

Temperature-jump solution X-ray scattering reveals distinct motions in a dynamic enzyme

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Correlated motions of proteins are critical to function, but these features are difficult to resolve using traditional structure determination techniques. Time-resolved X-ray methods hold promise for addressing this challenge, but have relied on the exploitation of exotic protein photoactivity, and are therefore not generalizable. Temperature jumps, through thermal excitation of the solvent, have been utilized to study protein dynamics using spectroscopic techniques, but their implementation in X-ray scattering experiments has been limited. Here, we perform temperature-jump small- and wide-angle X-ray scattering measurements on a dynamic enzyme, cyclophilin A, demonstrating that these experiments are able to capture functional intramolecular protein dynamics on the microsecond timescale. We show that cyclophilin A displays rich dynamics following a temperature jump, and use the resulting time-resolved signal to assess the kinetics of conformational changes. Two relaxation processes are resolved: a fast process is related to surface loop motions, and a slower process is related to motions in the core of the protein that are critical for catalytic turnover.

Protein motions are critical for functions such as enzyme catalysis and allosteric signal transduction¹, but it remains challenging to study excursions away from the most populated conformations². Traditional methods that utilize X-rays for structural characterization of biological macromolecules, such as crystallography and solution scattering, provide high-quality structural information, but this information is both spatially and temporally averaged because the measurements are performed on large ensembles of molecules and are typically slower than the timescales of molecular motion^{2,3}. To some extent, the spatial averaging inherent to X-ray experiments is advantageous, because it reveals the alternative local conformations of a molecule that are significantly populated at equilibrium; however, structural states that are not significantly populated at equilibrium, such as intermediates along a conformational transition pathway, are effectively invisible. The temporal averaging inherent to X-ray experiments also results in a loss of information about how transitions between local alternative conformational states are coupled to one another. To gain kinetic information about molecular motion, researchers often turn to spectroscopic methods, but it can be difficult to correlate spectroscopic observables with high-resolution structural models.

Time-resolved X-ray scattering and diffraction can overcome the limitations of traditional structure determination for studying the dynamics of biomolecules^{4–7}. In these experiments, a fast perturbation is applied to the sample to remove it from conformational equilibrium and synchronize conformational changes in a significant fraction of the molecules. Ultrafast X-ray pulses, which are short relative to motions of interest, are then used to perform structural measurements in real time as the system relaxes to a new equilibrium, providing simultaneous structural and kinetic information at high spatial and temporal resolution. Time-resolved X-ray experiments can identify transiently populated structural states along a conformational transition pathway, and reveal kinetic couplings

between conformations⁸. Despite this potential to provide a wealth of information, especially when combined with molecular dynamics simulation^{9–12}, time-resolved experiments have not been broadly applied by structural biologists. To date, systems that have been most rigorously studied are those in which a protein conformational change is coupled to excitation of a photoactive chromophore molecule, because the conformational change can be initiated with an ultrafast optical laser pulse (for example, refs. ^{13–18}). Unfortunately, the number of proteins that undergo specific photochemistry as part of their functional cycle is small, and there is a fundamental need to develop generalized methods that can be used to synchronously excite conformational transitions in any protein molecule and expand the utility of time-resolved structural experiments^{19–27}.

Protein structural dynamics are intimately coupled to the thermal fluctuations of the surrounding solvent ('solvent slaving'^{28,29}), and thermal excitation of the solvent by infrared laser temperature jump has been used in numerous pump–probe experiments. These experiments work on the principle that absorption of infrared photons excites the O–H stretching modes of water molecules, and the increased vibrational energy is dissipated through increased rotation and translation of the solvent molecules, effectively converting electromagnetic energy into kinetic (thermal) energy. Because this process of solvent heating and subsequent heat transfer to the protein is much faster than the large-scale molecular motions that define protein conformational changes, the sudden temperature jump removes conformational ensembles of protein molecules from their thermal equilibrium so that their structural dynamics can be measured using relaxation methods (Fig. 1a). For example, temperature-jump perturbations have been coupled to ultrafast spectroscopic methods, including Fourier-transform infrared spectroscopy^{30,31}, NMR^{32–34} and various forms of fluorescence spectroscopy^{35,36}, for the study of protein folding and enzyme dynamics. While these methods provide detailed kinetic information, they

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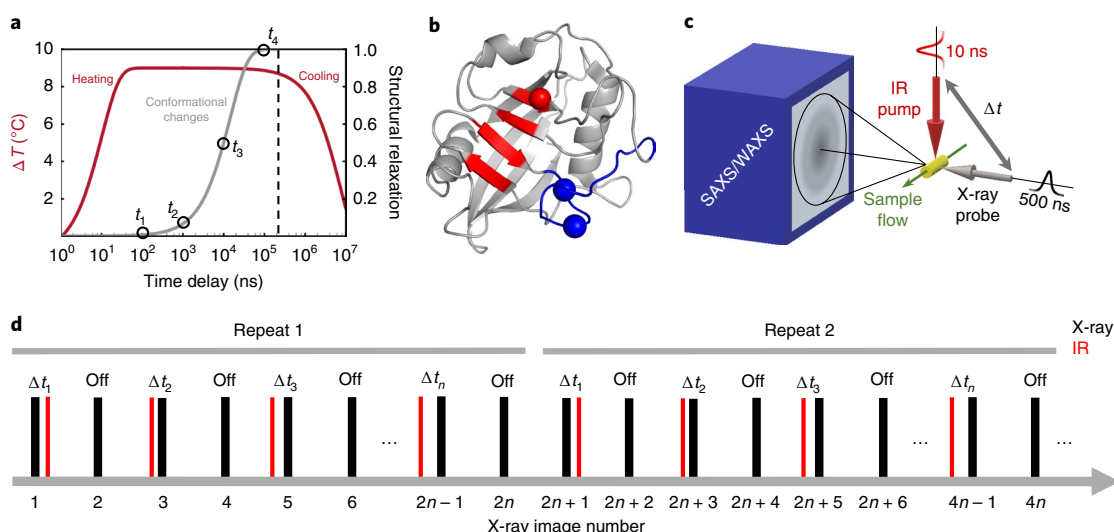


