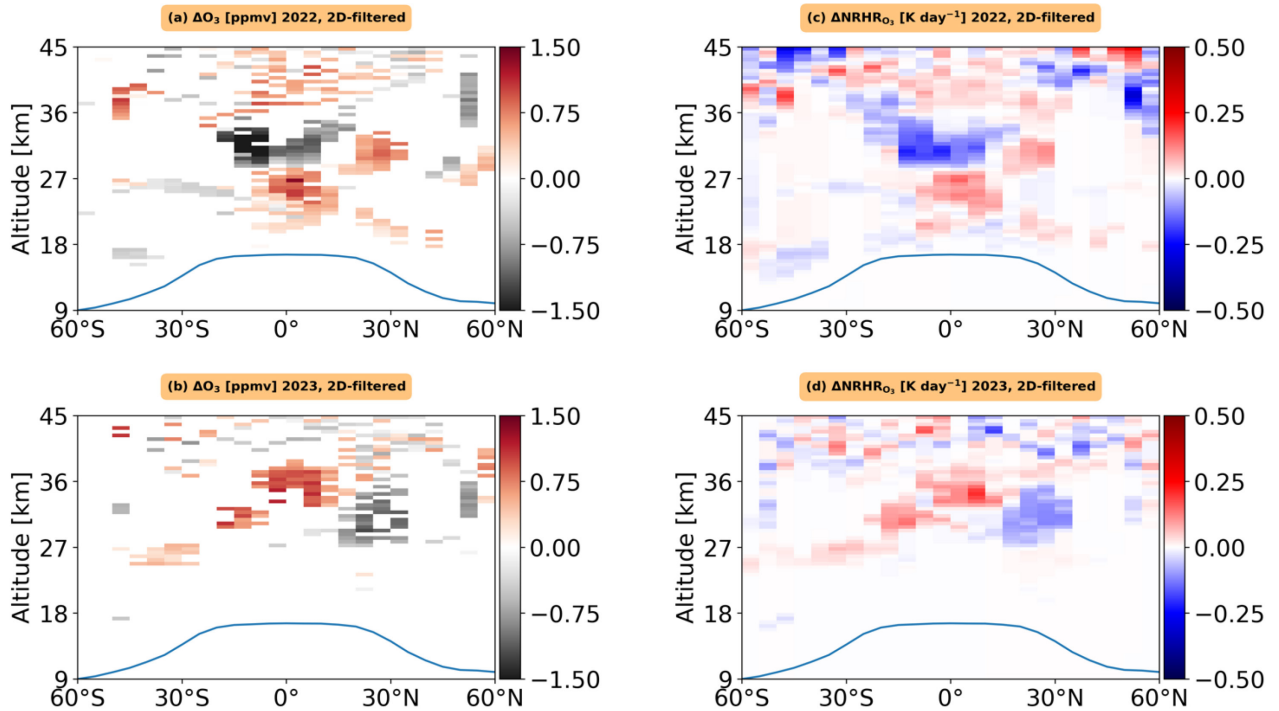

434

**Figure 4: Observed 2D-filtered changes in stratospheric ozone mixing ratios and the
associated simulated instantaneous radiative forcing during 2022 and 2023 following the
Hunga eruption.** *Panels (a) and (b) show the zonal- and annual-mean latitude–altitude
variation of the changes in stratospheric ozone mixing ratios for 2022 and 2023, respectively.*

439 *Panels (c) and (d) present the latitude–altitude distribution of the differences in net (longwave +*
440 *shortwave) radiative heating rates (ΔNRHR) for stratospheric ozone for 2022 and 2023,*
441 *respectively.*

442
443 **Changes in stratospheric ozone after the Hunga eruption: 2D-Filtered technique**
444
445   Following the methodologies of Wilmouth[35] and Santee[37], we assessed changes in
446 the ozone mixing ratio in 2022 and 2023 relative to background mean values using both SAGE
447 III[47] and MLS[58] datasets. To better quantify the impact of the Hunga eruption, we applied a 2D-
448 filtered (2DF) technique (see Methods) to distinguish changes in the zonal-mean mixing ratio.
449 Using the 2DF technique, the extent of the Hunga eruption's influence on ozone fluctuations be-
450 comes discernible, as the observed changes following the eruption exceed the typical year-to-
451 year variability seen before the event (Fig. 2c).

452
453   We find a slightly stronger negative ozone mixing ratio anomaly, with values of around -
454 0.5 ppmv at 27 km and -1 ppmv at 32 km in the SH during the summer of 2022 (Fig. 4a; Fig. 2c;
455 Supplementary Fig. 9c). In 2023, we also observe an increase in the ozone mixing ratio anomaly,
456 with values exceeding 0.5 ppmv between 22 km and 36 km altitude near equatorial regions and
457 in the SH (Fig. 4b; Fig. 2c; Supplementary Fig. 9i, j).

