

### Key Points:

- The 2022 Sumatra event ruptured up-dip of the 2007 Bengkulu event on a near-trench shallowly dipping fault, producing a moderate tsunami
- Joint modeling of seismic and tsunami data resolves shallow slip of up to 4.1 m in low rigidity material, typical of tsunami earthquakes
- Global observations show that up-dip regions not ruptured in great megathrust events can host shallow damaging tsunami earthquakes

### Supporting Information:

Supporting Information may be found in the online version of this article.

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### Citation:

Xia, T., Ye, L., Bai, Y., Lay, T., Xu, S., Kanamori, H., et al. (2024). The 2022  $M_w$  7.3 southern Sumatra tsunami Earthquake: Rupture up-dip of the 2007  $M_w$  8.4 Bengkulu event. *Journal of Geophysical Research: Solid Earth*, 129, e2024JB030284. <https://doi.org/10.1029/2024JB030284>

Received 3 SEP 2024

Accepted 30 NOV 2024

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## The 2022 $M_w$ 7.3 Southern Sumatra Tsunami Earthquake: Rupture Up-Dip of the 2007 $M_w$ 8.4 Bengkulu Event

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**Abstract** On 18 November 2022, a large earthquake struck offshore southern Sumatra, generating a tsunami with 25 cm peak amplitude recorded at tide gauge station SBLT. Our  $W$ -phase solution indicates a shallow dip of  $6.2^\circ$ , compatible with long-period surface wave radiation patterns. Inversion of teleseismic body waves indicates a shallow slip distribution extending from about 10 km deep to near the trench with maximum slip of  $\sim 4.1$  m and seismic moment of  $1.05 \times 10^{20}$  Nm ( $M_w$  7.3). Joint modeling of seismic and tsunami data indicates a shallow rigidity of  $\sim 23$  GPa. We find a low moment-scaled radiated energy of  $4.15 \times 10^{-6}$ , similar to that of the 2010  $M_w$  7.8 Mentawai event ( $3.1 \times 10^{-6}$ ) and other tsunami earthquakes. These characteristics indicate that the 2022 event should be designated as a smaller moment magnitude tsunami earthquake compared to the other 12 well-documented global occurrences since 1896. The 2022 event ruptured up-dip of the 2007  $M_w$  8.4 Bengkulu earthquake, demonstrating shallow seismogenic capability of a megathrust that had experienced both a deeper seismic event and adjacent shallow aseismic afterslip. We consider seismogenic behavior of shallow megathrusts and concern for future tsunami earthquakes in subduction zones globally, noting a correlation between tsunami earthquake occurrence and subducting seafloor covered with siliceous pelagic sediments. We suggest that the combination of pelagic clay and siliceous sediments and rough seafloor topography near the trench play important roles in controlling the genesis of tsunami earthquakes along Sumatra and other regions, rather than the subduction tectonic framework of accretionary or erosive margin.

**Plain Language Summary** Tsunami earthquakes are shallow subduction zone events that occur near the trench, usually on the plate boundary, producing larger than typical tsunamis relative to their surface and body wave earthquake magnitudes. These events produce weak ground shaking, leaving coastal populations unaware of the event until a large tsunami arrives. They are also important indicators of the frictional state of the shallow megathrust fault. We use seismic and tsunami observations to determine the rupture process of the 18 November 2022 Sumatra event, which occurred up-dip of the 2007 magnitude 8.4 Bengkulu rupture. We find that the 2022 moment-magnitude 7.3 event shares common rupture characteristics with other recognized tsunami earthquakes and we designate it as a tsunami earthquake. We compare the tsunami earthquakes that have struck in different subduction zones globally, noting examples involving spontaneous failures of stress accumulations at shallow depths or events triggered by stress and strain perturbation from deeper great earthquakes. Additional endeavors to estimate the seismic coupling of shallow megathrusts are needed to evaluate the potential for tsunami earthquakes.

### 1. Introduction

Earthquakes that rupture the near-trench portion of megathrusts are rare but can result in large tsunamis relative to their body and surface wave seismic magnitudes, leading to their designation as tsunami earthquakes based on characteristics of their source spectra (e.g., Kanamori, 1972, 2014; Lay & Bilek, 2007; Okal & Newman, 2001; Polet & Kanamori, 2000; Ye et al., 2016a). Tsunami earthquakes have occurred in the shallow megathrusts of several subduction zones, including the 1896 Sanriku, 1907 Sumatra, 1946 Aleutian Islands, 1960 and 1996 Peru, 1963 and 1975 Kuril, 1992 Nicaragua, 1994 and 2006 Java, 2010 Mentawai (Sumatra), and 2012 El Salvador events. The 2021  $M_w$  8.2 South Sandwich Island earthquake has a large slow rupture component located between deeper fast components (Jia et al., 2022), so it can be viewed as a tsunami earthquake with broader depth extent

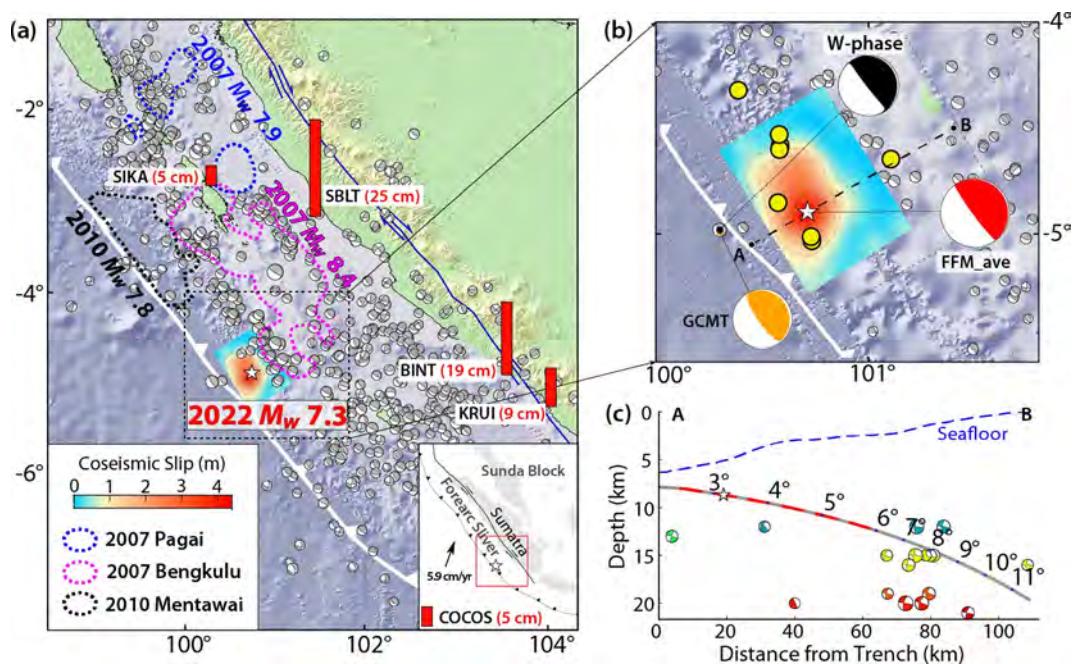
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than other tsunami earthquakes. These tsunami earthquakes exhibit distinctive rupture characteristics, including source spectra depleted in short period energy, low stress drop, slow rupture speed, enhanced slip at shallow depth with attendant strong excitation of tsunami waves, and aftershocks with normal faulting mechanisms (e.g., Kanamori, 2014; Lay & Bilek, 2007; Ye et al., 2016a). Various mechanisms have been proposed to explain one or more of these characteristics, such as subduction of a seamount (e.g., Abercrombie et al., 2001), variable frictional properties on the plate interface (e.g., Bilek & Lay, 2002), translation of sloping seafloor (e.g., Tanioka & Satake, 1996), anelastic deformation of the overlying sedimentary wedge (e.g., Du et al., 2021; Ma, 2012), low rigidity and low shear velocity of the shallow megathrust (e.g., Cheung et al., 2022; Polet & Kanamori, 2000; Sallarès & Ranero, 2019; Sallarés et al., 2021), release of gravitational potential energy (McKenzie & Jackson, 2012), free surface interactions (Gabuchian et al., 2017), plate bending (Oryan & Buck, 2020), and so on. However, the robustness of these candidate mechanisms still needs to be validated by further observations, as there are many regions where tsunami earthquakes could potentially strike.

The capacity of the shallow megathrust to store strain and rupture spontaneously depends on its physical characteristics and the stress state along the fault and within the outer accretionary wedge. These factors are intricately linked to the consolidation state and mineral content of accreted and underthrust pelagic sediments and to the deformation structures and width of the outer wedge of the frontal accretionary prism (Qiu & Barbot, 2022). Overlying sediments on top of the wedge can also influence shallow slip as a result of dynamic rupture interactions with energy radiated up from the megathrust (Li et al., 2023). The subducting sediments may govern the up-dip extent of seismic rupture during large megathrust events due to their unstable properties, thereby influencing the potential for devastating tsunami generation (e.g., Ruff & Kanamori, 1980). Specifically, thermally controlled chemical transformations including the smectite to illite transition (e.g., Pytte & Reynolds, 1989; Saffer et al., 2008), the opal to quartz transition (e.g., Spinelli & Underwood, 2004), and carbonate and zeolite cementation (e.g., Moore & Saffer, 2001), may occur within sediments, altering the frictional properties dependent on sediment mineral content, sediment consolidation, and fluid release (e.g., Ikari et al., 2018; Vrolijk, 1990). Some of these processes were inferred to influence shallow coseismic rupture during the 2004  $M_W$  9.2 Sumatra-Andaman (e.g., Dean et al., 2010; Geersen et al., 2013; Gulick et al., 2011; Hüpers et al., 2017; Stevens et al., 2021) and 2011  $M_W$  9.1 Tohoku (e.g., Fujie et al., 2020; Kameda et al., 2015; Ujiie et al., 2013) megathrust earthquakes, for which up-dip rupture that extended to the trench enhanced destructive tsunamis, essentially as co-seismic tsunami earthquake component of the total ruptures. In addition, the low shear velocity (or rigidity) near the toe of the wedge likely enhances the shallow slip for tsunami generation (Prada et al., 2021; Sallarès & Ranero, 2019; Sallarès et al., 2021), and the nature of the upper plate sedimentary wedge may control the volume above the underthrust slab in which low rigidity affects rupture properties of the tsunami earthquake. Qiu and Barbot (2022) establish that there is a correlation between maximum tsunami run-up for tsunami earthquakes and width of the outer wedge.

Subduction morphological condition is another key factor that appears to influence the occurrence of tsunami earthquakes. Kanamori and Kikuchi (1993) propose that a non-accreting (erosive) margin facilitates the rupture extending all the way to the trench, with examples being the 1960  $M_W$  7.6 Peru, 1963  $M_W$  7.8 and 1975  $M_W$  7.5 Kuril, and 1992  $M_W$  7.6 Nicaragua tsunami earthquakes. Minor incoming sediments are totally subducted along the plate-boundary interface in these regions, allowing slip to extend to the trench through a relatively weak and shallow plate-boundary interface covered by sediments. Polet and Kanamori (2000) argue that some sediment-starved trenches with small accretionary prisms and thin sediment can also support rupture nucleation and up-dip propagation in the shallow megathrust due to their rough horst-and-graben structure. Based on marine acoustic data, Geersen (2019) infers that tsunami earthquakes tend to occur in trenches characterized as sediment-starved with rough subducting plates.

Over the past decades, the Sunda megathrust has experienced several significant earthquakes, notably including the 2004  $M_W$  9.2 Sumatra-Andaman and 2005  $M_W$  8.6 Nias events, which ruptured the zone from 0°N to 14°N (e.g., Ammon et al., 2005; Hsu et al., 2006; Lay et al., 2005). The 1907  $M_W$  7.8 Sumatra earthquake located seaward of the 2005 Nias coseismic rupture, is regarded as a tsunami earthquake due to its shallow rupture, low rupture velocity, relatively long duration, and extensive tsunami impact (e.g., Kanamori et al., 2010; Martin et al., 2019). Kanamori et al. (2010) conclude, after reviewing historical seismograms of the 1907 earthquake, that it possibly originated as a thrust event on the subduction interface at a depth of ~30 km and then propagated up-dip to the trench, resulting in an extensive tsunami. Martin et al. (2019) derive the rupture area of the 1907 event located up-dip of the coseismic slip region of the 2005  $M_W$  8.6 Nias event and find that it had a maximum slip of



**Figure 1.** (a) Regional tectonic setting around the 2022  $M_w$  7.3 Sumatra earthquake near the Sunda trench. The white star indicates the epicenter of the earthquake from the U.S. Geological Survey National Earthquake Information Center (USGS-NEIC) catalog. Magenta, blue, and black outlined patches represent the rupture areas with slip  $>1$  m for the 12 September 2007  $M_w$  8.4 and  $M_w$  7.9 events (Konca et al., 2008) and the 25 October 2010 Mentawai  $M_w$  7.8 event (Yue, Lay, Rivera, Bai, et al., 2014), respectively. Focal mechanisms show seismicity with magnitude  $\geq 5.0$  from the Global Centroid Moment Tensor (GCMT) catalog since 1976. The red bars show five tide gauges with observed tsunami heights. The inset map shows the active tectonic framework in which the Australian plate subducts beneath the Sunda Block of the Eurasian plate along the Sunda trench. Black arrows represent the relative plate motion referenced to the Sunda Block (Demets et al., 2010). The barbed solid line represents the Sunda trench. (b) Coseismic slip distribution of the 2022 Sumatra earthquake, along with the average moment tensor of our preferred slip model (red), our W-phase solution (black), and the GCMT solution (orange). Yellow circles show 1-month aftershocks from the USGS-NEIC catalog. (c) Depth profile of the megathrust interface geometry (gray curve) with depth-varying dip from the Slab 2.0 model (Hayes et al., 2018) and varying seafloor bathymetry from BATNAS data (blue dashed curve). The red dashed curve indicates the interface extent used for slip inversion for the 2022 Sumatra earthquake. Focal mechanisms are for M5+ events in the dashed box area in panel (b) from the GCMT catalog, color-coded with the source depth.

$\sim 21$  m and extended  $\sim 220$  km southeastward along the shallow Sunda trench, based on the near- and far-field tsunami observations. Qiu and Barbot (2022) estimate a 30 km width to the outer wedge, with a large influx of trench-filled sediment in this accretionary region.

The plate interface in the Sumatra subduction zone from  $2^{\circ}$ S to  $5^{\circ}$ S near the Pagai Islands experienced several major ruptures in 2007 and 2010 (Figure 1) and previous great events in 1797  $M_w \sim 8.5$ – $8.7$  and 1833  $M_w \sim 8.6$ – $8.9$ . The estimated rupture zones of these events indicate strong interplate coupling beneath the Mentawai Islands from geodetic and paleogeodetic observations (e.g., Chlieh et al., 2008; Natawidjaja et al., 2006; Yue, Lay, Rivera, Bai, et al., 2014). The 2007  $M_w$  8.4 Bengkulu and 2007  $M_w$  7.9 Pagai events ruptured the central and deep portion of the megathrust, partially overlapping the slip zones of the 1797 and 1833 events, and the Pagai event may be correlated with accumulated slip deficit exceeding the slip that occurred during the 1797 earthquake (Chlieh et al., 2008). On 25 October 2010, the  $M_w$  7.8 Mentawai tsunami earthquake ruptured the shallow portion of the Sunda megathrust, up-dip of the northwestern region of the 2007  $M_w$  8.4 Bengkulu event, producing a stronger tsunami than the 2007 events (e.g., Borrero et al., 2009; Hill et al., 2012; Lay et al., 2011; Satake et al., 2013; Yue, Lay, Rivera, Bai, et al., 2014). The outer trench wedge width in this region is estimated as 35 km wide, with multiple splay faults in the wedge (e.g., Qiu & Barbot, 2024). The 2010 Mentawai event defied the common assumption that the up-dip region of prior great earthquakes that do not rupture to the trench is aseismic. It demonstrated that local frictional properties in the shallow megathrust could be capable of accumulating sufficient stress or receiving sufficient strain perturbation following great earthquakes downdip to fail

synchronously as for the 2011  $M_W$  9.1 Tohoku earthquake or fail subsequently in a tsunami earthquake as for the 2010 Mentawai earthquake following the 2007  $M_W$  8.4 Bengkulu event. Distributed faulting in the outer trench region likely reduces the loading rate for the décollement fault (Qiu & Barbot, 2024), resulting in longer intervals between tsunami events, but strain accumulation is indicated by the larger earthquakes that do occur.

The Sumatra earthquake of 18 November 2022 occurred on the shallow Sunda megathrust off southern Sumatra Island, Indonesia, at 13:37:08 UTC (Figure 1). The U.S. Geological Survey National Earthquake Information Center (USGS-NEIC) hypocenter for the 2022 Sumatra event is 4.90°S, 100.79°E, 25.0 km deep, roughly 204 km southwest of Bengkulu city (<https://earthquake.usgs.gov/earthquakes/eventpage/us7000iqpn/executive>). The short-period body wave magnitude was 6.1 and the 20-s period surface wave magnitude was 7.2 in the USGS-NEIC catalog. The 2022 Sumatra event ruptured the up-dip megathrust region of the southeastern region of the 2007  $M_W$  8.4 Bengkulu event and generated a moderate tsunami, which was recorded at four near-field tide gauge stations (Sriyanto et al., 2023) and one far-field tide gauge station (see Data and Resources). It occurred near the trench, providing an opportunity to further explore the seismogenic capability of the shallow megathrust of a subduction zone that previously experienced a great earthquake deeper on the megathrust. Here, we determine the source process of the 2022 Sumatra event using inversion of seismic data and joint modeling of tsunami waveforms, estimate the diagnostic source parameters of tsunami earthquakes, and then discuss diverse slip modes in the shallow megathrust of different subduction zones globally that have experienced major tsunami earthquakes.

## 2. Methods and Results

### 2.1. Long-Period Point-Source Solutions for the 2022 Sumatra Earthquake

Routine long-period moment tensor solutions from the USGS-NEIC (W-Phase) and Global Centroid Moment Tensor (GCMT) indicate shallow thrust faulting (Table S1 in Supporting Information S1) for the 2022 Sumatra event. Our W-phase moment-tensor solution (Kanamori & Rivera, 2008) of 157 three-component 5–10 mHz observations from 95 global broadband network stations indicates a predominantly double-couple point-source solution with  $M_0 = 4.39 \times 10^{19}$  Nm ( $M_W$  7.03) at a shallow centroid depth of 9.5 km,  $\varphi = 346.6^\circ$ ,  $\delta = 6.2^\circ$ ,  $\lambda = 114.4^\circ$ , and centroid time of 15.0 s. The observed long-period surface wave radiation patterns (Figure S1 in Supporting Information S1) have less-pronounced Love wave nodes near 150° and 330°, favoring the dip to be less than 10°, compatible with the W-phase solutions. The bootstrap analysis of the W-phase inversion demonstrates a well-resolved solution with a very shallow dip for the 2022 Sumatra event (Figure S2 in Supporting Information S1).

All of the long-period point-source solutions, summarized in Table S1 in Supporting Information S1, consistently indicate underthrusting focal mechanisms with the strike along the trench, relatively small dip angles ranging from 6.0° to 11.0°, and shallow source depths from 9.5 to 19.5 km. The estimated seismic moment is inversely proportional to the dip angle given that  $M_0 \sin 2\delta$  is roughly constant for shallow thrust events (Kanamori & Stewart, 1976). These long-period models all assume the Preliminary Reference Earth Model (PREM) velocity model for the source structure, thus the resolutions of dip angle and shallow source depth are limited. Also, pre-calculated Green's functions are usually fixed to depths of 9.5 km and deeper for the W-phase inversions.