Fig. 1 | Overview of temperature-jump SAXS/WAXS experiments. **a**, During a temperature-jump experiment, an infrared (IR) laser pulse, several nanoseconds in duration, vibrationally excites the water O–H stretch and rapidly heats an aqueous solution of protein molecules (red curve). Heating is fast, but kinetic barriers to protein motions cause a lag in the structural relaxation to a new thermal equilibrium of conformational states (grey curve). **b**, Ribbon model depicting a single CypA molecule, with the ‘core’ dynamic residues (linked to catalysis) coloured red and the ‘loop’ region adjacent to the active site (determines substrate specificity) coloured blue. Sites of key mutations (S99T) are identified by spheres at the corresponding α positions (red: S99T; blue: D66N/R69H). **c**, Schematic of the temperature-jump SAXS/WAXS instrumentation, with key features highlighted. Liquid sample flows through the interaction region, where it interacts with mutually perpendicular IR pump and X-ray probe beams. Both pump and probe sources are pulsed, with a defined time delay between arrivals at the sample. SAXS/WAXS patterns are recorded on a single detector panel. **d**, Data collection sequence used for the temperature-jump experiments. For each pump–probe time delay, a pump–probe measurement (‘laser on’) was performed, followed by a measurement with no application of the pump laser (‘laser off’). On–off pairs with increasing pump–probe time delays were measured in succession until all of the desired delay times were acquired, then the sequence was repeated as many as 50 times to improve the signal-to-noise ratio of the data. Note that the first measurement within each repeat is a control measurement, wherein the probe pulse arrives at the sample before the pump pulse (negative time delay).

yield only very limited structural information about the underlying atomic ensemble. In contrast, the application of temperature jumps to time-resolved X-ray scattering and diffraction has been very limited. Nearly two decades ago, Hori et al.³⁷ used temperature-jump Laue crystallography to study the initial unfolding step of 3-isopropylmalate dehydrogenase. That study explored only a single pump–probe time delay, which allowed them to observe laser-induced structural changes but precluded kinetic analysis. Within the past two years, the laser temperature-jump method has been paired with X-ray solution scattering to explore the oligomerization of insulin under non-physiological conditions^{25,26} and haemoglobin²⁷. The results and analysis we present here expand the role of the temperature-jump method in structural biology, by showing that temperature-jump X-ray scattering experiments can be used as a general method to explore the functional, internal dynamics of proteins in solution. Additionally, we provide a detailed outline of a data reduction and analysis procedure suitable for temperature-jump small-angle X-ray scattering (SAXS)/wide-angle X-ray scattering (WAXS) experiments.

The temperature-jump SAXS/WAXS experiments we describe here used human cyclophilin A (CypA)—a proline isomerase enzyme that functions as a protein-folding chaperone and as a modulator of intracellular signalling pathways. CypA has been the subject of many NMR experiments that have identified two primary dynamic features of interest (Fig. 1b). First, the active site-adjacent loops (covering approximately residues 60–80 and hereafter referred to as the ‘loops’ region) are mobile on the millisecond timescale³⁸. This region is especially interesting because evolutionarily selected mutations within these loops perturb the dynamics of the loop³⁹, alter the binding specificity of CypA for substrates such as immunodeficiency virus capsids⁴⁰, and restrict the host range of these

viruses^{41,42}. Second, a group of residues that extends from the active site into the core of the protein (hereafter referred to as the ‘core’ region) has also been shown to be mobile on a millisecond timescale³⁸. Subsequent work incorporating multi-temperature X-ray crystallography⁴³, mutagenesis⁴⁴ and further NMR experiments⁴⁵ has established a relationship between catalysis and conformational dynamics of a group of side chains in this region. Motivated by the sensitivity of the conformational state of the active site–core network to temperature⁴³, we performed infrared laser-driven temperature jumps on buffered aqueous solutions of CypA and measured subsequent, time-dependent changes in SAXS/WAXS. While our measurements provide only low-resolution structural information, we were able to measure the kinetics of protein conformational changes in CypA. We identified two relaxation processes, and by performing temperature-jump experiments at a range of different temperatures, were able to calculate thermodynamic properties of the transition states for the underlying conformational transitions. Specific mutants in the ‘loops’ or the ‘core’ regions of CypA show that the two processes are independent, each representing a distinct and uncoupled reaction coordinate on a complex conformational landscape. Collectively, our measurements and analysis show that a wealth of information about a protein’s conformational landscape can be obtained by pairing laser-induced temperature jump with time-resolved X-ray scattering.

Results

A method for simultaneous measurement of structural and kinetic details of intrinsic protein dynamics. To measure protein structural dynamics, we utilized a pump–probe method that pairs an infrared laser-induced temperature jump with global measurement of protein structure via X-ray solution scattering (Fig. 1c).