458
459   Wilmouth[35], Wang[36], and Santee[37] conducted a thorough analysis of various trace gas
460 species and also used global climate model simulations to determine the mechanism behind the
461 potential ozone loss in 2022. The results of their study demonstrate that the eruption impacted
462 ozone by influencing both stratospheric dynamics and ozone chemistry. We note that to exactly
463 pinpoint the physical mechanism related to the effects of the Hunga eruption on ozone gain in
464 2023, detailed modeling studies coupling stratospheric dynamics and chemistry are needed
465 which is out of the scope of this study. Therefore, our primary results are focused on assessing
466 the influence of radiative forcing of water vapor, SAOD and ozone (2D-Filtered technique).
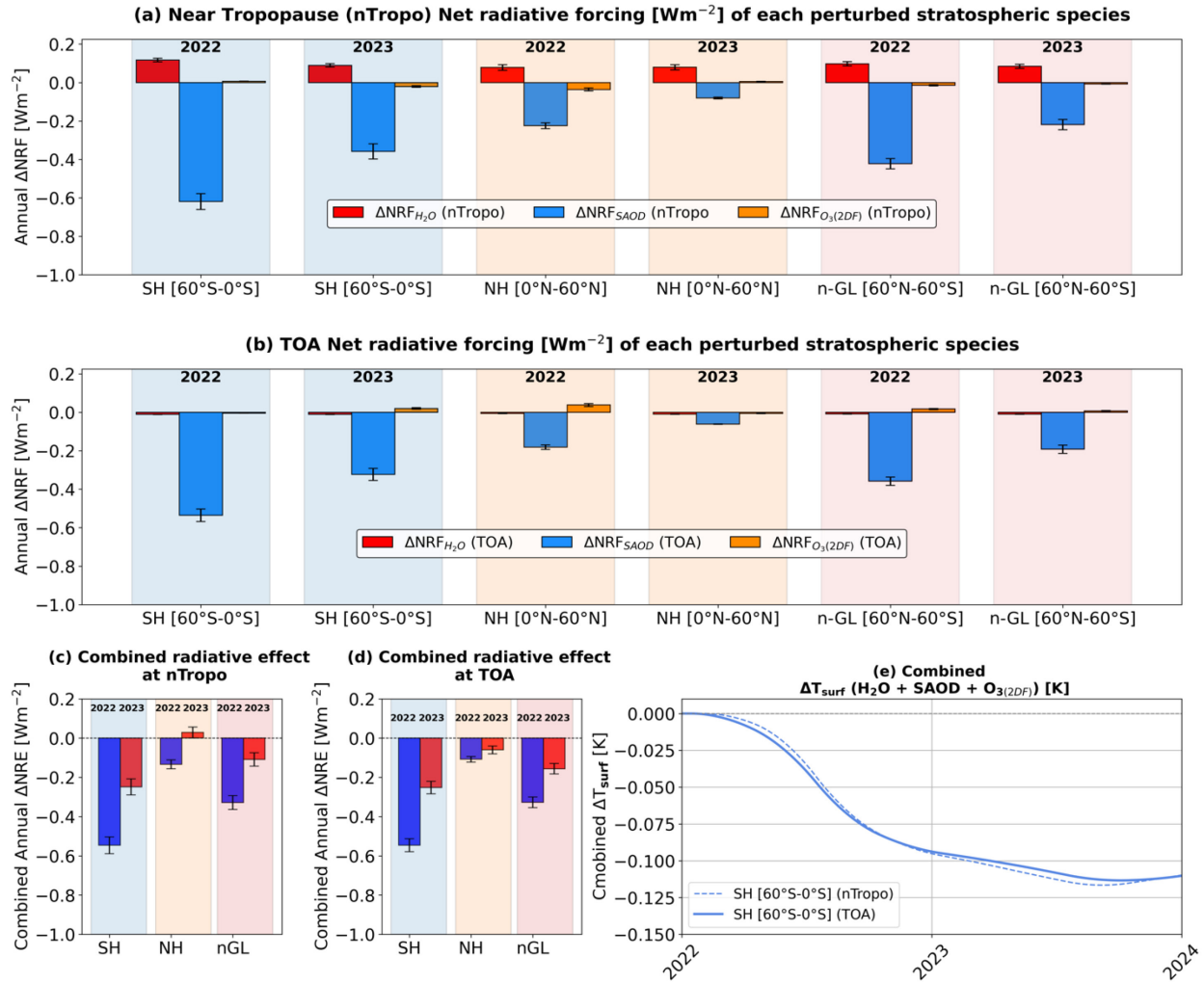467

Figure 5: Net radiative forcing during 2022 and 2023 due to three radiatively important stratospheric species perturbed by the Hunga eruption. *Panel (a) shows the clear-sky net (longwave + shortwave) instantaneous radiative forcing near the tropopause due to perturbations in stratospheric H_2O (red bars), aerosols (blue bars), and O_3 (orange bars) for the Southern Hemisphere (SH), Northern Hemisphere (NH), and near-global (60°S–60°N) regions. Panel (b) presents the clear-sky net instantaneous radiative forcing at the top-of-the-atmosphere (TOA) for the same species and regions. Panels (c) and (d) show the combined ΔNRF near the tropopause and at the TOA, respectively, averaged for the SH, NH, and near-global regions during 2022 and 2023. Panel (e) shows the surface temperature changes in the SH, estimated using a two-layer energy balance FaIR model based on the perturbed instantaneous energy balance at the TOA and near the tropopause. These estimates represent only a first-order approximation of the actual surface*

482

483 Net radiative heating rate and radiative forcing of the Hunga-perturbed stratospheric

484 ozone (2D-Filtered): 2022 versus 2023

485   Because ozone molecules absorb a substantial amount of solar radiation, a decrease in

486 their concentration would lead to reduced warming within the stratosphere. This, in turn, would

487 decrease the upwelling solar radiation from the uppermost layer of the stratosphere, resulting in

488 a relative reduction in the net radiative flux at Earth's TOA[33]. The opposite effect occurs in the

489 case of ozone gain.

