



The half-volume spectrum of a manifold

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Abstract

We define the half-volume spectrum $\{\tilde{\omega}_p\}_{p \in \mathbb{N}}$ of a closed manifold (M^{n+1}, g) . This is analogous to the usual volume spectrum of M , except that we restrict to p -sweepouts whose slices each enclose half the volume of M . We prove that the Weyl law continues to hold for the half-volume spectrum. We define an analogous half-volume spectrum $\tilde{c}(p)$ in the phase transition setting. Moreover, for $3 \leq n+1 \leq 7$, we use the Allen–Cahn min–max theory to show that each $\tilde{c}(p)$ is achieved by a constant mean curvature surface enclosing half the volume of M plus a (possibly empty) collection of minimal surfaces with even multiplicities.

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1 Introduction

The spectrum of the Laplacian is an important invariant of a closed Riemannian manifold (M^{n+1}, g) . A number λ is called an eigenvalue of the Laplacian provided there is a function $u: M \rightarrow \mathbb{R}$ such that $\Delta u + \lambda u = 0$. It is well-known that the eigenvalues form a discrete sequence $0 = \lambda_0 < \lambda_1 \leq \lambda_2 \leq \dots$ and $\lambda_p \rightarrow \infty$ as $p \rightarrow \infty$. In fact, the eigenvalues of Laplacian are characterized by the min–max formula

$$\lambda_p = \inf_{(p+1)\text{-planes } P \subset W^{1,2}(M)} \left[\sup_{u \in P \setminus \{0\}} \frac{\int_M |\nabla u|^2}{\int_M u^2} \right],$$

and they satisfy the Weyl law

$$\lambda_p \sim 4\pi^2 \text{Vol}(B)^{-\frac{2}{n+1}} \text{Vol}(M)^{-\frac{2}{n+1}} p^{\frac{2}{n+1}}$$

as $p \rightarrow \infty$. Here B is the unit ball in \mathbb{R}^{n+1} .

In [11], Gromov proposed a non-linear analog of the spectrum of the Laplacian. Roughly speaking, he defines a p -sweepout of M to be a family X of hypersurfaces with the following

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property: given any p points in M , there is a hypersurface Σ belonging to the family X which passes through all p of these points. Then he defines the p -widths

$$\omega_p = \inf_{p\text{-sweepouts } X} \left[\sup_{\Sigma \in X} \text{Area}(\Sigma) \right].$$

See Sect. 2 for precise definitions. The sequence $\{\omega_p\}_{p \in \mathbb{N}}$ is called the volume spectrum of M .

Gromov [12] and Guth [14] proved that the volume spectrum satisfies sublinear growth bounds. Namely, there are constants C_1 and C_2 depending on M such that

$$C_1 p^{\frac{1}{n+1}} \leq \omega_p \leq C_2 p^{\frac{1}{n+1}}.$$

Later, Liokumovich, Marques, and Neves [17] showed that the volume spectrum satisfies a Weyl law. That is, there is a universal constant a_n depending only on the dimension such that

$$\omega_p \sim a_n \text{Vol}(M)^{\frac{n}{n+1}} p^{\frac{1}{n+1}}$$

as $p \rightarrow \infty$; see Chodosh-Mantoulidis [6] for the calculation of a_n when $n = 2$. The Weyl law for the volume spectrum has been instrumental in the proof of many results on the existence of minimal surfaces in Riemannian manifolds.

In the early 1980s, Almgren [2], Pitts [21], and Schoen-Simon [22] developed a min-max theory for the area functional on closed Riemannian manifolds. Their combined work implies that every closed Riemannian manifold of dimension $3 \leq n+1 \leq 7$, contains a closed, smooth, embedded minimal surface. Around the same time, Yau [25] conjectured that every closed manifold should contain infinitely many minimal surfaces. Marques and Neves devised a program to prove Yau's conjecture by using the Almgren-Pitts min-max theory to find a minimal surface with area ω_p for each $p \in \mathbb{N}$.

This program has now been successfully carried out. Fix a closed Riemannian manifold (M^{n+1}, g) with $3 \leq n+1 \leq 7$. Irie, Marques, and Neves [16] showed that, for a generic metric g , the union of all minimal surfaces in M is dense in M . In particular, this proved Yau's conjecture for generic metrics. Later Marques, Neves, and Song [20] refined this result to show that, for a generic metric g , there is a sequence of minimal surfaces in M which becomes equidistributed in M . The Weyl law for the volume spectrum was a key ingredient in the proof of both of these results. Following the second named author's proof of the Multiplicity One Conjecture [26], Marques and Neves [19] showed that, for a generic metric g , there is a sequence of minimal surfaces Σ_p with index p and $\text{Area}(\Sigma_p) = \omega_p$. Song [23] proved Yau's conjecture for arbitrary metrics g . We refer to the survey articles [18, 27] for more detailed history on this exciting field.

The Almgren-Pitts min-max theory relies heavily on tools from geometric measure theory. There is a parallel min-max theory for finding minimal surfaces based on the theory of phase transitions. This theory relies on the Allen–Cahn PDE and the varifold regularity theory of Wickramasekera. Gaspar and Guaraco [9] defined a phase transition spectrum $\{c(p)\}_{p \in \mathbb{N}}$ associated to a Riemannian manifold via the Allen–Cahn PDE. They showed that each $c(p)$ is achieved by a collection of minimal surfaces with multiplicities. Chodosh and Mantoulidis [5] proved the Multiplicity One Conjecture in the phase transition setting in ambient dimension three. Thus, for generic metrics on M^3 , they obtained the existence of a sequence of minimal surfaces Σ_p with index p and $\text{Area}(\Sigma_p) = c(p)$.

Gaspar and Guaraco [10] showed that the phase transition spectrum also satisfies a Weyl law. That is, there is a constant τ_n depending only on the dimension such that

$$c(p) \sim \tau_n \text{Vol}(M)^{\frac{n}{n+1}} p^{\frac{1}{n+1}}.$$

Dey [7] proved that actually $\omega_p = c(p)$ for all $p \in \mathbb{N}$ and thus the Almgren-Pitts volume spectrum and the phase transition volume spectrum coincide. In particular, the constants a_n and τ_n appearing in the two Weyl laws are equal.

In this paper, we define a “half-volume” spectrum associated to a Riemannian manifold. In the Almgren-Pitts setting, we restrict to p -sweepouts by families of hypersurfaces that each enclose half the volume of M . Then we define

$$\tilde{\omega}_p = \inf_{\text{half-volume } p\text{-sweepouts } X} \left[\sup_{\Sigma \in X} \text{Area}(\Sigma) \right].$$

The sequence $\{\tilde{\omega}_p\}_{p \in \mathbb{N}}$ is called the half-volume spectrum of M . In the phase transition setting, we define an analogous half-volume spectrum $\tilde{c}(p)$ by looking at critical points of the Allen–Cahn energy subject to the volume constraint $\int_M u = 0$. In both cases, we show that the Weyl law continues to hold. This gives the following theorems.

Theorem 1 *The Almgren-Pitts half-volume spectrum satisfies*

$$\tilde{\omega}_p \sim a_n \text{Vol}(M)^{\frac{n}{n+1}} p^{\frac{1}{n+1}}, \quad \text{as } p \rightarrow \infty.$$

Theorem 2 *The phase transition half-volume spectrum satisfies*

$$\tilde{c}(p) \sim \tau_n \text{Vol}(M)^{\frac{n}{n+1}} p^{\frac{1}{n+1}}, \quad \text{as } p \rightarrow \infty.$$

In the Allen–Cahn setting, we are able to use the results of Bellettini and Wickramasekera [3] to find varifolds achieving each $\tilde{c}(p)$. In the following theorem, a hypersurface Σ is called almost-embedded if near each point in M either Σ is embedded or Σ decomposes into an ordered union of embedded sheets.

Theorem 3 *Let (M^{n+1}, g) be a closed Riemannian manifold with $3 \leq n+1 \leq 7$. Fix a number $p \in \mathbb{N}$. There are*

- (i) *a Caccioppoli set $\Omega \subset M$ with $\text{Vol}(\Omega) = \frac{1}{2} \text{Vol}(M)$ whose boundary is smooth and almost-embedded with constant mean curvature,*
- (ii) *a (possibly empty) collection of smooth, disjoint minimal surfaces $\Sigma_1, \dots, \Sigma_k \subset M \setminus \Omega$,*
- (iii) *and positive integers $\theta_0 \in \mathbb{Z}$ and $\theta_1, \dots, \theta_k \in 2\mathbb{Z}$*

such that $\tilde{c}(p) = \theta_0 \text{Area}(\partial\Omega) + \theta_1 \text{Area}(\Sigma_1) + \dots + \theta_k \text{Area}(\Sigma_k)$. Moreover, $\theta_0 = 1$ unless $\partial\Omega$ is also a minimal surface.

Note that Theorem 3 produces a constant mean curvature surface that encloses half the volume of M . Previously, the second author and Zhu [28] developed a min-max theory in the Almgren-Pitts setting capable of finding surfaces of constant mean curvature c ; see also [29]. However, there is no control over the volume enclosed by the surface. Likewise, Bellettini and Wickramasekera [3] developed a min-max theory in the Allen–Cahn setting capable of finding surfaces of constant mean curvature c . Again there is no control over the volume enclosed by the surface. Thus there is a trade off. Theorem 3 produces constant mean curvature surfaces enclosing half the volume of M , but at the expense of losing control over the exact value of the mean curvature.

We conclude the introduction with some open problems. First, we conjecture that $\tilde{\omega}_p = \tilde{c}(p)$ for all $p \in \mathbb{N}$. Second, we conjecture that, for a generic metric g , the phase transition half-volume spectrum is achieved by multiplicity one constant mean curvature surfaces enclosing half the volume of M . In other words, in Theorem 3 the collection $\Sigma_1, \dots, \Sigma_k$ is empty and $\theta_0 = 1$ for every p . In particular, we conjecture that generically there are infinitely many

constant mean curvature surfaces enclosing half the volume of M . Finally, it is interesting to know whether one can find surfaces achieving $\tilde{\omega}_p$ by applying the Almgren-Pitts min-max theory with a volume constraint. This seems to be a difficult task. Already it is not obvious how to define a suitable pull-tight on the space of half-volume cycles.