We performed solvent heating in aqueous protein solutions by exciting the water O–H stretch with mid-IR laser pulses (1,443 nm; 7 ns duration). At regularly defined time delays following the infrared heating pulse (from 562 ns to 1 ms), we probed the sample with high-brilliance synchrotron X-ray pulses from a pink-beam undulator (3% bandwidth at 12 keV; Supplementary Fig. 1) that were approximately 500 ns in duration, and measured X-ray scattering using a large charge-coupled device detector that was capable of capturing small and wide scattering angles on a single panel. Because the duration of the infrared pump pulse was sufficiently short compared with the duration of the X-ray probe, the heating was effectively instantaneous with respect to the relaxation processes we were able to observe. Data were collected as interleaved ‘laser on’ and ‘laser off’ X-ray scattering images, so that each pump–probe measurement could be paired to a measurement made immediately before application of the pump laser (Fig. 1d). We measured 27 unique pump–probe time delays across four decades of time spanning from 562 ns to 1 ms, performing 50 repeat measurements for each time delay. For each detector image, the individual pixel values were azimuthally averaged as a function of the scattering vector magnitude, q , to give one-dimensional scattering intensity profiles ($I(q)$ curves). All scattering profiles were scaled to a single reference, and the data were analysed as described below. These pump–probe measurements allowed us to monitor structural changes within the ensemble of heated molecules in real time as the system relaxed to a new thermal equilibrium following temperature jump (Fig. 1a). Additional details about the instrumentation used for our experiments are available in the Supplementary Methods.

Calibrating the magnitude of the temperature jump. Because the isothermal compressibility of liquid water is highly temperature dependent⁴⁶, X-ray scattering from the bulk solvent acts as an exquisitely sensitive thermometer that can be used to calibrate the magnitude of the temperature jump in our experiments^{9,25,27,47,48}. Our instrument configuration allowed us to record low-angle protein scattering (SAXS) and high-angle solvent scattering (WAXS) on the same detector image, enabling simultaneous measurement of CypA structure and water temperature. A method based on singular value decomposition of a matrix constructed from the scattering curves provided a simple way to measure the magnitude of the laser-induced temperature jump (Supplementary Data). The singular value decomposition (SVD) analysis showed that the average temperature jump produced by our infrared heating pulse was approximately 10.7 °C, and allowed us to judge when cooling of the system became significant, so that we could identify the maximum pump–probe time delay that was valid for our relaxation analysis. Additionally, the SVD analysis revealed that following laser temperature jump, the solvent reaches a new thermal equilibrium faster than the measurement dead time of our experiment (562 ns). This observation is consistent with other work, in which changes to the structure of bulk solvent following laser temperature jump have been shown to equilibrate within roughly 200 ns⁴⁹.

Temperature jump produces changes in the X-ray scattering profile of CypA. To determine the effect of the temperature jump, we initially averaged all data for a given time delay (Supplementary Fig. 2), examined the scattering profiles for differences, and observed a small laser-induced change in the low- q region of the scattering profiles. Next, we exploited the interleaved data collection scheme (Fig. 1d) to increase the sensitivity of the experiment. We calculated the ‘on–off difference’ between the paired ‘laser on’ and ‘laser off’ scattering profiles. Following subtraction, we binned the ‘on–off difference’ scattering curves according to the associated pump–probe time delay, performed an iterative chi-squared test to remove outliers, and averaged the calculated differences for all repeat measurements (Supplementary Methods). This subtraction and averaging

removed systematic errors that accumulated over long experiments, resulting in accurate measurements of difference signals as a function of the pump–probe time delay. Scattering differences at high q (1.0–4.2 Å^{−1}) were used to calibrate the final sample temperature after laser illumination (Supplementary Data), while differences at low q (0.03–0.3 Å^{−1}) were analysed in the context of the average physical dimensions and scattering density of the CypA ‘protein particle’. Here, we use the phrase ‘protein particle’ to describe the protein molecule plus the ordered solvent bound to its surface, since both the protein and its hydration layer have electron densities that differ from bulk solvent, and therefore contribute to the observed X-ray scattering by the CypA solution. Identical temperature-jump measurements were performed on protein samples and on samples consisting of buffer only without protein, and an additional step of scaling and subtraction isolated the signal changes at low q due only to the protein (Supplementary Fig. 2).

Time-resolved changes in small-angle X-ray scattering.

Comparison of difference scattering curves calculated for 27 unique time delays revealed a time-dependent change in X-ray scattering by the protein, demonstrating that the modest temperature jump we introduced was capable of exciting protein dynamics that could be observed in real time. The difference curves calculated from our data, showing the contribution of the protein to time-resolved changes in the SAXS/WAXS signal, have features in the low- q ($q=0.03$ – 0.3 Å^{−1}) region (Fig. 2a). The time-resolved differences are approximately the same in magnitude and direction as differences between static temperature measurements performed on samples equilibrated to temperatures that differ by 10 °C, and correspond to the ‘laser on’ and ‘laser off’ temperatures (Supplementary Fig. 3). Qualitatively, the time-resolved on–off differences show that the overall low-angle scattering and extrapolated value of $I(0)$ are reduced within the measurement dead time of our experiment (562 ns), then begin to increase slightly over the next few microseconds before decreasing further at longer pump–probe time delays out to 562 μs. The observed ‘laser on–off’ difference in $I(0)$ is approximately 3% of the total observed signal, with one-third of that signal change occurring in the time regime that can be resolved by our measurements. Additional static measurements as a function of both CypA concentration and temperature allowed us to characterize and correct for the effect of interparticle interactions, and to determine that their contribution to the X-ray scattering is temperature independent over the relevant temperature range (Supplementary Data).