490

491   We find that the ozone mixing ratio, derived from the 2D-filtered technique, decreased in

492 the stratosphere over 20°S–40°S between 24–27 km altitude and over 10°S–30°S between 30–35

493 km altitude during the first year after the eruption (Fig. 4a). This ozone loss contributed to strato-

494 spheric cooling of more than -0.15 K day^{-1} (Fig. 4c). Additionally, we observed an increase in

495 ozone in the equatorial regions (15°S–15°N) in the lower stratosphere below 27 km (Fig. 4a),

496 leading to stratospheric warming of more than +0.1 K day^{-1} (Fig. 4c). We observed a strong in-

497 crease in ozone mixing ratio from the middle stratosphere to lower stratosphere in 2023, possibly

498 influenced by dynamical processes associated with the easterly phase of the Quasi-Biennial Os-

499 cillation and secondary stratospheric circulation[15,32-36], contributing to a stratospheric heating

500 rate of approximately +0.15 to +0.2 K day^{-1}.

501

502   During 2022, we found that the reduction in O₃ concentration was primarily located in

503 the SH (Table 1). This decline in stratospheric ozone concentration by -0.06 DU in the SH

504 contributed to a small negative net radiative flux of -0.009 W m^{-2} at the TOA and a positive net

505 radiative flux of +0.018 W m^{-2} near the tropopause (Table 1). During 2023, an increase in ozone

506 concentration by 1.38 DU in the NH resulted in a net positive radiative flux change of +0.02 ±

507 0.01 W m^{-2} at the TOA and -0.05 ± 0.01 W m^{-2} near the tropopause (Table 1).

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Hunga eruption-perturbed species	SH		NH		nGL	
	2022	2023	2022	2023	2022	2023
$r_{eff_{HTHH}}[\mu m]$	0.30	0.29	0.26	0.27	0.28	0.28
$MEE_{HTHH}[m^2 g^{-1}]$	4.21	4.24	4.22	4.17	4.21	4.21
$M_{SO_{2_{HTHH}}}[Tg]$	0.46	0.31	0.26	0.19	0.71	0.50
$SAOD_{HTHH}[Unitless]$	0.017	0.013	0.010	0.007	0.013	0.010
$\Delta M_{SO_{2_{HTHH}}}[Tg]$	0.35	0.18	0.13	0.06	0.48	0.24
$\Delta SAOD_{HTHH}$ $[Unitless]$	0.014	0.007	0.005	0.002	0.008	0.005
ΔNRF_{SAOD} $[Wm^{-2}]$ (near tropopause)	**-0.66**	**-0.31**	**-0.18**	**-0.06**	**-0.41**	**-0.19**
ΔNRF_{SAOD} $[Wm^{-2}]$ (TOA)	-0.53	-0.26	-0.14	-0.05	-0.34	-0.15
ΔH_2O_{Total} $[Tg]$	86	72	59	66	145	138
ΔH_2O $[gm^{-2}]$	0.37	0.32	0.24	0.30	0.32	0.31
$\Delta NRF_{H_2O}[Wm^{-2}]$ (near tropopause)	**0.12**	**0.09**	**0.08**	**0.08**	**0.10**	**0.08**
$\Delta NRF_{H_2O}[Wm^{-2}]$ (TOA)	-0.010	-0.01	-0.005	-0.009	-0.008	-0.009
ΔO_3 $[DU]$ (2DF with H_2O)	-0.06	0.45	1.38	-0.36	0.66	0.05
$\Delta NRF_{O_3}[Wm^{-2}]$ (2DF; near tropopause)	**0.018**	**-0.028**	**-0.05**	**0.02**	**-0.018**	**-0.004**
$\Delta NRF_{O_3}[Wm^{-2}]$ (2DF; TOA)	-0.009	0.01	0.02	-0.006	0.007	0.002
Near tropopause $\Delta NRF_{H_2O+SAOD+O_{3,2DF}}$	-0.52	-0.25	-0.15	+0.04	-0.33	-0.11
At TOA $\Delta NRF_{H_2O+SAOD+O_{3,2DF}}$	-0.55	-0.26	-0.13	-0.07	-0.34	-0.16

Table 1: Properties and hemispherical mean net radiative forcing of radiatively important stratospheric species perturbed by the 2022 Hunga eruption. The blue rows list the mean

552 *effective radius (r_{eff}; in μm) of retrieved sulfate aerosols, the retrieved vertically integrated SO_2*

553 *mass (MSO_2; in Tg), the retrieved mass extinction efficiency (MEE; in $m^2 g^{-1}$) per gram of SO_2,*

554 *and the SAGE-III/ISS observed SAOD for Hunga for the Southern Hemisphere (SH), Northern*

555 *Hemisphere (NH), and near-global (n-GL) domains for 2022 and 2023. The light orange rows*

556 *present the retrieved absolute (Hunga–CLIM) change in sulfate aerosols (ΔMSO_{2Hunga}; Tg), the*