### 2.2. Finite-Fault Inversion With Moment Constraint and Slab Geometry

We invert teleseismic  $P$  and  $S$  wave data for space-time slip models utilizing a least squares kinematic multi-time window finite-fault inversion method with variable rakes for specified fault model geometry and rupture speed (e.g., Hartzell & Heaton, 1983; Kikuchi & Kanamori, 1991; Ye et al., 2016b). We apply the inversion procedure with a seismic moment constraint, which minimizes the difference between the inverted total seismic moment and a reference value based on long-period waves (e.g., W-phase or GCMT). This stabilizes the inversion process and provides flexibility in the parameterization of the subfault source time functions with a number of overlapping triangles, and use of subfaults with relatively small grid size, as required for smaller events (Ye et al., 2016c). The inversion problem with spatial smoothing and moment constraint with corresponding normalized weights of  $\beta_1$  and  $\beta_2$  can be written as,

$$\begin{bmatrix} G \\ \lambda_1 G_1 \\ \lambda_2 G_2 \\ \lambda_3 G_3 \end{bmatrix} [m] = \begin{bmatrix} d \\ 0 \\ \lambda_2 M_0^1 \\ \lambda_3 M_0^2 \end{bmatrix} \quad (1)$$

where  $[m]$  and  $d$  are model parameters and data, respectively;  $G$  is the Green's function matrix;  $G_1$ ,  $G_2$ , and  $G_3$  are the coefficient matrix for spatial smoothing and one-row matrices for moment constraints, respectively.  $M_0^1$ ,  $M_0^2$  are decomposed components of a given seismic moment  $M_0$  in two directions (Ye et al., 2016c). The weight  $\lambda_1$ ,  $\lambda_2$ , and  $\lambda_3$  are given by

$$\lambda_1^2 = \beta_1 \frac{|G|^2}{|G_1|^2}, \lambda_2^2 = \beta_2 \frac{|G|^2}{|G_2|^2}, \lambda_3^2 = \beta_2 \frac{|G|^2}{|G_3|^2}, \quad (2)$$

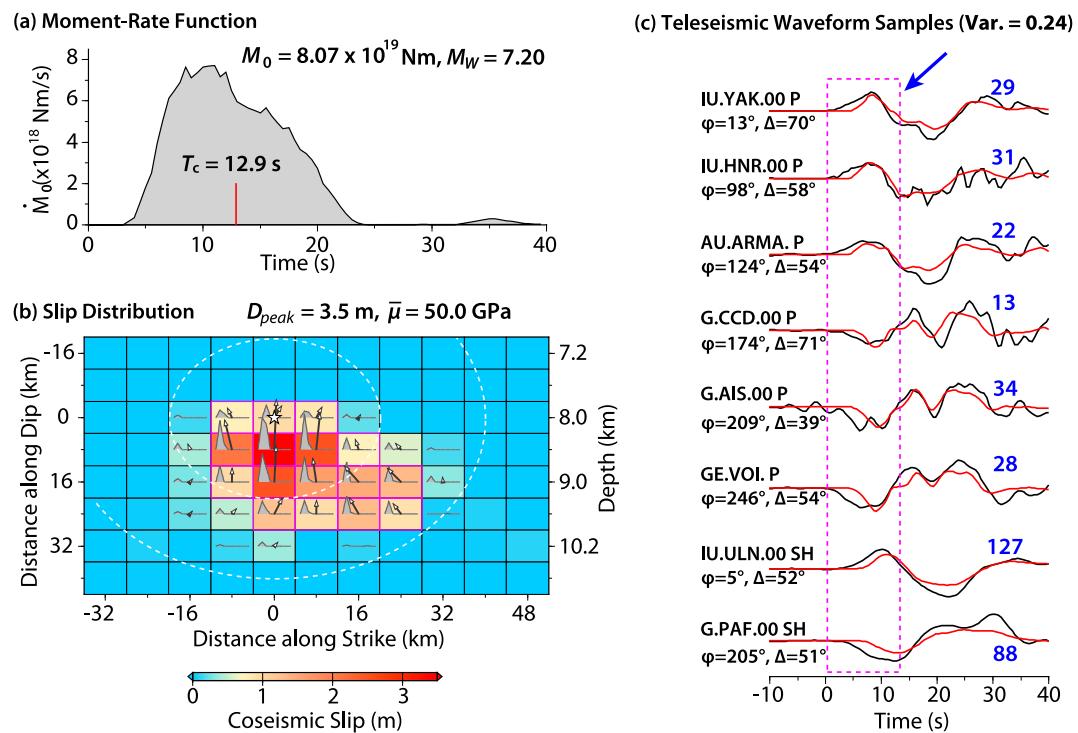
where  $\|$  is the quadratic norm of the matrix. We fix the normalized spatial smoothing weight  $\beta_1$  as 0.3, and vary the normalized seismic moment constraint weight  $\beta_2$  from 0.05 to 0.001 in the following iterative inversion section. We use the seismic moment  $M_0$  of  $3.84 \times 10^{19}$  Nm, which is the average of the  $W$ -phase solutions, as a reference value for the total seismic moment, and set  $M_0^1 = M_0 \cos(\lambda + 45^\circ)$ ,  $M_0^2 = M_0 \cos(\lambda - 45^\circ)$  with the rake angle  $\lambda = 114.4^\circ$  from our  $W$ -phase inversion.

We specify a 2.5-D dip-varying fault plane with 8 km by 8 km subfaults, megathrust dip ranging from  $2.8^\circ$  to  $5.5^\circ$  (Figure 1c) and an initial uniform strike of  $\varphi = 346.6^\circ$  indicated by the Slab 2.0 model (Hayes et al., 2018), long-period surface wave radiation patterns, and our  $W$ -phase solution. We adopt Model Crust 1.0 (Laske et al., 2013) as a local source velocity structure to calculate teleseismic Green's functions for a variable bathymetry interpolated from BATNAS (<https://tanahair.indonesia.go.id/portal-web/unduh/batnas>). Details of the local crustal structure are not known, and shallow velocities are varied in the analysis below. The subfault source time windows are parameterized by 8 overlapping 2.0 s rise-time symmetric triangles for a maximum possible duration of 18 s to account for higher-frequency details in the waveforms, and the peak rupture expansion speed is limited to 2.0 km/s (a range of speeds was considered, but it is not well resolved by the teleseismic ground displacement data). The hypocentral depth is set at 8.0 km constrained by the Slab 2.0 model and our  $W$ -phase solution, with the epicenter slightly adjusted from the USGS-NEIC location to  $(4.9^\circ\text{S}, 100.7^\circ\text{N})$  based on a grid search solution using regional arrival times (Figure S3 in Supporting Information S1).

### 2.3. Iterative Inversion and Modeling With Seismic and Tsunami Data

We employ seismic and tsunami observations to resolve the space-time history of the 2022 rupture. The seismic data set includes 85 vertical component  $P$  and 30 horizontal component  $SH$  teleseismic ground displacement recordings with well-distributed azimuthal coverage from global broadband stations, acquired from the Incorporated Research Institutions for Seismology (IRIS). We filtered the displacement recordings in the 0.005–0.9 Hz passband. The tsunami data include four tide gauge recordings from nearby stations along Sumatra (Sriyanto et al., 2023), as well as one far-field tide gauge recording at Cocos Island station from the European Commission Joint Research Centre, ISPRA, Space, Security and Migration Directorate (see Data and Resources).

Tsunami simulation is sensitive to the vertical seafloor deformation amplitude and area, which are proportional to the potency  $[\varphi \equiv \int D(x, y) dS]$  of the slip distribution  $D(x, y)$  (Heaton & Heaton, 1989). Accurately recovering reasonable slip amplitude and absolute fault location, which both affect tsunami excitation, is challenging in finite-fault inversions based only on teleseismic waves. Here, we apply an iterative inversion strategy, which has been applied previously to constrain self-consistent slip models using seismic, geodetic, and tsunami observations for large earthquakes (e.g., Yamazaki et al., 2011, 2013, 2018; Ye et al., 2016b, 2022; Yue et al., 2014, 2014a). We seek a preferred space-time slip model by inverting the seismic data and modeling the tsunami data, perturbing the structure and parameters in the inversion to iteratively seek a self-consistent model for both data sets. The iterative approach is used because of the non-linearity of changes in structural and inversion parameters needed to reconcile the data sets. We use NEOWAVE, a well-established forward computation code for the



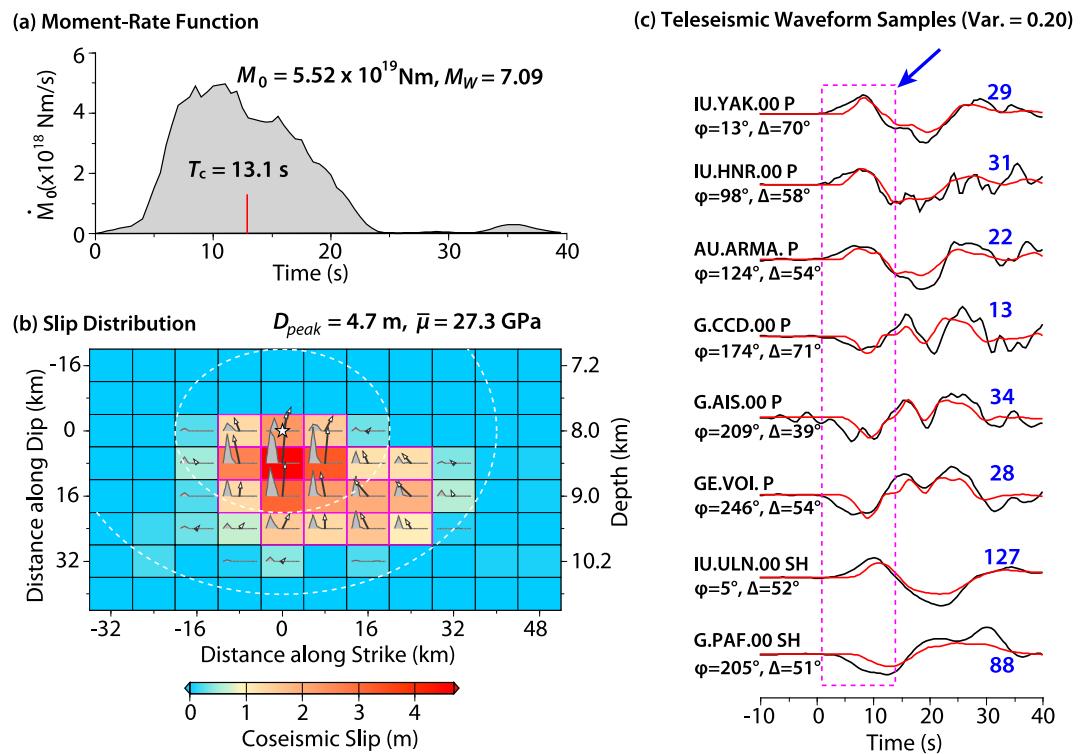
**Figure 2.** Slip model M01 using the Crust 1.0 source velocity structure with a modest seismic moment constraint ( $\beta_2 = 0.05$ ). (a) The moment rate function with a red tick at the centroid time  $T_c = 12.9 \text{ s}$ . (b) Slip distribution with arrows showing the magnitude and direction of slip (hanging-wall relative to footwall) and subfaults color-coded by peak slip. The subfault source time functions are shown within each subfault by gray polygons. White-dashed circles indicate the position of the rupture front in 10 s intervals. (c) Comparisons of the observed (black) and predicted (red) ground displacement waveforms for a subset of seismic waves. The purple dashed rectangular indicates the imprecise wave fits within the initial 15 s.

tsunami simulations (e.g., Bai et al., 2018; Liu et al., 2023; Ye et al., 2021, 2022). In this section, we discuss how we derive our preferred model through a sequence of representative steps.

The initial slip model M01 (Figure 2) is obtained by fixing the normalized moment constraint weight  $\beta_2$  at a relatively small value of 0.05 because the data fit degrades as  $\beta_2$  increases. Model M01 has an average rigidity  $\bar{\mu}$  of 50.0 GPa, a peak slip  $D_{\text{peak}}$  of 3.5 m, a potency  $\varphi$  of  $\sim 2.1 \text{ km}^3$  over the total fault area. It has an average slip  $\bar{D}_{0.15}$  of 1.4 m over a summed area of  $1088 \text{ km}^2$  for patches where the slip is greater than 15% of the  $D_{\text{peak}}$ , that is slip  $\geq \sim 0.5 \text{ m}$ , along with the corresponding potency  $\varphi_{0.15}$  of  $1.5 \text{ km}^3$  for that area (Table S2 in Supporting Information S1). The estimated seismic moment  $M_0$  of M01 is  $8.07 \times 10^{19} \text{ Nm}$ , substantially larger than that derived from the  $W$ -phase solutions due to the relatively small  $\beta_2$ . Model M01 fails to fit the seismic waveforms within the initial 15 s (Figure 2), and tsunami simulations for M01 are also unable to fit the first tsunami wave arrivals, showing relatively lower tsunami amplitudes (Figure S4a in Supporting Information S1).

A common, but imprecise practice for scaling slip models to match tsunami amplitudes is to simply enhance the slip in proportion to an assumed rigidity reduction while keeping the seismic moment fixed, but without using a corresponding seismic velocity model to reinvert the seismic data. We gradually increase the overall slip relative to that in model M01 to 180% (in 20% intervals) to determine how much stronger seafloor motion resulting from higher peak and average slip increases is needed to match the tsunami amplitudes. We find that a perturbed Test model with  $D_{\text{peak}}$  of 5.6 m and  $\bar{D}_{0.15}$  of 2.2 m from increasing the overall slip in model M01 by 160% is able to match the tsunami signal amplitudes at SBLT and BINT well (Table S2 in Supporting Information S1), but this model is intrinsically not consistent with the seismic inversion.

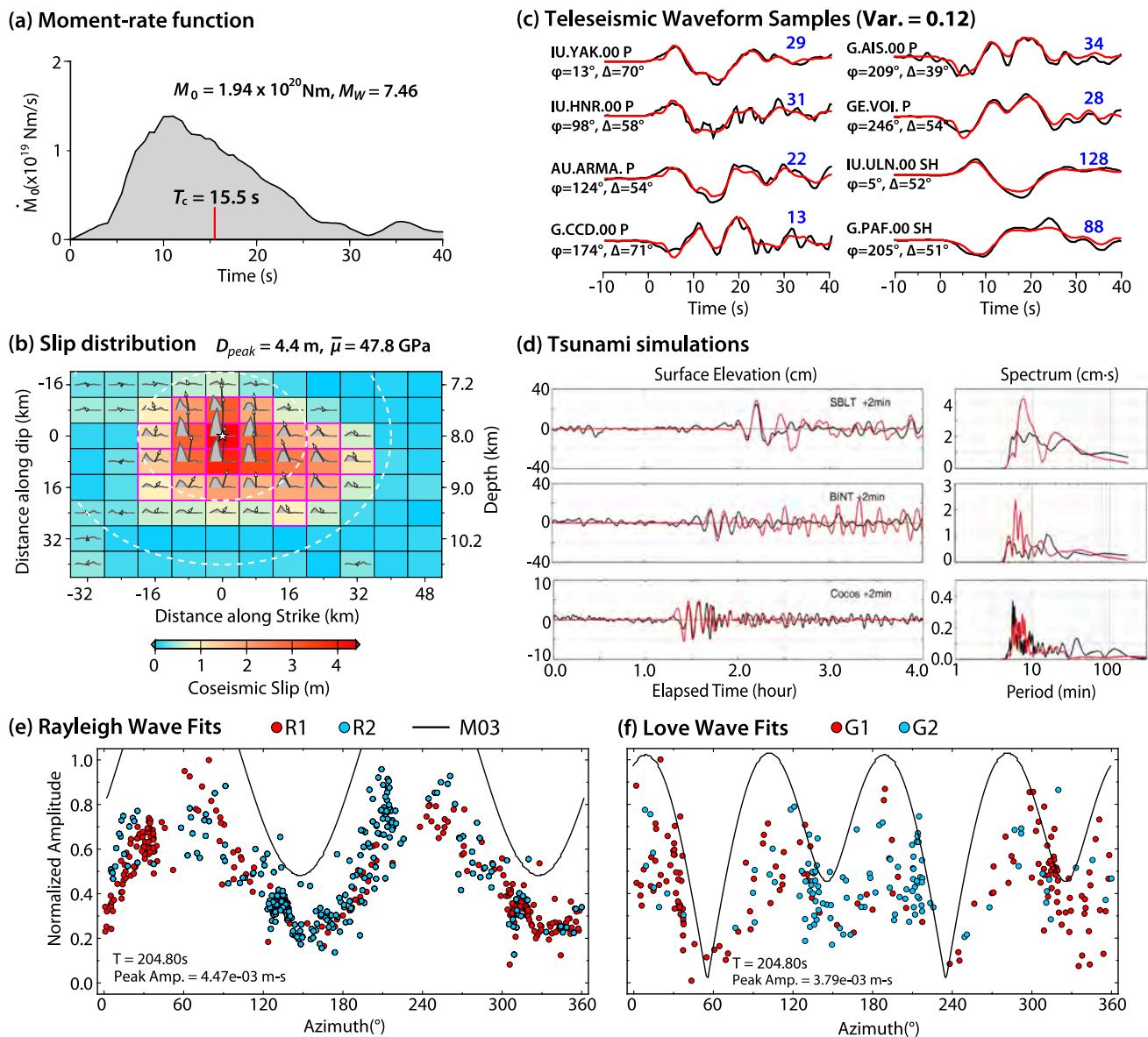
To seek a self-consistent model, we invert for slip model M02 (Figure 3) which has enhanced peak slip and average slip close to those for the Test model as a result of reducing the Crust 1.0 crustal seismic velocities by a factor of 0.79. This lowers the average rigidity in the source region (Table S3 in Supporting Information S1) for



**Figure 3.** Slip model M02 with lower source velocity and rigidity obtained by multiplying  $V_p$  and  $V_s$  of Crust 1.0 model by a factor of 0.79 and using a modest seismic moment constraint ( $\beta_2 = 0.05$ ). (a) The moment rate function with a red tick at the centroid time  $T_c = 13.1 \text{ s}$ . (b) Slip distribution with arrows showing the magnitude and direction of slip (hanging-wall relative to footwall) and subfaults color-coded by peak slip. The subfault source time functions are shown within each subfault by gray polygons. White-dashed circles indicate the position of the rupture front in 10 s intervals. (c) Comparisons of the observed (black) and predicted (red) ground displacement waveforms for a subset of seismic waves. The purple dashed rectangular indicates the imprecise wave fits within the initial 15 s.

fixed density and unchanged layer thicknesses. Compared with M01, M02 has a lower  $\bar{\mu}$  of 27.3 GPa, higher  $D_{\text{peak}}$  of 4.7 m and  $\bar{D}_{0.15}$  of 1.7 m, slightly higher  $\varphi$  of 2.8  $\text{km}^3$  and  $\varphi_{0.15}$  of 2.0  $\text{km}^3$ , along with a lower  $M_0$  of  $5.52 \times 10^{19} \text{ Nm}$  (Table S2 in Supporting Information S1). However, as apparent in Figure 3c, model M02 still fails to adequately fit the early portion of the seismic waves. The modest 30% increase of the potency in the model M02, which is less than the 60% increase for the Test model, proves insufficient to address this issue well. This is attributed to the weak excitation of seismic waves in very low-velocity material for a very shallow fault dip, which requires a larger moment to match the far-field seismic waves. We note that the moment constraint employed in M01 and M02 inhibits the increase in seismic moment and potency.

