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# Effects of Photopatterning Conditions on Azimuthal Surface Anchoring Strength

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**Abstract:** Spatially-varying alignment of liquid crystals is essential for research and applications. One widely used method is based on the photopatterning of thin layers of azo-dye molecules, such as Brilliant Yellow (BY), that serve as an aligning substrate for a liquid crystal. In this study, we examine how photopatterning conditions, such as BY layer thickness ( $b$ ), light intensity ( $I$ ), irradiation dose, and age affect the alignment quality and the strength of the azimuthal surface anchoring. The azimuthal surface anchoring coefficient,  $W$ , is determined by analyzing the splitting of integer disclinations into half-integer disclinations at prepatterned substrates. The strongest anchoring is achieved for  $b$  in the range of 5–8 nm.  $W$  increases with the dose, and within the same dose,  $W$  increases with  $I$ . Aging of a non-irradiated BY coating above 15 days reduces  $W$ . Our study also demonstrates that sealed photopatterned cells filled with a conventional nematic preserve their alignment quality for up to four weeks, after which time  $W$  decreases. This work suggests the optimization pathways for photoalignment of nematic liquid crystals.

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**Keywords:** Azimuthal anchoring, Surface anchoring, Liquid crystals, Photopatterning conditions, Alignment stability, Brilliant Yellow

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## 1. Introduction

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Alignment of liquid crystals (LCs) is crucial for their applications. While the unidirectional alignment can be achieved by mechanical rubbing [1–5], alignment with spatially-varying “easy axis” requires a more sophisticated approach. Spatially-varying alignment becomes exceedingly important in many academic and applied projects, such as the fabrication of planar optics elements [6–12], LC elastomer coatings with pre-designed topography [13–19], orientationally ordered environments that control collective and individual dynamics of microswimmers [20–25], and substrates that align living tissues [26–29]. The most popular approach to achieving spatially-varying alignment of LCs is photoalignment [13,30–38]. Photoalignment allows one to design complex director patterns with high spatial resolution [34,35,39], controllable surface anchoring [40], dynamic director alignment [41], and the capability to pattern the alignment on flexible and curved substrates [42–44]. Photoalignment does not induce impurities, electric charges, or mechanical damage to the treated surfaces, unlike conventional rubbing [45].

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Brilliant Yellow is one of the azo-dye materials used for photoalignment [46–48]. Extensive studies by Yaroshchuk et al. [49] brought the conclusion that BY provides “excellent photoalignment of nematic, smectic ferroelectric and reactive liquid crystals” and “shows extraordinarily high photo and thermal stability.” The study also noted that the double C=C bond in the core of BY molecules results in an absorption peak at 432 nm [49]. This peak, being in the lower part of the visible spectrum, might result in reduced stability

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when exposed to visible light. Nevertheless, the study by Yaroshchuk et al. [49] demonstrated that BY has excellent stability when irradiated with ultraviolet light and kept at an elevated temperature of 150 °C for one hour. Since BY allows one to produce high quality alignment in a relatively easy process, it continues to attract interest in research and applications, see, for example, some recent publications in Refs. [33,49–54].

Azo-dye molecules undergo photoinduced reorientation when exposed to light [50,55,56]. For irradiation with a linearly polarized beam, the probability of absorption is  $P \propto \cos^2 \beta$  where  $\beta$  is the angle between the long axis of the molecule in its *trans* state and the light polarization direction. The absorption-reorientation process repeats itself until the dye molecule aligns perpendicularly to the light polarization,  $\beta = \pi/2$ . Studies by Wang et al. [57] and Shi et al. [58] show that alignment quality strongly depends on exposure to humidity and only slightly depends on the type of surface used for depositing BY. Our previous work showed that a longer light exposure produces a stronger in-plane (azimuthal) surface anchoring [40]. However, the effect of other important factors remained unexplored.

This study investigates the effect of photopatterning conditions, such as BY layer thickness ( $b$ ), light intensity ( $I$ ), irradiation dose, and age of non-irradiated substrate, on the alignment quality and the strength of azimuthal anchoring expressed by the anchoring coefficient  $W$  in the surface potential  $\frac{1}{2}W\sin^2\alpha$ , where  $\alpha$  is the angle between the alignment direction imposed by the patterned BY layer and the actual director specifying the local orientation of the liquid crystal. The coefficient  $W$  is determined by analyzing the splitting of integer disclinations into half-integer disclinations at photopatterned substrates [26,40]. We find that  $b$  in the range 5–8 nm yields the strongest azimuthal anchoring.  $W$  increases with the dose and with  $I$  when the dose is fixed. BY-coated substrates, which are photopatterned within 15 days of substrate preparation (after BY spin coating and baking, but before irradiation), show no significant change in  $W$ . However, substrates aged for more than 15 days before irradiation exhibit a decline in  $W$ . Our study also demonstrates that photopatterned cells filled with a conventional nematic LC and sealed with epoxy glue preserve their alignment strength for about 4 weeks, but further aging of the filled cell leads to a reduction of  $W$ . The results facilitate the optimization of BY photoalignment for liquid crystal applications.