2 The Almgren–Pitts half-volume spectrum

In this section, we investigate the topology of the space of half-volume cycles in a given manifold, and then define the Almgren–Pitts half-volume spectrum. Let (M^{n+1}, g) be a closed Riemannian manifold. We will use the following notation.

- Let $\mathfrak{h} = \frac{1}{2} \text{Vol } M$.
- Let $\mathcal{C}(M)$ denote the collection of all Caccioppoli sets in M .
- Let $\mathcal{C}_{\mathfrak{h}}(M)$, $\mathcal{C}_{\geq \mathfrak{h}}(M)$, and $\mathcal{C}_{\leq \mathfrak{h}}(M)$ denote the space of Caccioppoli sets with volume equal to \mathfrak{h} , greater than or equal to \mathfrak{h} , and less than or equal to \mathfrak{h} , respectively.
- Let $\mathcal{Z}(M, \mathbb{Z}_2)$ denote the set of all n -dimensional flat chains mod 2 in M .
- Let $\mathcal{B}(M, \mathbb{Z}_2)$ denote the set of all $T \in \mathcal{Z}(M, \mathbb{Z}_2)$ such that $T = \partial\Omega$ for some $\Omega \in \mathcal{C}(M)$. This is the connected component of the empty set in $\mathcal{Z}(M, \mathbb{Z}_2)$ in the flat topology.
- Let $\mathcal{H}(M, \mathbb{Z}_2)$ be the set of all $T \in \mathcal{Z}(M, \mathbb{Z}_2)$ such that $T = \partial\Omega$ for some $\Omega \in \mathcal{C}_{\mathfrak{h}}(M)$. This is the space of “half-volume cycles.”
- We use \mathcal{F} to denote the flat topology, \mathbf{F} to denote the \mathbf{F} -topology, and \mathbf{M} to denote the mass topology. All spaces are assumed to be equipped with the flat topology except where otherwise noted.
- We will abuse notation and write $\text{Vol}(\Omega)$ and $\text{Area}(T)$ instead of $\mathbf{M}(\Omega)$ and $\mathbf{M}(T)$ for $\Omega \in \mathcal{C}(M)$ and $T \in \mathcal{Z}(M, \mathbb{Z}_2)$, respectively.

We will show that $\mathcal{H}(M, \mathbb{Z}_2)$ is weakly homotopy equivalent to $\mathbb{R}\mathbb{P}^\infty$. The first step is to show that the double cover $\mathcal{C}_{\mathfrak{h}}(M)$ is contractible.

Proposition 4 *The space $\mathcal{C}_{\leq \mathfrak{h}}(M)$ deformation retracts to $\mathcal{C}_{\mathfrak{h}}(M)$.*

Proof The union of two Caccioppoli sets is a Caccioppoli set. Choose a Morse function $f: M \rightarrow \mathbb{R}$. For $s \in [0, \mathfrak{h}]$, let B_s be the sublevel set of f with volume equal to s . Note that each B_s is a Caccioppoli set. For each $\Omega \in \mathcal{C}_{\leq \mathfrak{h}}(M)$ and $t \in [0, 1]$, choose a number $s(\Omega, t)$ so that

$$\text{Vol}(\Omega \cup B_{s(\Omega, t)}) = \text{Vol}(\Omega) + t(\mathfrak{h} - \text{Vol}(\Omega)).$$

Note that there is not necessarily a unique choice for $s(\Omega, t)$, and the mapping $(\Omega, t) \mapsto s(\Omega, t)$ may not be continuous. Nevertheless, we claim that the map $\phi: \mathcal{C}_{\leq \mathfrak{h}}(M) \times [0, 1] \rightarrow \mathcal{C}_{\leq \mathfrak{h}}(M)$ defined by

$$\phi(\Omega, t) = \Omega \cup B_{s(\Omega, t)}$$

is continuous in the flat topology. Given this claim, ϕ is the required deformation retraction. Indeed, $\phi(\Omega, t) = \Omega$ for all $\Omega \in \mathcal{C}_{\mathfrak{h}}(M)$ and all $t \in [0, 1]$, and moreover, $\phi(\Omega, 1) \in \mathcal{C}_{\mathfrak{h}}(M)$ for all $\Omega \in \mathcal{C}_{\leq \mathfrak{h}}(M)$.

To see that ϕ is continuous, let $\varepsilon, \eta > 0$ be small positive numbers. Assume that $\Omega, \Theta \in \mathcal{C}_{\leq \mathfrak{h}}(M)$ satisfy $\text{Vol}(\Omega \Delta \Theta) < \varepsilon$ and that $t, r \in [0, 1]$ satisfy $|t - r| < \eta$. We need to check that $\text{Vol}(\phi(\Omega, t) \Delta \phi(\Theta, r))$ is small. First observe that

$$\text{Vol}(\phi(\Omega, t) \Delta \phi(\Theta, r)) = \text{Vol}(\phi(\Omega, t)) + \text{Vol}(\phi(\Theta, r)) - 2 \text{Vol}(\phi(\Omega, t) \cap \phi(\Theta, r)).$$

Now, without loss of generality, we can suppose that $s(\Omega, t) \leq s(\Theta, r)$. Then

$$\begin{aligned}\text{Vol}(\phi(\Omega, t) \cap \phi(\Theta, r)) &\geq \text{Vol}((\Omega \cup B_{s(\Omega, t)}) \cap (\Theta \cup B_{s(\Omega, t)})) \\ &= \text{Vol}((\Omega \cap \Theta) \cup B_{s(\Omega, t)}) \\ &\geq \text{Vol}(\Omega \cup B_{s(\Omega, t)}) - \varepsilon \\ &= \text{Vol}(\phi(\Omega, t)) - \varepsilon.\end{aligned}$$

Therefore we have

$$\text{Vol}(\phi(\Omega, t) \Delta \phi(\Theta, r)) \leq \text{Vol}(\phi(\Theta, r)) - \text{Vol}(\phi(\Omega, t)) + 2\varepsilon.$$

It remains to note that

$$\begin{aligned}|\text{Vol}(\phi(\Theta, r)) - \text{Vol}(\phi(\Omega, t))| &= |\text{Vol}(\Theta) + r(\mathfrak{h} - \text{Vol}(\Theta)) - \text{Vol}(\Omega) - t(\mathfrak{h} - \text{Vol}(\Omega))| \\ &\leq |\text{Vol}(\Theta) - \text{Vol}(\Omega)| + |r - t|\mathfrak{h} + |r\text{Vol}(\Theta) - t\text{Vol}(\Omega)| \\ &\leq \varepsilon + \eta\mathfrak{h} + r|\text{Vol}(\Theta) - \text{Vol}(\Omega)| + |r - t|\text{Vol}(\Omega) \\ &\leq 2\varepsilon + 2\eta\mathfrak{h}.\end{aligned}$$

Therefore we obtain

$$\text{Vol}(\phi(\Omega, t) \Delta \phi(\Theta, r)) \leq 4\varepsilon + 2\eta\mathfrak{h}$$

and the continuity of ϕ follows. \square

Proposition 5 *The space $\mathcal{C}(M)$ deformation retracts to $\mathcal{C}_{\leq \mathfrak{h}}(M)$, and the space $\mathcal{B}(M, \mathbb{Z}_2)$ deformation retracts to $\mathcal{H}(M, \mathbb{Z}_2)$.*

Proof Consider the deformation retraction $\phi: \mathcal{C}_{\leq \mathfrak{h}}(M) \times [0, 1] \rightarrow \mathcal{C}_{\leq \mathfrak{h}}(M)$ from the previous proposition. We can extend ϕ to an odd map $\psi: \mathcal{C}(M) \times [0, 1] \rightarrow \mathcal{C}(M)$ by the formula

$$\psi(\Omega, t) = \begin{cases} \phi(\Omega, t), & \text{if } \text{Vol}(\Omega) \leq \mathfrak{h} \\ M \setminus \phi(M \setminus \Omega, t), & \text{if } \text{Vol}(\Omega) \geq \mathfrak{h}. \end{cases}$$

Then ψ is a deformation retraction of $\mathcal{C}(M)$ onto $\mathcal{C}_{\leq \mathfrak{h}}(M)$. Moreover, since ψ is odd, this descends to a map $\theta: \mathcal{B}(M, \mathbb{Z}_2) \times [0, 1] \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$. This is the required deformation retraction of $\mathcal{B}(M, \mathbb{Z}_2)$ onto $\mathcal{H}(M, \mathbb{Z}_2)$. \square

Proposition 6 *Let K be the maximal area of a level set of the Morse function f used in the proof of Proposition 4. Let θ be the deformation retraction from Proposition 5. Then $\text{Area}(\theta(T, t)) \leq \text{Area}(T) + K$ for all $T \in \mathcal{B}(M, \mathbb{Z}_2)$ and all $t \in [0, 1]$.*

Proof Fix some $T \in \mathcal{B}(M, \mathbb{Z}_2)$ and some $t \in [0, 1]$. Choose a set $\Omega \in \mathcal{C}_{\leq \mathfrak{h}}(M)$ such that $\partial\Omega = T$. Let ϕ be the deformation retraction from Proposition 4. Then

$$\theta(T, t) = \partial\phi(\Omega, t) = \partial(\Omega \cup B_{s(\Omega, t)}).$$

Note that $\partial(\Omega \cup B_{s(\Omega, t)}) \subset \partial\Omega \cup \partial B_{s(\Omega, t)}$ and therefore $\text{Area}(\theta(T, t)) \leq \text{Area}(\partial\Omega) + K = \text{Area}(T) + K$, as needed. \square

The homotopy groups of the cycle spaces were originally computed by Almgren [1]. Later, Marques and Neves [19] gave a simplified proof in the case of codimension 1 cycles.