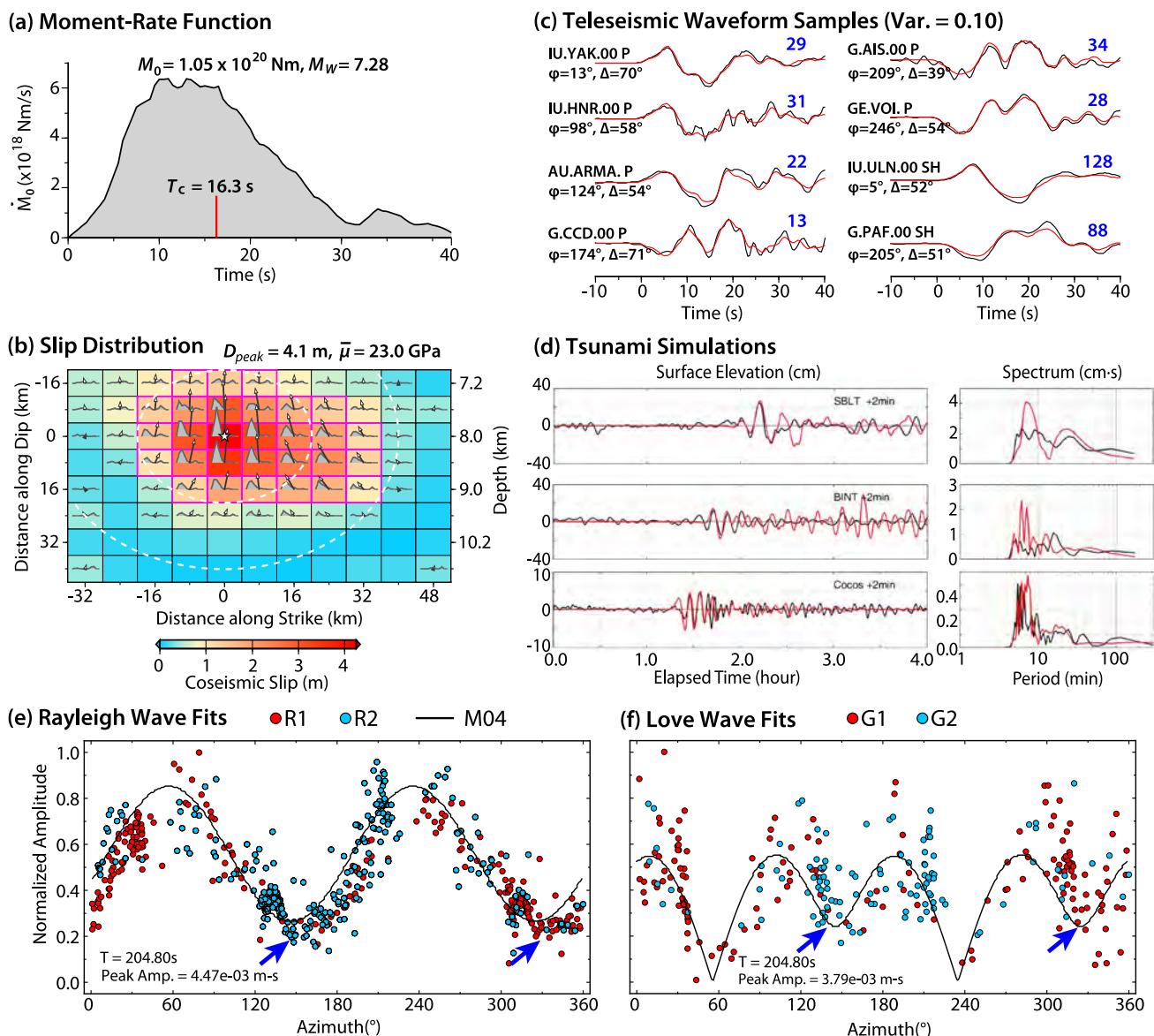
We next relax the total moment constraint by fixing  $\beta_2$  at a negligible value of 0.001 to invert for slip model M03 (Figure 4), using the same Crust 1.0 velocity model as M01. We adjust the strike from  $346.6^\circ$  to  $328.0^\circ$  to better match the tsunami signals. Model M03 has an  $\bar{\mu}$  of 47.8 GPa, a  $D_{\text{peak}}$  of 4.4 m,  $\varphi$  of 4.8  $\text{km}^3$  for total area, a  $\bar{D}_{0.15}$  of 1.8 m, and  $\varphi_{0.15}$  of 3.5  $\text{km}^3$ , along with  $M_0$  of  $1.94 \times 10^{20} \text{ Nm}$ , 2.5 times larger than for M01 (Table S2 in Supporting Information S1). While the peak slip of M03 is lower than for M01 and M02, it significantly improves the match to the early seismic waveforms (Figure 4c) and the initial tsunami wave arrivals at SBLT and BINT stations (Figure 4d). This notable improvement in fit is attributed to the significantly increased potency, which is  $\sim 2.3$  times larger than the value for M01, arising from the effects by lowering the moment constraint (Table S2 and Figure S5 in Supporting Information S1). We additionally validate model M03 using simplified point-source forward modeling of the long-period Rayleigh and Love wave spectral amplitudes (Figures 4e and 4f). We follow the spectral method developed by Kanamori and Given (1981) to compare measured long-period surface wave radiation patterns at a period of  $T = 204.8 \text{ s}$  with calculations using the PREM velocity model for the GCMT solution, our  $W$ -phase solution, model M03, and the following M04 model. We find that model M03 fails to match the long-period surface wave spectra, yielding significantly higher spectral amplitude compared to both GCMT



**Figure 4.** Slip model M03 using the Crust 1.0 source velocity model and a strongly reduced seismic moment constraint ( $\beta_2 = 0.001$ ). (a) The moment rate function with a red tick at the centroid time  $T_c = 15.5 \text{ s}$ . (b) Slip distribution with arrows showing the magnitude and direction of slip (hanging-wall relative to footwall) and subfaults color-coded by peak slip. The subfault source time functions are shown within each subfault by gray polygons. White-dashed circles indicate the position of the rupture front in 10 s intervals. (c) Comparisons of the observed (black) and predicted (red) ground displacement waveforms for a subset of seismic waves. (d) Comparison of recorded (black) and computed (red) sea level time series (left) and amplitude spectra (right) at SBLT, BINT, and Cocos tide gauges. The computed time series have been shifted by +2 min for each station, correcting for errors in average path bathymetry and approximations in the modeling. (e)–(f) Observed (dots) and predicted (black curves) Rayleigh wave and Love wave source spectral amplitudes from the model M03.

and our  $W$ -phase solutions (Figures 4e and 4f, and Figure S1 in Supporting Information S1). This discrepancy is attributed to the considerably larger seismic moment in a high-rigidity source velocity structure for Model M03.

Consequently, we reduce the average source region rigidity to 23.0 GPa by decreasing  $V_p$  and  $V_s$  in the source region by a factor of 0.7 (models for a suite of factors were considered) relative to model M01 and obtain slip model M04 (Figure 5). Model M04 has a seismic moment of  $1.05 \times 10^{20} \text{ Nm}$  ( $M_W = 7.3$ ), peak slip  $D_{\text{peak}}$  of 4.1 m, total potency  $\varphi$  of  $5.5 \text{ km}^3$  and truncated potency  $\varphi_{0.15}$  of  $4.0 \text{ km}^3$  (Table S2 in Supporting Information S1). Because of the enhanced potency and the decreased seismic moment relative to M03, model M04 model now fits both the initial 15 s of the seismic waves (Figure 5c) and tsunami observations (Figure 5d) well, while satisfactorily matching the

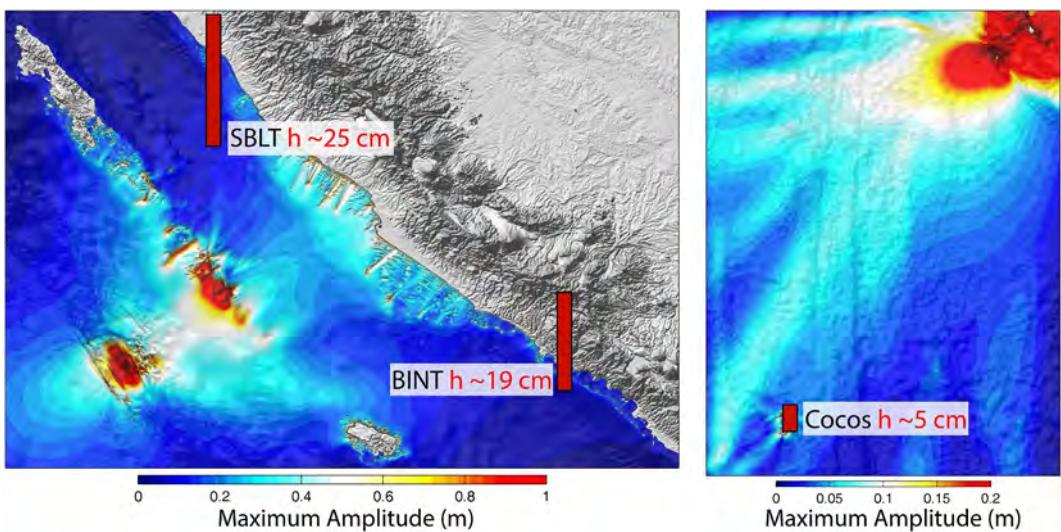


**Figure 5.** Preferred slip model M04 with Crust 1.0 source velocities being multiplied by a factor of 0.7 and strong reduction of the seismic moment constraint ( $\beta_2 = 0.001$ ). The average rigidity above 10 km depth is 23 GPa. (a) The moment rate function with a red tick at the centroid time  $T_c = 16.3$  s. (b) Slip distribution with arrows showing the magnitude and direction of slip (hanging-wall relative to footwall) and subfaults color-coded by peak slip. The subfault source time functions are shown within each subfault by gray polygons. White-dashed circles indicate the position of the rupture front in 10 s intervals. (c) Comparisons of the observed (black) and predicted (red) ground displacement waveforms of a subset of seismic waves. (d) Comparison of recorded (black) and computed (red) sea level time series (left) and amplitude spectra (right) at SBLT, BINT, and Cocos tide gauges. The computed time series have been shifted by 2 min for each station, approximately correcting for errors in average path bathymetry and approximations in the modeling. (e)–(f) Observed (dots) and predicted (black curves) Rayleigh wave and Love wave source spectral amplitudes from the model M04. The blue arrows indicate the symmetry in the Rayleigh wave radiation nodes (e) and the asymmetry in the Love wave radiation nodes (f).

long-period surface wave spectra with a point-source approximation (Figures 5e and 6f). We therefore prefer model M04 as a self-consistent representation of the source for the 2022 Sumatra earthquake.

#### 2.4. Characteristics of the Preferred Slip Model

Our preferred slip model M04 for the 2022 Sumatra earthquake (Figure 5) has an average rake of  $94.4^\circ$  and a rough upper bound on the rupture speed of  $2.0$  km/s. The moment rate function (Figure 5a) has a centroid time of  $16.3$  s. The seismic moment of  $1.05 \times 10^{20}$  Nm ( $M_W 7.28$ ) exceeds those estimated by the GCMT



**Figure 6.** Tsunami simulation results. Maximum sea surface elevation maps near the source region predicted by our preferred finite-fault model M04. The dark red bars of different sizes represent the tsunami heights observed at SBLT, BINT, and Cocos tide gauge stations.

( $2.81 \times 10^{19}$  Nm,  $M_W$  6.90), the USGS-NEIC *W*-phase ( $3.29 \times 10^{19}$  Nm,  $M_W$  6.94), and our *W*-phase ( $4.39 \times 10^{19}$  Nm,  $M_W$  7.03). This difference is caused by a combination of the very small dip ( $<4^\circ$ ) in the region where most slip is concentrated (Figure 5b) (even lower dip than the point-source model values), the shallow depth into rock (Figure 1c), and the low velocities and rigidity in the crustal model (the source model is listed in Table S1 in Supporting Information S1). The slip model has a centroid depth of 8.3 km below sea level, with an average water depth of about 4 km over the region of large slip. The average rigidity at the depths with large slip is 23 GPa, which reconciles the seismic and tsunami wave excitation. The large slip has a shallow bilateral concentration extending  $\sim$ 54 km along the strike, extending seaward to near the trench along the dip, with a maximum slip of  $\sim$ 4.1 m. All well-resolved slip is at depths shallower than 10 km.

The preferred slip model M04 fits the complete set of *P*-wave and *SH*-wave data well, accounting for approximately 91% of the power in the observations (Figure S6 in Supporting Information S1). The waveforms are relatively smooth and simple, but have long *P*-coda waves associated with deep water *pwP* reverberations. Our preferred slip model fits some of the *P*-coda and depth phases well due to the use of a 2.5D dip-varying fault model with a gradient in overlying water depth. A fully 3D modeling effort could possibly fit the later coda, but the signals confirm the concentration of slip below deep water. We find similar waveform fits and slip distributions for models with maximum rupture speeds of 1.0–4.0 km/s, indicating limited resolution of rupture expansion speed. This is typical of bilateral ruptures with moderate finiteness. The basic depth range of slip and absolute fault placement is partially constrained by the joint modeling of the tsunami signals. However, high-frequency *P*-wave motions are not all well predicted, probably due to limitations of the simplified source and receiver velocity models and Green's functions that we use for our ground displacement seismograms.

Maximum sea surface elevation maps near the source region and Cocos Island predicted for model M04 are shown in Figure 6, and comparisons between tsunami predictions and observations at SBLT, BINT, and Cocos Island tide gauge stations are shown in Figure 5d. Station SBLT is located closest to the source and recorded a peak tsunami wave first arrival amplitude of 25 cm, while station BINT located on the southeastern coast of Sumatra Island recorded a maximum amplitude of 19 cm. The Cocos Island station, located far from the source, reported a maximum amplitude of 5 cm. The arrival time of the first modeled peak at these three stations is delayed by 2 minutes to align with the observations reasonably, with this shift approximately correcting for errors in the overall bathymetric structure. The fits to these signals confirm the reliability of the seafloor deformation pattern (Figure S5 in Supporting Information S1) produced by our preferred slip model. We are unable to predict the tsunami signal at station SIKA because of the inaccurate available bathymetry information (Figure S7 in Supporting Information S1), and we cannot model the tsunami signal at KRUI well either, due to its low signal-noise ratio (Figure S8 in Supporting Information S1). Overall, the tsunami signals at SBLT, BINT, and Cocos

Island indicate compatibility with the relatively large seismic moment and shallow slip of our preferred model. We also validate our preferred slip model M04 using forward modeling of long-period surface wave spectral amplitude for Rayleigh and Love waves, as discussed in the last section. These comparisons (Figures 5e and 5f) indicate that our preferred model accounts for the asymmetry in the long-period Love wave nodes near 60°/240°, and 150°/330° sufficiently, somewhat better than the point-source moment tensor solutions. The spectral fits for the enhanced moment and reduced dip of the preferred model geometry confirm the viability of our estimated  $M_W = 7.3$ .

## 2.5. Tsunami Earthquake Designation

We estimate a radiated energy of  $E_R = 4.36 \times 10^{14}$  J for the 2022 event by combining contribution from the moment rate function of our preferred slip model M04 for periods longer than 20 s and from stacked broadband  $P$  wave spectra for shorter periods corrected for radiation pattern and propagation effects (e.g., Ye et al., 2016a). The estimated source spectra with the seismic moment from our preferred slip model and  $W$ -phase seismic moment (Figure 7a) both indicate a depletion of high-frequency seismic wave radiation for frequencies greater than 0.1 Hz, which is a consistent characteristic of tsunami earthquakes. Figure 7b shows a comparison of the moment-scaled radiated energy for major interplate earthquakes deeper on the megathrust (Ye et al., 2016c) and shallow interplate tsunami earthquakes. The estimated moment-scaled radiated energy  $E_R/M_0 = 4.15 \times 10^{-6}$  for the 2022 Sumatra earthquake is close to the values for recognized large well-recorded tsunami earthquakes which range from  $\sim 1 \times 10^{-6}$  to  $4 \times 10^{-6}$  (Ye et al., 2016c). Additionally, we note that the low average rigidity of 23 GPa found for our preferred slip model is consistent with the source having a low rupture speed and depleted high-frequency seismic radiation. The static slip-weighted stress drop  $\Delta\sigma_E = 1.41$  MPa, and the factor of 0.15 trimmed-slip circular stress drop  $\Delta\sigma_{0.15} = 2.30$  MPa (following Ye et al., 2016c), are both smaller than the average for deeper megathrust events ( $\sim 3.4$ –4.6 MPa, Ye et al., 2016c). Thus, we find that the 2022 Sumatra event, albeit a moderate-size rupture, involved large shallow slip in low-rigidity material, and its rupture shares several common features with tsunami earthquakes. This includes relatively strong excitation of tsunami waves by large shallow slip, depletion in high-frequency seismic wave radiation, low moment-scaled radiated energy, low rigidity, and low stress drop (e.g., Lay & Bilek, 2007; Ye et al., 2016c). While not a devastating tsunami event, the 2022 event is similar in size to the 2012 El Salvador earthquake which produced a moderate tsunami (Ye et al., 2013), and thus both events can be identified as tsunami earthquakes.