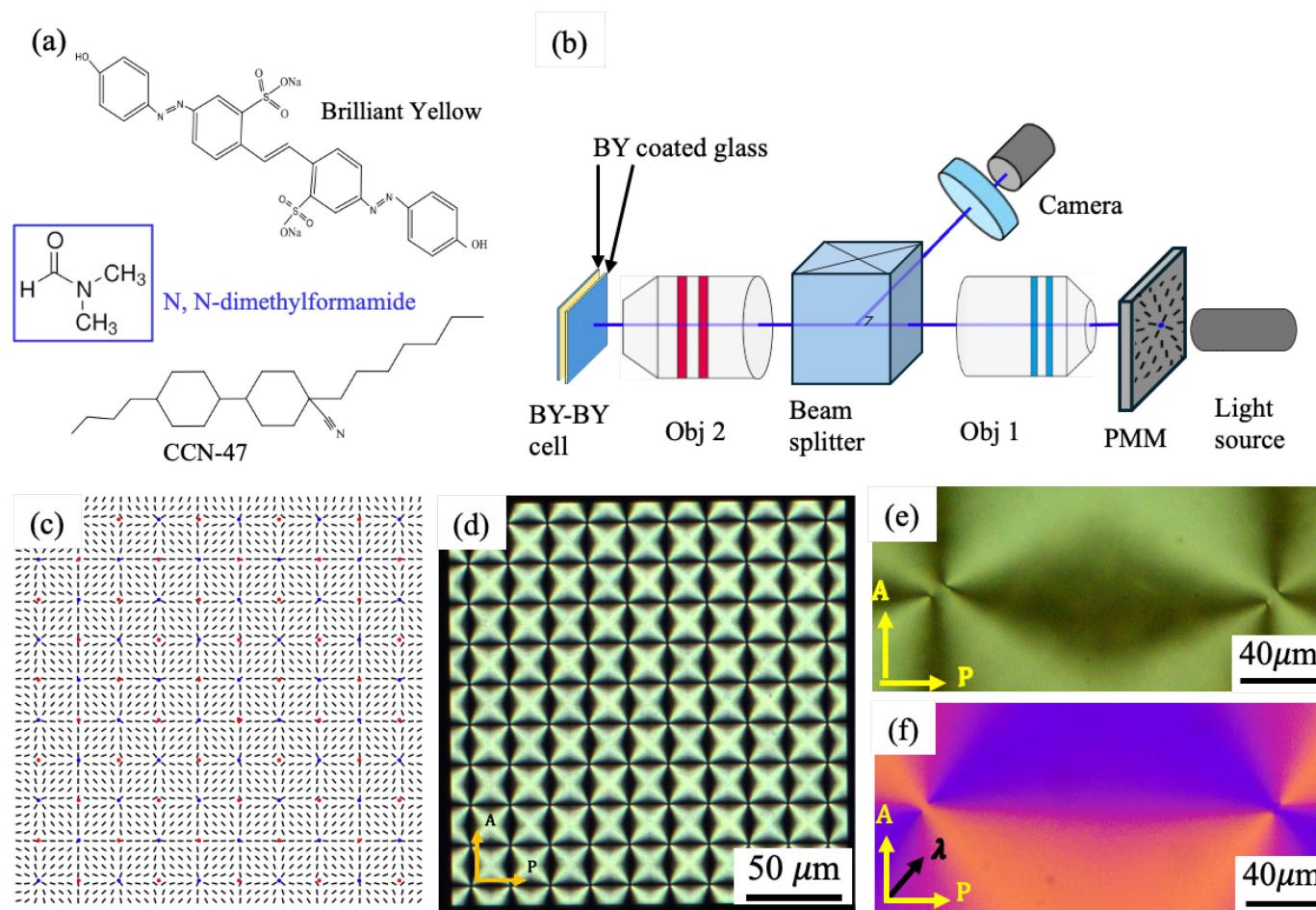
## 2. Materials and Methods

### 2.1. Cell Preparations

Indium tin oxide (ITO) coated glass plates are sonicated in water with a small amount of detergent at 60 °C for 15 min. Although the studies do not require an application of an electric field, the choice of ITO-coated plates is justified by the fact that most applications of liquid crystals involve electro-optic effects thus the presence of the ITO electrodes is often a necessity. The plates are rinsed with isopropanol, dried in an oven at 80 °C for 15 min, and exposed to UV in an ozone chamber for 15 min. The plates are spin-coated with a solution of azo-dye BY in N, N-dimethylformamide (DMF), Figure 1(a) (both purchased from Sigma Aldrich) at 3000 rpm for 30 seconds, and baked at 80 °C for 30 min. We use DMF solutions with various BY concentrations, 0.2, 0.4, 0.5, 0.6, 0.8, 1.0, 2.0, and 4.0 wt%, in order to vary the thickness  $b$  of the resulting BY layer and to explore its effect on the anchoring strength. All other experiments are performed with substrates coated with a 0.5 wt% BY solution, resulting in  $b = 7.6$  nm, which is in the optimal thickness range of 5–8 nm that yields the strongest  $W$ . To avoid detrimental effects of humidity on BY alignment [57], we control the relative humidity (RH) of the environment at less than 20% during the spin coating and baking, and at 20–35% during substrate storage, cell assembly, and photopatterning. During imaging, the sealed LC filled cells are kept at RH of less than 50% and a constant temperature of 45 °C.

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Cells are assembled from two glass substrates with the BY-coated surfaces facing each other; these are called BY-BY cells throughout the text. The gap is fixed using epoxy glue NOA 65, without spacers, to achieve a thickness of  $\sim 1 \mu\text{m}$ . The thickness  $h$  of the gap between two plates is measured by an interferometric technique using a UV/VIS spectrometer Lambda 18 (Perkin Elmer).



**Figure 1. Chemical structures of materials and photopatterned defect array.** (a) Photoresponsive azo-dye Brilliant Yellow, liquid crystal CCN-47, and solvent N, N-dimethylformamide. (b) Schematic of the photopatterning setup. Obj: objective, PMM: plasmonic metmask. (c) Director pattern of the defects array. The blue dots represent the cores of +1 defects, and the red dots represent the cores of -1 defects. (d) A polarizing optical microscope texture of the plasmonic meta mask with  $\pm 1$  radial defect array. (e) Polarized optical microscope images of a portion of a photopatterned cell with a +1 radial defect on the left and a -1 defect on the right. The cell thickness is  $1.1 \mu\text{m}$ . (f) The same polarizing microscopy with a full-wavelength optical compensator with the slow axis  $\lambda$ . P and A represent the polarizer and analyzer, respectively.

## 2.2 Photopatterning

We use the plasmonic metmask (PMM) technique introduced by Guo et al. [33,34] to pattern the substrates with an array of +1 and -1 defects (these defects should not be confused with “defects” induced by mishandling of the samples or with dust particles). PMM is an aluminum film with a thickness of 150 nm containing an array of nanoslits, each with a length of 220 nm and a width of 100 nm. When an unpolarized light beam passes through a nanoslit, the transmitted light becomes polarized along the short axis of

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the nanoslit. The degree of polarization of the transmitted light depends on the wavelength. Guo et al. [34] demonstrated experimentally and through numerical simulations that the polarization contrast ratio exceeds 7 dB for wavelengths ranging from 400 to 800 nm. The transmitted light beam irradiates the cells, as shown in Figure 1(b). We use two types of cells in this work: BY-BY and BY-PS cells. The BY-PS cells are assembled with one BY-coated substrate and one polystyrene (PS) coated substrate, with the coated surfaces facing each other. Preparation of the PS-coated substrates is explained later in the text, section 3.1. During photoalignment, BY-PS cells are positioned so that the BY-coated substrate is closer to the light source. Irradiated light aligns azobenzene molecules perpendicularly to the light polarization direction. As a result, a desired pattern, which replicates the pattern of nanoslits in the PMM, is produced in the azo-dye layer. The BY molecule exhibits an absorption range of 350–500 nm, with a peak at 432 nm [49]. We use a light source EXFO X-Cite with a wavelength range of (320-750) nm, which covers fully the absorption spectrum of BY. The light beam propagates along the normal to the PMM and the cell, Figure 1(b). The light intensity  $I$  is measured at the point of cell incidence using a power and energy meter console PM 100D (THORLABS).

A periodic square lattice of defects with strength +1 and -1 is designed by using a superposition rule for the in-plane director,  $\hat{\mathbf{n}}_{\text{BY}} = (n_x, n_y, 0) = (\cos \varphi, \sin \varphi, 0)$ , where  $\varphi = \sum_{i=1}^p \sum_{j=1}^q (-1)^{i+j} \arctan \left( \frac{y-jb}{x-ia} \right)$ ,  $x$  and  $y$  are the Cartesian coordinates,  $p$  and  $q$  are the numbers of defects in rows and columns, respectively,  $p = q = 10$ ;  $a = b = 200 \mu\text{m}$  is the distance between the defects along the  $x$  and  $y$  directions, respectively. The +1 defects in the pattern are of a radial type, so that the director around them experiences mostly splay, Figure 1(c).

### 2.3 Nematic material

The photopatterned cells are filled with 4-butyl-4-heptyl-bicyclohexyl-4-carbononitrile (CCN-47), Figure 1(a), by capillary action in the isotropic state at the temperature 70 °C. The material exhibits the following phase transitions upon heating: Smectic A 29.9 °C Nematic 58.5 °C Isotropic. After filling the cells, they are kept at 45 °C during the experiments using a Linkam hot stage. At 45 °C, the elastic constants  $K_1$  of splay and  $K_3$  of bend of CCN-47 are equal,  $K_1 = K_3 = K = 8 \text{ pN}$  [59,60], which allows one to use the superposition rule for the director field and to analyze the elastic properties of the patterns in the so-called one-constant approximation [61].

### 2.4 Optical microscopy characterization

Optical textures of photopatterned LC cells are recorded using a polarized optical microscope (Olympus BX51) equipped with a Basler (acA1920-155um) digital color camera, Figure 1(e,f). A full-wavelength (530 nm) optical compensator is used to reconstruct the director field, Figure 1(f). Regions where the director aligns parallel to the slow axis of the compensator, exhibit a blue interference color, while areas where the director is perpendicular to the slow axis appear yellow. The left defect in Figure 1(e,f) is a +1 radial defect, while the right is a -1 defect. The separation distance between half integer defects,  $d$ , is measured using the open-source software package Fiji/ImageJ.

### 2.5 Theoretical background

Defects of strength  $\pm 1$  tend to split into pairs of  $\pm 1/2$  to reduce their elastic energy [61]. In the so-called one-constant approximation, the elastic repulsive potential of two  $+1/2$  or two  $-1/2$  defects is weakly dependent on their separation  $d$ :  $F_E = -\frac{\pi K h}{2} \ln \frac{d}{2r_c}$ , where  $K$  is the average Frank elastic modulus,  $h$  and  $r_c$  are the cell thickness and the radius of the disclination core, respectively [26,61]. In a photopatterned cell, the separation

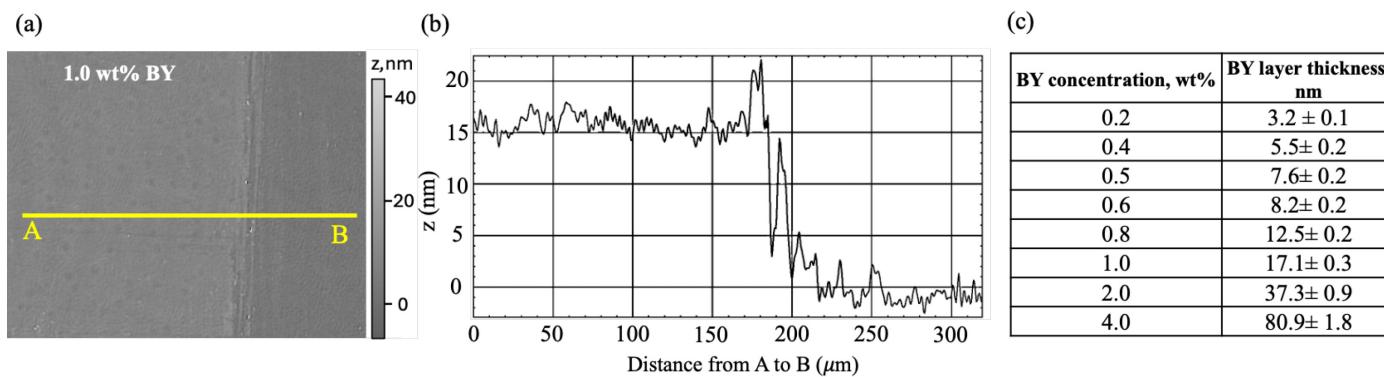
of the defects is resisted by surface anchoring that tends to enforce the patterned  $\pm 1$  defects. The elasticity-anchoring balance determines the equilibrium separation distance  $d$  of the semi-integer cores. The surface anchoring energy of a patterned cell can be found by integrating the Rapini-Papoulier potential,  $F_S = 2 \int_0^{2\pi} \int_0^d \frac{1}{2} W [1 - (\hat{\mathbf{n}}_{BY} \cdot \hat{\mathbf{n}}_{LC})^2] r dr d\varphi$  which yields  $F_S = 2\alpha W d^2$ ; here  $\hat{\mathbf{n}}_{LC}$  is the actual director field of a split defect pair,  $\hat{\mathbf{n}}_{BY}$  is the ideal radial pattern of the easy axis at the substrates,  $\alpha \approx 0.184$  is a numerical coefficient, and the factor 2 reflects anchoring at both plates [26,40,61]. The equilibrium value of  $d$  allows one to calculate the anchoring strength as  $W = \pi K h / (8\alpha d^2)$ . For cells assembled with one photopatterned surface and the second plate providing degenerate in-plane anchoring,  $W = \pi K h / (4\alpha d^2)$ .