Theorem 7 (Marques-Neves) *The map $\partial: \mathcal{C}(M) \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$ is a double cover. The space $\mathcal{C}(M)$ is contractible, and $\mathcal{B}(M, \mathbb{Z}_2)$ is weakly homotopy equivalent to $\mathbb{R}\mathbb{P}^\infty$.*

Combined with the previous propositions, this yields the following corollary.

Corollary 8 *The map $\partial: \mathcal{C}_{\mathfrak{h}}(M) \rightarrow \mathcal{H}(M, \mathbb{Z}_2)$ is a double cover. The space $\mathcal{C}_{\mathfrak{h}}(M)$ is contractible and $\mathcal{H}(M, \mathbb{Z}_2)$ is weakly homotopy equivalent to $\mathbb{R}\mathbb{P}^\infty$. The inclusion map $\mathcal{H}(M, \mathbb{Z}_2) \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$ is a homotopy equivalence.*

We now recall the notion of sweepouts. Since $\mathcal{B}(M, \mathbb{Z}_2)$ is weakly homotopy equivalent to $\mathbb{R}\mathbb{P}^\infty$, it follows that the cohomology ring of $\mathcal{B}(M, \mathbb{Z}_2)$ with \mathbb{Z}_2 coefficients is $\mathbb{Z}_2[\lambda]$, where the generator λ is of degree 1. Let X be a cubical complex.

Definition 9 A flat continuous map $\Phi: X \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$ is called a p -sweepout if $\Phi^*\lambda^p \neq 0$ in $H^p(X, \mathbb{Z}_2)$.

Definition 10 A map $\Phi: X \rightarrow \mathcal{Z}(M, \mathbb{Z}_2)$ is said to have *no concentration of mass* provided

$$\lim_{r \rightarrow 0} \left[\sup_{q \in M} \sup_{x \in X} \text{Area}(\Phi(x) \llcorner B(q, r)) \right] = 0.$$

Definition 11 Let $\mathcal{P}_p(M)$ denote the collection of all p -sweepouts of M with no concentration of mass. Note that different p -sweepouts may have different domains.

Definition 12 The p -width of M is

$$\omega_p = \inf_{\Phi \in \mathcal{P}_p(M)} \left[\sup_{x \in \text{dom}(\Phi)} \text{Area}(\Phi(x)) \right].$$

Remark 13 In [17], the authors state that the cohomology ring of $\mathcal{Z}(M, \mathbb{Z}_2)$ is isomorphic to $\mathbb{Z}_2[\lambda]$. Then they define a p -sweepout as a map $\Phi: X \rightarrow \mathcal{Z}(M, \mathbb{Z}_2)$ such that $\Phi^*(\lambda^p) \neq 0$. However, the cohomology ring of $\mathcal{Z}(M, \mathbb{Z}_2)$ is actually $\bigoplus_i \mathbb{Z}_2[\lambda_i]$ where the direct sum is taken over the connected components of $\mathcal{Z}(M, \mathbb{Z}_2)$. These connected components are in bijection with homology classes in $H_n(M, \mathbb{Z}_2)$. Given this, there are several possible ways to define a p -sweepout. The simplest, which we shall adopt, is to replace the space $\mathcal{Z}(M, \mathbb{Z}_2)$ with $\mathcal{B}(M, \mathbb{Z}_2)$ as in Definition 9 so that the cohomology ring is indeed $\mathbb{Z}_2[\lambda]$. Alternatively, one could define a p -sweepout as a map $\Phi: X \rightarrow \mathcal{Z}(M, \mathbb{Z}_2)$ such that $\Phi^*(\lambda_i^p) \neq 0$ for some i . In either case, it is straightforward to see that one still obtains a Weyl law for the resulting p -widths.

We can now introduce the central object of the paper. By Corollary 8, the cohomology ring of $\mathcal{H}(M, \mathbb{Z}_2)$ with \mathbb{Z}_2 coefficients is also $\mathbb{Z}_2[\lambda]$. Again let X be a cubical complex.

Definition 14 A flat continuous map $\Phi: X \rightarrow \mathcal{H}(M, \mathbb{Z}_2)$ is called a half-volume p -sweepout if $\Phi^*\lambda^p \neq 0$ in $H^p(X, \mathbb{Z}_2)$.

Definition 15 Let $\mathcal{Q}_p(M)$ denote the collection of all half-volume p -sweepouts of M with no concentration of mass.

Definition 16 The half-volume p -width of M is

$$\tilde{\omega}_p = \inf_{\Phi \in \mathcal{Q}_p(M)} \left[\sup_{x \in \text{dom}(\Phi)} \text{Area}(\Phi(x)) \right].$$

We will call the sequence $\{\tilde{\omega}_p\}_{p \in \mathbb{N}}$ the half-volume spectrum of M .

Liokumovich, Marques, and Neves [17] showed that the p -widths of M satisfy a Weyl law.

Theorem 17 (Liokumovich, Marques, Neves) *There is a universal constant a_n such that $\omega_p \sim a_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}$ as $p \rightarrow \infty$.*

Next, we will show that the half-volume spectrum also satisfies a Weyl law. It is possible to prove this directly. However, this is not the approach we will take. Rather, we will show that the Weyl law for the half-volume spectrum follows from Theorem 17, together with the fact that every p -sweepout is homotopic to a p -sweepout by half-volume cycles.

Proposition 18 *The half-volume spectrum satisfies $\omega_p \leq \tilde{\omega}_p$ for all $p \in \mathbb{N}$.*

Proof Notice that any half-volume p -sweepout with no concentration of mass automatically belongs to $\mathcal{P}_p(M)$. Therefore, the proposition follows immediately from the definitions of ω_p and $\tilde{\omega}_p$. \square

Proposition 19 *There is a constant K depending only on M such that $\tilde{\omega}_p \leq \omega_p + K + 1$ for all $p \in \mathbb{N}$.*

Proof Choose a p -sweepout $\Phi: X \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$ in $\mathcal{P}_p(M)$ with

$$\sup_{x \in X} \text{Area}(\Phi(x)) \leq \omega_p + 1.$$

Let $\theta: \mathcal{B}(M, \mathbb{Z}_2) \times [0, 1] \rightarrow \mathcal{B}(M, \mathbb{Z}_2)$ be the deformation retraction constructed in Proposition 5. By Proposition 6, there is a constant K such that

$$\text{Area}(\theta(T, t)) \leq \text{Area}(T) + K$$

for all $T \in \mathcal{B}(M, \mathbb{Z}_2)$ and all $t \in [0, 1]$. Therefore, the map $\Psi: X \rightarrow \mathcal{H}(M, \mathbb{Z}_2)$ given by $\Psi(x) = \theta(\Phi(x), 1)$ is a half-volume p -sweepout with

$$\sup_{x \in X} \text{Area}(\Psi(x)) \leq \sup_{x \in X} \text{Area}(\Phi(x)) + K.$$

Moreover, it is straightforward to check that Ψ has no concentration of mass. Indeed, for any $q \in M$ and $r > 0$, we have

$$\sup_{x \in X} \text{Area}(\Psi(x) \llcorner B(q, r)) \leq \sup_{x \in X} \text{Area}(\Phi(x) \llcorner B(q, r)) + \sup_{s \in [0, \hbar]} \text{Area}(\partial B_s \llcorner B(q, r)),$$

where B_s is defined as in Proposition 4. Since Φ has no concentration of mass, and the level sets of a Morse function also have no concentration of mass, it follows that Ψ has no concentration of mass. This proves that $\tilde{\omega}_p \leq \omega_p + K + 1$. \square

We are now able to prove Theorem 1.

Theorem 1 *The Weyl law holds for the half-volume spectrum. In other words, we have $\tilde{\omega}_p \sim a_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}$ as $p \rightarrow \infty$.*

Proof This follows from Proposition 18 and Proposition 19. Indeed, we have

$$\omega_p \leq \tilde{\omega}_p \leq \omega_p + K + 1.$$

Theorem 17 implies that

$$\lim_{p \rightarrow \infty} \frac{\omega_p}{a_n \operatorname{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1 \quad \text{and} \quad \lim_{p \rightarrow \infty} \frac{\omega_p + K + 1}{a_n \operatorname{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1,$$

and it follows that

$$\lim_{p \rightarrow \infty} \frac{\tilde{\omega}_p}{a_n \operatorname{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1$$

as well. \square

3 The phase transition half-volume spectrum

There is also an analogous half-volume spectrum in the Allen–Cahn setting. Let $W: \mathbb{R} \rightarrow \mathbb{R}$ be an even double-well potential. This means that

- (i) W is smooth and non-negative,
- (ii) $W(x) = W(-x)$ for all $x \in \mathbb{R}$,
- (iii) W has non-degenerate minima $W(\pm 1) = 0$,
- (iv) W has a non-degenerate maximum $W(0) > 0$,
- (v) W is increasing on $(-1, 0)$ and $(1, \infty)$ and decreasing on $(0, 1)$ and $(-\infty, -1)$,
- (vi) there are constants $\kappa > 0$ and $\alpha \in (0, 1)$ such that $W''(x) \geq \kappa$ for all $|x| \geq \alpha$.

Define the constant

$$\sigma = \int_{-1}^1 \sqrt{W(s)/2} \, ds.$$

Let $u: M \rightarrow \mathbb{R}$ be an $W^{1,2}$ function. For $\varepsilon > 0$ define the Allen–Cahn energy

$$E_\varepsilon(u) = \int_M \frac{\varepsilon}{2} |\nabla u|^2 + \frac{W(u)}{\varepsilon}.$$

In [9], Gaspar and Guaraco define a phase-transition spectrum associated to M via the Allen–Cahn energy.