Changes in low-angle scattering generally reflect changes in the overall size and shape of the particles in solution, and changes in $I(0)$ result from changes to the scattering density of the particle (that is, the number of excess electrons in the particle relative to the bulk solvent that it displaces). The observed reduction of both quantities indicated shrinkage of the scattering particles, and a decrease in their scattering density. The on–off difference for our shortest pump–probe time delay (562 ns) is significantly different from 0 at low scattering angles (a reduction in $I(0)$ of approximately 2%), showing the existence of structural changes that are faster than the measurement dead time of our experiments. We hypothesize that the physical basis for this fast signal change is a combination of thermal disorder and thermal expansion effects (Supplementary Data), possibly including a temperature-induced loss (or ‘melting’) of ordered solvent within the protein hydration shell. Previous work, such as measurement of ‘protein-quake’ motions in photoactive systems^{9,50}, has shown that vibrational energy transfer is fast in proteins, and we expect that the onset of thermal disorder following temperature jump happens on a similar timescale, within hundreds of picoseconds. Furthermore, thermal expansion of solvent following temperature jump is well known to equilibrate within approximately 200 ns^{25,47,48}, and it is reasonable to assume that protein thermal expansion may occur on a

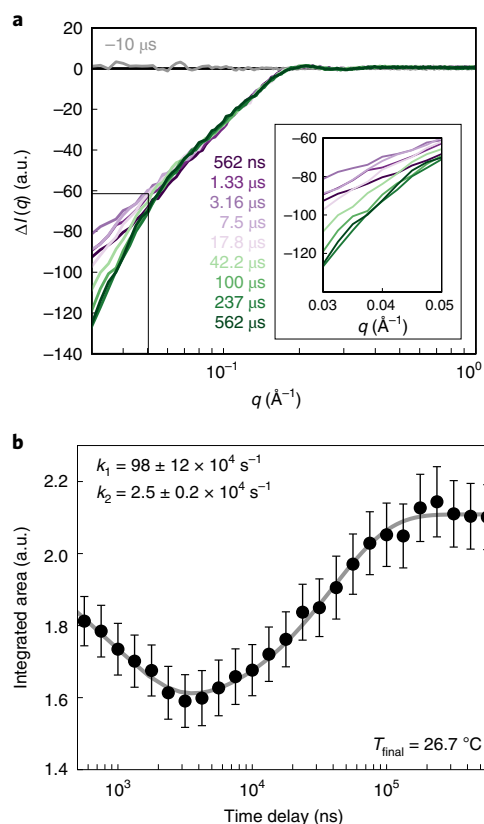


Fig. 2 | Temperature-jump data allow kinetic modelling of conformational dynamics. **a**, A series of averaged time-resolved difference X-ray scattering curves is shown for a subset of our temperature-jump data (10 out of 26 unique time delays). Data at low q are plotted on a linear q scale in the inset. **b**, The area under the averaged difference scattering curve in the $q = 0.03\text{--}0.05\text{ \AA}^{-1}$ region was integrated for all measured pump-probe time delays, and the resulting absolute values and their standard deviations (error bars) are plotted as a function of the pump-probe time delay. The plotted data suggest the existence of multiple relaxation processes, and we used a two-step model of relaxation kinetics to fit the observations (grey line). The rates calculated from the kinetic fit are provided, along with their standard deviations.

similar timescale. The notion that thermal disorder and expansion occur rapidly in our system is experimentally supported by Kratky plots created from our static and temperature-jump data, which suggest a slight increase in protein flexibility without unfolding (Supplementary Fig. 4). The effect is temperature dependent, but not time dependent over the pump-probe time delays we explored, confirming that the process is faster than the measurement dead time of our experiment. The physical details and kinetics of these fast processes, although interesting, are invisible to our experiment, and it is impossible to disentangle and quantify the exact contribution of each phenomenon to the fast signal decrease that we observe. Therefore, our analysis focused on structural dynamics that occur in the microsecond regime, which is resolved by our measurements. In contrast, ultrafast single pulses from synchrotrons (pulse duration: ~ 100 ps) or X-ray free-electron lasers (pulse duration: ~ 10 fs) offer the opportunity to study fast timescale processes using temperature jump; however, a major consideration for temperature-jump SAXS/WAXS experiments on faster timescales is the additional complexity in signal interpretation resulting from the complex dynamics of the solvent that immediately follow laser heating^{25,47,48} but equilibrate within the dead time of the experiments reported here.

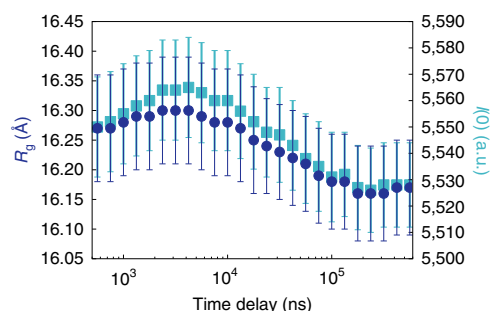


Fig. 3 | Time-resolved Guinier analysis. Changes in the average physical parameters of CypA protein particles in the ensemble were estimated by performing Guinier analysis on the time-resolved scattering data. Consistent with our kinetic analysis, the radius of gyration (R_g) of the average CypA particle in solution first increases and then decreases as a function of time following the temperature jump. Additionally, the value of $I(0)$ extrapolated from the Guinier analysis shows an analogous increase and decrease, suggesting that the change in the particle size is coupled to a change in its average electron density, which is probably due to the acquisition and loss of water molecules from the solvation shell as the protein swells and then shrinks. Error bars represent standard deviations of the calculated quantities based on propagation of measurement standard deviations.