We calculate  $W$  by measuring  $d$  and  $h$  and using the known values of  $K$  and  $\alpha$ , Figures 3, 5-7. The resolution of the optical microscope is sufficient for precise measurements, as the experiments show that the parameter  $d$  is higher than 3  $\mu\text{m}$  even for the highest achieved anchoring. The distance between neighboring photopatterned  $+1$  and  $-1$  defects is set to be  $\sim 200 \mu\text{m}$ , which is sufficiently large to prevent any interactions between them. The experiments are designed with thin cells to ensure core splitting is the prevailing director structure as opposed to the escape in the third dimension [40,62]. Each data point for  $d$  and  $W$  represents the average value obtained from 50 defects of the same sign within the array, with the errors calculated as the standard deviation [63].

### 3. Results

#### 3.1 Effect of BY layer thickness on photoalignment

The thickness  $b$  of BY coatings affects the anchoring strength. Therefore, BY coatings with different thicknesses are produced by spin coating solutions with different BY concentrations. The thickness  $b$  is measured using a digital holographic microscope (Lyncée Tec.) in a reflection mode with a vertical resolution better than 1 nm. A portion of the BY coating is removed by wiping the substrate with water, followed by isopropanol using a cotton bud, exposing the glass surface to serve as the reference for thickness measurement, Figure 2(a). After imaging the surface, the height difference between the coating and the glass substrate is measured using the open-source software package Fiji/ImageJ, Figure 2(b). Thickness is measured at 10 different locations on the coated surface, and the average of these measurements is used as  $b$ . Changing the BY concentrations (0.2, 0.4, 0.5, 0.6, 0.8, 1.0, 2.0, and 4.0 wt%) in DMF results in different values of  $b$ , ranging from 3.2–80.9 nm, Figure 2(c). The thickness  $b$  was also verified by scratching the BY-coated glass with a sharp blade to create a groove; the depth  $b$  of the groove is measured by a digital holographic microscope. The  $b$  values obtained by the two methods are in good agreement with each other, being within the measurement error.

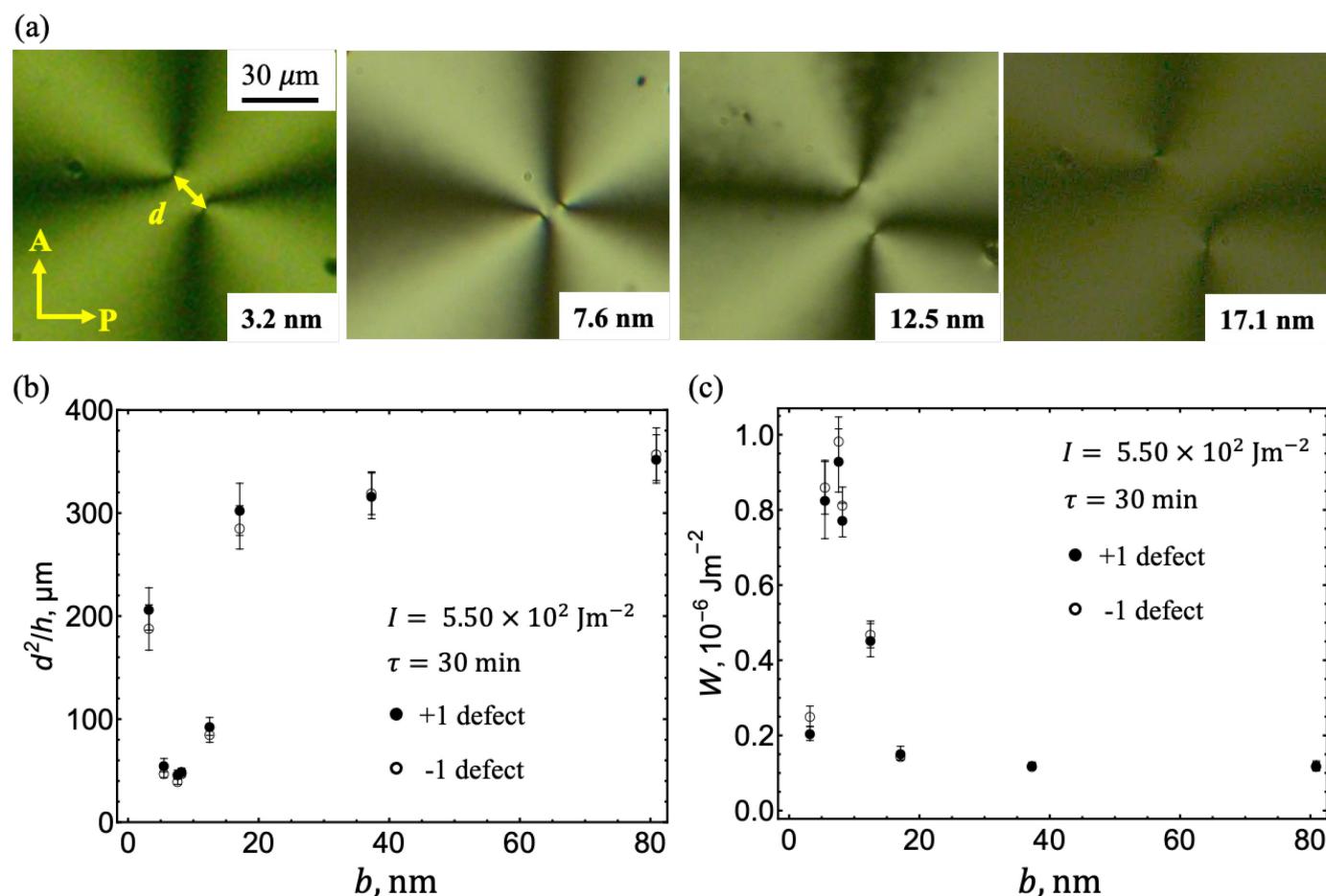


**Figure 2. Brilliant Yellow layer thickness measurement.** (a) An image of a 1.0 wt% BY-coated, partially wiped surface taken using a digital holographic microscope. The left side

of the image shows the BY coating, while the right side shows the glass surface after wiping off the coating. (b) The surface profile plot along the line AB marked in (a), and (c) BY layer thicknesses for different BY concentrations in the spin-coated solutions.

A larger  $b$  reduces the intensity of light that reaches the second plate that is further away from the source due to absorption by the BY layer on the front plate. To explore the effect, we prepare the BY-PS cells with one BY-coated substrate as the front plate and a polystyrene (PS) coated glass plate as the back substrate. During photoirradiation, the pattern is focused on the BY-coated front plate. PS-coated plates yield a negligibly weak azimuthal anchoring,  $W_{PS} \sim 10^{-10} \text{ J/m}^2 \ll W$  [64]. Thus, it is the patterned substrate that dictates the director orientation and the separation distance between the defects. The anchoring coefficient for these BY-PS cells is calculated as  $W = (\pi K h) / (4 \alpha d^2)$ . The PS-coated plates are prepared by spin-coating (3000 rpm, 30 s) a solution of 0.5 wt% PS in chloroform (Sigma-Aldrich, >98%) on clean glass substrates. The plates are kept at 80 °C for 30 min to evaporate the solvent. All cell preparation steps and observations are done on the same day. The cells are photopatterned with  $I = 5.50 \times 10^2 \text{ W m}^{-2}$  for 30 min.