557 *SAGE-III/ISS observed absolute (Hunga–CLIM) changes in the vertically integrated SAOD*

558 *($\Delta SAOD$), and the resulting simulated mean net radiative forcing (ΔNRF_{SAOD}; $W m^{-2}$) near the*

559 *tropopause and at the top of the atmosphere (TOA). The dark orange rows show the SAGE-*

560 *III/ISS observed total injected H_2O mass (ΔH_2O; Tg) and its corresponding simulated mean net*

561 *radiative forcing (ΔNRF_{H_2O}; $W m^{-2}$) near the tropopause and at the TOA. The light magenta*

562 *rows display the absolute changes in ozone derived from the 2D-filtered technique (ΔO_3; DU;*

563 *2DF with water vapor) and their simulated mean net radiative forcing (ΔNRF_{O_3}; $W m^{-2}$; 2DF)*

564 *near the tropopause and at the TOA. The final grey rows show the combined net radiative forc-*

565 *ing of the three species ($NRF_{H_2O+SAOD+O_{3,2DF}}$; $W m^{-2}$) near the tropopause and at the TOA.*

566

567

568 **The combined net radiative effects of the three perturbed stratospheric species**

569   After determining the domain-averaged net instantaneous radiative effects of

570 enhanced stratospheric water vapor, stratospheric sulfate aerosol, and ozone using the 2D-filtered

571 technique, we combined these three radiative perturbations to estimate the net TOA radiative

572 forcing and the associated surface temperature changes due to the eruption (Fig. 5; Table 1).

573 Using a simple climate emulator (FaIR)[7-8], we find that the Hunga eruption induced a mean

574 surface cooling effect of approximately -0.10 ± 0.01 K by the end of 2022 (Fig. 5e). This cooling

575 resulted from a combined instantaneous radiative forcing of ~-0.55 ± 0.05 W m⁻² at the TOA,

576 concentrated almost entirely in the SH (Fig. 4c; Table 1). This cooling effect was primarily

577 driven by sulfate aerosols. During 2023, in the SH, the combined effects of SH_2O,

578 SAOD, and ozone (O_3-2DF) resulted in a TOA radiative forcing of -0.26 ± 0.04 W m⁻², leading

579 to a cooling of -0.10 ± 0.01 K by the end of 2023 (Fig. 5d). Thus, the Hunga-associated

580 cooling (and any potential future warming) is likely very small compared to natural climate

581 variability and will be difficult to detect observationally.

582

583   The stronger surface cooling observed in both 2022 and 2023, despite variations in the
584 magnitude of the combined net radiative flux at the TOA and near the tropopause altitude,
585 suggests a cumulative effect of radiative forcing over time, influenced by thermal inertia and
586 feedback processes. This implies that, although radiative forcing in 2023 was relatively small,
587 the temperature change in 2023 also reflects the residual effects of the stronger forcing from
588 2022, leading to a similar overall temperature response.

589

590   Also, given that the Pinatubo eruption injected approximately 20 Tg of sulfur, resulting in
591 a maximum cooling of about –0.5 K, a simple linear scaling suggests that the HTHH event,
592 which injected less than 1 Tg of sulfur, would induce a cooling effect of approximately –0.025
593 K[15,22]. Such a small temperature change is likely to be masked by natural variability and may not
594 be directly detectable. Given the modest magnitude of the expected cooling, further investigation
595 using comprehensive climate model simulations is necessary to determine whether these
596 gradients could trigger any measurable atmospheric responses.

597

598   In conclusion, our findings indicate that the Hunga eruption did not cause warming in
599 the SH or globally between 2022 and 2023. Instead, it had a cooling effect in the SH.
600 The efficient conversion of SO_2 into sulfate aerosols in a water-rich stratosphere in the SH likely
601 contributed to the net cooling effect observed at Earth's TOA in 2022 and 2023.

602

603   However, in 2023, we find a slight increase in net radiative flux near the tropopause in
604 the NH, primarily due to increased SH_2O and reduced ozone, with minimal perturbation of
605 SAOD. Given that the stratospheric lifetime of SH_2O is much longer than that of SAOD (around
606 2.5 years[64-65]), this warming trend in the NH may persist and spread over the years until
607 the perturbed SH_2O is fully depleted. Consequently, depending on ozone perturbations,
608 this longer-lasting SH_2O effect could eventually outweigh the cooling effect of SAOD in the SH.

609

610   Our findings show that the climate impact of shallow submarine eruptions (water-rich
611 eruptions with a moderate amount of SO_2) differs from that of subaerial
612 eruptions (e.g., Pinatubo[22]), depending on the injected altitude and composition of the volcanic
613 plume.

Methods

This study aimed to assess the radiative effects of the Hunga eruption. To achieve this, we analyzed observed changes in the three main radiatively important species in the stratosphere that were substantially perturbed by the eruption: water vapor, aerosols, and ozone (see Supplementary Information). We then used the observed changes in these three stratospheric species before and after the eruption to simulate the radiative effects of the Hunga eruption using the LibRadtran radiative transfer model (Supplementary Figure 1).