Kinetic modelling of structural dynamics from time-resolved scattering differences. Our time-resolved measurements of scattering differences allowed us to model the kinetics of global structural changes induced by the temperature jump (Supplementary Methods). We integrated the area under each of our time-resolved difference curves in the $q = 0.03\text{--}0.05\text{ \AA}^{-1}$ region and plotted the absolute value of the area as a function of the associated pump-probe time delay (Fig. 2b). Based on the apparent shape of the area versus time delay plot, we reasoned that a two-step kinetic model would be needed to fit the data, since the area first decreases and then increases as a function of time delay. We fit the observed data to a two-step model of relaxation kinetics (independent steps) using a nonlinear least-squares curve-fitting algorithm, and calculated rates of $98 \times 10^4 \pm 12 \times 10^4\text{ s}^{-1}$ for the fast process (k_1) and $2.5 \times 10^4 \pm 0.2 \times 10^4\text{ s}^{-1}$ for the slow process (k_2) at 26.7°C (299.7 K). The errors calculated for these rates were determined using a bootstrap analysis⁵¹. We note that while other experimenters have employed SVD for kinetic analysis of time-resolved scattering, in our case, implementation of this method produced comparable results to the integration analysis, as one singular vector dominated the time-dependent signal identified by the SVD (Supplementary Fig. 5).

In addition to performing kinetic analysis, we also utilized the on-off difference curves to generate $I(q)$ scattering curves, which represent the time-dependent X-ray scattering from the CypA sample, but are generated in a manner that makes use of the paired ‘laser on’ and ‘laser off’ measurements to reduce the effects of systematic error (Supplementary Methods). We subsequently used these scattering curves for Guinier analysis (Supplementary Methods) to determine how the radius of gyration (R_g) of the average CypA particle in the conformational ensemble changes as a function of time following the temperature jump (Fig. 3). This analysis showed that after the temperature jump, the average CypA particle shrinks in the dead time of our experiment; however, within a few microseconds of the temperature jump, a fast structural transition (described by k_1 in our kinetic analysis) causes the average CypA particle to expand subtly. While the increase in the calculated radius of gyration is small relative to the error on the Guinier fit for any single data point, our conclusion that the particle is expanding is supported by multiple time

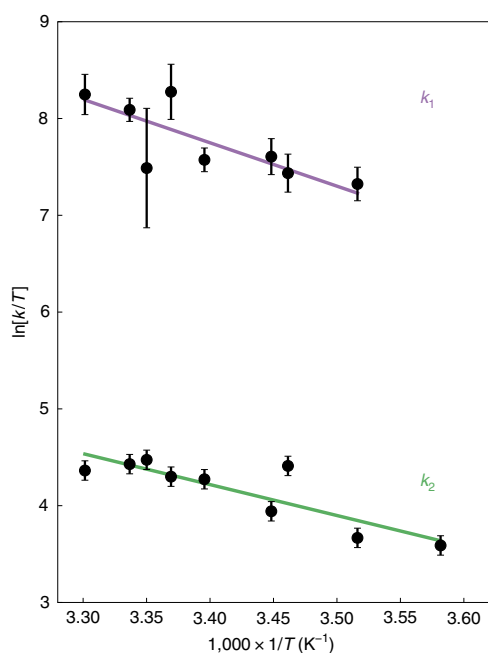


Fig. 4 | Eyring analysis of transition-state thermodynamics. Linear Eyring plots are shown for each of the two relaxation processes observed in our temperature-jump experiments with CypA. Both processes show Arrhenius behaviour over the temperature range we explored, and the apparent linear relationship between $1/T$ and $\ln[k/T]$ justifies the use of the linear Eyring equation (equation (1)) to fit the data and characterize the thermodynamics of the transition states. Data points represent rates determined from the kinetic analysis, and error bars represent the standard deviations of the rates based on propagation of the measurement standard deviations and bootstrap analysis (see Supplementary Methods). The fit to the linear Eyring equation for the fast process (k_1) is shown in purple, and the fit for the slow process (k_2) is shown in green.

points and kinetic analysis of the integrated area under difference scattering curves. Following this fast increase in R_s , a second, slower process (described by k_2 in our kinetic analysis) reverses this trend, causing the average CypA particle to shrink again.

To learn more about the conformational transitions in CypA that are excited by the temperature jump, we repeated the experiment at multiple different jumped (final) temperatures ranging from 6.2–29.9 °C (279.2–302.9 K). We modelled the kinetics of the SAXS/WAXS signal changes to observe how the relaxation rates changed as a function of temperature. The calculated rates (k_1 and k_2) for all temperatures are provided in Supplementary Table 1. We analysed the temperature dependence of these rates, which provided insight into the thermodynamics of the transition states for the two processes. We plotted $\ln[k/T]$ versus $1/T$ (Fig. 4), and noted that for both processes the relationships appeared to be linear. Therefore, we calculated enthalpies and entropies of activation for each of the two processes by fitting the data to the linearized Eyring equation:

$$\ln \left[\frac{k}{T} \right] = \left(\frac{\Delta H^\ddagger}{R} \times \frac{1}{T} \right) + \frac{\Delta S^\ddagger}{R} + \ln \left[\frac{k_B}{h} \right] \quad (1)$$

where R is the gas constant, k_B is the Boltzmann constant and h is Planck's constant. The enthalpies of activation ($\Delta H^\ddagger$) and entropies of activation ($\Delta S^\ddagger$) and their standard deviations are given in Table 1. The fast process (k_1) has a large, positive enthalpy of activation, but this is partially offset by a slightly positive entropy of activation. Formation of the transition state during the slow process (k_2) has a smaller enthalpic cost, but is also entropically disfavoured. We note

Table 1 | Thermodynamic properties of transition states

	$\Delta H^\ddagger$ (10^3 J mol $^{-1}$)	$\Delta S^\ddagger$ (J mol $^{-1}$ K $^{-1}$)
Fast process (k_1)	37 ± 8	32 ± 26
Slow process (k_2)	27 ± 3	-34 ± 11

Enthalpies ($\Delta H^\ddagger$) and entropies ($\Delta S^\ddagger$) of activation, and their standard deviations based on propagation of measurement standard deviations through all analysis steps, for the fast (k_1) and slow (k_2) processes observed for wild-type CypA, calculated from Eyring analysis.

that the lowest temperature measurement (279.2 K) was not used in the Eyring analysis of the fast process (k_1) because the error on the measured rate was large due to the low magnitude of the overall time-resolved signal changes at this temperature.