The anchoring coefficient is measured for BY-PS cells with different  $b$ , Figure 3(a-c). For a thin BY layer of 3.2 nm,  $W$  is weak  $\sim 0.2 \times 10^{-6} \text{ J m}^{-2}$ , Figure 3(c). Anchoring increases sharply with  $b$  and reaches a maximum of  $0.98 \times 10^{-6} \text{ J m}^{-2}$  for  $b$  ranging from 5.5 nm to 8.2 nm, Figure 3(c). For layers with  $b = (12.5 - 37.3) \text{ nm}$ ,  $W$  sharply decreases to around  $10^{-7} \text{ J m}^{-2}$  and remains constant for thicker BY layers, Figure 3(c).

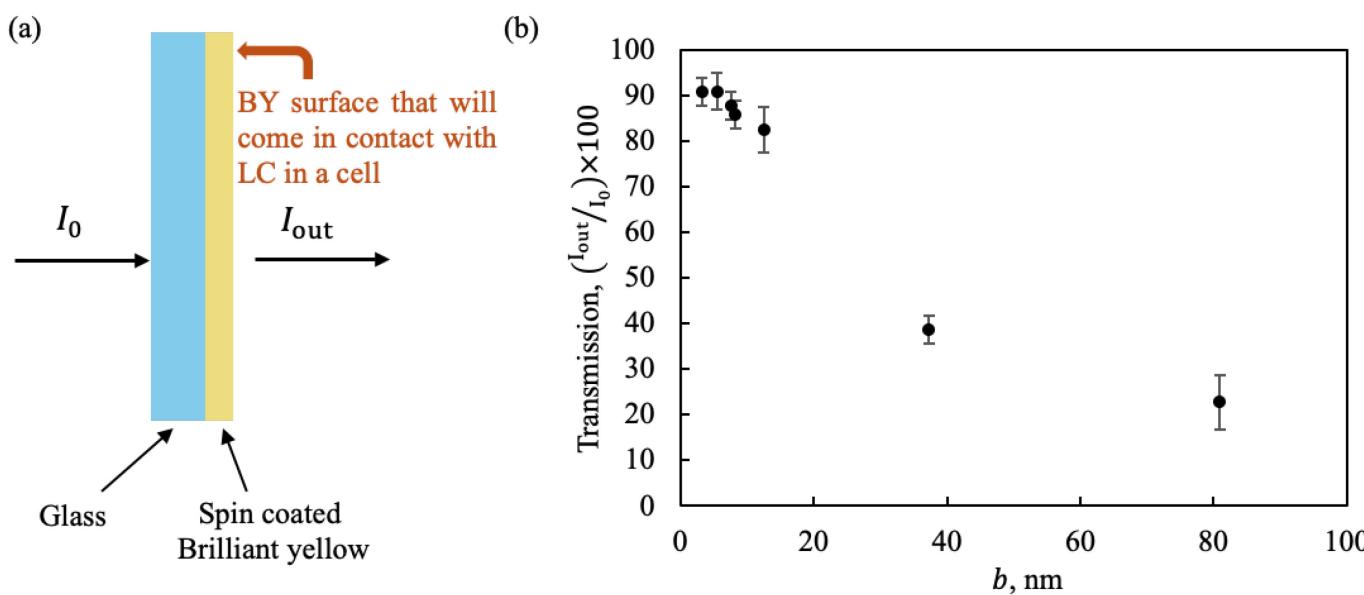


**Figure 3. Effect of Brilliant Yellow coating thickness on photoalignment in BY-PS cells.**  
 (a) Optical microscopy textures of photopatterned cells with different BY layer thicknesses,  $b$ .  $h = 2.2 \pm 0.2 \mu\text{m}$ . Cells are photopatterned with  $I = 5.50 \times 10^2 \text{ W m}^{-2}$  and a dose of  $9.90 \times 10^5 \text{ J m}^{-2}$ , corresponding to an irradiation time  $\tau = 30 \text{ min}$ . Cell

preparation and observation steps are performed on the same day. (b)  $d^2/h'$  and (c)  $W$ , as functions of  $b$ . Each data point represents the average value obtained from 50 defects of the same sign within the array, with the errors calculated as the standard deviation.

The weak  $W$  at the thinnest BY coatings ( $h = 3.2$  nm) is most likely caused by the surface roughness of the glass, ITO, and the coating itself. As clear from Fig. 2b, the thickness 3.2 nm is within the range of the measured variations of the coating's surface, which implies that in some places there might not be enough BY molecules. Another potential reason for weak anchoring at thin coatings is a formation of hydrogen bonds between the BY molecules and the hydrophilic UV/O<sub>3</sub> treated ITO surface. As noted by Wang et al. [57], BY films of a thickness 3 nm on a hydrophilic substrate of polyvinyl alcohol (PVA) demonstrated a low degree of orientational order as compared to thicker films. The effect was attributed to the formation of hydrogen bonds between the BY molecules and the hydroxide OH group of PVA which hinders the *trans-cis* isomerization of BY molecules needed for light-induced alignment [57,65].

The weak  $W$  at thick BY layers is attributed to weaker light intensity received by BY molecules at the interface with the LC, due to absorption by BY along the light path. The light transmittance through BY layers of varying thicknesses is measured at wavelength 410 nm using a UV/VIS spectrometer Lambda 18 (Perkin Elmer). The data show 90% of incident light is transmitted by the BY layers when  $b \leq 6$  nm, but the transmission is reduced to 22% for  $b = 80.9$  nm, Figure 4(b). As a result, the light intensity received by BY molecules on the surface that will meet the LC is substantially reduced, Figure 4(a,b). As the effective intensity decreases, the effective dose also decreases; both factors yield a weaker  $W$ . Since BY layers with thicknesses less than 8 nm transmit over 85% of the incident light, and the thickness of 5–8 nm results in stronger azimuthal anchoring, a 7.6 nm thick BY coating layer (spin coated with a 0.5 wt% BY solution) is used for the rest of experiments, presented in Figures 5,6,7. For a BY-BY cell, each with a 7.6 nm thick BY layer, the substrate that is further from the light source receives 87% of the incident light intensity during photoalignment. In such a cell, both BY surfaces that will later meet the LC are aligned with the same light intensity, 87% of the incident light. As shown in Figure 3(c), this intensity is sufficient to achieve the strongest anchoring. Therefore, all subsequent experiments are performed using BY-BY cells assembled with two BY-coated glass substrates.



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**Figure 4. Transmittance measurements of different Brilliant Yellow layer thicknesses.**

(a) Schematic illustration of BY-coated glass. Here,  $I_0$  is the incident light intensity and  $I_{\text{out}}$  is the intensity of light transmitted through BY layer. (b) Transmission as a function of BY coating thickness. The wavelength of light is 410 nm.

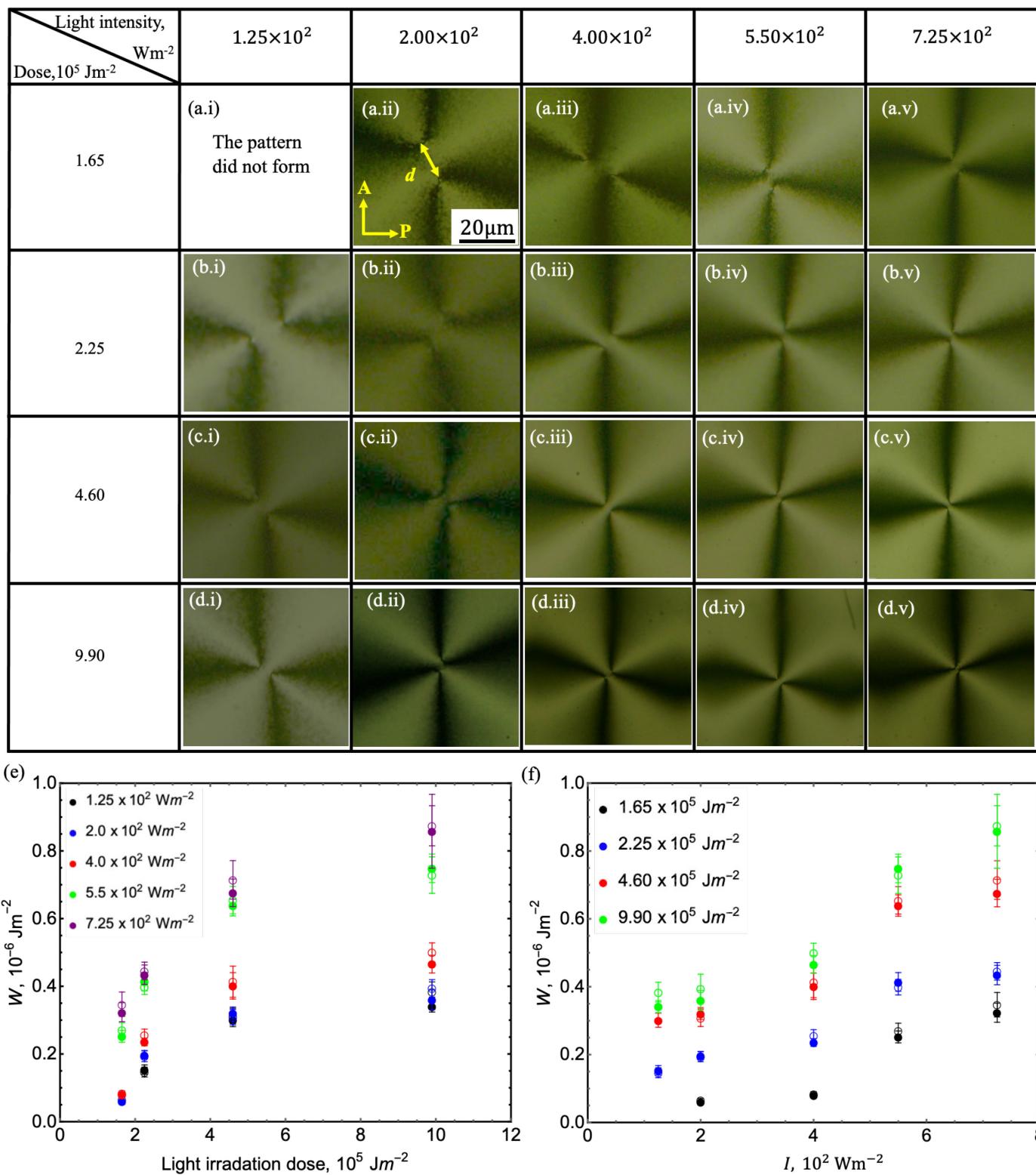
### 3.2 Effect of light dose and intensity on photoalignment