Descriptions of data sets

SAGE-III/ISS observations

We analyzed SAGE III/ISS data from 7 June 2017 to 31 December 2023, covering latitudes from 70°N to 70°S, depending on variations in solar insolation, solar zenith angle, and the coverage range of the ISS instrument. Specifically, we used solar occultation data from SAGE III on the ISS (SAGE III/ISS; version 5.3, accessed on 28 October 2023; Algorithm Theoretical Basis Document (ATBD)[70-71]), which provided approximately 30 atmospheric profiles per day of H_2O, O_3, and aerosol extinction coefficient at sunrise and sunset across the globe.

Compared to estimates of stratospheric temperature and composition from passive limb sounding instruments (e.g., Aura Microwave Limb Sounder; MLS), the solar occultation technique of SAGE III/ISS provides much higher vertical resolution measurements. While MLS offers daily near-global observations (82°S–82°N; Supplementary Figures 11-12), its lower vertical resolution makes it challenging to fully resolve vertical variations in SH_2O and ozone, as well as the associated radiative forcings[49-50]. From the perspective of higher vertical resolution, accurate estimation of atmospheric heating rates and top-of-atmosphere radiative forcing requires detailed profiling of radiatively active species, such as water vapor, ozone, and aerosols[49]. Therefore, in this study, we primarily use multi-wavelength SAGE III/ISS data as input for most of our radiative forcing analysis. However, we also compare our results with other satellite products, including MLS for H_2O and O_3 (Supplementary Figs. 3, 11–14) and SAOD measurements from multiple instruments, such as OMPS-SASK, OMPS-NASA, and

644 GloSSAC, to ensure consistency and to assess and highlight uncertainties across different
645 datasets (Supplementary Fig. 2; Table 1).

646

647   We analyzed approximately 70,000 profiles over our study period, conducting the
648 analysis on a seasonal basis to ensure sufficient global observations while accounting
649 for temporal changes in solar angle and other variables when calculating the radiative effects of
650 the eruption (Supplementary Methods). We computed the observed changes in SH_2O, multi-
651 wavelength aerosol extinction coefficient, and ozone as a function of altitude (binned at 0.5 km),
652 latitude (binned at 5° to ensure robust sampling statistics per bin), and longitude (binned at
653 10°). Additionally, we estimated the seasonal variations in a near-global map, depicting
654 the vertical integral of water vapor, aerosol extinction coefficient, and ozone concentration from
655 the lower to upper stratosphere.

656

657 We divided the background period before the Hunga eruption (CLIM period) into
658 four seasons: DJF (December–January–February), MAM (March–April–May), JJA (June–July–
659 August), and SON (September–October–November), using data from 7 June 2017 to 9
660 December 2021 (data was unavailable between 10–18 December 2021). We included all
661 available data from December 2021 before the initial phase of the Hunga eruption began on 19
662 December 2021. Similarly, we divided the Hunga eruption period into: Hunga-2022: JF (15
663 January – 28 February 2022), MAM (1 March – 31 May 2022), JJA (1 June – 31 August
664 2022), and SON (1 September – 30 November 2022). Hunga-2023: DJF (December 2022 – 28
665 February 2023), MAM (1 March – 31 May 2023), JJA (1 June – 31 August 2023), and SON (1
666 September – 30 November 2023).

667

668   For our radiative calculations, we applied the nearest-neighbor interpolation technique (K-
669 dimensional tree interpolation) when necessary to achieve complete latitude-longitude
670 coverage between 70°S and 70°N. Finally, we computed the weighted mean, accounting
671 for differences in the number of days across the four seasonal analysis periods during both
672 the Hunga and CLIM periods. This approach enabled us to assess annual-mean changes in
673 the radiation budget, thereby quantifying the annual radiative effect of Hunga-2022 and Hunga-
674 2023.

Ozone perturbations analysis following the Hunga Eruption: 2D-filtered anomaly

To assess ozone perturbations associated with the Hunga eruption during 2022 and 2023, we employed a structured methodology based on the approach adopted by Wilmouth[35]:

(a) **Evaluation of ozone mixing ratio anomalies exceeding 2σ**: We analyzed the zonal-mean and seasonal-mean ozone (O_3) mixing ratio in 2022 and 2023 to identify deviations exceeding ±2 standard deviations from the 2005–2021 MLS and 2017–2021 SAGE III seasonal means.

(b) **Detection of ozone anomalies exceeding 2σ, synchronized with >2σ water vapor anomalies**: The Hunga eruption strongly perturbed water vapor, meaning that any ozone changes caused by the eruption should co-occur with water vapor anomalies (Wilmouth[35]). Therefore, regions where water vapor was perturbed served as a reference to determine when and where the eruption also affected ozone.

Following Wilmouth[35], we identified concurrent anomalies in ozone and water vapor. Specifically, we assessed ozone changes as a function of season, latitude bin (5°), and pressure level for MLS (and 0.5 km altitude bin for SAGE III), ensuring that they coincided with water vapor anomalies exceeding >2σ above the MLS 16-year mean and SAGE III 5-year mean.

(c) **Removal of pre-existing anomaly values from 2022 and 2023**: To quantify ozone anomalies caused by the Hunga eruption, we examined whether the 2022–2023 anomalies had occurred in previous years and whether they fell within background variability, independent of the eruption. Using abnormally high water vapor levels in 2022 and 2023 as a reference, we identified and excluded zonal-mean ozone anomaly values (classified by season, latitude, and altitude) that predated these exceptional conditions.