In order to state the definition of the spectrum, we shall need some further background. A paracompact topological space X is called a \mathbb{Z}_2 -space if it admits a free \mathbb{Z}_2 -action. Given such a space, there is always a quotient space $T = X/\mathbb{Z}_2$ and the natural map $X \rightarrow T$ is a principal \mathbb{Z}_2 -bundle. Any such bundle arises as a pullback of the universal bundle $S^\infty \rightarrow \mathbb{R}\mathbb{P}^\infty$. More precisely, there is a classifying map $f: T \rightarrow \mathbb{R}\mathbb{P}^\infty$ such that $X \rightarrow T$ is the pullback of $S^\infty \rightarrow \mathbb{R}\mathbb{P}^\infty$ via f . The Alexander–Spanier cohomology ring of $\mathbb{R}\mathbb{P}^\infty$ with \mathbb{Z}_2 coefficients is $\mathbb{Z}_2[\mu]$ where the generator μ is in degree one. The map f is unique up to homotopy, and therefore the cohomology classes $f^* \mu^p$ are well-defined in the Alexander–Spanier cohomology ring $H^*(T, \mathbb{Z}_2)$. The \mathbb{Z}_2 -index of X is defined to be the largest p such that $f^* \mu^{p-1} \neq 0$ in $H^*(T, \mathbb{Z}_2)$. A subspace A of X is called invariant if it is closed under the \mathbb{Z}_2 -action.

The \mathbb{Z}_2 -index enjoys the following useful properties. See Fadell and Rabinowitz [8] for more details.

- (i) (Monotonicity) If X_1 and X_2 are \mathbb{Z}_2 -spaces and there is a continuous equivariant map $X_1 \rightarrow X_2$ then $\operatorname{ind}_{\mathbb{Z}_2}(X_1) \leq \operatorname{ind}_{\mathbb{Z}_2}(X_2)$.
- (ii) (Subadditivity) If X is a \mathbb{Z}_2 -space and A_1 and A_2 are closed, invariant subsets with $A_1 \cup A_2 = X$ then $\operatorname{ind}_{\mathbb{Z}_2}(X) \leq \operatorname{ind}_{\mathbb{Z}_2}(A_1) + \operatorname{ind}_{\mathbb{Z}_2}(A_2)$.

(iii) (Continuity) If X is a \mathbb{Z}_2 -space and A is a closed, invariant subset of X then there is an invariant neighborhood V of A in X such that $\text{ind}_{\mathbb{Z}_2}(A) = \text{ind}_{\mathbb{Z}_2}(\overline{V})$.

The space $W^{1,2}(M) \setminus \{0\}$ is paracompact since it is a metric space. Moreover, it admits a natural \mathbb{Z}_2 action $u \mapsto -u$. Note that E_ε respects this action since $E_\varepsilon(u) = E_\varepsilon(-u)$. This uses the fact that W is even. A set $A \subset W^{1,2}(M) \setminus \{0\}$ is called invariant provided $u \in A$ if and only if $-u \in A$. The \mathbb{Z}_2 -action on $W^{1,2}(M) \setminus \{0\}$ descends to any such A . Define the families

$$\mathcal{F}_p = \{A \subset W^{1,2}(M) \setminus \{0\} : A \text{ is compact and invariant with } \text{ind}_{\mathbb{Z}_2}(A) \geq p + 1\}.$$

Gaspar and Guaraco define the min-max values

$$c(\varepsilon, p) = \frac{1}{2\sigma} \inf_{A \in \mathcal{F}_p} \left[\sup_{u \in A} E_\varepsilon(u) \right].$$

Then they set $c(p) = \liminf_{\varepsilon \rightarrow 0} c(\varepsilon, p)$. The sequence $\{c(p)\}_{p \in \mathbb{N}}$ is the phase transition spectrum of M .

Gaspar and Guaraco [10] showed that the Weyl law also holds for the phase transition spectrum.

Theorem 20 (Gaspar and Guaraco) *There is a universal constant τ_n such that $c(p) \sim \tau_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}$ as $p \rightarrow \infty$.*

Dey [7] proved that $\omega_p = c(p)$ for all $p \in \mathbb{N}$. In particular, this implies that the constant τ_n is equal to the constant a_n .

Remark 21 Gaspar and Guaraco do not include the normalization constant $\frac{1}{2\sigma}$ in the definition of $c(\varepsilon, p)$ and $c(p)$. We have chosen to include it so that one has $\omega_p = c(p)$.

It is also possible to define a half-volume spectrum in the phase transition setting. Define

$$Y = \{u \in W^{1,2}(M) : \int_M u = 0\}.$$

Note that Y is a closed subspace of $W^{1,2}(M)$ and so Y is also a Hilbert space. We can run essentially the same construction using Y in place of $W^{1,2}(M)$. For each $p \in \mathbb{N}$, define

$$\mathcal{G}_p = \{A \subset Y \setminus \{0\} : A \text{ is compact and invariant with } \text{ind}_{\mathbb{Z}_2}(A) \geq p + 1\},$$

and then set

$$\tilde{c}(\varepsilon, p) = \frac{1}{2\sigma} \inf_{A \in \mathcal{G}_p} \left[\sup_{u \in A} E_\varepsilon(u) \right].$$

Taking the limit as $\varepsilon \rightarrow 0$ gives the phase-transition half volume spectrum.

Definition 22 For each $p \in \mathbb{N}$, let $\tilde{c}(p) = \liminf_{\varepsilon \rightarrow 0} \tilde{c}(\varepsilon, p)$. The phase transition half volume spectrum of M is the sequence $\{\tilde{c}(p)\}_{p \in \mathbb{N}}$.

Proposition 23 *The phase transition half-volume spectrum satisfies $c(p) \leq \tilde{c}(p)$ for all $p \in \mathbb{N}$.*

Proof Note that $\mathcal{G}_p \subset \mathcal{F}_p$ for every $p \in \mathbb{N}$. Therefore, for every $\varepsilon > 0$, it holds that $c(\varepsilon, p) \leq \tilde{c}(\varepsilon, p)$. The result then follows by sending $\varepsilon \rightarrow 0$. \square

Proposition 24 *The phase transition half-volume spectrum satisfies $\tilde{c}(p) \leq c(p+1)$ for all $p \in \mathbb{N}$.*

Proof Fix an $\varepsilon > 0$. Select a set $A \in \mathcal{F}_{p+1}$ with

$$\sup_{u \in A} E_\varepsilon(u) \leq 2\sigma [c(\varepsilon, p+1) + \varepsilon].$$

Define the set $B = \{u \in A : \int_M u = 0\}$ and note that B is closed and invariant. We claim that $\text{ind}_{\mathbb{Z}_2}(B) \geq p+1$ so that $B \in \mathcal{G}_p$. Given this, we obtain that $\tilde{c}(\varepsilon, p) \leq c(\varepsilon, p+1) + \varepsilon$, and the result follows upon sending $\varepsilon \rightarrow 0$.

It remains to prove the claim. By the continuity of the index, there is a neighborhood V of B in A such that $\text{ind}_{\mathbb{Z}_2}(B) = \text{ind}_{\mathbb{Z}_2}(\overline{V})$. There is an $\eta > 0$ such that

$$\{u \in A : -\eta < \int_M u < \eta\} \subset V.$$

Indeed, if not, then there is a sequence u_k in $A \setminus V$ with $\int_M u_k \rightarrow 0$. Since $A \setminus V$ is compact, we can find a subsequence u_{k_j} that converges to a limit u in $A \setminus V$. But u satisfies $\int_M u = 0$ and therefore $u \in B \subset V$ and this is a contradiction. Therefore, such an η exists.

Let $K = \{u \in A : |\int_M u| \geq \frac{\eta}{2}\}$. Then K is a closed invariant subset of A and $K \cup \overline{V} = A$. Define a map $K \rightarrow S^0$ by sending u to 1 if $\int_M u > 0$ and sending u to -1 if $\int_M u < 0$. This map is continuous and equivariant and so by the monotonicity of the index we have

$$\text{ind}_{\mathbb{Z}_2}(K) \leq \text{ind}_{\mathbb{Z}_2}(S^0) = 1.$$

Hence by the subadditivity of the index, we get

$$p+2 \leq \text{ind}_{\mathbb{Z}_2}(A) \leq \text{ind}_{\mathbb{Z}_2}(K) + \text{ind}_{\mathbb{Z}_2}(\overline{V}) \leq \text{ind}_{\mathbb{Z}_2}(\overline{V}) + 1.$$

This implies that $\text{ind}_{\mathbb{Z}_2}(B) = \text{ind}_{\mathbb{Z}_2}(\overline{V}) \geq p+1$, and so $B \in \mathcal{G}_p$ as needed. \square

Remark 25 The upper bound $\tilde{c}(p) \leq c(p+1)$ proven above is more natural than the upper bound $\tilde{\omega}_p \leq \omega_p + K + 1$ proven in Proposition 19 in the Almgren-Pitts setting. We believe it should also be true that $\tilde{\omega}_p \leq \omega_{p+1}$. In other words, every $(p+1)$ -sweepout contains a half-volume p -sweepout. In principle, it should be possible to prove this by using the relative cup product to adapt the above argument. However, there are technical difficulties in this approach, relating to the fact that restriction is not well-defined on cycles; c.f. [17, Lemma 2.15] where Liokumovich, Marques, and Neves encountered and overcame a similar difficulty. We found it simpler to argue as in Proposition 19, rather than adapting [17, Lemma 2.15] to the half-volume setting.

We can now prove Theorem 2.