The BY-BY cells are photopatterned with four different light irradiation doses. The dose is defined as a product of irradiation time ( $\tau$ ) and the light intensity  $I$ .  $\tau$  is changed to vary the dose while  $I$  is kept constant. When studying the effect of  $I$ , the dose is kept constant by varying  $\tau$ . The cell preparation and characterization are completed within a maximum of two consecutive days to minimize the effects of aging.

Increasing the dose from  $1.65 \times 10^5 \text{ Jm}^{-2}$  to  $9.90 \times 10^5 \text{ Jm}^{-2}$  results in an increase of  $W$ , Figure 5(a-d rows) and Figure 5(e);  $W$  saturates at doses higher than  $4.6 \times 10^5 \text{ Jm}^{-2}$ . This behavior is consistent with the previous study on the effect of photopatterning time on  $W$  [40].

Increasing  $I$  from  $1.25 \times 10^2 \text{ Wm}^{-2}$  to  $7.25 \times 10^2 \text{ Wm}^{-2}$  at a fixed dose increases  $W$  and creates better alignment, Figure 5(i-v columns) and Figure 5(f). It is notable that although dose and  $I$  are interrelated, both must be set properly to achieve a high quality patterning. For instance, photopatterning with a high dose of  $9.90 \times 10^5 \text{ Jm}^{-2}$  and a low  $I$  of  $1.25 \times 10^2 \text{ Wm}^{-2}$  results in a poor alignment (Figure 5(d,i)), and a weak  $W = 0.36 \pm 0.02 \times 10^{-6} \text{ Jm}^{-2}$ , Figure 5(e,f). However, a moderate dose of  $2.25 \times 10^5 \text{ Jm}^{-2}$  at  $I > 5.50 \times 10^2 \text{ Wm}^{-2}$  results in a better alignment with  $W \geq 0.4 \times 10^{-6} \text{ Jm}^{-2}$ . Using  $W$  values, one can establish criteria for setting irradiation conditions to achieve a good alignment. The results also show that  $W$  can be tuned by adjusting the dose and  $I$ .

At a constant dose, photopatterning with high  $I$  for a short  $\tau$  is more efficient than low intensity irradiation over a prolonged  $\tau$  (compare columns i and v in Figure 5), Figure 5(f). This indicates that the total number of photons is not the single decisive factor defining the anchoring strength. The intensity-dependent behavior of  $W$  suggests that photoisomerization occurs as a collective process, where the isomerization of individual BY molecules depends on the isomerization probability of neighboring molecules. Schönhoff et al. [66] observed a similar dependency of molecular photoreorientation on intensity for 4-(4'-N-octadecylamino)phenylazocyanobenzene (amino azobenzene) films irradiated by polarized light with varying intensity but a constant dose and concluded that photoreorientation is a pronounced collective effect.



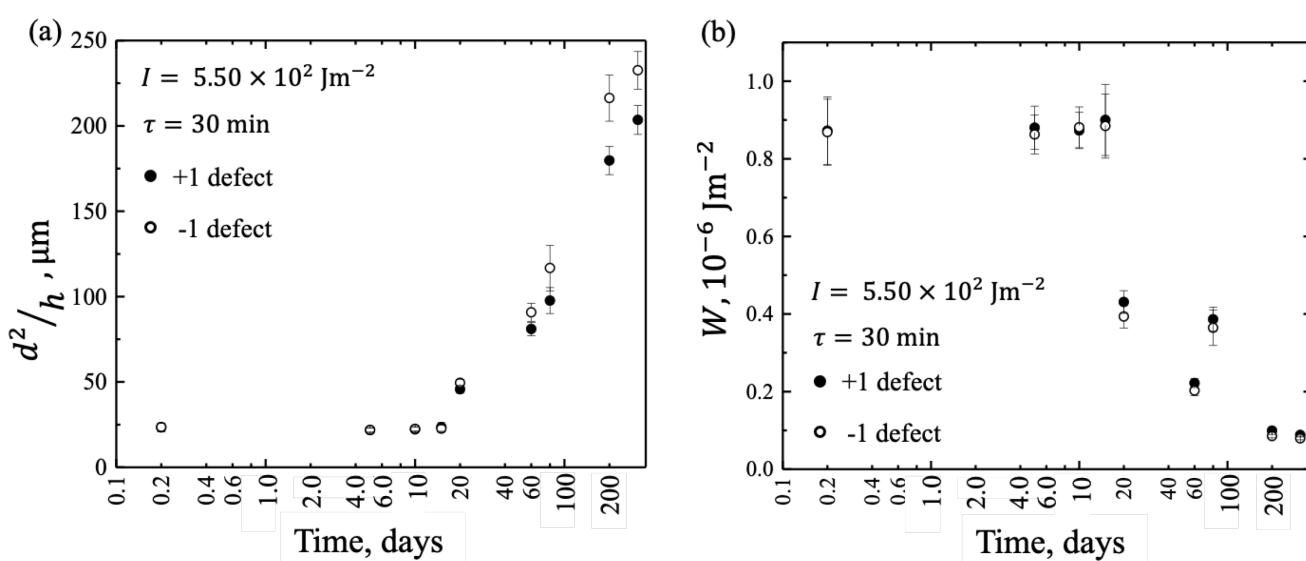
**Figure 5. Effect of dose and light intensity on photoalignment in BY-BY cells.** (a-d) Optical microscopy textures of photopatterned cells at different irradiation doses and light intensities. The images show a pair of  $+1/2$  defects split from a photopatterned  $+1$  defect.  $h = 1.3 \pm 0.3 \mu m$  and  $b = 7.6 nm$ . Cell preparation and characterization are performed within two consecutive days. (e) Azimuthal surface anchoring coefficient,  $W$ , as a function of dose for various light intensities. (f)  $W$  as a function of  $I$  for different doses. Colors indicate different intensities and doses in (e) and (f), respectively. Filled circles represent

+1 defects, and open circles represent -1 defects. Each data point for  $W$  represents the average value obtained from 50 defects in the array, with the errors calculated as the standard deviation.

### 3.3 Effect of aging of non-irradiated BY coatings.

The glass substrates, which were spin-coated with BY and then baked, are stored in a Humidity and Temperature-Controlled Cabinet (SIRUI HC series) for up to 310 days at relative humidity (RH) 25–35%, and 23 °C. Then BY-BY cells are assembled using two aged, non-irradiated BY-coated plates. After a predetermined aging time, the BY coatings are photopatterned with an  $I$  of  $5.50 \times 10^2 \text{ Wm}^{-2}$  for 30 min. The cells are filled with the nematic material, sealed and analyzed under the optical microscope

Aging of the non-irradiated BY layer impacts photopatterning in two phases. For BY layers irradiated within 15 days after the layers were spin-coated, dried and baked,  $W$  remains constant at  $\sim(0.88 \pm 0.01) \times 10^{-6} \text{ Jm}^{-2}$ , Figure 6. This suggests that BY-coated substrates can be safely stored at controlled humidity conditions (at RH 25–35%) for around two weeks without a noticeable effect on their photopatterning quality. However, for older BY layers,  $W$  decreases continuously with age and reduces to  $0.08 \times 10^{-6} \text{ Jm}^{-2}$  for substrates photopatterned 310 days after preparation, Figure 6. A possible reason could be the water absorption by the BY layers during storage. As already noted, we control the RH of the environment at less than 20% during the spin coating and baking, but during the storage, the RH is at higher levels 20–35%. Absorption of water during the prolonged storage can result in aggregation of BY molecules into J-structures, a process called by Shi et al. [58] “hydrogen-bond-assisted self-assembly of BY molecules with water molecule insertion”. The molecules in J-aggregates are less likely to undergo an efficient trans-to-cis isomerization. Accumulation of water at the substrate-BY interface with a suppression of isomerization through hydrogen bonds might also contribute to the effect of aging.