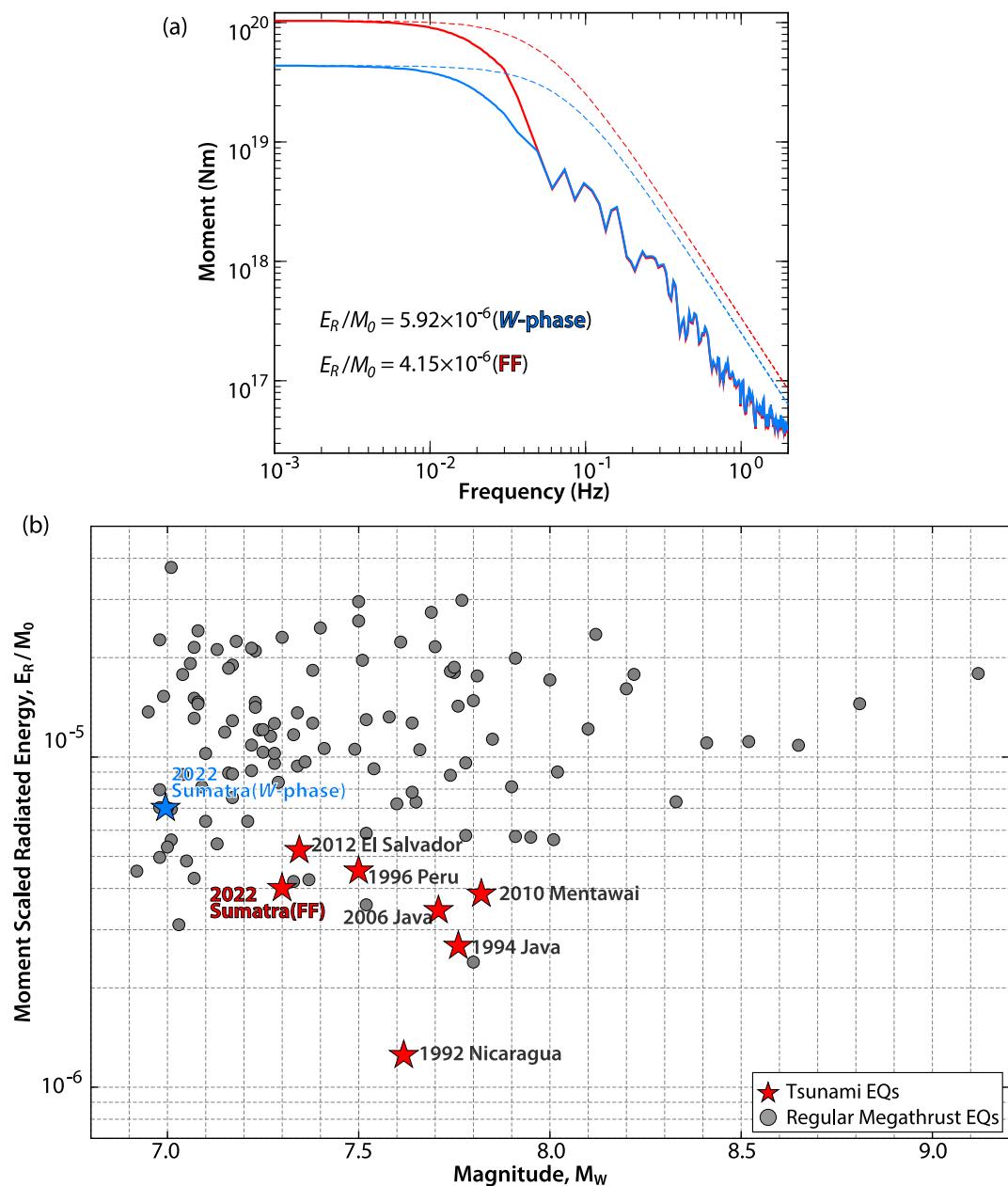
## 2.6. Coulomb Failure Stress Change From the Previous Earthquakes

The Mentawai segment of the Sunda trench has experienced several great megathrust earthquakes in its history (Figure 8a) and is prone to be relatively strongly locked as inferred from local GPS data (Chlieh et al., 2008). During the period from the 12 September 2007  $M_W 8.4$  Bengkulu earthquake to the 18 November 2022 Sumatra  $M_W 7.3$  earthquake, most aftershocks of the 2007  $M_W 8.4$  Bengkulu and the 2007  $M_W 7.9$  earthquakes were distributed within and around the boundary of the large-slip areas for the 2007 mainshocks and did not extend to the shallow megathrust, consistent with the absence of shallow rupture. There was very sparse seismicity (M4.5+) within the coseismic slip area of the 2022  $M_W 7.3$  Sumatra event and the northern Mentawai segment near Siberut island (Figure 8). There has not been a great earthquake in the northern Mentawai segment since 1797, so Haridhi et al. (2018) designate it as the North Mentawai Seismic Gap. We explore the stress perturbation in the area of the 2022  $M_W 7.3$  Sumatra earthquake and the North Mentawai Seismic Gap from the 2007  $M_W 8.4$  Bengkulu, 2007  $M_W 7.9$  Pagai, and 2010  $M_W 7.8$  Mentawai earthquakes by evaluating the change of Coulomb failure stress ( $\Delta\text{CFS}$ ).

The  $\Delta\text{CFS}$  involves projecting stress perturbation on a receiver fault resulting from earlier coseismic slip on a nearby fault, which can be written as,

$$\Delta\text{CFS} = \Delta\tau + \mu'\Delta\sigma_n$$

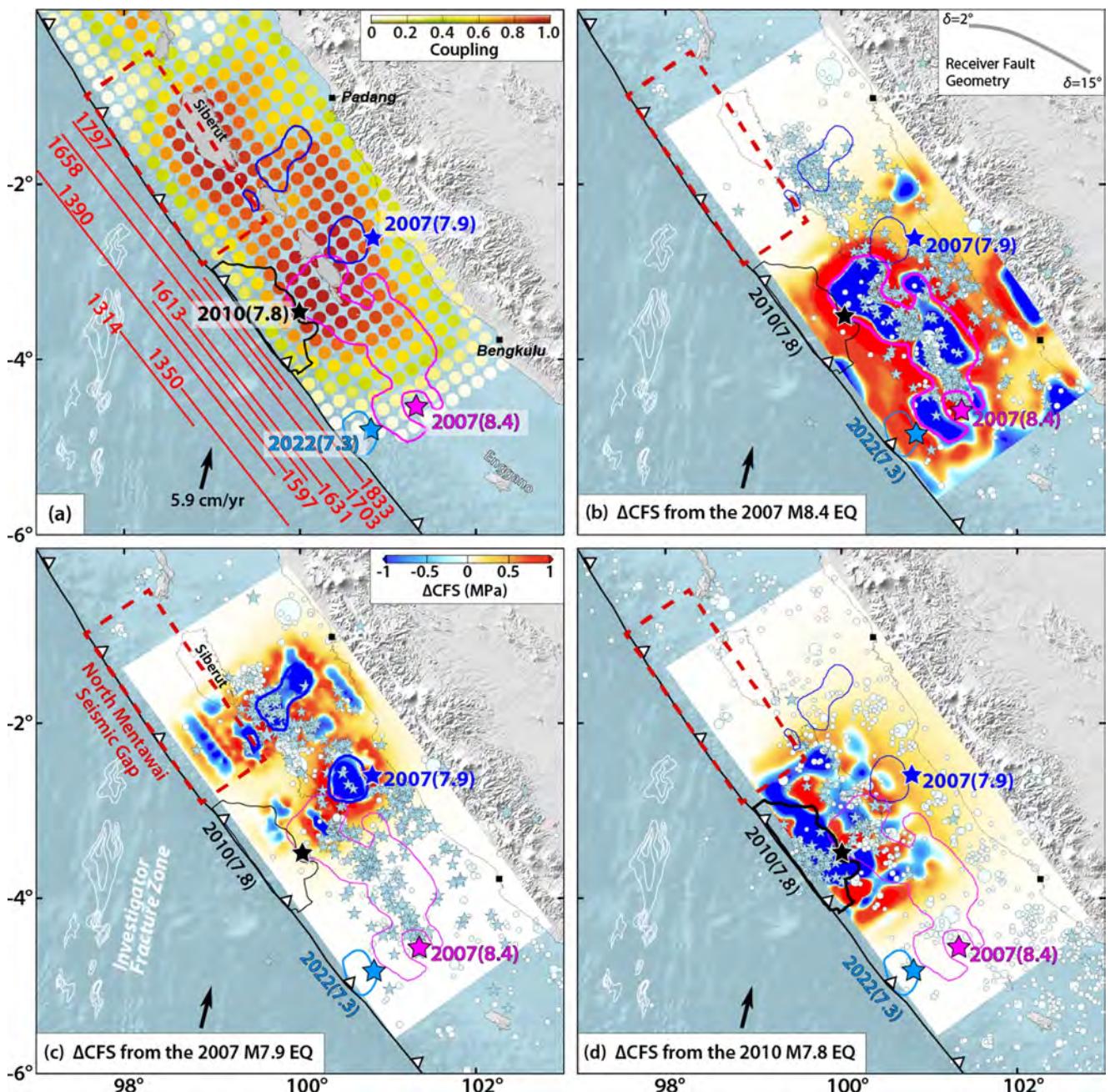
where  $\Delta\tau$  is the shear stress change (positive in the slip direction),  $\Delta\sigma_n$  is the normal stress change (positive for unclamped fault),  $\mu'$  is the effective frictional coefficient and set to be 0.4 in this study. We calculate  $\Delta\text{CFS}$  from the 2007  $M_W 8.4$  Bengkulu earthquake (Konca et al., 2008) on a receiver fault along the slab interface using the Slab 2.0 model. Our results show an increase of the Coulomb stress in the 2022  $M_W 7.3$  rupture zone of  $\sim 0.5$ –



**Figure 7.** (a) Source spectra for the 2022 Sumatra event obtained from teleseismic  $P$  spectra and the moment-rate spectrum with different corrections of the seismic moment from the  $W$ -phase solution (bold dashed blue curve) and our model M04 (red curve), along with corresponding reference  $\omega^{-2}$  source spectrum for 3 MPa stress drop (fine dashed curves). (b) Comparisons of moment-scaled radiated energy ( $E_R/M_0$ ) for major tsunami earthquakes (stars) and other major megathrust earthquakes from 1990 to 2016 (gray dots; Ye et al., 2016a).

1.5 MPa, as was also the case for the 2010  $M_W$  7.8 Mentawai tsunami earthquake rupture zone (Figure 8b). We calculate the Coulomb failure stress change caused by the 2007  $M_W$  7.9 Pagai and 2010  $M_W$  7.8 Mentawai earthquakes along the slab interface, showing no significant increase in the Coulomb failure stress change in the 2022  $M_W$  7.3 rupture zone (Figures 8c and 8d).

Thus, considering the cumulative contribution of stress changes due to the 2007  $M_W$  8.4 Bengkulu coseismic slip and the ensuing shallow afterslip in the shallow portion of the megathrust adjacent to the 2010  $M_W$  7.8 Mentawai and 2022  $M_W$  7.3 Sumatra earthquake rupture areas, it is evident that the 2022  $M_W$  7.3 Sumatra earthquake was



**Figure 8.** (a) Distribution of interseismic coupling (Chlieh et al., 2008) and historical great earthquakes in the Mentawai segment along the Sunda megathrust. Bold magenta, blue, black, and sky blue contours represent the rupture areas of the 2007  $M_W$  8.4 Bengkulu (Konca et al., 2008), 2007  $M_W$  7.9 Pagai (Konca et al., 2008), 2010  $M_W$  7.8 Mentawai (Yue, Lay, Rivera, Bai, et al., 2014), and 2022  $M_W$  7.3 Sumatra (this study) earthquakes with stars showing corresponding epicenters locations. Red bars show along-strike extension for historical great earthquakes from Philibosian et al. (2017). (b)–(d) Coulomb stress changes along the megathrust from the coseismic slip of the 2007  $M_W$  8.4 Bengkulu (Konca et al., 2008), 2007  $M_W$  7.9 Pagai (Konca et al., 2008), and 2010  $M_W$  7.8 Mentawai (Yue, Lay, Rivera, Bai, et al., 2014) earthquakes, respectively. The receiver fault is along the slab interface using the Slab 2.0 model, having strike  $\varphi = 324.0^\circ$ , rake  $\lambda = 99.0^\circ$  and a varying dip  $\delta$  from 2° to 15.0°. Dark red dashed rectangle shows the North Mentawai Seismic Gap (Haridhi et al., 2018). White circles in (b) and (c) represent  $M_{4.5+}$  seismicity from the 2007 Bengkulu earthquake to the 2010 Mentawai earthquake. White circles in (d) represent  $M_{4.5+}$  seismicity from the 2010 Mentawai earthquake to the 2022 Sumatra earthquake. Light blue stars show  $M_{4.5+}$  aftershocks within 3 months after each mainshock.

promoted by the stress transfer from the deeper 2007  $M_W$  8.4 Bengkulu event. The downdip 2007  $M_W$  7.9 Pagai rupture exerts positive and negative Coulomb failure stress change in the North Mentawai Seismic Gap of  $\sim 0.5$ –1.0 MPa (Figure 8c), but the region was not as strongly or uniformly perturbed as the 2010 and 2022 rupture regions were by the 2007  $M_W$  8.4 Bengkulu event.

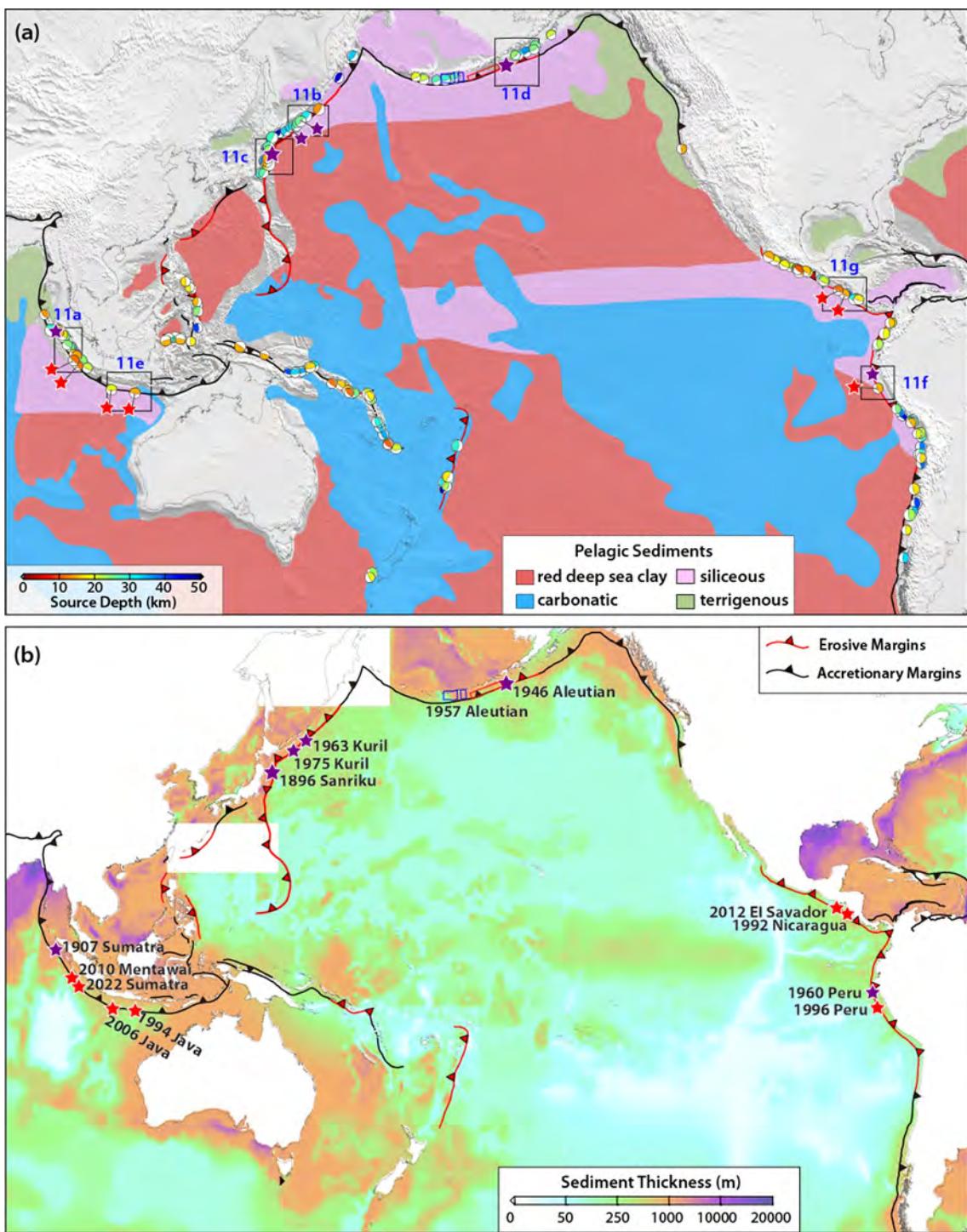
### 3. Discussion

Given that the near-trench megathrust was previously widely considered to be aseismic (e.g., Byrne et al., 1988; Scholz, 1998), the mechanism for shallow tsunami earthquakes remains unclear. A critical issue is whether there are local frictional heterogeneities in the shallow megathrust capable of accumulating sufficient stress to fail spontaneously in a tsunami earthquake. To that end, we consider the seismogenic behavior of the well-documented tsunami earthquakes since 1896 with good coverage and historical seismograms that have struck in different subduction zones globally.

We summarize the relatively sparse sampling provided by documented tsunami earthquakes, and their occurrence in both continental and island arc settings (Figure 9). There is a wide range of subducting lithosphere ages among the subduction zones that have hosted tsunami earthquakes. Figure 9 highlights one first-order correlation among almost all tsunami earthquakes which is that they have occurred in regions where siliceous pelagic sediments carpet the oceanic plates, but of course, only limited portions of such regions have produced recent tsunami earthquakes, so any causal relationship is unclear. Some destructive tsunamigenic earthquakes including shallow slip on the megathrust covered by siliceous pelagic sediments have been studied, such as the 2004  $M_W$  9.2 Sumatra-Andaman earthquake, for which a portion of the rupture extended close to the trench (e.g., Dean et al., 2010; Lay et al., 2005). Stevens et al. (2021) compile core samples from the International Ocean Discovery Program and the multichannel seismic reflection data, suggesting that the smectite-illite transformation and complete dehydration of silicates within thick basal sediments before subduction are essential factors contributing to the enhancement of the shallow slip. They found these processes produced substantial, dehydrated bulk sediment, along with fluid trapped by the low-permeability siliciclastic-fan and basal pelagic sediments, resulting in velocity-weakening behavior close to the trench in the southern part of the rupture area. Some regions like northern Honshu, where the 1896 Sanriku tsunami earthquake occurred, appear to have the red deep sea clay on the plate, and the same holds for the shallow rupture component of the 2011 Tohoku earthquake to the south of that event. Ujiiie et al. (2013) suggest the abundant red deep sea clay retrieved from the plate-boundary décollement near the Japan trench, which consists primarily of  $\sim 78\%$  smectite, exhibits very low shear strength under high-velocity friction experiments, leading to dynamically enhanced shallow slip under fluid-saturated and impermeable conditions. The large-slip region of the 2010  $M_W$  8.8 Maule (Chile) earthquake included two patches of slip that extended to the trench (e.g., Yue, Lay, Rivera, An, et al., 2014), where red deep sea clay is also found on the incoming Nazca Plate (Figure 9). The 2015  $M_W$  8.3 Illapel (Chile) earthquake to the north also involved coseismic slip extending up-dip to near the trench (e.g., Li et al., 2016). A 600-km long shallow very large-slip region for the 1957 Aleutian earthquake (Yamazaki et al., 2024) extended along the shallow megathrust in the eastern Aleutian Islands in a region where siliceous pelagic sediments are found on the incoming plate (Figure 9).

However, it is difficult to conclude whether the widespread occurrence of sediments can govern the genesis of tsunami earthquakes individually due to the very limited tsunami earthquake sampling. Here, we discuss the possibility that combined effects of the trench pelagic sediment and rough morphological condition of subduction zones may influence seismic behavior in the shallow megathrust. Figure 9b illustrates that eight tsunami earthquakes occurred in subduction zones with an erosive margin (Kuril, northeastern Japan, Alaska-Aleutian, Chile-Peru, Middle America), while the remainder occurred in zones with an accretionary margin (Sunda, Java). These tsunami earthquake occurrences in both erosive and accretionary subduction zones suggest a lack of clear systematic correlation between tsunami earthquake occurrence and subduction tectonic framework. Excluding the 1907 Sumatra event which has some location uncertainty, the area of the 2010 Mentawai and 2022 Sumatra earthquakes has a maximum sediment thickness of  $\sim 1100$  m, while other tsunami earthquake regions host little or no sediment in the trench.