Theorem 2 *The phase transition half-volume spectrum satisfies the Weyl law. In other words, we have $\tilde{c}(p) \sim \tau_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}$ as $p \rightarrow \infty$.*

Proof By Propositions 23 and 24 we have

$$c(p) \leq \tilde{c}(p) \leq c(p+1).$$

By the Weyl law for the phase-transition spectrum, we have

$$\lim_{p \rightarrow \infty} \frac{c(p)}{\tau_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1 \quad \text{and} \quad \lim_{p \rightarrow \infty} \frac{c(p+1)}{\tau_n \text{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1$$

and therefore

$$\lim_{p \rightarrow \infty} \frac{\tilde{c}(p)}{\tau_n \operatorname{Vol}(M)^{n/(n+1)} p^{1/(n+1)}} = 1$$

as well. \square

4 Surfaces associated to the half-volume spectrum

In this section, we use the Allen–Cahn min–max theory to construct surfaces associated to the phase transition half-volume spectrum. The goal is to prove Theorem 3. Fix a closed Riemannian manifold (M^{n+1}, g) with $3 \leq n+1 \leq 7$. Fix a number $p \in \mathbb{N}$. In this section, we require the following additional hypothesis on the double-well potential W .

(vii) There are constants $0 < C_1 < C_2$ and $\beta > 1$ and $2 < q < \frac{11}{5}$ such that

$$C_1|x|^q \leq W(x) \leq C_2|x|^q \quad \text{and} \quad C_1|x|^{q-1} \leq |W'(x)| \leq C_2|x|^{q-1}$$

for all $|x| \geq \beta$.

The first step of the proof is to construct, for each small enough $\varepsilon > 0$, a critical point u_ε of E_ε subject to the volume constraint

$$\int_M u_\varepsilon = 0.$$

Given such a u_ε , there is a Lagrange multiplier $\lambda_\varepsilon \in \mathbb{R}$ such that u_ε is a critical point of

$$F_{\varepsilon, \lambda_\varepsilon}(v) = E_\varepsilon(v) + \lambda_\varepsilon \int_M v$$

on all of $W^{1,2}$. The construction of u_ε is similar to that of Gaspar and Guaraco [9] in the unconstrained case.

Remark 26 The assumption (vii) is not new, and has been considered before by Hutchinson and Tonegawa (c.f. [15, Section 6.1]). For us, there are two purposes for imposing the growth condition (vii):

- First, the bound $|W'| \leq C_2|x|^{q-1}$, is used to verify the Palais–Smale condition with the volume constraint. In the unconstrained case, one has

$$E_\varepsilon(\max(\min(u, 1), -1)) \leq E_\varepsilon(u)$$

and so by a truncation argument it is enough to verify the Palais–Smale condition along Palais–Smale sequences which are bounded in L^∞ . See [9] for more details. However, truncation may not preserve the volume constraint. In the volume constrained case, we instead rely on the estimate $|W'| \leq C_2|x|^{q-1}$ to show that $W^{1,2}$ bounds on u imply L^2 bounds on $W'(u)$. Then we show that these L^2 bounds on $W'(u)$ still suffice to check the Palais–Smale condition.

- Second, assumption (vii) is needed to get uniform L^∞ bounds on critical points of $F_{\varepsilon, \lambda}$, and to get uniform bounds on the Lagrange multipliers λ . See Proposition 32 in the Appendix, which is based on [15, Section 6.1] and work of X. Chen [4]. More precisely, the upper bound $|W'| \leq C_2|x|^{q-1}$ is used in the elliptic estimates at the beginning of the proof of Proposition 32, the lower bound $|W| \geq C_1|x|^q$ is used in Lemma 34 which is

needed to bound the Lagrange multiplier, and the lower bound $|W'| \geq C_1|x|^{q-1}$ is used at the end of the proof of Proposition 32 to obtain uniform L^∞ bounds on the critical points.

The upper bound $|W| \leq C_2|x|^q$ is not explicitly used, but follows from $|W'| \leq C_2|x|^{q-1}$. Also, it is not essential that the exponent in the upper and lower bounds are the same, but we have chosen to assume this for simplicity.

We recall (see Proposition 4.4 in [13]) that the first variation of E_ε is given by

$$DE_\varepsilon(u)(v) = \int_M \frac{\varepsilon}{2} \nabla u \cdot \nabla v + \frac{W'(u)}{\varepsilon} v.$$

Fix a number $\varepsilon > 0$. A sequence A_k in \mathcal{G}_p is called a critical sequence if

$$\lim_{k \rightarrow \infty} \left[\sup_{u \in A_k} E_\varepsilon(u) \right] = 2\sigma\tilde{c}(\varepsilon, p).$$

A sequence $u_k \in A_k$ is called a min-max sequence provided $\lim_{k \rightarrow \infty} E_\varepsilon(u_k) = 2\sigma\tilde{c}(\varepsilon, p)$.

In the unconstrained case, it is not necessarily true that every min-max sequence is bounded in $W^{1,2}$. However, one can obtain the existence of a bounded min-max sequence via a truncation argument. See, for example, the remarks before Proposition 4.5 in [13]. We cannot employ truncation because it doesn't preserve the volume constraint. Fortunately, in the volume constrained case, every min-max sequence is automatically bounded in $W^{1,2}$.

Proposition 27 *Any min-max sequence u_k is uniformly bounded in $W^{1,2}(M)$.*

Proof Assume that $u \in Y$ satisfies $E_\varepsilon(u) \leq K$. Since $W \geq 0$, it follows immediately that

$$\int_M |\nabla u|^2 \leq \frac{2K}{\varepsilon}.$$

Since u has average 0, the Poincaré inequality implies that $\|u\|_{W^{1,2}} \leq CK/\varepsilon$. This proves the result. \square

Proposition 28 *Assume that u_k is a sequence uniformly bounded in $W^{1,2}$. Then $W'(u_k)$ is uniformly bounded in L^2 .*

Proof By assumption the sequence u_k is uniformly bounded in $W^{1,2}$. As $3 \leq n+1 \leq 7$, the Sobolev embedding theorem implies that u_k is uniformly bounded in $L^{12/5}$. Now $|W'(u_k)| \leq C|u_k|^{q-1} \leq C|u_k|^{6/5}$ whenever $|u_k| \geq \beta$. Therefore

$$\begin{aligned} \int_M W'(u_k)^2 &= \int_{|u_k| \leq \beta} W'(u_k)^2 + \int_{|u_k| > \beta} W'(u_k)^2 \\ &\leq C \operatorname{Vol}(M) + C \int_M |u_k|^{12/5}, \end{aligned}$$

and it follows that $W'(u_k)$ is uniformly bounded in L^2 . \square

The functional $E_\varepsilon|_Y$ satisfies the Palais-Smale condition. See [13] Proposition 4.4 for the proof without a volume constraint.

Proposition 29 *The functional $E_\varepsilon|_Y$ satisfies the Palais-Smale condition. More precisely, assume that u_k is a bounded sequence in Y and that $\|DE_\varepsilon|_Y(u_k)\| \rightarrow 0$. Then a subsequence of u_k converges strongly to a limit $u \in Y$.*

Proof Assume that u_k is a bounded sequence in Y such that $\|DE_\varepsilon|_Y(u_k)\| \rightarrow 0$. We need to show that some subsequence of u_k converges strongly to a limit $u \in Y$. Note that Y is closed and convex in $W^{1,2}(M)$ and so Y is weakly closed. Thus, passing to a subsequence, we can assume that u_k converges weakly in $W^{1,2}$ and strongly in $L^{12/5}$ to a point $u \in Y$.

Observe that

$$DE_\varepsilon|_Y(u)(u_k - u) = \int_M \varepsilon \nabla u \cdot \nabla(u_k - u) + \int_M \frac{W'(u)}{\varepsilon}(u_k - u).$$

The first term on the right hand side goes to 0 by the weak convergence $u_k \rightharpoonup u$. Note that $W'(u) \in L^2$ since $u \in L^{12/5}$. Therefore the second term on the right hand side also goes to 0 since $u_k \rightarrow u$ in L^2 . Thus we obtain

$$DE_\varepsilon|_Y(u)(u_k - u) \rightarrow 0, \quad \text{as } k \rightarrow \infty.$$

Also note that $DE_\varepsilon|_Y(u_k)(u_k - u) \rightarrow 0$ since $\|DE_\varepsilon|_Y(u_k)\| \rightarrow 0$ and $u_k - u$ is uniformly bounded in $W^{1,2}$. On the other hand,

$$DE_\varepsilon|_Y(u_k)(u_k - u) = \int_M \varepsilon \nabla u_k \cdot \nabla(u_k - u) + \int_M \frac{W'(u_k)}{\varepsilon}(u_k - u).$$

The second term on the right hand side goes to 0 as $W'(u_k)$ is uniformly bounded in L^2 and $u_k - u \rightarrow 0$ in L^2 .

Now observe that

$$\begin{aligned} & DE_\varepsilon|_Y(u_k)(u_k - u) - DE_\varepsilon|_Y(u)(u_k - u) \\ &= \int_M \varepsilon |\nabla u_k - \nabla u|^2 + \int_M \frac{W'(u_k)}{\varepsilon}(u_k - u) - \int_M \frac{W'(u)}{\varepsilon}(u_k - u). \end{aligned}$$

We have already seen that every term in this formula goes to 0 except $\int_M \varepsilon |\nabla u_k - \nabla u|^2$, and therefore $\int_M \varepsilon |\nabla u_k - \nabla u|^2$ goes to 0 as well. This proves that $u_k \rightarrow u$ strongly in $W^{1,2}$, as needed. \square

According to Gaspar and Guaraco [9], for each given p , we have $2\sigma c(\varepsilon, p+1) < E_\varepsilon(0)$ provided ε is small enough. Therefore, we also have $2\sigma \tilde{c}(\varepsilon, p) < E_\varepsilon(0)$ provided ε is small enough. Hence, for ε small enough, any min-max sequence remains bounded away from 0. By the classical theory for functionals satisfying the Palais-Smale condition (see [24]), we get the following existence result for critical points of $E_\varepsilon|_Y$. See Theorem 3.3 in [9] for the corresponding result in the unconstrained case.