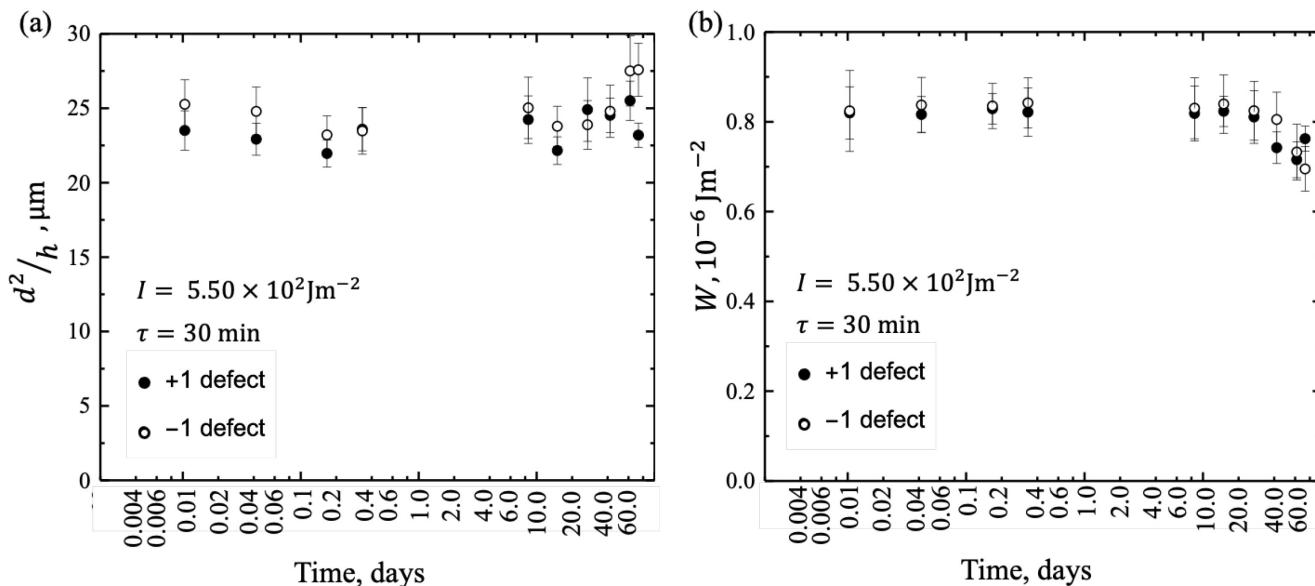


**Figure 6. Effect of non-irradiated BY-coated layer aging on photoalignment in BY-BY cells.** (a)  $d^2/h$  and (b)  $W$ , as a function of non-irradiated BY layer aging time. Each data point represents the average value obtained from 50, +1 and 50, -1 defects within the array, with the errors calculated as the standard deviation.  $h = 1.1 \pm 0.2 \mu\text{m}$  and  $b = 7.6 \text{ nm}$ . Cell assembly, photopatterning, and characterization are performed on the same day. Cells are photopatterned with  $I = 5.50 \times 10^2 \text{ Wm}^{-2}$  and dose of  $9.90 \times 10^5 \text{ Jm}^{-2}$ .

### 3.4 Surface patterning stability of aged LC-filled cells.

A photopatterned BY-BY cell is filled with CCN-47, and the edges of the cell are sealed with epoxy glue after the filling. Photopatterning is performed with  $I = 5.50 \times 10^2 \text{ Wm}^{-2}$  for 30 min. The LC-filled BY-BY cell is maintained in an environment with RH <50% at 45 °C for 72 days. While maintaining the temperature at 45 °C to stabilize the nematic phase, optical textures of the cell are recorded over time as the cell ages.

The distance  $d$  remains at around  $10.4 \pm 0.3 \mu\text{m}$  for 26 days, corresponding to  $W = (0.82 \pm 0.01) \times 10^{-6} \text{ Jm}^{-2}$ , Figure 7(a,b). Maintaining the cell for a longer period (around 72 days) results in a slight decrease in  $W$  to  $0.76 \times 10^{-6} \text{ Jm}^{-2}$ , Figure 7(b). This result suggests that filling the photopatterned cell with LC and sealing it helps to preserve the quality of photopatterning. The observed decline could be due to the gradual dissolution of LC into the BY coating [67,68] or the dissolving of epoxy glue into the LC over time [69].



**Figure 7. Surface patterning stability of aged LC-filled BY-BY cells.** (a)  $d^2/h$ , and (b)  $W$ , as functions of the age of the LC filled cell. Each data point represents the average value obtained from 50 defects of the same sign within the array, with the errors calculated as the standard deviation.  $h = 1.5 \pm 0.02 \mu\text{m}$  and  $b = 7.6 \text{ nm}$ . The cell is photopatterned with  $I = 5.50 \times 10^2 \text{ Wm}^{-2}$  and a dose of  $9.90 \times 10^5 \text{ Jm}^{-2}$ . Cell preparation steps are performed on the same day and observation is done as the cell ages.

#### 4. Conclusions

We have demonstrated how various photopatterning conditions, including the thickness  $b$  of the azobenzene alignment layers the intensity  $I$  of the light used for photopatterning, the irradiation dose, and substrate aging, affect the anchoring strength of a photopatterned nematic LC. Our results show that BY layers of thickness  $b = 5\text{--}8 \text{ nm}$  produce the strongest anchoring coefficient  $W$ . Moreover,  $W$  increases with both irradiation dose and  $I$ . Also, aging of non-irradiated substrates beyond 15 days significantly reduces  $W$ . However, if the cell is filled with a LC immediately after photopatterning of the BY layer,  $W$  remains constant for up to 4 weeks. This work provides practical strategies for enhancing the azimuthal strength of the photopatterned anchoring of nematics. The results offer guidelines for optimizing BY photoalignment parameters and storage.

The method to measure the azimuthal anchoring coefficient described in this paper is based on the properties of topological defects produced by photopatterning; it does not require one to use a second plate, for example, a rubbed polyimide plate, with an anchoring much stronger than the photoalignment anchoring [47,51,70]. It also does not require

one to use any external fields [71-73] or to prepare wedge samples of varying thickness [74-76]. Within a broader prospectus, our approach to measuring the azimuthal anchoring coefficient can be extended to other photoalignment materials, such as SD-1 that does not feature absorption peaks in the visible spectral range [49] and thus might be better suited for applications that require a long-term stability.

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## References

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1. Cognard, J. Alignment of nematic liquid crystals and their mixtures. *Mol. Cryst. Liq. Cryst.* **1982**, *78*, 1.
2. Lee, K.W.; Paek, S.H.; Lien, A.; Durning, C.; Fukuro, H. Microscopic molecular reorientation of alignment layer polymer surfaces induced by rubbing and its effects on LC pretilt angles. *Macromolecules* **1996**, *29*, 8894-8899.
3. Chen, J.; Cranton, W.; Fihn, M. *Handbook of visual display technology*; Springer Publishing Company, Incorporated: 2011.
4. Babakhanova, G.; Lavrentovich, O.D. The techniques of surface alignment of liquid crystals. In Proceedings of the Modern Problems of the Physics of Liquid Systems: Selected Reviews from the 8th International Conference “Physics of Liquid Matter: Modern Problems”, Kyiv, Ukraine, May 18-22, 2018 8, 2019; pp. 165-197.
5. Muravsky, A.; Murauski, A. Q&A of liquid crystal alignment: theory and practice. *Frontiers in Soft Matter* **2024**, *4*, 1382925.
6. Huang, Y.H.; Ko, S.W.; Li, M.S.; Chu, S.C.; Fuh, A.Y.G. Modulation of shape and polarization of beam using a liquid crystal q-plate that is fabricated via photo-alignment. *Optics Express* **2013**, *21*, 10954-10961.
7. Jiang, M.; Yu, H.; Feng, X.; Guo, Y.; Chaganava, I.; Turiv, T.; Lavrentovich, O.D.; Wei, Q.H. Liquid crystal Pancharatnam - Berry micro - optical elements for laser beam shaping. *Advanced Optical Materials* **2018**, *6*, 1800961.
8. Chen, P.; Wei, B.Y.; Hu, W.; Lu, Y.Q. Liquid - crystal - mediated geometric phase: from transmissive to broadband reflective planar optics. *Advanced Materials* **2020**, *32*, 1903665.
9. Tabiryan, N.V.; Roberts, D.E.; Liao, Z.; Hwang, J.Y.; Moran, M.; Ouskova, O.; Pshenichnyi, A.; Sigley, J.; Tabirian, A.; Vergara, R. Advances in transparent planar optics: enabling large aperture, ultrathin lenses. *Advanced Optical Materials* **2021**, *9*, 2001692.
10. Tang, D.; Shao, Z.; Zhou, Y.; Lei, Y.; Chen, L.; Xie, J.; Zhang, X.; Xie, X.; Fan, F.; Liao, L. Simultaneous surface display and holography enabled by flat liquid crystal elements. *Laser & Photonics Reviews* **2022**, *16*, 2100491.
11. Chen, F.; Zheng, J.; Xing, C.; Sang, J.; Shen, T. Applications of liquid crystal planer optical elements based on photoalignment technology in display and photonic devices. *Displays* **2024**, *102632*.
12. Tabiryan, N.V.; Pshenichnyi, A.; Ouskova, O.; Gerosa, K.; Crenshaw, E.; Carter, M.J.; McConney, M.; Urbas, A.; Bunning, T.J.; Slagle, J. Liquid crystalline geometrical phase Alvarez-Lohmann lenses. *Optical Materials Express* **2024**, *14*, 2786-2799.
13. McConney, M.E.; Martinez, A.; Tondiglia, V.P.; Lee, K.M.; Langley, D.; Smalyukh, I.I.; White, T.J. Topography from topology: photoinduced surface features generated in liquid crystal polymer networks. *Advanced Materials* **2013**, *25*, 5880-5885.