Figure 10 shows cross-sections with relocated hypocenters and slip regions for seven well-recorded tsunami earthquakes since 1990. These events occurred in regions with either thick sediment on the incoming plate (2022 Sumatra, 2010 Mentawai, 2006 Java, 1994 Java) or rough bathymetry with minor sediment (2012 El Salvador, 1992 Nicaragua, 1996 Peru). Rough seafloor topography, such as horst-and-graben structures and subducted



**Figure 9.** The distribution of well-known tsunami earthquakes globally along with pelagic sediment and sediment thickness. Focal mechanisms in (a) show 172 M6.9+ megathrust events since 1976 from the Global Centroid Moment Tensor catalog, coded by the centroid depth. Red stars indicate the well-documented tsunami earthquakes with broadband source spectrum for the 2022 Sumatra earthquake (this study) and six others since 1990 (Ye et al., 2016b). Six purple stars show the inferred tsunami earthquakes of the 1896 Sanriku (Kanamori, 1972), 1907 Sumatra (Kanamori et al., 2010), 1946 Aleutian (Kanamori, 1972), 1960 Peru (Pelayo & Wiens, 1992), 1963 and 1975 Kuril (Fukao, 1979; Pelayo & Wiens, 1992), plotted at the USGS-NEIC location. Three boxes show main slip patches for the 1957 Aleutian earthquake with the red box indicating a 600-km long shallow very large-slip region (Yamazaki et al., 2024). The pelagic sediment distribution is modified from Jenkyns (1986), which was adapted from Davies and Gorsline (1976). The sediment thickness data is from Whittaker et al. (2013). Barbed lines represent subducting boundary types (black: accretionary; red: erosive), modified from Noda (2016). Subduction zones with tsunami earthquakes displayed in Figure 11 are outlined by black boxes (a–g).

seamounts, is common in regions with tsunami earthquake occurrences (Figures 9 and 10), as is strong deformation of the outer wedge (Qiu & Barbot, 2022). The rough seafloor topography may foster thick, fragmented, and faulted plate-boundary shear zones when subducted, enhancing contact area with the overriding plate and accumulating patchy stress in the shallow megathrust (e.g., Polet & Kanamori, 2000; Tanioka et al., 1997). The presence of incoming sediment deposits likely reduces the interface frictional strength, as in the case of abundant smectite observed in the 2011 Tohoku rupture area, favoring the possibility of shallow tsunami earthquakes (e.g., Geersen, 2019; Moore et al., 2015; Polet & Kanamori, 2000; Ujiie et al., 2013). Given that tsunami earthquakes have often involved large slip with correspondingly long recurrence intervals, the available sampling from recent seismic history has to be considered with caution, as we proceed to discuss some other aspects of these events.

As shown in this study, the 2022 Sumatra tsunami earthquake ruptured ~50 km along strike of the shallow portion of the Sunda trench, with all significant slip locating shallower than ~10 km (from 2 to 8 km into rock), which has been defined as Domain A in the depth-varying segmentation model (Lay et al., 2012). The rupture is located up-dip of the Domain B (15–35 km deep) 2007  $M_W$  8.4 Bengkulu event. This geometry is similar to that of the 2010  $M_W$  7.8 Mentawai tsunami earthquake to the northwest and that of the 1907  $M_W$  7.8 Sumatra tsunami earthquake (Figure 11a and Figure S15 in Supporting Information S1). In the shallow megathrust between the 2010 Mentawai and 2022 Sumatra events, Tsang et al. (2016) suggests that ~1.5 m of aseismic afterslip of the 2007 events occurred, consistent with our Coulomb failure stress change pattern. Thus, it appears that there is little likelihood of another tsunami earthquake occurring up-dip of the Bengkulu event in the near future. The patchy distribution of seismic and aseismic slip along the shallow Sumatra subduction zone suggests a role for small-scale structural heterogeneity influencing the conditions for localized tsunami earthquake occurrence amidst adjacent regions of slow slip.

Similarly, the 1963  $M_W$  7.8 and 1975  $M_W$  7.7 Kuril earthquakes ruptured near the Kuril trench (Fukao, 1979; Pelayo & Wiens, 1992; Pérez, 2000), where the deep megathrust has experienced two great megathrust Domain B events, including the 1963  $M_W$  8.5 and 1969  $M_W$  8.2 Kuril events (Figure 11b and Figure S16 in Supporting Information S1). Notably, the 1963  $M_W$  7.8 earthquake is an aftershock one week after the 1963  $M_W$  8.5 earthquake, probably resulting from stress transfer from the deeper rupture zone to the shallower megathrust, or possibly onto a splay fault. The outer wedge is only about 15 km wide in this region (Qiu & Barbot, 2022).

Two great tsunami earthquakes: the 1896  $M_W$  8.1 Sanriku and the shallow component of the 2011  $M_W$  9.1 Tohoku events, ruptured the Domain A segment of the Japan trench along Honshu, which has also experienced numerous ruptures confined to Domains B and C (35–50 km deep) (Figure 11c and Figure S17 in Supporting Information S1). The shallow plate interface in this region is very irregular and there are multiple deformation structures in the ~30 km wide outer wedge (Qiu & Barbot, 2022). The coseismic rupture extent spanning Domains A, B, and C for the 2011 event coincides with the region of estimated strong megathrust coupling area and afterslip in the deeper zone (Ozawa et al., 2011), indicating that balanced slip accommodation exists between the shallower and deeper portion of this region. The 1896 event represents spontaneous rupture of Domain A without deeper coseismic slip, and the down-dip region appears to have little slip deficit (e.g., Ye et al., 2011), so it is possible that stress shadowing by the deeper well-locked region to the south prevented a similar Domain A rupture prior to the 2011 event. This illustrates a major challenge of evaluating whether tsunami earthquakes can occur in those circum-Pacific regions without historical examples.

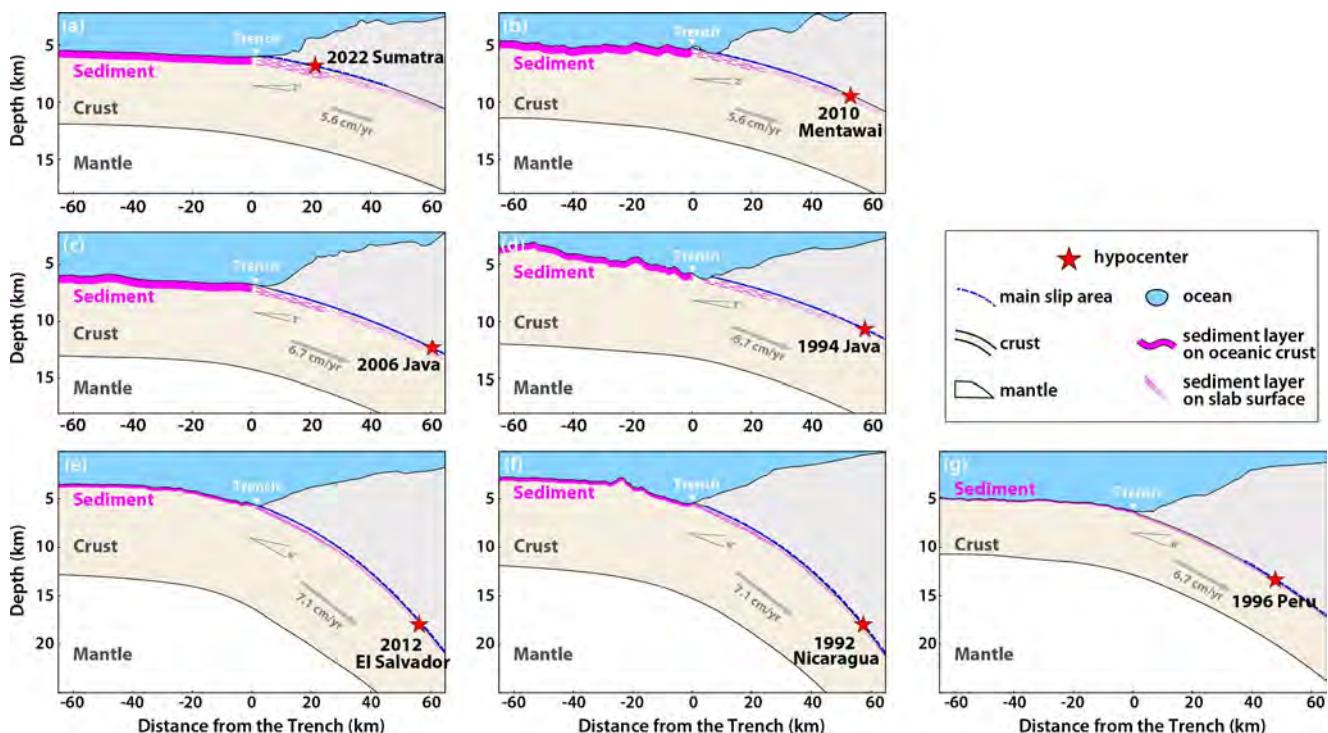
Another tsunami earthquake that spontaneously ruptured Domain A is the 1946  $M_W$  8.6 Aleutian Islands event, which occurred adjacent to the Shumagin Island region of the Aleutian-Alaska trench. Qiu and Barbot (2022) infer a very wide (~50 km) outer trench wedge in this region. There is an along-strike offset in the large-slip areas of earthquakes between the 1946 event and deeper (Domains B–C) 2020  $M_W$  7.8 Shumagin and the 2021  $M_W$  8.2 Chignik earthquakes (Figure 11d and Figure S18 in Supporting Information S1). Geodetic slip deficit appears to be very low along the deeper megathrust near the 1946 event (Freymueller & Beavan, 1999), so this may be a region where only Domain A ruptures occur, and the 1957 rupture zone to the west had 600 km long rupture of Domain A (Yamazaki et al., 2024). Aftershocks of the 2021 event distribute in the shallow portion and overlap its afterslip area at shallower depths (Brooks et al., 2023), also suggesting some strain balancing across different depths. The shallow region up-dip of the 2020 Shumagin earthquake did not have coseismic slip, but did show some slip deficit accumulation prior to the event (DeSanto et al., 2023). Some uncertainty remains regarding the seismic potential of the shallow part of the former Shumagin gap and the Chignik region, due to the lack of

shallow seismicity and limited resolution of interplate coupling along the dip in this region (e.g., Liu et al., 2023; Xiao et al., 2021; Ye et al., 2021, 2022).

Some other tsunami earthquakes have occurred in regions that experienced no historic downdip Domain B or Domain C events. The 1994  $M_W$  7.8 and 2006  $M_W$  7.7 Java events ruptured the Java trench, where the deeper megathrust region is generally regarded as weakly coupled (e.g., Scholz & Campos, 2012) with no history of great megathrust events (Figure 11e and Figure S17 in Supporting Information S1). Both events have low moment-scaled radiated energy, low apparent stress, and large slip in the depth range of 6–22 km (Ye et al., 2016a). They also produced predominantly normal-faulting aftershocks including several seaward of the large-slip zones (e.g., Abercrombie et al., 2001; Ammon et al., 2006; Bilek & Engdahl, 2007), which has been related to shallow ruptures (Sladen & Trevisan, 2018). The 1994 and 2006 Java events probably included rupture of localized unstable patches with strong coupling possibly associated with subducted seamounts or bathymetry (e.g., Abercrombie et al., 2001; Ammon et al., 2006; Bilek & Engdahl, 2007). The outer trench wedge along this accretionary boundary is estimated to be about 40 km wide in this region (Qiu & Barbot, 2022).

The 1960  $M_W$  7.6 and 1996  $M_W$  7.5 Peru events ruptured the northern Peru trench, which is an erosive boundary that has low seismic coupling in general (e.g., Scholz & Campos, 1995) without any recorded deep megathrust events in Domains B and C (Figure 11f and Figure S20 in Supporting Information S1). The rupture extent of the 1960 event is correlated with a locally strong shallow coupling region, but the 1996 event cannot be assessed due to the low-resolution coupling model (e.g., Villegas-Lanza et al., 2016). The outer wedge width is estimated as about 20 km near the 1996 event by Qiu and Barbot (2022). The 1992  $M_W$  7.6 Nicaragua and the 2012  $M_W$  7.3 El Salvador events ruptured patchy areas in the shallow portion of the Middle American megathrust along Nicaragua and El Salvador (Figure 11g and Figure S21 in Supporting Information S1). These are along an erosive margin as well, with ~20 km wide outer wedge (Qiu & Barbot, 2022). The down-dip regions don't appear to be strongly coupled, but the presence of an upper plate sliver complicates this inference (Ye et al., 2013). There are other moderate ruptures showing regular seismic behavior along the strike in the Middle American megathrust, such as the 2012  $M_W$  7.6 Costa Rica event in Domain B and the 2012  $M_W$  7.4 Guatemala event in Domain C (e.g., Ye et al., 2013).

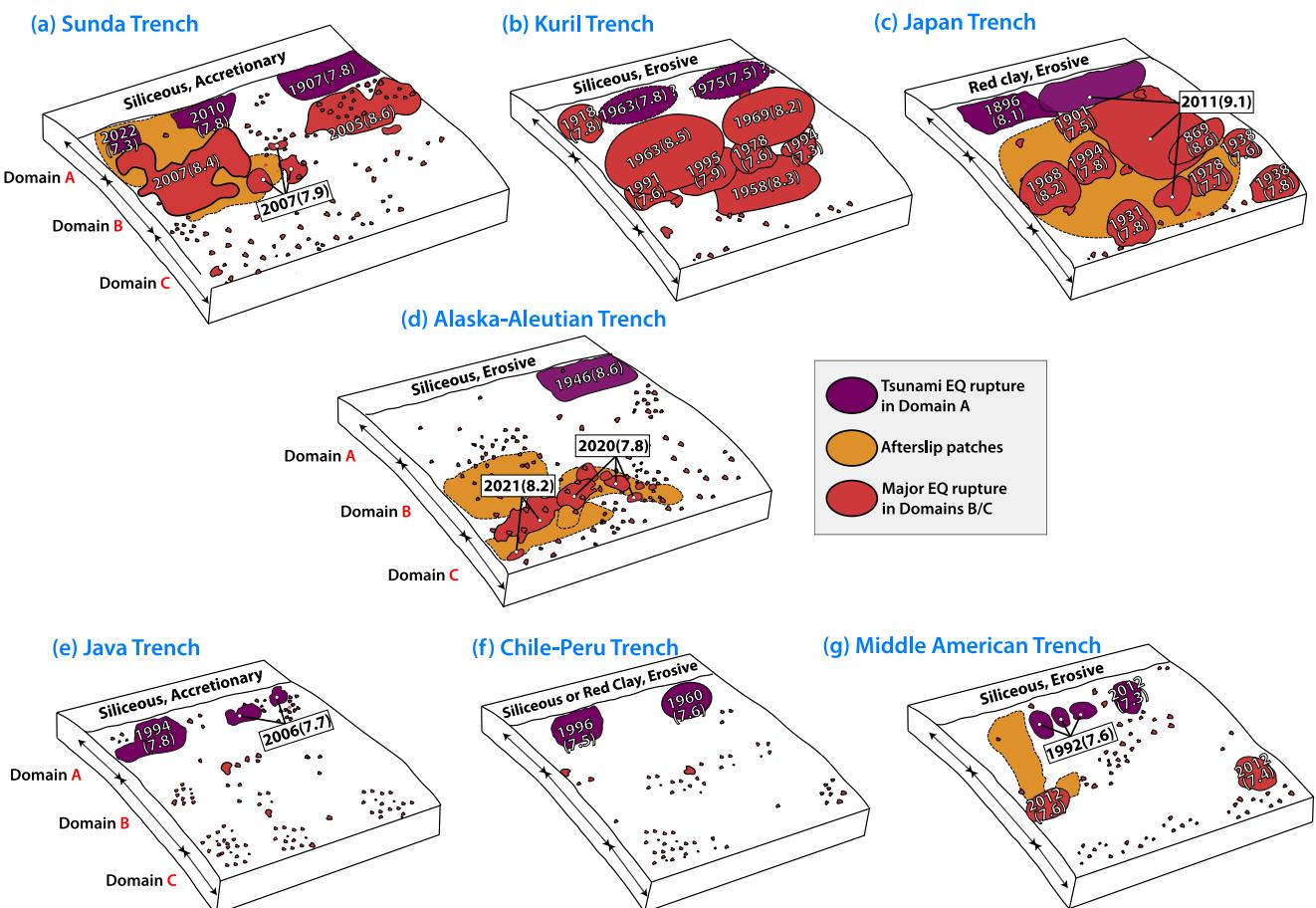
The documented tsunami earthquakes occurred in diverse megathrust environments, which can be roughly divided into two categories. One type includes regions where tsunami earthquake rupture up-dip of major or great megathrust events at deeper depths, with the deeper ruptures likely promoting triggering (i.e., clock-advance) of the ruptures at shallow depths, if not coseismic rupture as for the 2011 Tohoku event. The deeper portion of megathrust, when locked, likely casts a stress shadow for the shallow depth megathrust, regardless of the frictional properties at shallow depth (Lindsey et al., 2021). The heterogeneous coupling and seismogenic capability in the shallow Domain A region indicated by this type of tsunami earthquake may be attributed to the presence of fragmented and faulted shear zones resulting from the combined effects of sediments and rough topography, as discussed previously. It is unclear why immediate rupture of the shallow region does not always take place when deeper ruptures occur, as was the case for the 3 and 15 years delays for the 2010 Mentawai and 2022 Sumatra events following the 2007 Bengkulu earthquake; sluggish poroelastic and viscoelastic effects and additional stress increments from afterslip may be required to push the patchy shallow interface to subsequent failure. The presence of low-velocity sedimentary layers deposited on the wedge may cause dynamic interactions with the rupture that allow deeper ruptures to have more shallow slip (Li et al., 2023), but such superimposed layers are not evident in the steeply dipping outer wedges where isolated tsunami earthquake occur (Qiu & Barbot, 2022). Another type of tsunami earthquake includes regions where tsunami earthquakes occur in Domain A without having large events in Domains B and C. In these regions, localized asperities in the shallow megathrust are able to accumulate stress for a long time and fail spontaneously. The occurrence of such ruptures does not preclude some degree of stress shadowing from slip deficit accumulation and eventual rupture of deeper Domains B and C; the earthquake catalog may simply be too short to provide a full picture of the process. In these regions, long-term seafloor geodetic monitoring is required to assess the potential for tsunami earthquakes as well as to detect strain accumulation that may result in deeper ruptures.



**Figure 10.** Profiles for seven regions that experienced well-documented tsunami earthquakes since 1990, including (a) 2022  $M_W$  7.3 Sumatra, (b) 2010  $M_W$  7.8 Mentawai, (c) 2006  $M_W$  7.7 Java, (d) 1994  $M_W$  7.8 Java, (e) 2012  $M_W$  7.3 El Salvador, (f) 1992  $M_W$  7.6 Nicaragua, and (g) 1996  $M_W$  7.6 Peru. Red stars show relocated hypocenters (Details in Table S4 and Figures S9–S14 in Supporting Information S1), with blue dashed curves indicating the main slip area. Each profile is constructed using GEBCO 2023 bathymetry data, sediment thickness data (Whittaker et al., 2013), and the Slab 2.0 model. The crustal thickness is derived from the Crust 1.0 model. The amplitude of the seafloor bathymetry is amplified by a factor of 3. Depths in the Slab 2.0 model are manually adjusted to match the trench depth from the bathymetry data.

#### 4. Conclusion

The 2022  $M_W$  7.3 Sumatra underthrusting earthquake ruptured the megathrust region up-dip of the 2007  $M_W$  8.4 Bengkulu rupture on a shallow, sub-horizontal fault plane. The major slip patch had a peak slip of 4.1 m, located at depths shallower than 10 km, adjacent to the shallow afterslip zone of the 2007 major events. It exhibits common features of tsunami earthquakes, including a depletion in high-frequency seismic wave radiation, low moment-scaled radiated energy of  $4.15 \times 10^{-6}$ , and low source region rigidity of 23 GPa. Calculation of the Coulomb failure stress suggests that the 2022  $M_W$  7.3 Sumatra event was probably promoted by stress transfer resulting from the coseismic slip and afterslip of the 2007 Bengkulu event. We consider tsunami earthquakes that occurred in the Sumatra, offshore Honshu, Alaska-Aleutian, Peru, Java, and Middle American subduction zones using the depth-varying megathrust subdivision of Lay et al. (2012) with the combined effects of sediments and rough morphological. The diverse slip modes in different subduction zones reveal substantial complexity in the shallow megathrust, including ruptures in Domain A that may be induced by deeper ruptures in Domains B and C, or may occur independently due to the long-term stress accumulation. The presence of tsunami earthquakes in both erosive and accretionary subduction zones suggests a lack of distinct systematic correlation between their occurrence and subduction tectonic framework. The width of the strongly deformed outer wedge of the accretionary prism appears to correlate with tsunami excitation, with coseismic splay faulting within the wedge possibly contributing to tsunami excitation (Qiu & Barbot, 2022), but the shallow megathrust supports strain accumulation that releases in tsunami earthquakes. Furthermore, the combination of pelagic clay and siliceous sediments on the subducting ocean floor, which are subducted or accumulate in the outer sedimentary wedge, and rough seafloor topography may result in the heterogeneous coupling in the Domain A region, which appears to play important roles in controlling the genesis of tsunami earthquake along Sumatra and other regions. However, the variability of structures among limited recorded tsunami earthquakes obscures the factors that control tsunami earthquake occurrence.