(d) **Application of the 2D-Filtered (2DF) technique**: The 2D-Filtered (2DF) technique was applied to ozone anomalies as a function of season, latitude, and altitude (or pressure level) if they met the criteria from Steps (a) to (c) (Fig. 2c; Fig. 4a,b; Supplementary Figure 10, 13, 14). When the 2DF technique was applied without water vapor constraints, it was referred to as 2D-Filtered (2DF) without water vapor (see Supplementary Fig. 10). Ozone anomalies that did not meet these criteria were classified as unfiltered conditions (Supplementary Fig. 3e-h). Note that one

702 limitation of this filtering approach is that it may underestimate ozone anomalies influenced by
703 secondary circulation effects that extend beyond the core of the water vapor anomalies.

Mass extinction efficiency of sulfate aerosol

705   Mass extinction efficiency (MEE; $m^2 g^{-1}$) of sulfate aerosols indicates how efficiently of
706 aerosols scatter or absorb of solar radiation per unit mass. It can be expressed as:

707

$$MEE[m, 521nm, D_s] = \frac{\int_{D_{smin}}^{D_{smax}} \frac{3}{2\rho_s D_s} Q_{ext}(D_s, n, k, 521nm) n_M(D_s) dD_s}{\int_{D_{smin}}^{D_{smax}} n_M(D_s) dD_s} \tag{1}$$

708

709 The mass size distribution, $n_M(D_s)$, is given by:

$$n_M(D_s) = \frac{\pi}{6} \rho_s D_s^3 n(D_s)$$

711 where $n(D_s)$ represents the log-normal number size distribution at geometric standard deviation
712 of 1.2. ρ_s is the SO_2 density. Based on the retrieved effective radius (r_{eff}) of sulfate aerosol (D_s
713 = $2r_{eff}$) and the complex refractive index of sulfate aerosol at 521 nm (real part: n = 1.431;
714 imaginary part: k = 1.0 x 10^{-8}), we determined the extinction efficiency factor (Q_{ext} at 521 nm)
715 using Mie theory (Supplementary Figure 7, 8,11). Details of the expression of
716 $Q_{ext}(r_{eff}, n, k, 521nm)$ are provided in the Supplementary Methods. The refractive index of
717 sulfate aerosol is obtained from GEISA database.

718

719 The hemispherical mean value of MEE is also estimated using:

720

$$MEE = \frac{SAOD * \text{Surface area}[m^2]}{\text{mass of } SO_2[g] * \left(\frac{\frac{\text{Molar Mass of } H_2SO_4}{\text{Molar mass of } SO_2} * 1}{wt} \right)} \tag{2}$$

721

722 Here, wt is the 75% H_2SO_4 weight.

723

Relative and absolute changes of three stratospheric species: Hunga versus CLIM

725   We determined the relative and absolute changes in the three stratospheric species (Y) for

726   the Hunga and CLIM periods using:

727

$$Absolute\ change = \Delta Y = Y[HTHH] - Y[CLIM] \tag{3}$$

728

$$\%\ Relative\ change = 100 \times \frac{\Delta Y}{Y[CLIM]} \tag{4}$$

729

730   **Retrieval of optical and microphysical properties of stratospheric sulfate aerosol: Hunga**

731   **versus CLIM**

732   We utilized near-global coverage of the spatial and temporal distribution of aerosol

733   extinction [km^{-1}] (SAGE-III; ATBD, 2002) at nine wavelengths (384.224 nm, 448.511 nm,

734   520.513 nm, 601.583 nm, 676.037 nm, 755.979 nm, 869.178 nm, 1021.20 nm, and 1543.92 nm).

735   It is important to note that uncertainty in aerosol extinction varies with wavelength[70]. For

736   instance, at 407 nm and 1089 nm, the uncertainty can exceed 25%, whereas at 450 nm, it

737   decreases to 10%, and at 521 nm, it is less than 5%. To maximize accuracy, we jointly

738   utilized all nine wavelengths to retrieve stratospheric sulfate aerosol optical and microphysical

739   properties using a pre-calculated look-up table (LUT) based on Mie theory[66,71-74].

740

741   **Mie-LUT based retrieval of stratospheric sulfate aerosol**

742   An accurate assessment of the size distribution of stratospheric aerosols is crucial

743   for accurately quantifying their radiative effects before and after the eruption. When

744   analyzing background stratospheric aerosols, studies have shown that their size distribution can

745   typically be characterized by a unimodal lognormal distribution[75]. These aerosols

746   are predominantly spherical and consist of approximately 75% sulfuric acid and 25% water

747   vapor[68-73].

748   Ground-based stratospheric aerosol observations from AERONET indicate that during

749   the early stages of aerosol growth following the Hunga eruption, bi-modal and tri-modal log-

750   normal size distributions better represented the aerosol characteristics, particularly the coarse

751   mode of aerosol particles[16,17]. However, within a few months after the eruption, the data showed

752   that fine-mode stratospheric aerosols became dominant in aerosol extinction[16,67]. Thus, we

753   adopted a unimodal log-normal size distribution to represent the size distribution of sulfate

754   aerosols resulting from the Hunga eruption, with a geometric standard deviation of 1.2. This is

755   broadly consistent with the more comprehensive size distribution analysis by Knepp[66],

756   Khaykin[66] and Duchamp[16]. This size distribution parameterization has previously been used to

757   assess stratospheric sulfate aerosol growth one year after the Pinatubo eruption[22].