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14. Ware, T.H.; McConney, M.E.; Wie, J.J.; Tondiglia, V.P.; White, T.J. Voxelated liquid crystal elastomers. *Science* **2015**, *347*, 982-984. 443

15. Mostajeran, C.; Warner, M.; Ware, T.H.; White, T.J. Encoding Gaussian curvature in glassy and elastomeric liquid crystal solids. *Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences* **2016**, *472*, 20160112. 445

16. Babakhanova, G.; Turiv, T.; Guo, Y.; Hendrikx, M.; Wei, Q.H.; Schenning, A.P.; Broer, D.J.; Lavrentovich, O.D. Liquid crystal elastomer coatings with programmed response of surface profile. *Nature communications* **2018**, *9*, 456. 447

17. Feng, W.; Liu, D.; Broer, D.J. Functional liquid crystal polymer surfaces with switchable topographies. *Small Structures* **2021**, *2*, 2000107. 449

18. Chen, J.; Akomolafe, O.I.; Jiang, J.; Peng, C. Light-actuated liquid crystal elastomer prepared by projection display. *Materials* **2021**, *14*, 7245. 451

19. Chen, J.; Johnson, A.S.; Weber, J.; Akomolafe, O.I.; Jiang, J.; Peng, C. Programmable light - driven liquid crystal elastomer kirigami with controlled molecular orientations. *Advanced Intelligent Systems* **2022**, *4*, 2100233. 453

20. Peng, C.; Turiv, T.; Guo, Y.; Wei, Q.H.; Lavrentovich, O.D. Command of active matter by topological defects and patterns. *Science* **2016**, *354*, 882-885. 455

21. Straube, A.V.; Pagès, J.M.; Ortiz-Ambriz, A.; Tierno, P.; Ignés-Mullol, J.; Sagués, F. Assembly and transport of nematic colloidal swarms above photo-patterned defects and surfaces. *New Journal of Physics* **2018**, *20*, 075006. 457

22. Turiv, T.; Koizumi, R.; Thijssen, K.; Genkin, M.M.; Yu, H.; Peng, C.; Wei, Q.H.; Yeomans, J.M.; Aranson, I.S.; Doostmohammadi, A. Polar jets of swimming bacteria condensed by a patterned liquid crystal. *Nature Physics* **2020**, *16*, 481-487. 459

23. Dhakal, N.P.; Jiang, J.; Guo, Y.; Peng, C. Self-assembly of aqueous soft matter patterned by liquid-crystal polymer networks for controlling the dynamics of bacteria. *ACS applied materials & interfaces* **2020**, *12*, 13680-13685. 462

24. Koizumi, R.; Turiv, T.; Genkin, M.M.; Lastowski, R.J.; Yu, H.; Chaganava, I.; Wei, Q.H.; Aranson, I.S.; Lavrentovich, O.D. Control of microswimmers by spiral nematic vortices: Transition from individual to collective motion and contraction, expansion, and stable circulation of bacterial swirls. *Physical Review Research* **2020**, *2*, 033060. 464

25. Ma, L.L.; Liu, C.; Wu, S.B.; Chen, P.; Chen, Q.M.; Qian, J.X.; Ge, S.J.; Wu, Y.H.; Hu, W.; Lu, Y.Q. Programmable self-propelling actuators enabled by a dynamic helical medium. *Science Advances* **2021**, *7*, eabh3505. 467

26. Turiv, T.; Krieger, J.; Babakhanova, G.; Yu, H.; Shiyanovskii, S.V.; Wei, Q.H.; Kim, M.H.; Lavrentovich, O.D. Topology control of human fibroblast cells monolayer by liquid crystal elastomer. *Science Advances* **2020**, *6*, eaaz6485. 469

27. Dhakal, N.P.; Jiang, J.; Guo, Y.; Peng, C. Photopatterning DNA structures with topological defects and arbitrary patterns through multiple length scales. *Physical Review Applied* **2020**, *13*, 014026. 471

28. Jiang, J.; Dhakal, N.P.; Guo, Y.; Andre, C.; Thompson, L.; Skalli, O.; Peng, C. Controlled dynamics of neural tumor cells by templated liquid crystalline polymer networks. *Advanced Healthcare Materials* **2020**, *9*, 2000487. 473

29. Zhao, Z.; Li, H.; Yao, Y.; Zhao, Y.; Serra, F.; Kawaguchi, K.; Zhang, H.; Chate, H.; Sano, M. Integer topological defects provide a new way to quantify and classify cell sheets. *bioRxiv* **2024**, 2024.2008.2028.610106. 475

30. Culbreath, C.; Glazar, N.; Yokoyama, H. Note: Automated maskless micro-multidomain photoalignment. *Review of Scientific Instruments* **2011**, *82*. 477

31. Shteyner, E.A.; Srivastava, A.K.; Chigrinov, V.G.; Kwok, H.S.; Afanasyev, A.D. Submicron-scale liquid crystal photoalignment. *Soft Matter* **2013**, *9*, 5160-5165. 479

32. Chigrinov, V. Liquid crystal devices based on photoalignment and photopatterning materials. In Proceedings of the Emerging Liquid Crystal Technologies IX, 2014; pp. 38-47. 481

33. Guo, Y.; Jiang, M.; Peng, C.; Sun, K.; Yaroshchuk, O.; Lavrentovich, O.D.; Wei, Q.H. Designs of plasmonic metamasks for photopatterning molecular orientations in liquid crystals. *Crystals* **2016**, *7*, 8. 483

484

34. Guo, Y.; Jiang, M.; Peng, C.; Sun, K.; Yaroshchuk, O.; Lavrentovich, O.; Wei, Q.H. High Resolution and High Throughput Plasmonic Photopatterning of Complex Molecular Orientations in Liquid Crystals. *Advanced Materials* **2016**, *28*, 2353–2358. 485

35. Yun, H.; Jiang, S.; Chen, H.; Zhu, Y.; Xu, X.; Li, B.; Xi, P.; Jiang, M.; Wei, Q.H. Ultra-high spatial resolutions in photopatterning molecular orientations. *Optics Express* **2024**, *32*, 31107–31119. 488

36. Ropač, P.; Hsiao, Y.T.; Berteloot, B.; Ussembayev, Y.; Nys, I.; Ravnik, M.; Neyts, K. Liquid Crystal 3D Optical Waveguides Based on Photoalignment. *Advanced Optical Materials* **2024**, *24*, 2402174. 489

37. Kudreyko, A.; Chigrinov, V.; Neyts, K.; Chausov, D.; Perestoronina, A. Photonic Devices with Multi-Domain Liquid Crystal Structures. *Crystals* **2024**, *14*, 512. 491

38. Modin, A.; Leheny, R.L.; Serra, F. Spatial Photo - Patterning of Nematic Liquid Crystal Pretilt and its Application in Fabricating Flat Gradient - Index Lenses. *Advanced Materials* **2024**, *23*, 10083. 493

39. Liu, S.; Nys, I.; Neyts, K. Two - Step Photoalignment with High Resolution for the Alignment of Blue Phase Liquid Crystal. *Advanced Optical Materials* **2022**, *10*, 2200711. 494