CypA mutations with distinct effects on conformational dynamics alter time-resolved signal changes. The time-resolved signal changes that we attributed to wild-type CypA were observed only at low scattering angles; therefore, the resulting structural information had very limited resolution. To gain a better understanding of the structural transitions excited by the temperature jump, we next studied two specific CypA mutants, S99T (in the 'core' region; red in Fig. 1b) and NH (D66N/R69H in the 'loops'; blue in Fig. 1b). The conformational dynamics of these two variants of the enzyme each differ from the wild type in distinct ways: S99T is catalytically impaired due to a loss of rotameric exchange in a key network of residues, whereas NH alters the substrate specificity of CypA by enhancing the dynamics of the surface-exposed loops adjacent to the active site. Importantly, NMR relaxation measurements indicate that S99T perturbs the active site but not the loops⁴⁴, and that NH only perturbs the loops³⁹.

We observed that both S99T and NH mutants showed time-resolved SAXS signal changes that differed from the wild-type enzyme. Both mutants showed a fast signal change that occurred within the measurement dead time of the experiment, similar to what was observed for the wild-type enzyme, and consistent with these changes being largely due to thermal expansion or temperature-dependent changes to the solvation shell. Beyond the initial fast loss of scattering intensity that was observed (and nearly identical) for all three CypA variants we studied (wild type, S99T and NH), the evolution of the time-resolved signals for each of the two mutants differed substantially from the wild type and from one another. In the S99T mutant (Fig. 5a), we observed only the fast decrease (k_1) of the integrated area under the difference curve ($q=0.03$ – 0.05 Å $^{-1}$), and a striking absence of the subsequent increase (k_2) in the integrated area at longer time delays that was observed for the wild-type enzyme. In contrast, for NH, the plot of integrated area under the difference curve as a function of the time delay (Fig. 5b) appears to lack the initial fast decrease (k_1), but it does appear to retain the slower signal change (k_2) that results in an increase for this quantity at longer time delays. We found that for the mutants, the two-step kinetic model that was fit to the wild-type data yielded at least one rate with a large error. For the S99T mutant, the first step (k_1) was well fit but the second step (k_2) was poorly fit, while the opposite was true for the NH variant. After visual inspection, we chose to use a single-step kinetic model to fit the data for the S99T and NH mutants, and the calculated rates for the two mutants (k_1 for S99T and k_2 for NH) are also given in Supplementary Table 1. Plots of the residuals for these fits revealed no structure, suggesting that a single-step kinetic model is sufficient to explain the data for the CypA mutants. In contrast, fitting kinetic data collected for the wild-type enzyme using a single-step model results in residuals with exponential character, and a two-step kinetic model is needed to reduce the error in the fit (Supplementary Fig. 6).

Our measurements of the S99T and NH variants of CypA clearly show that mutations that are known to impinge on the activity and