Proposition 30 Fix $p \in \mathbb{N}$. For all small enough ε , there is a critical point $u_\varepsilon \in Y$ of $E_\varepsilon|_Y$ with $E_\varepsilon(u_\varepsilon) = \tilde{c}(\varepsilon, p)$. There is a number $\lambda_\varepsilon \in \mathbb{R}$ such that u_ε is a critical point of

$$v \mapsto F_{\varepsilon, \lambda_\varepsilon}(v) = E_\varepsilon(v) + \lambda_\varepsilon \int_M v$$

on all of $W^{1,2}$, and u_ε satisfies the PDE

$$-\varepsilon \Delta u_\varepsilon + \frac{W'(u_\varepsilon)}{\varepsilon} = \lambda_\varepsilon$$

in the weak sense. Moreover, we have $\int_M u_\varepsilon = 0$. The index of u_ε as a critical point of $E_\varepsilon|_Y$ is at most p .

Given the existence of u_ε , the second step in the proof is to study the convergence of u_ε as $\varepsilon \rightarrow 0$. Fortunately for us, Bellettini and Wickramasekera [3] have already studied the regularity of such limits. Let us recall the setup in [3]. For each $\varepsilon > 0$, suppose u_ε is a critical point of $F_{\varepsilon, \lambda_\varepsilon}$ and that $E_\varepsilon(u_\varepsilon)$ and $\text{ind}_{F_{\varepsilon, \lambda_\varepsilon}}(u_\varepsilon)$ and $\|u_\varepsilon\|_{L^\infty}$ are uniformly bounded. Choose a sequence $\varepsilon_j \rightarrow 0$ and assume that $\lambda_{\varepsilon_j} \rightarrow \lambda$. Passing to a subsequence if necessary, there exist a radon measure μ on M and a function $u_\infty \in BV(M)$ such that

$$\frac{1}{2\sigma} \left(\frac{\varepsilon_j}{2} |\nabla u_{\varepsilon_j}|^2 + \frac{W(u_{\varepsilon_j})}{\varepsilon_j} \right) \rightharpoonup \mu$$

and $u_{\varepsilon_j} \rightarrow u_\infty$ in L^1 . Moreover, u_∞ takes only the values ± 1 . Hutchinson and Tonegawa [15] proved that there is an integral varifold V on M such that $\|V\| = \mu$. The following is the special case of Theorem 4.1 in Bellettini and Wickramasekera [3] where the prescription functions are assumed to be constants and the ambient dimension is assumed to be between 3 and 7.

Theorem 31 (See Theorem 4.1 in [3]) *Let (M^{n+1}, g) be a closed Riemannian manifold with $3 \leq n+1 \leq 7$. Assume (u_{ε_j}) is a sequence as above, and assume that $\lambda > 0$. Let $\Omega = \text{int}(\{x \in M : u_\infty(x) = 1\})$. Then $V = V_0 + V_\lambda$ where*

- (i) *V_0 is induced by a collection of smooth, disjoint minimal surfaces equipped with even multiplicities. Moreover, $\text{spt}(\|V_0\|) \subset M \setminus \Omega$.*
- (ii) *If $\Omega = \emptyset$ then $V_\lambda = 0$. If $\Omega \neq \emptyset$ then $V_\lambda = |\partial^* \Omega| \neq 0$, and moreover, V_λ is induced by a smooth surface with constant mean curvature λ whose mean curvature vector points into Ω .*

The minimal surfaces may have tangential intersection with the CMC surface. Likewise, the CMC surface may have tangential intersections with itself but it never crosses itself.

We can now complete the proof of Theorem 3.

Theorem 3 *Let (M^{n+1}, g) be a closed Riemannian manifold with $3 \leq n+1 \leq 7$. Fix a number $p \in \mathbb{N}$. There are*

- (i) *a Caccioppoli set $\Omega \subset M$ with $\text{Vol}(\Omega) = \frac{1}{2} \text{Vol}(M)$ whose boundary is smooth and almost-embedded with constant mean curvature,*
- (ii) *a (possibly empty) collection of smooth, disjoint minimal surfaces $\Sigma_1, \dots, \Sigma_k \subset M \setminus \Omega$,*
- (iii) *and positive integers $\theta_0 \in \mathbb{Z}$ and $\theta_1, \dots, \theta_k \in 2\mathbb{Z}$*

such that $\tilde{c}(p) = \theta_0 \text{Area}(\partial\Omega) + \theta_1 \text{Area}(\Sigma_1) + \dots + \theta_k \text{Area}(\Sigma_k)$. Moreover, $\theta_0 = 1$ unless $\partial\Omega$ is also a minimal surface.

Proof Choose a sequence $\varepsilon_k \rightarrow 0$ so that $\tilde{c}(\varepsilon_k, p) \rightarrow \tilde{c}(p)$. Let u_{ε_k} be the critical points of $F_{\varepsilon_k, \lambda_{\varepsilon_k}}$ constructed in Proposition 30. Then $E_{\varepsilon_k}(u_{\varepsilon_k})$ and $\text{ind}_{F_{\varepsilon_k, \lambda_{\varepsilon_k}}}(u_{\varepsilon_k})$ are uniformly bounded. According to Hutchinson and Tonegawa [15] Section 6.1 and Lemma 3.4 in [4], the Lagrange multipliers λ_{ε_k} are also uniformly bounded and there is a constant $K > 0$ such that $\|u_{\varepsilon_k}\|_{L^\infty} \leq K$ for all k . For the interested reader, we include the details of the argument in the appendix. Applying Theorem 31 to the sequence u_{ε_k} now yields the result. \square

5 Appendix

The goal of the appendix is to prove the following proposition. We largely follow the sketch in [15, Section 6.1], giving details as appropriate. The proof that the Lagrange multipliers are bounded depends on work of Chen [4].

Proposition 32 *Assume that the potential W satisfies the growth condition (vii). Choose a sequence $\varepsilon_k \rightarrow 0$ and let u_{ε_k} be a critical point of $F_{\varepsilon_k, \lambda_{\varepsilon_k}}$. Assume that $E_{\varepsilon_k}(u_{\varepsilon_k})$ is uniformly bounded. Then the Lagrange multipliers λ_{ε_k} are uniformly bounded and $\|u_{\varepsilon_k}\|_{L^\infty}$ is also uniformly bounded.*

Proof The first step is to check that each u_{ε_k} is smooth. Recall that $|W'(u_{\varepsilon_k})| \leq C|u_{\varepsilon_k}|^{q-1}$ for $|u_{\varepsilon_k}| \geq \beta$. For simplicity, we will give the argument assuming $3 \leq n+1 \leq 5$. In this case, by the Sobolev embedding theorem, u_{ε_k} belongs to $L^{10/3}$. It follows that $W'(u_{\varepsilon_k}) \in L^{10/(3q-3)}$. Hence by elliptic regularity we have $u_{\varepsilon_k} \in W^{2,q_1}$ for $q_1 = 10/(3q-3)$. Note that

$$\frac{n+1}{q_1} \leq \frac{3}{2}(q-1) \leq \frac{18}{10} < 2.$$

Thus by the Sobolev inequalities we obtain that u_{ε_k} is Hölder continuous. Standard elliptic regularity then implies that u_{ε_k} is smooth. The cases $6 \leq n+1 \leq 7$ are handled similarly. One applies the Sobolev inequalities together with elliptic regularity several times. Each application improves the regularity of u_{ε_k} until eventually one obtains $u_{\varepsilon_k} \in W^{2,q_1}$ with $2 > (n+1)/q$. This gives Hölder continuity of u_{ε_k} , and standard elliptic regularity then implies that u_{ε_k} is smooth. Note at this point we do not have uniform L^∞ estimates on u_{ε_k} .

To prove that the Lagrange multipliers are bounded we follow [4] Lemma 3.4. Note that [4] Lemma 3.4 is proved for domains in Euclidean space. Some additional difficulties arise in adapting the mollifier arguments used in [4] to the case of a closed manifold. We need to prove a sequence of lemmas. In what follows, C denotes a positive constant that is allowed to change from line to line.

Lemma 33 *We have*

$$\int_M (|u_{\varepsilon_k}| - 1)^2 \rightarrow 0, \quad \text{as } k \rightarrow \infty.$$

Proof Define the sets

$$A_1 = \{|u_{\varepsilon_k}| \geq \alpha\}, \quad A_2 = \{|u_{\varepsilon_k}| < \alpha\},$$

where α is the constant appearing in assumption (vi) for the potential W . Then since $W(\pm 1) = 0$ and $W'(\pm 1) = 0$ and $W''(t) \geq \kappa$ for $|t| \geq \alpha$, it follows that

$$\frac{\kappa}{2} (|u_{\varepsilon_k}| - 1)^2 \leq W(u_{\varepsilon_k})$$

on A_1 . This gives the estimate

$$\int_{A_1} (|u_{\varepsilon_k}| - 1)^2 \leq \frac{2\varepsilon_k}{\kappa} E_{\varepsilon_k}(u_{\varepsilon_k}).$$

On the other hand, since $W(u_{\varepsilon_k}) \geq W(\alpha)$ on A_2 , it follows that

$$\int_{A_2} (|u_{\varepsilon_k}| - 1)^2 \leq \int_{A_2} 1 \leq \frac{1}{W(\alpha)} \int_{A_2} W(u_{\varepsilon_k}) \leq \frac{\varepsilon_k}{W(\alpha)} E_{\varepsilon_k}(u_{\varepsilon_k}).$$

The lemma now follows by combining the previous two estimates, and using the fact that the energy is uniformly bounded. \square

Define

$$\Phi(s) = \int_0^s \sqrt{W(t)/2} dt.$$

Let $w_{\varepsilon_k} = \Phi \circ u_{\varepsilon_k}$. It is easy to check that ∇w_{ε_k} is uniformly bounded in L^1 (see [15]).