**Figure 11.** Schematics for subduction zones with major tsunami earthquakes. (a) Sumatra subduction zone with the 1907  $M_w$  7.8 Sumatra (Kanamori et al., 2010; Martin et al., 2019), 2005  $M_w$  8.6 Nias (Hsu et al., 2006), 2007  $M_w$  8.4 Bengkulu (Konca et al., 2008), the 2007  $M_w$  7.9 Pagai (Konca et al., 2008), the 2010  $M_w$  7.8 Mentawai (Yue, Lay, Rivera, Bai, et al., 2014) and the 2022  $M_w$  7.3 Sumatra events (this study). (b) Kuril subduction zone with the 1963  $M_w$  7.8, 1975  $M_w$  7.7, 1969  $M_w$  8.2, 1963  $M_w$  8.5, and 1995  $M_w$  7.9 Kuril events (Fukao, 1979; Pelayo & Wiens, 1992; Pérez, 2000). (c) Offshore Honshu subduction zone with the 1896  $M_w$  8.1 Sanriku, the 1978  $M_w$  7.7 Miyagi-oki, and the 2011  $M_w$  9.1 Tohoku events (Ye et al., 2021). (d) Alaska-Aleutian subduction zone with the 1946  $M_w$  8.6 Alaska, the 2020  $M_w$  7.8 Shumagin (Liu et al., 2023), and the 2021  $M_w$  8.2 Chignik events (Ye et al., 2022). (e) Java subduction zone with the 1994  $M_w$  7.8 Java and the 2006  $M_w$  7.7 Java events (Bilek & Engdahl, 2007). (f) Chile-Peru subduction zone with the 1960  $M_w$  7.6 Peru events (Bilek, 2010; Pelayo & Wiens, 1992) and the 1996  $M_w$  7.5 Peru (Ihmé et al., 1998). (g) Middle American subduction zone with the 1992  $M_w$  7.6 Nicaragua, the 2012  $M_w$  7.3 El Salvador, the 2012  $M_w$  7.4 Guatemala (Ye et al., 2013), and the 2012  $M_w$  7.6 Costa Rica events (Yue et al., 2013). Red polygons indicate major earthquake ruptures in Domains B and C; Purple polygons indicate tsunami earthquakes in the shallow depth; and orange polygons indicate afterslip patches, such as for the 2007  $M_w$  8.4 Bengkulu (Tsang et al., 2016), the 2011  $M_w$  9.1 Tohoku (Ozawa et al., 2011), the 2020  $M_w$  7.8 Shumagin, the 2021  $M_w$  8.2 Chignik (Brooks et al., 2023), and 2012  $M_w$  7.6 Costa Rica (Hobbs et al., 2017) events.

## Conflict of Interest

The authors declare no conflicts of interest relevant to this study.

## Data Availability Statement

Global seismic data were downloaded from the Incorporated Research Institutions for Seismology (IRIS) data management center ([http://ds.iris.edu/wilber3/find\\_event](http://ds.iris.edu/wilber3/find_event)), including stations from Global Seismographic Network code II (Scripps Institution of Oceanography [SIO], 1986, available at <https://doi.org/10.7914/SN/II>), IU (ASL/USGS, 1988, available at <https://doi.org/10.7914/SN/IU>), and International Federation of Digital Seismic Networks (FDSN) code AU (Geoscience Australia, 2021, available at <https://doi.org/10.26186/144675>), AW (Alfred Wegener Institute For Polar And Marine Research [AWI], 1993, available at <https://doi.org/10.14470/NJ617293>), FN (Sodankylä Geophysical Observatory/University Of Oulu (Finland), 1980, available at <https://doi.org/10.14470/SA879454>), G (Institut de physique du globe de Paris (IPGP) & École et Observatoire des Sciences de la Terre de Strasbourg (EOST), 1982, available at <https://doi.org/10.18715/GEOSCOPE.G>), GE

(GEOFON Data Centre, 1993, available at <https://doi.org/10.14470/TR560404>), GT (Albuquerque Seismological Laboratory (ASL)/USGS, 1993, available at <https://doi.org/10.7914/SN/GT>), HU [Kövesligethy Radó Seismological Observatory (Geodetic And Geophysical Institute, Research Centre For Astronomy And Earth Sciences, Hungarian Academy Of Sciences), 1992, available at <https://doi.org/10.14470/UH028726>], IC (Albuquerque Seismological Laboratory (ASL)/USGS, 1992, available at <https://doi.org/10.7914/SN/IC>), JP, KC (Central Asian Institute for Applied Geosciences, 2008, available at <https://doi.org/10.7914/SN/KC>), MN (MedNet Project Partner Institutions, 1990, available at <https://doi.org/10.13127/sd/fbbbttd6q>) and PS. The tsunami observation at Cocos Island tide gauge station is downloaded from the European Commission Joint Research Centre, ISPRA, Space, Security and Migration Directorate (JRC, [https://webcritech.jrc.ec.europa.eu/tad\\_server/](https://webcritech.jrc.ec.europa.eu/tad_server/)). The high-resolution bathymetry BATNAS data set is downloaded from Indonesian Geospatial Information Agency (<https://tanahair.indonesia.go.id/portal-web/unduh/batnas>). We thank the facilities of IRIS Data Services, and specifically the IRIS Data Management Center, which were used for access to waveforms, related metadata, and/or derived products used in this study. Data Services are funded through the Seismological Facilities for the Advancement of Geoscience (SAGE) Award of the National Science Foundation under Cooperative Support Agreement EAR-1851048. Global Centroid Moment Tensor solutions are from <https://www.globalcmt.org/CMTsearch.html>. The USGS-NEIC catalog is from <https://earthquake.usgs.gov/earthquakes/search/>. All links were last accessed on 29 October 2024.

## Acknowledgments

We thank Prof. Bin Chen at SUSTech for the discussion about the pelagic sediment, Prof. Haipeng Luo at SUSTech for the discussion about the tsunami earthquakes. We thank editor Rachel Abercrombie, associate editor Yihe Huang, and two anonymous reviewers for their constructive comments. The work was supported in part by National Key Research and Development Program of China 2023YFF0803200 (L. Y.), National Natural Science Foundation of China 42376212 (Y. B.), National Natural Science Foundation of China 42074048 (S. X.), and National Science Foundation Grant EAR1802364 (T. L.).

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Supporting Information for

# The 2022 $M_W$ 7.3 Southern Sumatra Tsunami Earthquake: Rupture Up-dip of the 2007 $M_W$ 8.4 Bengkulu Event

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**Table S1.** Point-source solutions

	Lon.	Lat.	H (km)	Strike	Dip	Rake	$M_0$ (Nm)	$M_W$
USGS	100.74°	-4.96°	19.5	347.0°	6.0°	125.0°	3.29 x10 <sup>19</sup>	6.94
GCMT	100.31°	-4.98°	12.0	335.0°	11.0°	100.0°	2.81 x10 <sup>19</sup>	6.90
This study (W-Phase)	100.30°	-4.99°	9.5	346.6°	6.2°	114.4°	4.39 x10 <sup>19</sup>	7.03

**Table S2.** Parameters for slip models M01-M04

Models	$\bar{\mu}$ (GPa)	$\bar{\delta}$	$D_{\text{peak}}$ (m)	$\bar{D}_{0.15}$ (m)	$A_{0.15}$ (km <sup>2</sup> )	$\varphi_{0.15}$ (km <sup>3</sup> )	$\varphi$ (km <sup>3</sup> )	$M_0$ (Nm)	$\beta_2$	Data fits	
										Teleseismic	Tsunami
M01	50.0	3.8°	3.5	1.4	1088	1.5	2.0	8.07 x10 <sup>19</sup>	0.05	Poor	Poor
M02	27.3	3.7°	4.7	1.7	1152	2.0	2.5	5.52 x10 <sup>19</sup>	0.05	Poor	Poor
M03	47.8	3.5°	4.4	1.8	1920	3.5	5.0	19.40 x10 <sup>19</sup>	0.001	Good	Good
M04	23.0	3.5°	4.1	1.8	2240	4.0	5.6	10.50 x10 <sup>19</sup>	0.001	Good	Good

Notes:  $\bar{\mu}$  and  $\bar{\delta}$  are average rigidity and dip angle weighted by the slip distribution  $D(x, y)$  [ $\bar{\mu} = \int \mu(x, y)D(x, y)dS / \int D(x, y)dS$ ;  $\bar{\delta} = \int \delta(x, y)D(x, y)dS / \int D(x, y)dS$ ], respectively.  $D_{\text{peak}}$  is the peak slip.  $\bar{D}_{0.15}$ ,  $A_{0.15}$  and  $\varphi_{0.15}$  are average slip, total area and seismic potency for grids which has a slip value larger or equal to  $0.15D_{\text{peak}}$ , respectively.  $\varphi_{0.15} = A_{0.15} \times \bar{D}_{0.15}$ .  $\varphi$  is the total seismic potency [ $\varphi = \int D(x, y)dS$ ].  $\beta_2$  is the weighting factor of the moment constraint. Model M04 is our preferred model.

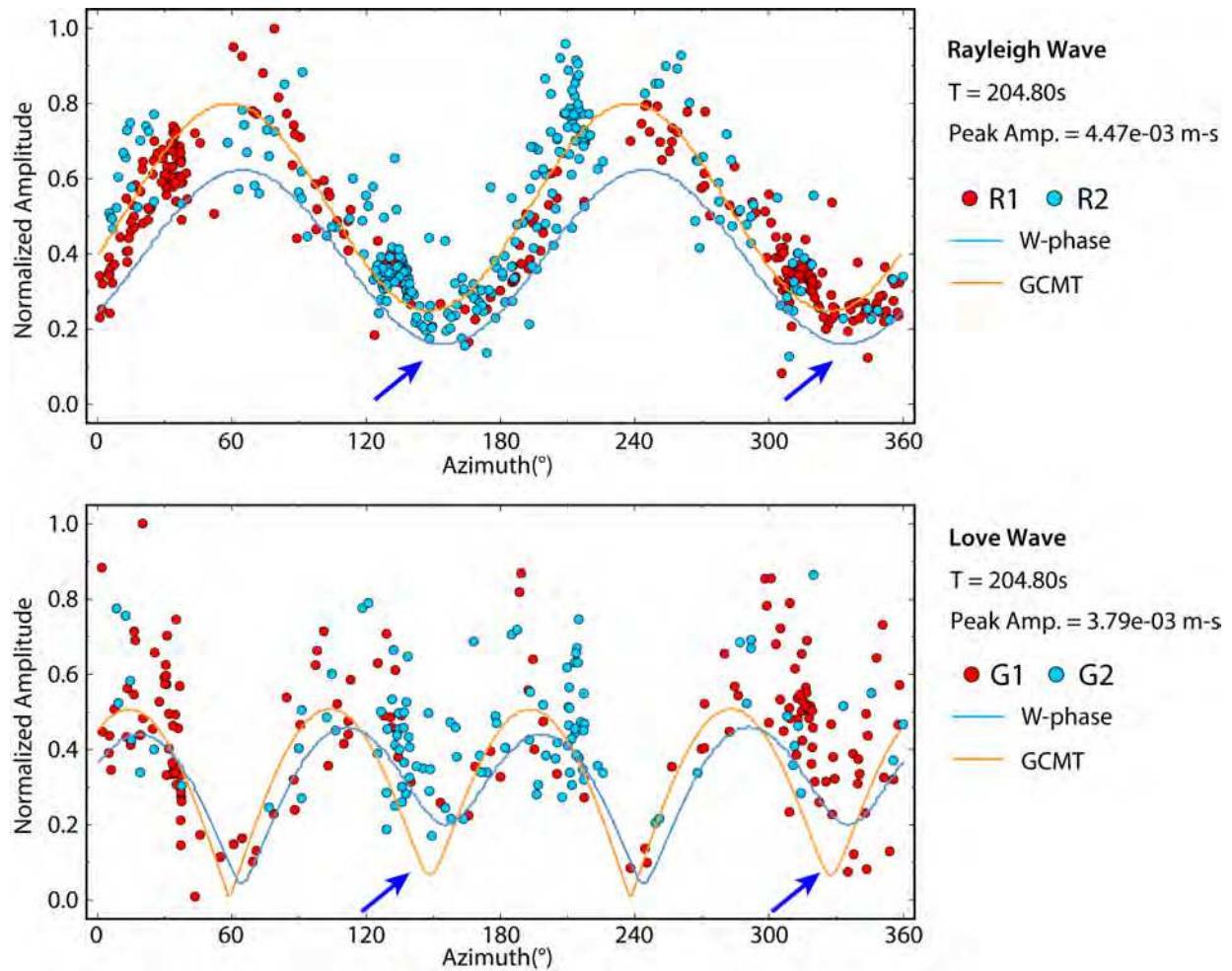
**Table S3.** Different velocity models tested in iterative finite-fault inversion.

Depth (km)	$V_P$ (km/s)			$V_S$ (km/s)			$\rho$ (g/cm <sup>3</sup> )	Rigidity $\mu$ (GPa)		
	M01/M03	M02	M04	M01/M03	M02	M04		M01/M03	M02	M04
0.00	1.50			0.00			1.02	0.00		
5.52	1.50			0.00			1.02	0.00		
5.52	1.75			0.34			1.82	0.21		
5.82	1.75			0.34			1.82	0.21		
5.82	5.00	3.95	3.50	2.70	2.13	1.89	2.55	18.59	11.57	9.11
6.49	5.00	3.95	3.50	2.70	2.13	1.89	2.55	18.59	11.57	9.11
6.49	6.50	5.10	4.55	3.70	2.92	2.59	2.85	39.02	24.30	19.12
7.97	6.50	5.10	4.55	3.70	2.92	2.59	2.85	39.02	24.30	19.12
7.97	7.10	5.60	4.97	4.05	3.20	2.84	3.05	50.03	31.23	24.60
12.55	7.10	5.60	4.97	4.05	3.20	2.84	3.05	50.03	31.23	24.60
12.55	7.80			4.40			3.50	67.76		
$+\infty$	7.80			4.40			3.50	67.76		

**Table S4. Relocation info. for 7 well-documented tsunami earthquakes since 1990.**

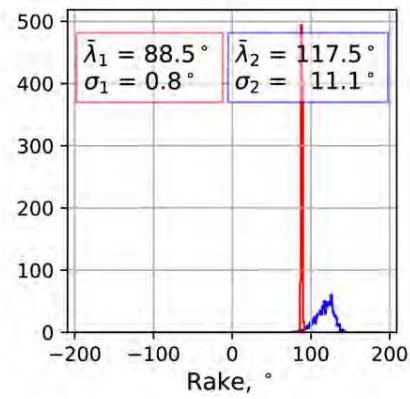
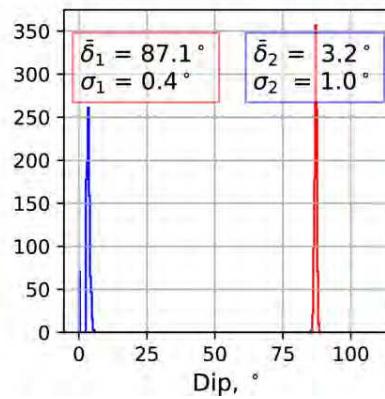
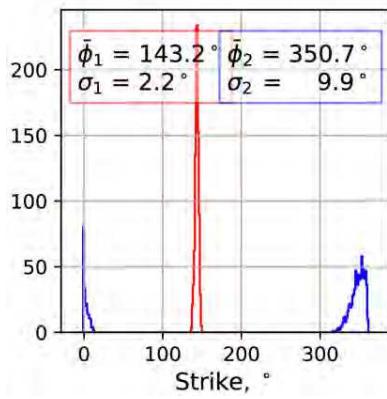
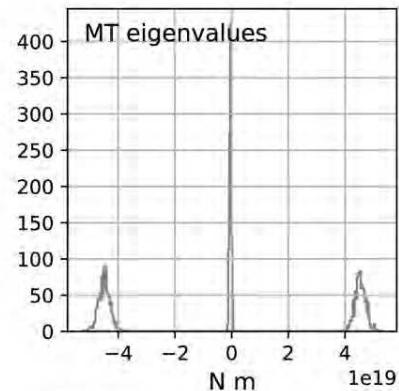
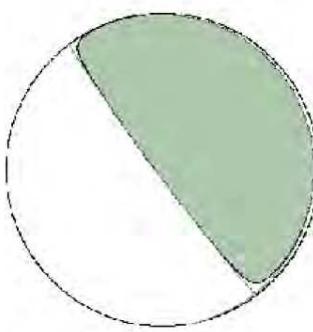
Events	USGS-NEIC			The optimal solutions in this study		
	Lat.	Lon.	Depth (km)	Lat.	Lon.	Depth (km)*
1992 $M_w$ 7.7 Nicaragua 1992-09-02T00:16:01	11.742°	-87.340°	44.8	11.602°	-87.305°	19.9
1994 $M_w$ 7.8 Java 1994-06-02T18:17:34	-10.477°	112.835°	18.4	-10.345°	112.928°	11.2
1996 $M_w$ 7.4 Peru 1996-02-21T12:51:01	-9.593°	-79.587°	10.0	-9.678°	-79.897°	16.9
2006 $M_w$ 7.7 Java 2006-07-17T08:19:26	-9.284°	107.419°	20.0	-9.184°	107.569°	14.2
2010 $M_w$ 7.8 Mentawai 2010-10-25T14:42:22	-3.487°	100.082°	20.1	-3.485°	100.184°	12.1
2012 $M_w$ 7.3 El Salvador 2012-08-27T04:37:19	12.139°	-88.590°	28.0	12.299°	-88.604°	19.9
2022 $M_w$ 7.3 Sumatra 2022-11-18T13:37:08	-4.904°	100.786°	25.0	-4.904°	100.736°	8.9

Velocity model: Crust 1.0 and PREM models.  
\*Source depth is obtained by projecting the epicenter to the Slab 2.0 model.