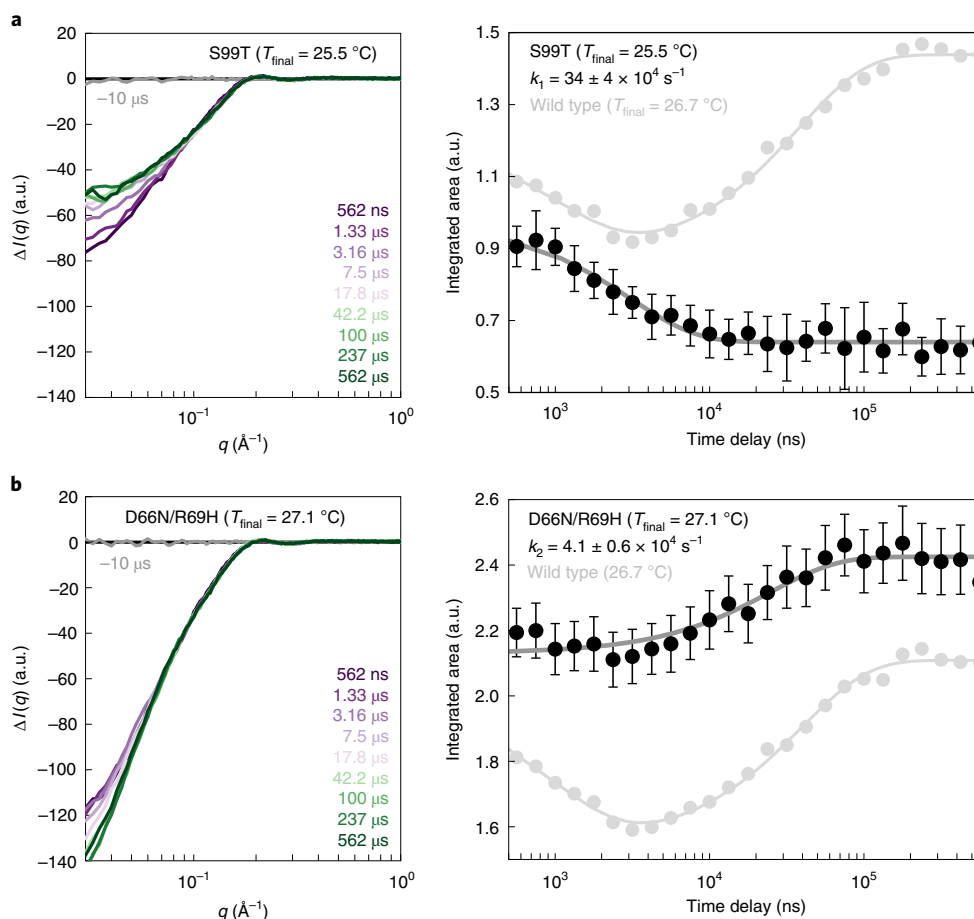


Fig. 5 | Kinetic analysis of CypA variants. **a,b**, Two CypA mutants (S99T (**a**) and D66N/R69H (**b**)), with distinct effects on the enzyme's function, show the link between the observed temperature-jump signal and functional dynamics. The data are presented in the same manner as for the wild-type enzyme shown in Fig. 4. Left panels show a subset of averaged difference scattering curves for several pump-probe time delays. Data points in the right panels represent the absolute value of the integrated area under the corresponding averaged difference scattering curve in the $q = 0.03\text{--}0.05\text{ \AA}^{-1}$ region, and the error bars represent standard deviations. In the plots of integrated area versus pump-probe time delay (right panels), the signal observed for the wild-type enzyme is shown in light grey for comparison. In **a**, the S99T mutant, which displays defective catalytic function, shows only the fast relaxation process (k_1) and lacks the slower process (k_2). In **b**, the D66N/R69H (NH) mutant, with altered substrate specificity, shows the slow relaxation process (k_2) and lacks the faster process (k_1).

specificity of the enzyme also perturb the observed time-resolved signal relative to the wild type in our temperature-jump experiments. Most notably, the slow relaxation process (modelled by k_2) is shared only by the catalytically competent wild-type and NH variants, and its absence from the S99T variant suggests that the underlying conformational change is related to the catalytically coupled motions that are arrested by the S99T mutation. These results indicate that temperature-jump experiments are capable of exciting and measuring functionally relevant, intramolecular structural dynamics of proteins, even when the data are limited to relatively low scattering angles.

Discussion

Time-resolved temperature-jump X-ray scattering experiments have the potential to be a powerful tool for understanding the complex dynamics of protein molecules, such as the model enzyme CypA, which has no intrinsic photoactivity. In our experiments, temperature jumps of 10–11 °C modified the CypA conformational ensemble, producing a clear, time-dependent change at low scattering angles. Because the time-resolved scattering changes could be observed only at low q , we demonstrated that interparticle spacings (quantified by the structure factor, $S(q)$) were temperature independent, and that

the changes were due to structural rearrangements within the protein particle. Guinier analysis of time-resolved scattering curves allowed us to track changes in the average radius of gyration (R_g) of the protein particles, which include the CypA molecules plus their ordered solvation shells. This signal change comprises an initial reduction in low-angle scattering that occurs within the measurement dead time of our experiment, followed by a small increase in low-angle scattering that equilibrates within a few microseconds ($k_1 = 98 \pm 12 \times 10^4\text{ s}^{-1}$ at 26.7 °C), and finally a further reduction in low-angle scattering that equilibrates within tens of microseconds ($k_2 = 2.5 \pm 0.2 \times 10^4\text{ s}^{-1}$ at 26.7 °C). High-angle scattering differences required for atomistic structural interpretation were not observed due to signal-to-noise considerations. We suspect that high-angle features in time-resolved difference scattering curves may be especially weak for proteins such as CypA, in which conformational motions involve correlated shifts of atoms that can preserve many properties of short-range structure.

This analysis suggests a model in which the scattering density of the CypA particle (protein and ordered solvent) first increases and then decreases after excitation by the temperature jump. By performing temperature-jump experiments over a range of temperatures, we discovered that both the fast and slow processes we observed could be described using Arrhenius kinetics. An Eyring analysis revealed

relatively large, positive enthalpies of activation for both processes, consistent with the idea that conformational changes generally require breakage of existing interactions in both the protein and the solvent. The activation enthalpy for the fast process (k_1) is larger, but the overall activation energy is lower because of a favourable activation entropy. The opposite is true for the slower process (k_2), which has a smaller overall activation enthalpy, but a disfavoured activation entropy. We used the S99T and NH variants of CypA to disentangle the nature of these processes and their associated functions. The S99T mutant is capable of undergoing the fast (k_1) expansion process, but does not experience the subsequent slow (k_2) shrinkage. NMR and crystallography have shown that this mutation arrests the conformational exchange of the 'core' catalytic network of residues in CypA by creating steric hindrance, strongly favouring a minor conformation of the wild-type enzyme⁴⁵. This interpretation suggests that the internal rearrangements are related to the k_2 process. However, there is a separation of timescales between the NMR results, which indicate millisecond dynamics in the 'core' region, and the temperature-jump SAXS results here, which indicate microsecond dynamics. This discrepancy may reflect coupled processes that are related by a population-shuffling mechanism⁵² and agree with a broad timescale range of side chain dynamics in CypA uncovered by molecular dynamics experiments⁵³. In contrast with the S99T mutant, the NH variant lacks the initial fast signal change (k_1) in our temperature-jump experiments, but clearly retains the slow (k_2) signal. NMR and crystallographic studies of the D66N/R69H (NH) double mutant showed that it maintains wild-type catalytic motions, but breakage of several hydrogen-bonding interactions results in enhanced flexibility of a surface loop adjacent to the active site and altered substrate specificity³⁹. NMR experiments with NH have shown that the loop motions still occur, but at an increased rate that renders them invisible to our experiments. Therefore, we hypothesize that the loop motions are responsible for the fast (k_1) signal in wild-type CypA and S99T, where these motions have been shown (by NMR) to be unperturbed⁴⁴.