Lemma 34 *There is a constant $C > 0$ such that $|x - y|^2 \leq C|\Phi(x) - \Phi(y)|$ for all $x, y \in \mathbb{R}$.*

Proof We check a number of cases. First suppose that $x, y \geq \beta$ and assume without loss of generality that $x \geq y$. Then

$$\begin{aligned} |\Phi(x) - \Phi(y)| &= \int_y^x \sqrt{W(s)/2} ds \geq \frac{1}{\sqrt{2}} \int_y^x s^{q/2} ds \\ &\geq \frac{1}{\sqrt{2}} \int_y^x s ds = \frac{1}{2\sqrt{2}}(x^2 - y^2) \geq \frac{1}{2\sqrt{2}}|x - y|^2, \end{aligned}$$

where the last inequality uses the fact that $x \geq y \geq 0$ and so $x + y \geq x - y$. The same argument works in the case where $x, y \leq -\beta$.

Now suppose that $x \geq \beta$ and $y \leq -\beta$. Let $a = \int_{-\beta}^{\beta} \sqrt{W(s)/2} ds > 0$. Then

$$\begin{aligned} |\Phi(x) - \Phi(y)| &= \int_{\beta}^x \sqrt{W(s)/2} ds + a + \int_y^{-\beta} \sqrt{W(s)/2} ds \\ &\geq \frac{1}{\sqrt{2}} \int_{\beta}^x s ds + a + \frac{1}{\sqrt{2}} \int_y^{-\beta} (-s) ds \\ &\geq \frac{1}{2\sqrt{2}}(x^2 - \beta^2) + a + \frac{1}{2\sqrt{2}}(y^2 - \beta^2). \end{aligned}$$

We have

$$|x - y|^2 \leq 2(x^2 + y^2) \leq 4\beta^2 + 2(x^2 - \beta^2) + 2(y^2 - \beta^2) \leq C|\Phi(x) - \Phi(y)|.$$

The same argument works if $x \leq -\beta$ and $y \geq \beta$.

It remains to handle the case when $-\beta \leq x, y \leq \beta$. Assume for contradiction there are two sequences $x_j, y_j \in [-\beta, \beta]$ such that

$$|x_j - y_j|^2 > j|\Phi(x_j) - \Phi(y_j)|. \quad (1)$$

Passing to a subsequence if necessary, we can assume that $x_j \rightarrow x \in [-\beta, \beta]$ and $y_j \rightarrow y \in [-\beta, \beta]$. It is clear from (1) that $x = y$. Passing to the limit in

$$|x_j - y_j| \geq j \frac{|\Phi(x_j) - \Phi(y_j)|}{|x_j - y_j|}$$

we obtain that $\Phi'(x) = 0$. Thus $W(x) = 0$ and so $x = \pm 1$. Without loss of generality assume that $x = 1$. Note that there is a constant $C > 0$ such that $W(s) \geq C|s - 1|^2$ for s close to 1. Thus

$$|\Phi(x_j) - \Phi(y_j)| = \left| \int_{y_j}^{x_j} \sqrt{W(s)/2} ds \right| \geq C \left| \int_{y_j}^{x_j} |s - 1| ds \right|.$$

If $x_j \geq y_j \geq 1$ then

$$\begin{aligned} \left| \int_{y_j}^{x_j} |s - 1| ds \right| &= \int_{y_j}^{x_j} s - 1 ds = \frac{1}{2} [(x_j - 1)^2 - (y_j - 1)^2] \\ &= \frac{1}{2} [(x_j + y_j - 2)(x_j - y_j)] \geq \frac{1}{2} (x_j - y_j)^2, \end{aligned}$$

which combined with the previous equation yields a contradiction. If $x_j \geq 1 \geq y_j$ then

$$\begin{aligned} \left| \int_{y_j}^{x_j} |s - 1| ds \right| &= \int_{y_j}^1 1 - s ds + \int_1^{x_j} s - 1 ds \\ &= \frac{1}{2} [(1 - y_j)^2 + (x_j - 1)^2] \geq \frac{1}{4} (x_j - y_j)^2 \end{aligned}$$

since either $x_j - 1 \geq \frac{1}{2}(x_j - y_j)$ or $1 - y_j \geq \frac{1}{2}(x_j - y_j)$. Again this gives a contradiction. The remaining possibilities likewise lead to contradiction and the lemma is proved. \square

For each $\eta \in (0, 1)$, let $u_{\varepsilon_k, \eta}$ be a mollified version of u_{ε_k} . More precisely, choose an isometric embedding of M into \mathbb{R}^m . Let N be a small tubular neighborhood of M where the nearest point projection $\Pi: N \rightarrow M$ is well-defined and a submersion. Let $B_r(x)$ denote the open ball of radius r centered at $x \in \mathbb{R}^m$. Also let B_1 denote the open unit ball centered at the origin in \mathbb{R}^m . Choose a non-negative smooth function $\rho: B_1 \rightarrow \mathbb{R}$ which is compactly supported in B_1 and satisfies

$$\int_{B_1} \rho(y) d\mathcal{L}^m(y) = 1.$$

Extend u_{ε_k} to a function v_{ε_k} on N by setting $v_{\varepsilon_k} = u_{\varepsilon_k} \circ \Pi$. Then let

$$u_{\varepsilon_k, \eta}(x) = \int_{B_1} \rho(y) v_{\varepsilon_k}(x - \eta y) d\mathcal{L}^m(y), \quad x \in M$$

be the mollified version of u_{ε_k} .

Lemma 35 *For each fixed η , there is a uniform bound $\|u_{\varepsilon_k, \eta}\|_{L^\infty(M)} \leq C(\eta)$.*

Proof Observe that

$$\begin{aligned} |u_{\varepsilon_k, \eta}(x)| &\leq \int_{B_1} \rho(y) |v_{\varepsilon_k}(x - \eta y)| d\mathcal{L}^m(y) \\ &\leq 1 + \int_{B_1} \rho(y) \left| |v_{\varepsilon_k}(x - \eta y)| - 1 \right| d\mathcal{L}^m(y) \\ &\leq 1 + C \eta^{-(n+1)} \int_{B_\eta(x)} \left| |v_{\varepsilon_k}(z)| - 1 \right| d\mathcal{L}^m(z). \end{aligned}$$

By the co-area formula we have

$$\begin{aligned}
& \int_{B_\eta(x)} \left| |v_{\varepsilon_k}(z)| - 1 \right| d\mathcal{L}^m(z) \\
& \leq C \int_{B_\eta(x)} \left| |v_{\varepsilon_k}(z)| - 1 \right| J\Pi(z) d\mathcal{L}^m(z) \\
& \leq C \int_{\Pi(B_\eta(x))} \int_{p \in \Pi^{-1}(q)} \left| |v_{\varepsilon_k}(p)| - 1 \right| d\mathcal{H}^{m-n-1}(p) d\mathcal{H}^{n+1}(q) \\
& \leq C \int_{\Pi(B_\eta(x))} \left| |u_{\varepsilon_k}(q)| - 1 \right| d\mathcal{H}^{n+1}(q).
\end{aligned}$$

Inserting this into the previous equation and using Lemma 33 gives the result. \square

Lemma 36 *For each fixed η , there is a uniform bound $\|\nabla u_{\varepsilon_k, \eta}\|_{L^\infty(M)} \leq C(\eta)$.*

Proof Fix an index $i \in \{1, \dots, m\}$. Note the formula for $u_{\varepsilon_k, \eta}(x)$ also makes sense for $x \in N$ so we can regard $u_{\varepsilon_k, \eta}$ as a function defined on N . For $x \in M$ we have

$$\partial_i u_{\varepsilon_k, \eta}(x) = \int_{B_1} \partial_i \rho(y) v_{\varepsilon_k}(x - \eta y) d\mathcal{L}^m(y).$$

Thus we have

$$\begin{aligned}
|\partial_i u_{\varepsilon_k, \eta}(x)| & \leq \int_{B_1} |\partial_i \rho(y)| |v_{\varepsilon_k}(x - \eta y)| d\mathcal{L}^m(y) \\
& \leq C + \int_{B_1} |\partial_i \rho(y)| (|v_{\varepsilon_k}(x - \eta y)| - 1) d\mathcal{L}^m(y) \\
& \leq C + C\eta^{-(n+1)} \int_{B_\eta(x)} \left| |v_{\varepsilon_k}(z)| - 1 \right| d\mathcal{L}^m(z).
\end{aligned}$$

Using the coarea formula as in the proof of the previous lemma now gives the result. \square

Lemma 37 *There is a uniform bound $\|u_{\varepsilon_k, \eta} - u_{\varepsilon_k}\|_{L^2(M)}^2 \leq C\eta$ for $k \geq K(\eta)$.*

Proof Observe that

$$\begin{aligned}
\int_M |u_{\varepsilon_k, \eta} - u_{\varepsilon_k}|^2 & = \int_M \left| \int_{B_1} \rho(y) v_{\varepsilon_k}(x - \eta y) d\mathcal{L}^m(y) - u_{\varepsilon_k}(x) \right|^2 d\mathcal{H}^{n+1}(x) \\
& \leq \int_M \int_{B_1} \rho(y) |v_{\varepsilon_k}(x - \eta y) - u_{\varepsilon_k}(x)|^2 d\mathcal{L}^m(y) d\mathcal{H}^{n+1}(x) \\
& \leq C \int_M \int_{B_1} \rho(y) |f(x - \eta y) - f(x)| d\mathcal{L}^m(y) d\mathcal{H}^{n+1}(x),
\end{aligned}$$

where $f = \Phi \circ u_{\varepsilon_k} \circ \Pi = w_{\varepsilon_k} \circ \Pi$ and we used Lemma 34 to get the last inequality. By Fubini's theorem we get

$$\begin{aligned}
& \int_M \int_{B_1} \rho(y) |f(x - \eta y) - f(x)| d\mathcal{L}^m(y) d\mathcal{H}^{n+1}(x) \\
& \leq \eta \int_M \int_{B_1} \int_0^1 \rho(y) |\nabla f(x - t\eta y)| dt d\mathcal{L}^m(y) d\mathcal{H}^{n+1}(x) \\
& = \eta \int_{B_1} \rho(y) \int_0^1 \int_M |\nabla f(x - t\eta y)| d\mathcal{H}^{n+1}(x) dt d\mathcal{L}^m(y).
\end{aligned}$$