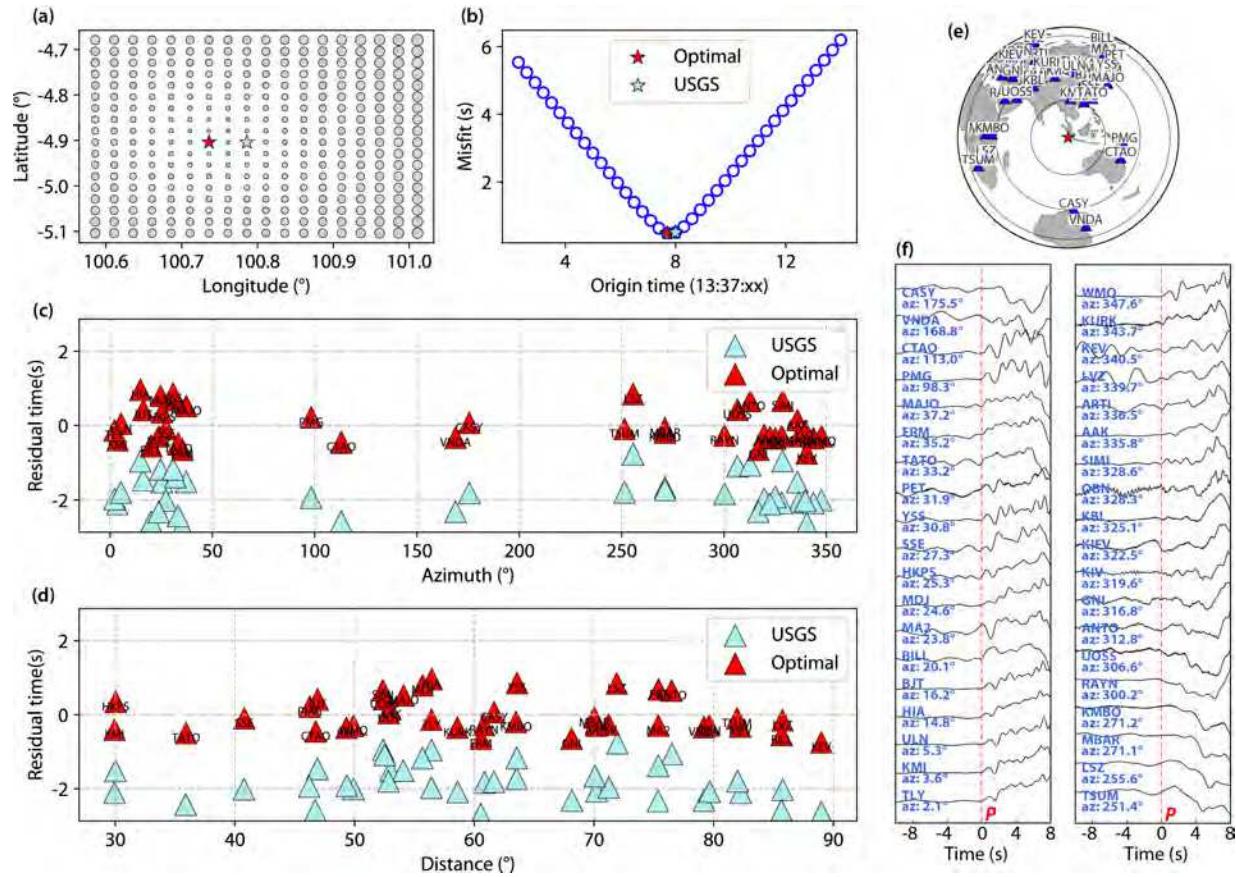


**Figure S1.** Observed and predicted Rayleigh wave (top) and Love wave (bottom) source spectral amplitudes for period  $T = 204.80$  s for the 18 November 2022 Sumatra earthquake. Red and cyan dots indicate short-arc (R1, G1) and long-arc (R2, G2) observations, respectively. Cyan and orange lines indicate theoretical amplitudes for our *W*-phase solution with strike  $346.6^\circ$ , dip  $6.2^\circ$ , rake  $114.4^\circ$ ,  $M_0 = 4.39 \times 10^{19}$  Nm and the GCMT solution with strike  $335^\circ$ , dip  $11^\circ$ , rake  $100^\circ$ ,  $M_0 = 2.81 \times 10^{19}$  Nm, respectively. The blue arrows indicated the symmetry in the Rayleigh wave radiation nodes (e) and the asymmetry in the Love wave radiation nodes (f).

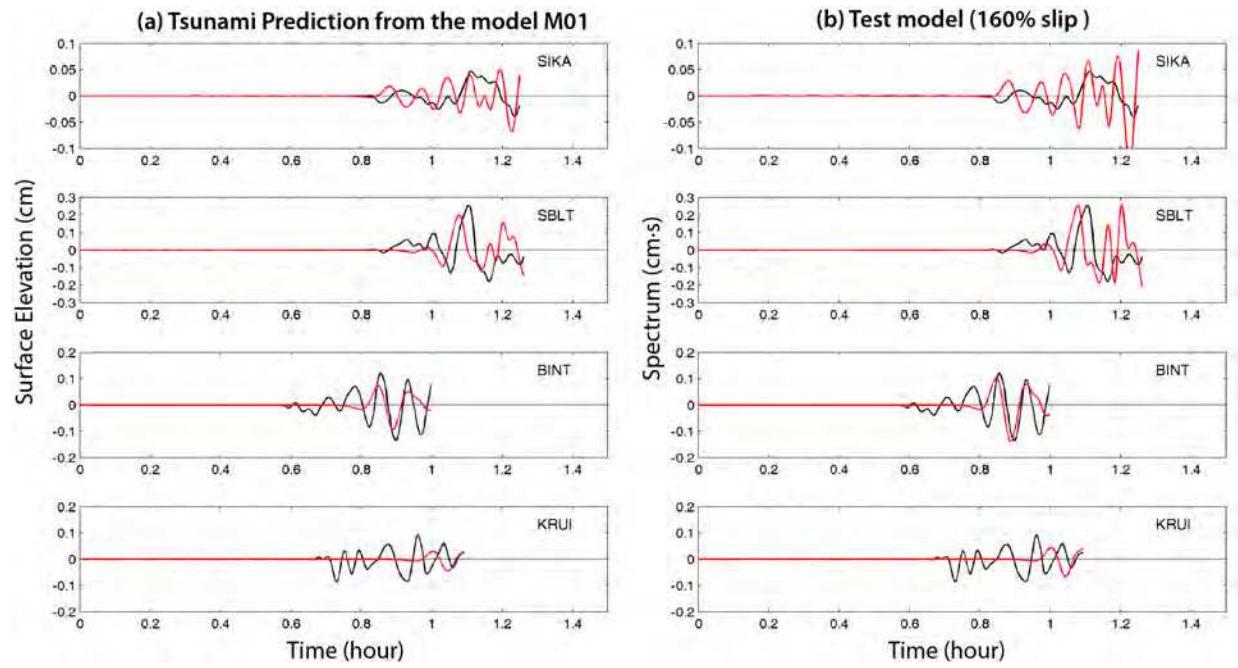
2022\_1118\_Sumatra  
 2022/11/18 13:37:08 + 15s  
 Centroid location:  
 Latitude: 4.992 S  
 Longitude: 100.301 E  
 Depth: 9.5 km  
 $M_0 = 4.52e+26 \text{ Nm}$  ( $M_w = 7.04$ )  
 Stations / Channels = 95 / 157  
 Distance range: 4.36° - 83.04°  
 Azimuthal gap = 29.7°  
 C# = 69  
 Screening: 5. 3. .9  
 Filter = 0.00500-0.01000 Hz  
 N\_bootstrap = 1000



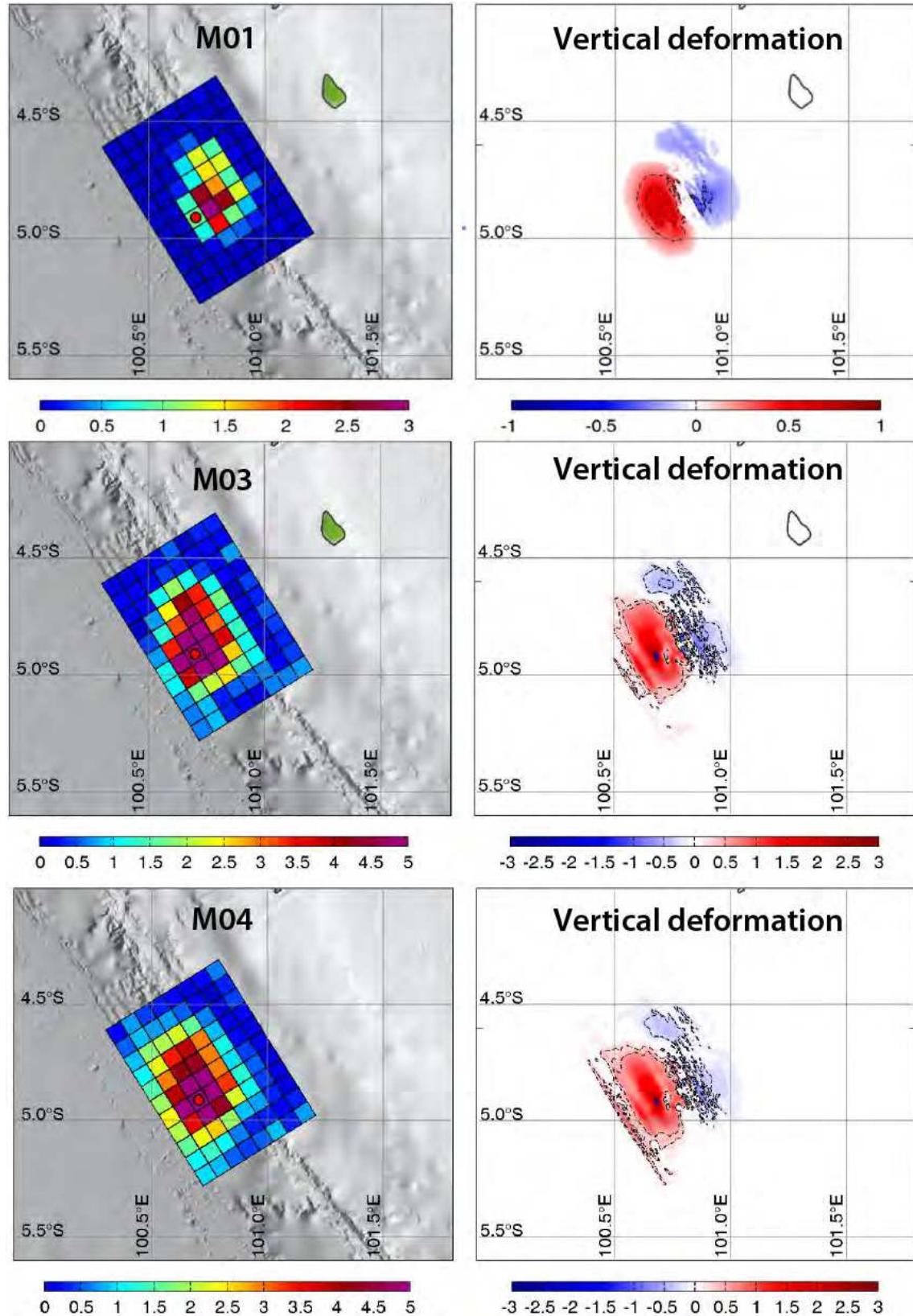
**Figure S2.** Bootstrap statistics of the  $W$ -phase solution. The box on the top left and focal mechanism gives the event details and a summary of the  $W$ -phase source inversion parameters, solution, and quality control parameters. The histograms show the output of 1,000 bootstrap inversion solutions. The initial dataset comprises 157 channels (Nch). Each bootstrap dataset is built by randomly drawing (with replacement) Nch channels from the initial dataset. Histograms for the eigenvalues of the moment tensors are depicted on the top right panel. Although not enforced during the inversion, for practical purposes, all the solutions are pure double couples. The bottom histograms correspond to the geometrical parameters: strike ( $\varphi$ ), dip ( $\delta$ ), and rake ( $\lambda$ ) of the double couples. Blue and red, respectively, indicate the shallow and the near-vertical planes. They are remarkably stable. For geometrical reasons, the strike and the rake of the sub-horizontal planes are more spread. The average and standard deviation ( $\sigma$ ) for each parameter are given. The dip of the shallowly dipping plane averages  $3.2^\circ$  with a standard deviation of  $1^\circ$ .



**Figure S3.** Relocation of the 2022 Sumatra earthquake. The source depth of our optimal solution is fixed at 8.9 km by projecting the corresponding epicenter onto the Slab 2.0 model. (a) and (b) represent misfit between observed and theoretical first arrival in both space and time, respectively. (c) and (d) Plots of arrival time residuals using the USGS-NEIC solution and our optimal location with station azimuth and distance. (e) and (f) Station distribution and initial waveforms of P-wave aligned with the predicted time from our optimal solution.

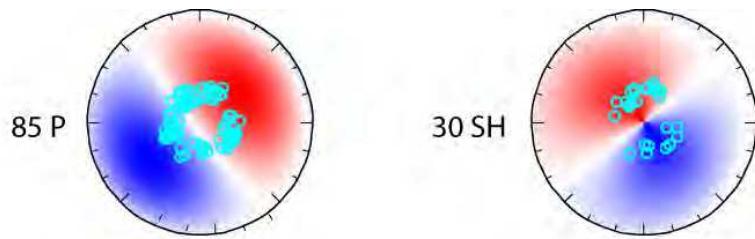


**Figure S4.** Comparison of recorded (black) and computed (red) sea level time series at SIKA, SBLT, BINT, and KRUI tide gauges for (a) slip model M01 and (b) the Test model that increases slip of all subfaults by a factor of 1.6 relative to M01.

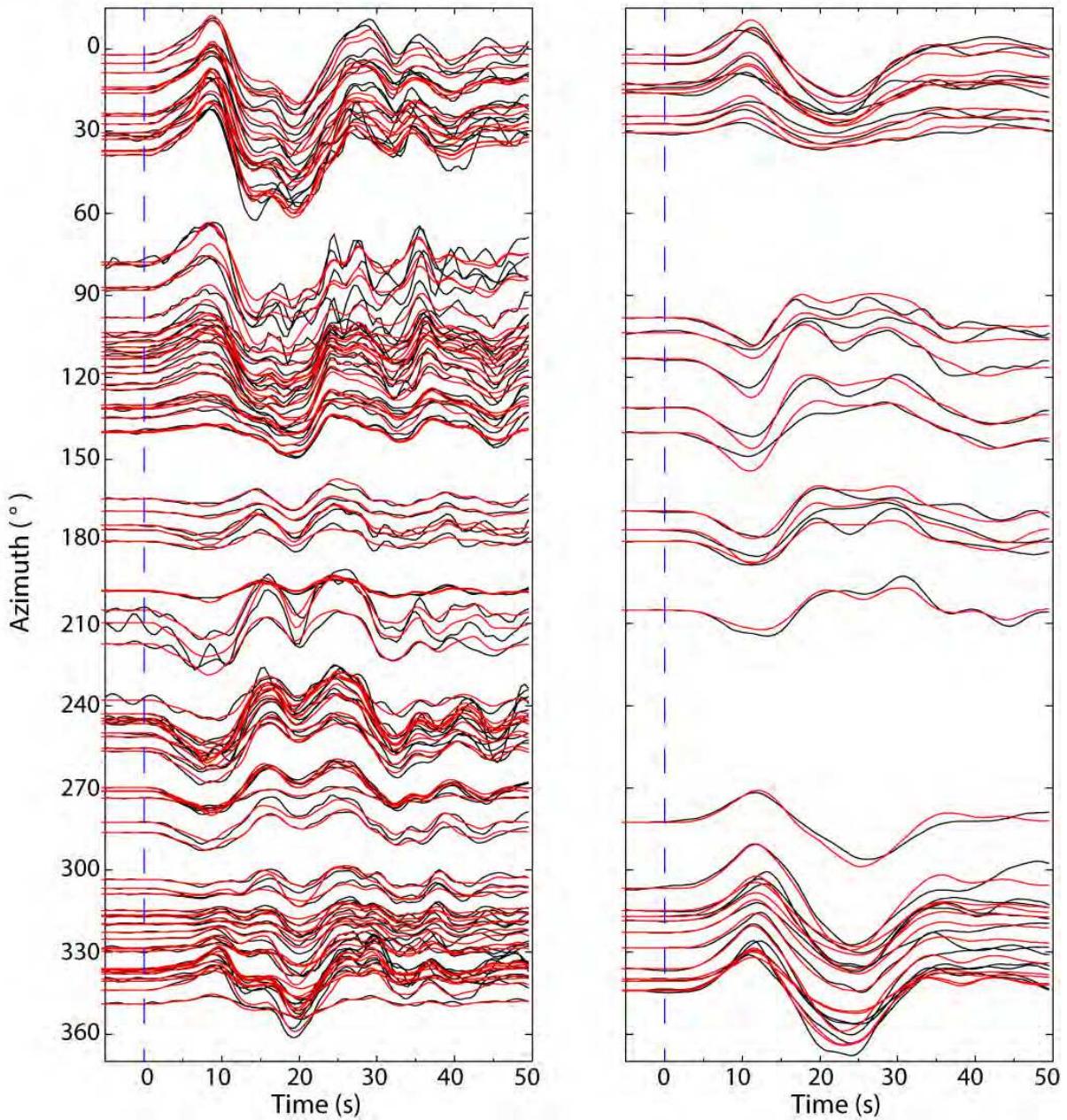


**Figure S5.** Map and seafloor vertical deformation for slip models M01, M03, and M04. The red dot indicates the epicenter. Model M04 is our preferred solution.