The assignment of the motions responsible for the experimentally observed R_g and $I(0)$ changes by mutational analysis is also consistent with X-ray crystal structures of CypA. Using a room temperature X-ray crystal structure of wild-type CypA (Protein Data Bank ID: 3K0N), we calculated the R_g of the enzyme with the core catalytic network (Arg55, Met61, Ser99 and Phe113) in both the major and minor conformational states, and found that the predicted R_g of the minor state is 0.07 Å smaller than the major state (14.09 versus 14.16 Å, for protein atoms only). In addition to decreasing the R_g of the protein, conversion from the major to the minor state also results in a small reduction in solvent-exposed surface area, which would necessarily reduce the size of the protein solvation layer. While the reduction of the protein's R_g makes it more compact, and therefore should increase $I(0)$ because the protein has become more dense, the loss of material from the solvation layer opposes this effect, probably leading to the observed decrease in the scattering density of the CypA particle and reduction in $I(0)$. This coupling of R_g changes with changes in the solvent-exposed surface area of the protein can potentially explain the observed correlation between R_g and $I(0)$ changes calculated from our time-resolved scattering data. Additionally, the increase in the average R_g and $I(0)$ during the faster process (k_1) is consistent with the loops sampling an expanded conformational ensemble (with increased surface area), as indicated by recent exact nuclear Overhauser effect NMR ensembles³⁴. Our kinetic modelling of the wild-type and mutant data suggests that two dynamic modes are observed, each of which is individually perturbed by different mutations. However, the extent to which the observed motions are coupled is unclear from our measurements. The S99T variant clearly shows that the fast motion can occur independent of the slower motion, because S99T is known to have arrested slow dynamics. In contrast, the fast motion we observe is

likely to be accelerated by the NH mutation, becoming too fast for us to observe rather than being impeded. Therefore, we are unable to determine whether the fast motion is a requisite first step that precedes the slower motion.

Time-resolved X-ray structural measurements are critical for decoupling the experimental signatures of conformational changes that can become convoluted by the spatial and temporal averaging that is inherent to traditional X-ray experiments. If one were to assess traditional, static SAXS data for CypA, they would find that increasing the temperature of the sample results in a decrease in the average particle size at equilibrium. These static measurements as a function of equilibrium temperature fail to capture that the temperature change actually perturbs two distinct protein motions, which have the opposite effect on the enzyme's global structural characteristics. This information can only be obtained through a time-resolved experiment, which is able to separate the effects of these two motions because they have substantially different rates. The ability to dissect individual conformational motions and measure their rates is important for understanding processes involving complex protein dynamics. Many of these dynamic processes, including allostery^{55–58} and enzyme catalysis^{59–63}, involve extensive reorganization of interactions between the protein and its ordered solvation shell, which are key contributors to the energetics that govern protein motions^{28,29,64–68}. Because X-ray solution scattering experiments report on the structure of a protein and the ordered solvent molecules that constitute its solvation shell^{69–72}, the widespread application of time-resolved SAXS/WAXS experiments will enhance our understanding of protein motions, including how they are driven by solvent dynamics, especially when they can be combined with longer and increasingly detailed molecular dynamics simulations to provide atomic-scale insight into the underlying structural changes^{9–12}.

For time-resolved experiments to enter the mainstream of structural biology, it is necessary to create general perturbations that can be applied universally (to any protein of interest), and our results establish that temperature jump can excite functional intramolecular protein motions for time-resolved X-ray structural measurements. The intrinsic motions of proteins and other biological macromolecules are naturally driven by thermal fluctuations, which makes temperature jump an ideal perturbation for studying their dynamics using time-resolved X-ray scattering. Consistent with this idea, a growing number of X-ray diffraction studies across multiple systems have shown that conformational equilibria, which are sensitive to temperature perturbation, are the same ones exploited by evolution to create new functions^{73,74}, by medicinal chemists in identifying novel small-molecule binding sites^{75–77}, and by enzymes during their catalytic cycles^{44,78,79}, paving the way for time-resolved measurements on a broad variety of biochemical systems. Temperature-jump solution scattering experiments are powerful tools for understanding protein dynamics, especially when they can be paired with atomistic simulations⁶⁹, or with crystal structures of alternative conformations and functional perturbations, as demonstrated here. In the future, functional perturbations such as the mutations that we used to study CypA, or binding of chemical ligands, could be paired with time-resolved structural data from temperature-jump experiments to provide insight into how disease alleles or drug molecules impinge on protein dynamics.

Data availability

Scattering data are deposited at NIH Figshare (<https://doi.org/10.35092/yhjc.9177143>). Additional information and files are available from the corresponding author upon reasonable request.

Code availability

All Python scripts used for analysis of integrated X-ray scattering curves are publicly available from GitHub (<https://github.com/>

fraser-lab/solution_scattering). A code release checkpoint containing the exact scripts used in this work is available via Zenodo (<https://doi.org/10.5281/zenodo.3355707>).

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Author contributions

M.C.T., P.A. and J.S.F. conceived and designed the experiments. M.C.T., B.A.B., A.M.W., H.S.C., F.S., D.M.C.S. and P.A. performed the experiments. M.C.T., B.A.B. and A.M.W. analysed the data. M.C.T., B.A.B., A.M.W., H.S.C., F.S. and P.A. contributed materials/analysis tools. M.C.T. and J.S.F. wrote the paper. All authors discussed the results and commented on the manuscript.

Competing interests

The authors declare no competing interests.

Additional information

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