758   Notably, Knepp[66] introduced a novel method based on SAGE III data, indicating a spectral width

759   of 1.2 to 1.3 post-Hunga eruption, which aligns with our choice of 1.2 for the geometric standard

760   deviation in the Hunga case.

761   We used the Mie LUT to retrieve the optical and microphysical properties of sulfate

762   aerosols from SAGE III/ISS stratospheric aerosol extinction profiles at nine wavelengths,

763   covering the period from 7 June 2017 to 31 December 2023 (Hunga-2022 vs. CLIM and Hunga-

764   2023 vs. CLIM periods). For this purpose, we followed these steps:

765

766   1. We calculate the aerosol extinction ratios at 521 nm as: 384.224 /520.513 nm,

767   448.511/520.513 nm, 520.513/520.513 nm, 601.583/520.513nm, 676.037/520.513

768   nm, 755.979/520.513 nm, 869.178/520.513 nm, 1021.20/520.513 nm, and

769   1543.92/520.513 nm) using SAGE-III/ISS observations.

770

771   2. We compare the SAGE-III aerosol extinction ratios (from point 1) with the Mie LUT

772   aerosol extinction ratios and determine the closest matching extinction ratios at each

773   given wavelength.

774   3. We calculate the squared difference between the eight aerosol extinction ratios from

775   SAGE-III/ISS and the Mie LUT for each effective radius, then sum these differences

776   to quantify the discrepancy between the two datasets.

777   Here, Eq. (5) summarizes above step 3:

$$p(r_{eff}) = \sum_{\lambda_i} \left[Q_{ratio_{SAGE}}(\lambda_i) - Q_{ratio_{LUT}}(\lambda_i, r_{eff}) \right]^2 \tag{5}$$

778   where $Q_{ratio,SAGE}$ and $Q_{ratio_{LUT}}$ represent the extinction ratio at 521 nm from SAGE-III

779   observations and LUT, respectively. Here λ_i sums over the eight wavelengths from 400

780   to 1450 nm (as listed above), determined for both SAGE-III and LUT table using

781   interpolation techniques.

782

783   4. The minimum sum of squared differences represents the best solution for the effective

784   radius, asymmetry factor, and single scattering albedo. This step identifies the closest

785   match between the observed and modeled data.

786

787   Simulations of radiative effects of three stratospheric species (Hunga versus CLIM periods)

788 | Columnar- & Zonal-mean

789   We utilized observational data from multiple instruments (SAGE III/ISS and Aura MLS)

790   to obtain atmospheric and aerosol extinction profiles. Specifically, we retrieved vertical

791   profiles of temperature, pressure, density, water vapor (H_2O), ozone (O_3), and multi-wavelength

792   aerosol extinction from SAGE III/ISS observations between 9 km and 100 km altitude.

793   For atmospheric profiles below 9 km, we incorporated data from other sources,

794   including MLS[58] and ERA5[78].

795

796   Using the input data described above, we conducted idealized radiative transfer

797   simulations to examine the 3D distribution of stratospheric radiative heating

798   rates and instantaneous radiative fluxes between two periods: Hunga-2022 and Hunga-

799   2023 versus the CLIM period, at the TOA and near the tropopause. A summary of these data

800   sources and radiative transfer model simulations is provided in Supplementary Figure 1. Using

801   the LibRadtran model with the DISORT solver[59,79-80] and eight streams, we simulated radiative

802   flux changes and associated heating rates by incorporating observed seasonal-mean vertical

803   profiles of atmospheric constituents, aerosol properties, and seasonal-mean daily-insolation

804   weighted solar zenith angle as a function of latitude and longitude. These simulations were

805   conducted separately for the shortwave (0.28–4 μm) and longwave (4.0–100 μm) components

806   of radiative flux changes and heating rates (see Supplementary Methods). We used the CLIM

807   period to define the atmospheric temperature profiles (per season) and applied these fixed

808   temperature profiles for the Hunga period analysis. Thus, our radiative calculations do not

809   account for any temperature adjustments in the stratosphere in the year following the eruption

810   due to net heating/cooling from different radiative species. This analysis design allows us

811   to clearly isolate the instantaneous radiative response of the Hunga eruption from eruption-

812   associated climate feedbacks (e.g., dynamical changes) that influence the observed stratospheric

813   and tropospheric temperature profiles over the past year.

814

815   Uncertainty analysis related to three stratospheric species

816   We estimated seasonal and yearly standard deviation values for three stratospheric

817   species (H_2O, sulfate aerosol, and O_3) using SAGE III/ISS observations from 2017 to 2021,

818   before the Hunga eruption. Based on these yearly standard deviation values of the perturbed

819   stratospheric species, we calculated the corresponding standard deviation associated with

820   their instantaneous radiative effects at the TOA and near the tropopause (Fig. 2). Additionally,

821   we determined the propagated standard deviation for the combined radiative effects of the

822   three stratospheric species (Fig. 2). It is important to note that the standard deviation associated

823   with surface temperature changes, estimated using the FaIR model[7-8], is derived from

824   the propagated standard deviation of the combined instantaneous radiative effects of the

825   three stratospheric species[66]. We presented the standard deviation values rounded to two and

826   three places. Other sources of uncertainty, such as dynamical variations, cloud effects, and biases

827   across different satellite products, may also influence temperature changes. However, these

828   factors are beyond the scope of this study and are therefore not included in our uncertainty

829   estimate[66].