Now for fixed y and t , we have

$$\int_M |\nabla f(x - t\eta y)| d\mathcal{H}^{n+1}(x) = \int_{M-t\eta y} |\nabla f(z)| d\mathcal{H}^{n+1}(z).$$

Provided η is small enough, the map $\Pi: M - t\eta y \rightarrow M$ is a diffeomorphism and so by the change of variables formula

$$\begin{aligned} \int_{M-t\eta y} |\nabla f(z)| d\mathcal{H}^{n+1}(z) &\leq \int_{M-t\eta y} |\nabla w_{\varepsilon_k}(\Pi(z))| d\mathcal{H}^{n+1}(z) \\ &\leq C \int_{M-t\eta y} |\nabla w_{\varepsilon_k}(\Pi(z))| J\Pi(z) d\mathcal{H}^{n+1}(z) \\ &= C \int_M |\nabla w_{\varepsilon_k}(q)| d\mathcal{H}^{n+1}(q). \end{aligned}$$

Thus we obtain

$$\begin{aligned} &\eta \int_{B_1} \rho(y) \int_0^1 \int_M |\nabla f(x - t\eta y)| d\mathcal{H}^{n+1}(x) dt d\mathcal{L}^m(y) \\ &\leq C\eta \int_{B_1} \rho(y) \|w_{\varepsilon_k}\|_{L^1(M)} d\mathcal{L}^m(y) \leq C\eta. \end{aligned}$$

Putting everything together we get

$$\int_M |u_{\varepsilon_k, \eta} - u_{\varepsilon_k}|^2 \leq C\eta,$$

and the result follows. \square

Lemma 38 *Let $\bar{u}_{\varepsilon_k, \eta}$ be the average of $u_{\varepsilon_k, \eta}$. There is a uniform bound $|\bar{u}_{\varepsilon_k, \eta}| \leq C\eta^{\frac{1}{2}}$ for $k \geq K(\eta)$.*

Proof Observe that

$$\bar{u}_{\varepsilon_k, \eta} = \frac{1}{\text{Vol}(M)} \int_M u_{\varepsilon_k, \eta} = \frac{1}{\text{Vol}(M)} \int_M (u_{\varepsilon_k, \eta} - u_{\varepsilon_k}).$$

Now the result follows from Lemma 37 and Hölder's inequality. \square

Now fix an η to be specified later. Given a large integer k , let ψ be the solution to

$$-\Delta\psi = u_{\varepsilon_k, \eta} - \bar{u}_{\varepsilon_k, \eta}, \quad \int_M \psi = 0.$$

By Lemmas 35, 36, and 38 the right hand side of the above PDE is uniformly bounded in C^1 . Therefore by elliptic regularity, ψ is bounded in C^2 by a constant that depends on η but not on k . Note that u_{ε_k} satisfies the PDE

$$-\varepsilon_k \Delta u_{\varepsilon_k} + \frac{W'(u_{\varepsilon_k})}{\varepsilon_k} = \lambda_{\varepsilon_k}.$$

Multiplying by $\nabla\psi \cdot \nabla u_{\varepsilon_k}$ and integrating yields

$$\lambda_{\varepsilon_k} \int_M \nabla\psi \cdot \nabla u_{\varepsilon_k} = \int_M \nabla\psi \cdot \nabla u_{\varepsilon_k} \left(-\varepsilon_k \Delta u_{\varepsilon_k} + \frac{W'(u_{\varepsilon_k})}{\varepsilon_k} \right). \quad (2)$$

Observe that

$$\int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} \frac{W'(u_{\varepsilon_k})}{\varepsilon_k} = \int_M \nabla \psi \cdot \nabla \left(\frac{W(u_{\varepsilon_k})}{\varepsilon_k} \right) = - \int_M \frac{W(u_{\varepsilon_k})}{\varepsilon_k} \Delta \psi.$$

Also, by the integration by parts formula for the Hessian, we have

$$\varepsilon_k \int_M D^2 \psi(\nabla u_{\varepsilon_k}, \nabla u_{\varepsilon_k}) = -\varepsilon_k \int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} \Delta u_{\varepsilon_k} + \frac{\varepsilon_k}{2} \int_M |\nabla u_{\varepsilon_k}|^2 \Delta \psi.$$

Thus we have the following formula for the right hand side of (2):

$$\begin{aligned} & \int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} \left(-\varepsilon_k \Delta u_{\varepsilon_k} + \frac{W'(u_{\varepsilon_k})}{\varepsilon_k} \right) \\ &= \varepsilon_k \int_M D^2 \psi(\nabla u_{\varepsilon_k}, \nabla u_{\varepsilon_k}) - \int_M \left(\frac{\varepsilon_k}{2} |\nabla u_{\varepsilon_k}|^2 + \frac{W(u_{\varepsilon_k})}{\varepsilon_k} \right) \Delta \psi. \end{aligned}$$

Since $\|\psi\|_{C^2} \leq C(\eta)$, this gives a bound

$$\left| \int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} \left(-\varepsilon_k \Delta u_{\varepsilon_k} + \frac{W'(u_{\varepsilon_k})}{\varepsilon_k} \right) \right| \leq C(\eta) E_{\varepsilon_k}(u_{\varepsilon_k}).$$

We now turn attention to the left hand side of (2). Integrating by parts, we have

$$\lambda_{\varepsilon_k} \int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} = -\lambda_{\varepsilon_k} \int_M u_{\varepsilon_k} \Delta \psi = \lambda_{\varepsilon_k} \int_M u_{\varepsilon_k} (u_{\varepsilon_k, \eta} - \bar{u}_{\varepsilon_k, \eta}).$$

Now observe that

$$\int_M u_{\varepsilon_k} (u_{\varepsilon_k, \eta} - \bar{u}_{\varepsilon_k, \eta}) = \int_M u_{\varepsilon_k} (u_{\varepsilon_k, \eta} - u_{\varepsilon_k}) + \int_M (u_{\varepsilon_k}^2 - 1) - \bar{u}_{\varepsilon_k, \eta} \int_M u_{\varepsilon_k} + \text{Vol}(M).$$

By Lemmas 33 and 38 and Hölder's inequality, we can select η small enough that

$$\left| \int_M u_{\varepsilon_k} (u_{\varepsilon_k, \eta} - u_{\varepsilon_k}) \right| \leq \frac{1}{4} \text{Vol}(M), \quad \left| \bar{u}_{\varepsilon_k, \eta} \int_M u_{\varepsilon_k} \right| \leq \frac{1}{4} \text{Vol}(M).$$

By Lemma 33 and Hölder's inequality, we have

$$\left| \int_M (u_{\varepsilon_k}^2 - 1) \right| \leq \int_M ||u_{\varepsilon_k}| - 1| (|u_{\varepsilon_k}| + 1) \leq \frac{1}{4} \text{Vol}(M)$$

for k large enough. It follows that

$$\int_M \nabla \psi \cdot \nabla u_{\varepsilon_k} \geq \frac{1}{4} \text{Vol}(M)$$

for large enough k . Using equation (2) then gives an upper bound on λ_{ε_k} which is independent of k .

To complete the proof of the Proposition 32, it remains to show that $\|u_{\varepsilon_k}\|_{L^\infty}$ is uniformly bounded. Let $M_{\varepsilon_k} = M/\varepsilon_k$. Define the rescaled functions $h_{\varepsilon_k} : M_{\varepsilon_k} \rightarrow \mathbb{R}$ by $h_{\varepsilon_k}(x) = u_{\varepsilon_k}(\varepsilon_k x)$. Then h_{ε_k} solves

$$-\Delta h_{\varepsilon_k} + W'(h_{\varepsilon_k}) = \varepsilon_k \lambda_{\varepsilon_k}.$$

Fix some $p \geq 1$. Multiplying the equation by $|h_{\varepsilon_k}|^{p-1} h_{\varepsilon_k}$ and integrating gives

$$\int_{M_{\varepsilon_k}} p |h_{\varepsilon_k}|^{p-1} |\nabla h_{\varepsilon_k}|^2 + \int_{M_{\varepsilon_k}} W'(h_{\varepsilon_k}) |h_{\varepsilon_k}|^{p-1} h_{\varepsilon_k} = \varepsilon_k \lambda_{\varepsilon_k} \int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^{p-1} h_{\varepsilon_k}. \quad (3)$$

Now for $|x| \geq \beta$ we have $W'(x)|x|^{p-1}x \geq C|x|^{p+q-1}$. Since $\varepsilon_k \lambda_{\varepsilon_k} \rightarrow 0$, it follows from (3) that

$$\int_{\{|h_{\varepsilon_k}| \geq \beta\}} |h_{\varepsilon_k}|^{p+q-1} \leq \frac{1}{2} \int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^p$$

assuming k is large enough. This implies that

$$\int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^{p+q-1} \leq \left(\beta^{q-1} + \frac{1}{2} \right) \int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^p.$$

By induction, for any positive integer r , this gives

$$\int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^{2+r(q-1)} \leq \left(\beta^{q-1} + \frac{1}{2} \right)^r \int_{M_{\varepsilon_k}} |h_{\varepsilon_k}|^2.$$

By Hölder's inequality, this implies

$$\|h_{\varepsilon_k}\|_{L^{2+r(q-1)}(M_{\varepsilon_k})}^{r(q-1)} \leq \left(\beta^{q-1} + \frac{1}{2} \right)^r.$$

Raising both sides to the power $(r(q-1))^{-1}$ and then sending $r \rightarrow \infty$ gives the bound

$$\|h_{\varepsilon_k}\|_{L^\infty(M_{\varepsilon_k})} \leq \left(\beta^{q-1} + \frac{1}{2} \right)^{\frac{1}{q-1}}$$

provided k is large enough. This implies that $\|u_{\varepsilon_k}\|_{L^\infty(M)}$ has a uniform L^∞ bound for all large k , as needed. \square

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