40. Padmini, H.N.; Rajabi, M.; Shiyanovskii, S.V.; Lavrentovich, O.D. Azimuthal anchoring strength in photopatterned alignment of a nematic. *Crystals* **2021**, *11*, 675. 497

41. Jiang, J.; Wang, X.; Akomolafe, O.I.; Tang, W.; Asilehan, Z.; Ranabhat, K.; Zhang, R.; Peng, C. Collective transport and reconfigurable assembly of nematic colloids by light-driven cooperative molecular reorientations. *Proceedings of the National Academy of Sciences* **2023**, *120*, e2221718120. 499

42. Chigrinov, V.G.; Kozenkov, V.M.; Kwok, H.S. *Photoalignment of liquid crystalline materials: physics and applications*; John Wiley & Sons: 2008. 502

43. Chigrinov, V.; Sun, J.; Wang, X. Photoaligning and photopatterning: New LC technology. *Crystals* **2020**, *10*, 323. 504

44. Chigrinov, V.; Kudreyko, A.; Sun, J. Flexible Optically Rewritable Electronic Paper. *Crystals* **2023**, *13*, 1283. 505

45. Schadt, M.; Schmitt, K.; Kozinkov, V.; Chigrinov, V. Surface-induced parallel alignment of liquid crystals by linearly polymerized photopolymers. *Japanese Journal of Applied Physics* **1992**, *31*, 2155. 506

46. West, J.L.; Su, L.; Reznikov, Y. Photo-alignment using adsorbed dichroic molecules. *Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals* **2001**, *364*, 199–210. 508

47. Yaroshchuk, O.; Reznikov, Y. Photoalignment of liquid crystals: basics and current trends. *Journal of Materials Chemistry* **2012**, *22*, 286–300. 510

48. McGinty, C.P.; Kołacz, J.; Spillmann, C.M. Large rewritable liquid crystal pretilt angle by in situ photoalignment of brilliant yellow films. *Applied Physics Letters* **2021**, *119*. 512

49. Yaroshchuk, O.; Gurumurthy, H.; Chigrinov, V.G.; Kwok, H.S.; Hasebe, H.; Takatsu, H. Photoalignment properties of brilliant yellow dye. In Proceedings of the Proc. IDW, 2007; pp. 1665–1668. 514

50. Yin, K.; Xiong, J.; He, Z.; Wu, S.T. Patterning liquid-crystal alignment for ultrathin flat optics. *ACS omega* **2020**, *5*, 31485–31489. 516

51. Folwill, Y.; Zeitouny, Z.; Lall, J.; Zappe, H. A practical guide to versatile photoalignment of azobenzenes. *Liquid Crystals* **2021**, *48*, 862–872. 518

52. Chigrinov, V.; Kudreyko, A.; Guo, Q. Patterned photoalignment in thin films: physics and applications. *Crystals* **2021**, *11*, 84. 520

53. Lall, J.; Zappe, H. In situ, spatially variable photoalignment of liquid crystals inside a glass cell using brilliant yellow. In Proceedings of the Photosensitive Materials and their Applications II, 2022; pp. 91–97. 522

54. Lovšin, M.; Petelin, A.; Berteloot, B.; Osterman, N.; Aya, S.; Huang, M.; Drevenšek-Olenik, I.; Mandle, R.; Neyts, K.; Mertelj, A. Patterning of 2D second harmonic generation active arrays in ferroelectric nematic fluids. *Giant* **2024**, *19*, 100315. 524

55. Fang, G.; Shi, Y.; MacLennan, J.E.; Clark, N.A.; Farrow, M.J.; Walba, D.M. Photo-reversible liquid crystal alignment using azobenzene-based self-assembled monolayers: comparison of the bare monolayer and liquid crystal reorientation dynamics. *Langmuir* **2010**, *26*, 17482–17488. 526

56. Fang, G.; MacLennan, J.; Yi, Y.; Glaser, M.; Farrow, M.; Korblova, E.; Walba, D.; Furtak, T.; Clark, N. Athermal photofluidization of glasses. *Nature communications* **2013**, *4*, 1521. 529

57. Wang, J.; McGinty, C.; West, J.; Bryant, D.; Finnemeyer, V.; Reich, R.; Berry, S.; Clark, H.; Yaroshchuk, O.; Bos, P. Effects of humidity and surface on photoalignment of brilliant yellow. *Liquid Crystals* **2017**, *44*, 863–872. 531

58. Shi, Y.; Zhao, C.; Ho, J.Y.L.; Song, F.; Chigrinov, V.G.; Luo, D.; Kwok, H.S.; Sun, X.W. High photoinduced ordering and controllable photostability of hydrophilic azobenzene material based on relative humidity. *Langmuir* **2018**, *34*, 4465–4472. 533

59. Dhara, S.; Madhusudana, N. Physical characterisation of 4'-butyl-4-heptyl-bicyclohexyl-4-carbonitrile. *Phase Transitions* **2008**, *81*, 561–569. 535

60. Zawadzki, A.; Walton, H. Measurements of the splay and bend elastic constants of 4'-butyl-4-heptyl-bicyclohexyl-4-carbononitrile, CCN47. *Molecular Crystals and Liquid Crystals* **2012**, *569*, 10–14. 537

61. Kleman, M.; Lavrentovich, O.D. *Soft Matter Physics: An Introduction*; Springer: 2003. 539

62. Chiccoli, C.; Feruli, I.; Lavrentovich, O.; Pasini, P.; Shiyanovskii, S.V.; Zannoni, C. Topological defects in schlieren textures of biaxial and uniaxial nematics. *Physical Review E* **2002**, *66*, 030701. 540

63. Taylor, J. *Introduction to error analysis, the study of uncertainties in physical measurements*, 2nd ed.; University Science Books: 1997. 542

64. 543

64. Ramdane, O.O.; Auroy, P.; Forget, S.; Raspaud, E.; Martinot-Lagarde, P.; Dozov, I. Memory-free conic anchoring of liquid crystals on a solid substrate. *Physical Review Letters* **2000**, *84*, 3871. 544

65. Chaplanova, Zh.D.; Murauski, A.A.; Rogachev, A.A.; Agabekov, V.E.; Gracheva, E.A. Multi-layered anisotropic films based on the azo dye brilliant yellow and organic polymers. *Journal of Applied Spectroscopy* **2013**, *80*, 658–662. 546

66. Schönhoff, M.; Mertesdorf, M.; Lösche, M. Mechanism of photoreorientation of azobenzene dyes in molecular films. *The Journal of Physical Chemistry* **1996**, *100*, 7558–7565. 548

67. Perny, S.; Barny, P.L.; Delaire, J.; Buffeteau, T.; Sourisseau, C.; Dozov, I.; Forget, S.; Martinot-Lagarde, P. Photoinduced orientation in poly (vinylcinnamate) and poly (7-methacryloyloxycoumarin) thin films and the consequences on liquid crystal alignment. *Liquid crystals* **2000**, *27*, 329–340. 550

68. Vilfan, M.; Olenik, I.D.; Mertelj, A.; Čopič, M. Aging of surface anchoring and surface viscosity of a nematic liquid crystal on photoaligning poly-(vinyl-cinnamate). *Physical Review E* **2001**, *63*, 061709. 553

69. Drzaic, P.S. *Liquid Crystal Dispersions*; World Scientific: 1995. 555

70. Andrienko, D.; Dyadyusha, A.; Iljin, A.; Kurioz, Y.; Reznikov, Y. Measurement of azimuthal anchoring energy of nematic liquid crystal on photoaligning polymer surface. *Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals* **1998**, *321*, 271–281. 556

71. Gu, D.-F.; Uran, S.; Rosenblatt, C. A simple and reliable method for measuring the liquid crystal anchoring strength coefficient. *Liquid Crystals* **1995**, *19*, 427–431. 559

72. Yang, F.; Cheng, H.; Gao, H.; Sambles, J. Technique for characterizing azimuthal anchoring of twisted nematic liquid crystals using half-leaky guided modes. *JOSA B* **2001**, *18*, 994–1002. 561

73. Zhang, B.; Sheng, P.; Kwok, H.S. Optical measurement of azimuthal anchoring strength in nematic liquid crystals. *Physical Review E* **2003**, *67*, 041713. 563

74. Polossat, E.; Dozov, I. New optical method for the measurement of the azimuthal anchoring energy of nematic liquid crystals. *Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals* **1996**, *282*, 223–233. 566

567

75. Fonseca, J.G.; Galerne, Y. Simple method for measuring the azimuthal anchoring strength of nematic liquid crystals. *Applied Physics Letters* **2001**, *79*, 2910-2912. 568

569

76. Vilfan, M.; Čopič, M. Azimuthal and zenithal anchoring of nematic liquid crystals. *Physical Review E* **2003**, *68*, 031704. 570

571