830   Limitations of Radiative transfer model

831   Even though we directly use the observed vertical profiles of aerosols, stratospheric water

832   vapor, and ozone as inputs to the LibRadtran model, our results are subject to several important

833   limitations.

834

835   First, the infrequent sampling of aerosol extinction profiles at multiple wavelengths by

836   SAGE III introduces uncertainties in SAOD retrievals, particularly in the lower stratosphere

837   following the Hunga eruption. Thus, cross-instrument comparisons are critical for improving

838   confidence in SAOD estimates and reducing uncertainties in radiative forcing assessments

839   following major volcanic events (Supplementary Figure 2 and Supplementary Table 1).

840

841 Second, the high signal-to-noise ratio SAGE-III/ISS dataset we used provides only dawn-
842 dusk solar stratospheric profiles. Consequently, we cannot capture the diurnal variations of
843 radiative fluxes in our calculations, which are limited to daytime periods (i.e., no diurnal
844 variability is considered). However, we anticipate that this limitation does not significantly affect
845 our results, as the impact of diurnal temperature and water vapor variations in the stratosphere is
846 relatively small[81-84].

847

848 Third, for computational simplicity, we assume a fixed, idealized clear-sky surface
849 albedo of 0.15 between 60°N and 60°S (Supplementary Figure 15). This assumption may
850 influence the net magnitude of the cooling effect, which could be even stronger if surface albedo
851 were higher. However, the overall sign of the radiative effect remains the same even if the clear-
852 sky surface albedo is assumed to be 0.2 over the same latitude range.

853

854 Fourth, our idealized radiation-only simulations do not account for the dynamical effects
855 of radiative components. Additionally, we do not incorporate measured temperature profiles
856 from the Hunga eruption period to estimate radiative forcing under a relaxed stratospheric state
857 (i.e., effective radiative forcing). Notably, Schoeberl[15] accounted for temperature anomalies and
858 found that while they are not always as pronounced as those associated with water vapor or
859 aerosols, they are non-negligible. To illustrate the potential impact of stratospheric temperature
860 anomalies (cooling), we provide idealized radiative forcing column calculations in
861 Supplementary Figures 15–19, showing that these anomalies primarily affect longwave radiation.
862 When accounting for stratospheric temperature changes, we find that the radiative forcing at both
863 the TOA and near the tropopause altitude due to stratospheric water vapor remains positive.
864 Furthermore, surface albedo should ideally be treated as a dynamic variable, varying with
865 surface type, land use, and seasonal changes. Nonetheless, despite these simplifications, our
866 primary conclusion—that sulfate aerosols dominate the stratospheric radiative cooling effect
867 among the three species considered—remains robust.

868

869 **The Finite-amplitude Impulse-Response (FaIR) model**
870 We use the Finite-Amplitude Impulse-Response (FaIR) model, an emulator of a two-
871 layer energy balance system, to estimate the first-order temperature response to the Hunga

872 eruption[7–8]. In this model, "finite-amplitude" refers to the response magnitude being determined
873 by the forcing strength, while "impulse-response" describes the system's reaction to a sudden
874 forcing change. FaIR is applied to assess the temperature response to Hunga-induced radiative
875 forcing from stratospheric water vapor, sulfate aerosols, and ozone. The model provides an
876 estimate of the resulting temperature change, particularly in the SH. However, as FaIR primarily
877 simulates global and hemispheric-scale surface temperature changes, it does not fully account for
878 regional variability or all Earth system feedbacks.

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882 **Data availability**
883 The SAGE-III/ISS v053 solar datasets (filename: g3bssp_53) are freely available from
884 (https://asdc.larc.nasa.gov/project/SAGE%20III-ISS/g3bssp_53). The Aura MLS dataset is also
885 freely available and can be obtained using https://search.earthdata.nasa.gov/. The ECMWF
886 ERA5 reanalysis pressure level datasets can be obtained from
887 https://cds.climate.copernicus.eu/cdsapp#!/search?type=dataset. Underlying data (Supplementary
888 Data 1-5) related to the manuscript is publicly available at Zenodo website
889 (https://zenodo.org/records/14955808 ; Gupta[85]). The data related to the refractive index of
890 sulfate aerosols is taken from https://geisa.aeris-data.fr/litms/. The Mie table for sulfate aerosol is
891 available here: https://github.com/matthew2e/easy-volcanic-
892 aerosol/blob/master/eva_Mie_lookuptables.nc[72].

893
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896 **Code availability**
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898 The LibRadtran model is available at http://www.libradtran.org/doku.php?id=download here. All
899 the Figures (including Supplementary Figures) were originally produced and plotted using
900 various open-source Python libraries (e.g., https://matplotlib.org/stable/). The LibRadtran model-
901 based processed data and input is also publicly available at the Zenodo
902 (https://zenodo.org/records/14955808; Gupta[85]).

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Acknowledgments

Author contributions

A.G. contributed to the conceptual development of this work, conducted the analysis, and produced

all figures and results. R.B. ran the FaIR model. A.G. drafted the original manuscript, which was

reviewed by T.M., J.F.K., R.B., and K.E.F.

Competing interests

The authors declare no competing interests.