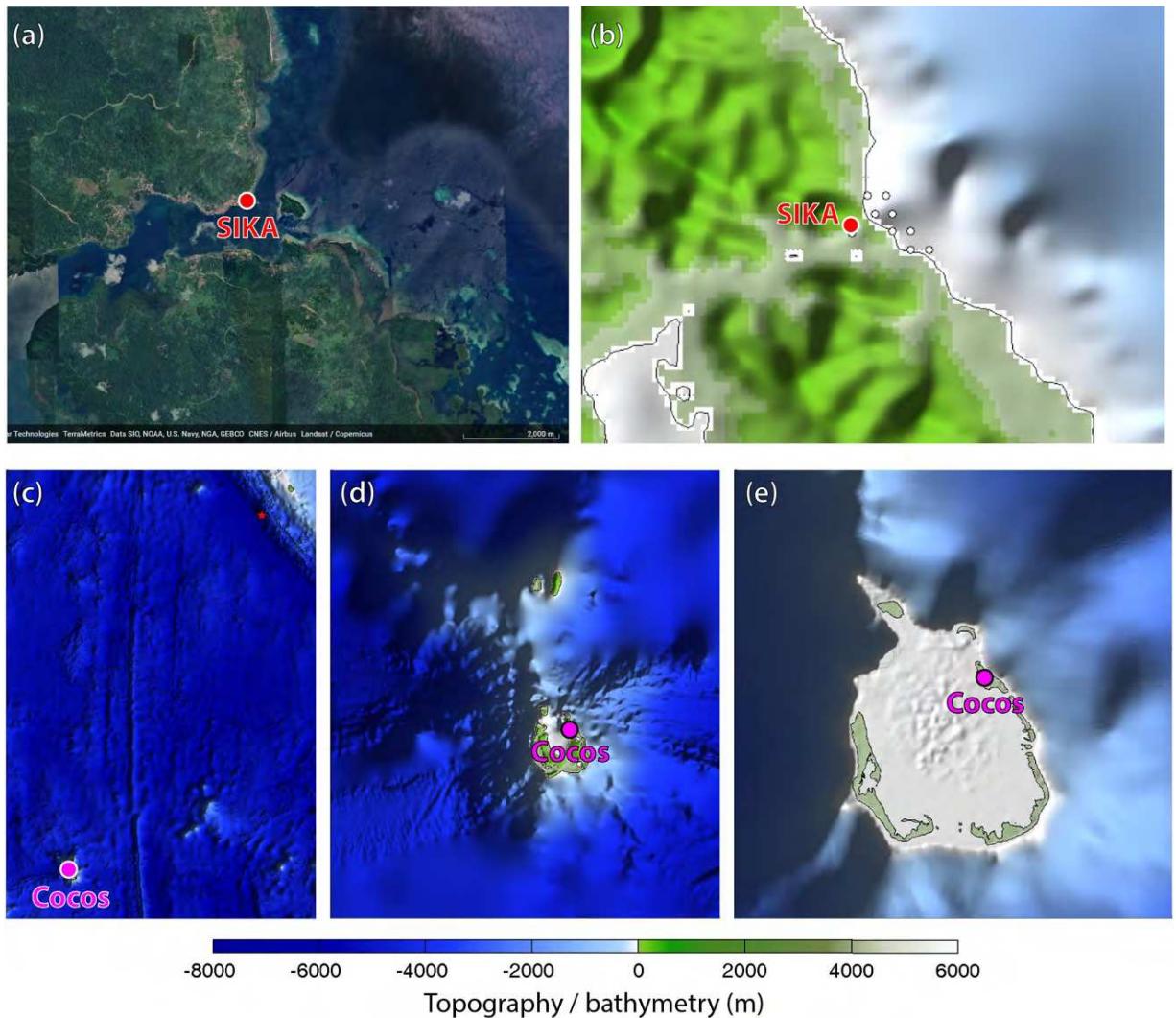
(a)



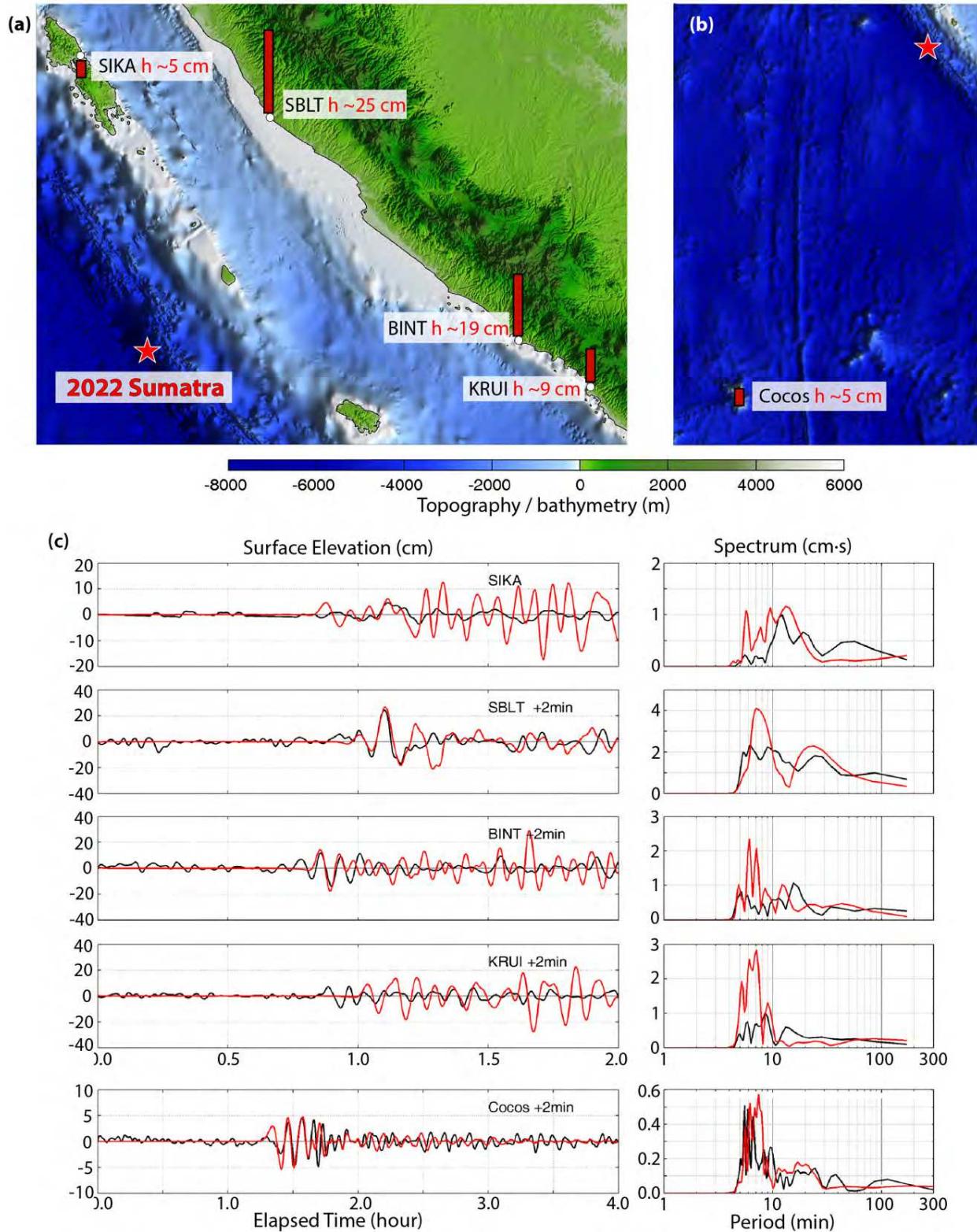
## (b) Telescismic Waveform Data and Synthetics



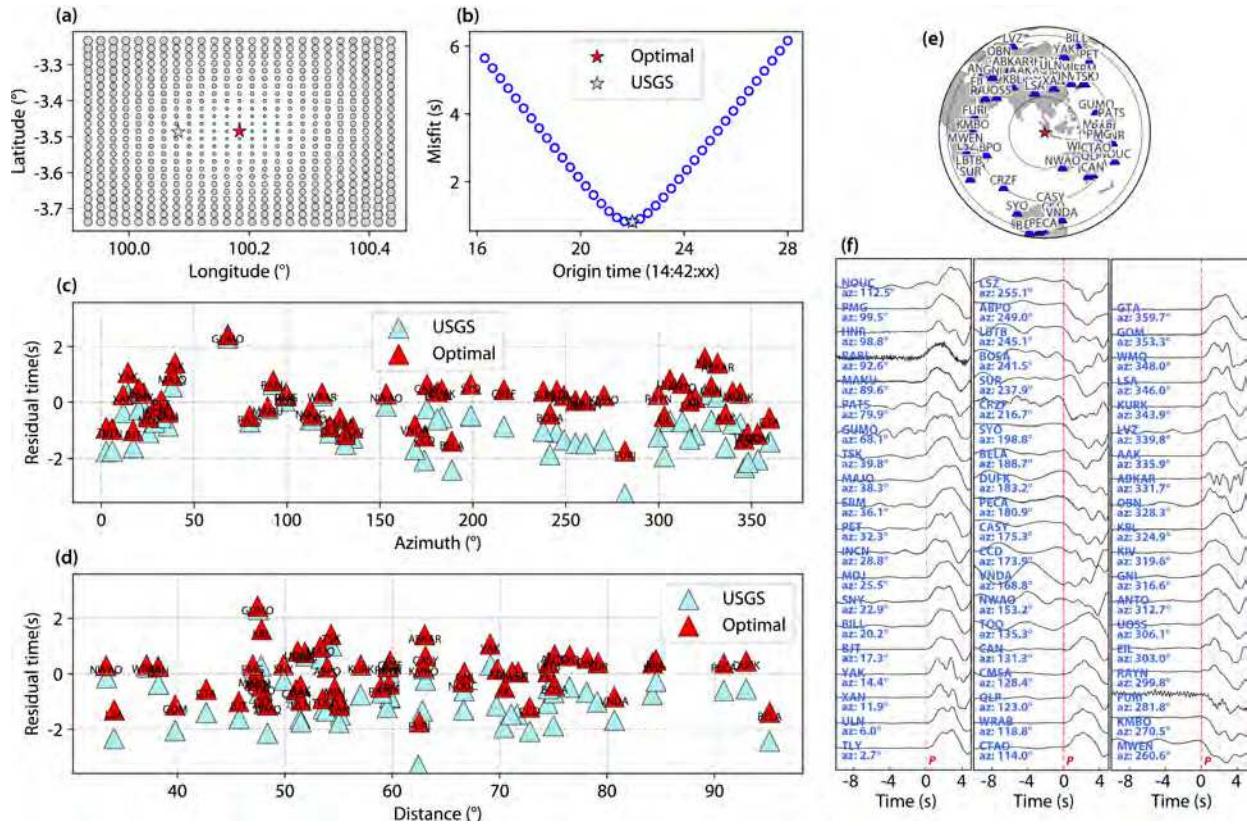
**Figure S6.** Telescismic data fits from the preferred slip model M04. (a) Lower-hemisphere stereographic projections of radiation patterns for the 85  $P$ -wave and 30  $SH$ -wave (cyan) ground motions used in the inversion. (b) Comparisons of the observed (black) and predicted (red) ground displacement waveforms of telescismic  $P$  (left) and  $SH$  (right) waves.



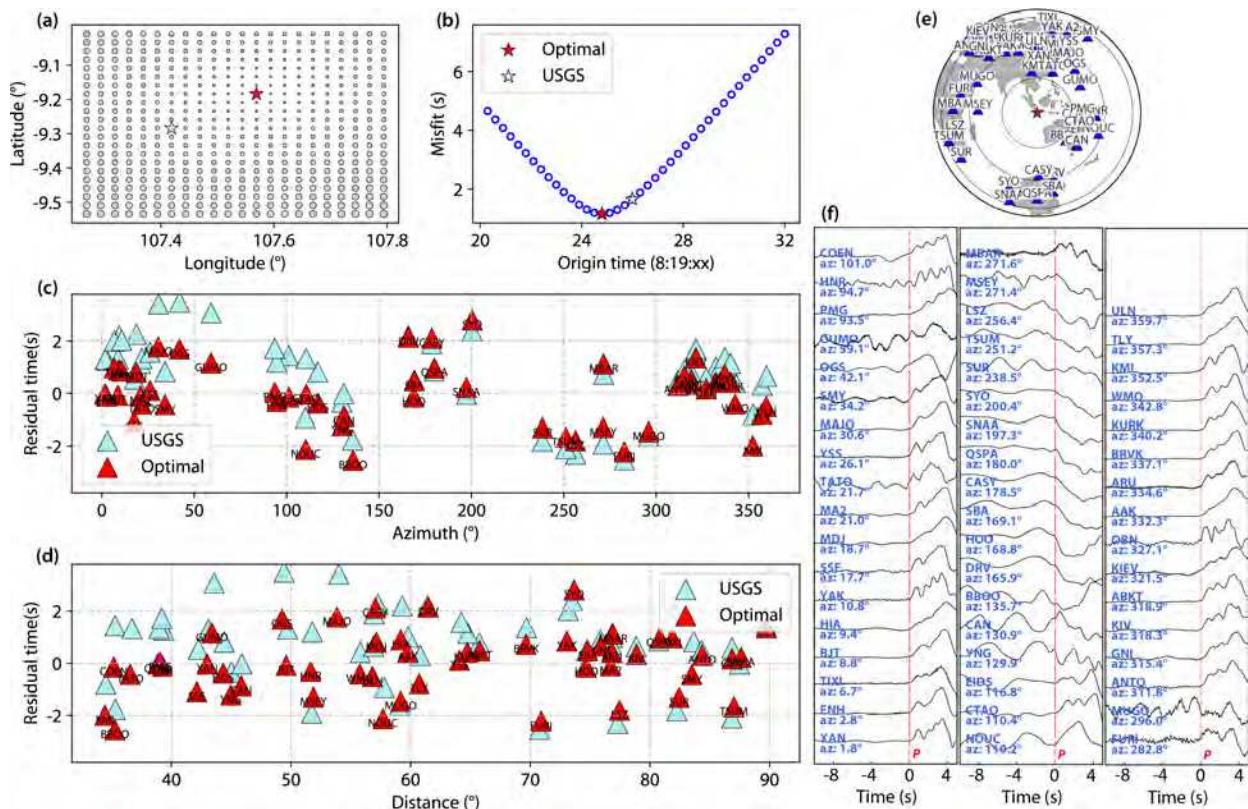
**Figure S7.** Comparison of computation domains for SIKA and Cocos stations. (a) Observed topography and bathymetry near SIKA Tide Gauge. (b) Computation domain for SIKA Gauge; the water inlet is not accurately represented in the bathymetry model, so virtual station locations (white dots) are used to estimate the response. (c-e) 12-arcsec, 6-arcsec and 1.5-arcsec resolution computation domains for Cocos Island, respectively.



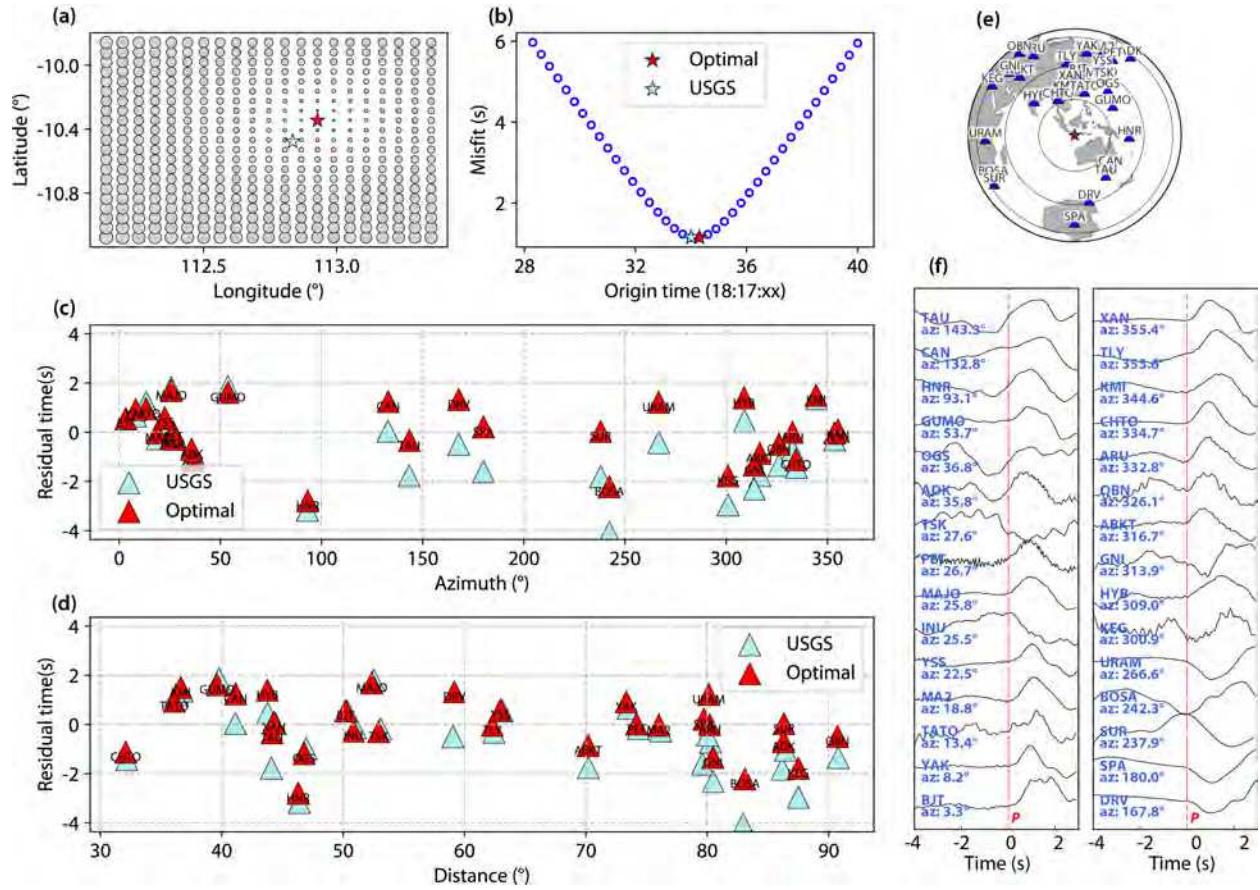
**Figure S8.** Tsunami simulation bathymetry and results for all five tide gauges from the preferred model (M04). (a-b) Location of the five tide gauges. The red bars represent the tsunami heights observed at each station. (c) Comparison of recorded (black) and computed (red) time series (left) for the preferred slip model and amplitude spectra (right) at each tide gauge.



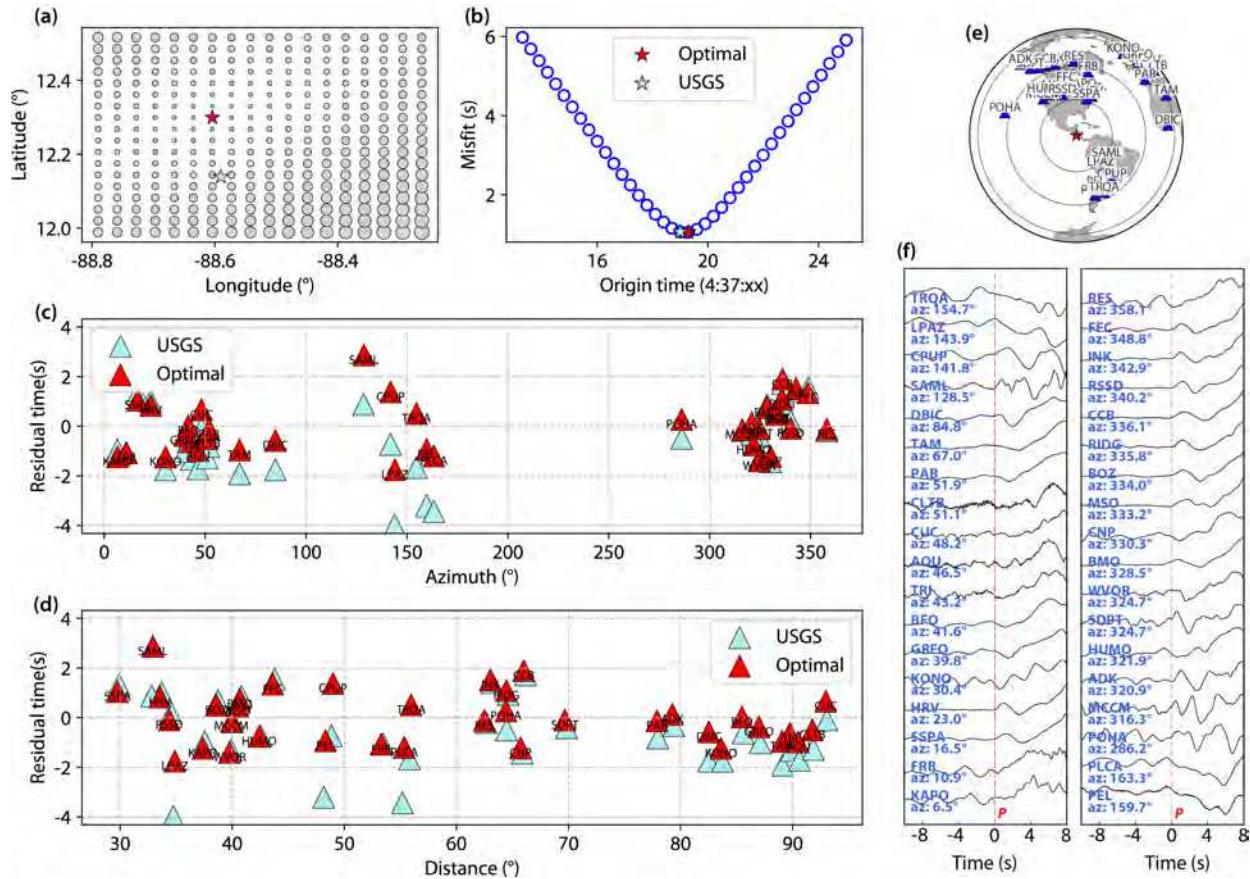
**Figure S9.** Relocation of the 2010  $M_w$  7.8 Mentawai earthquake. The depth of our optimal solution is fixed at 12.1 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



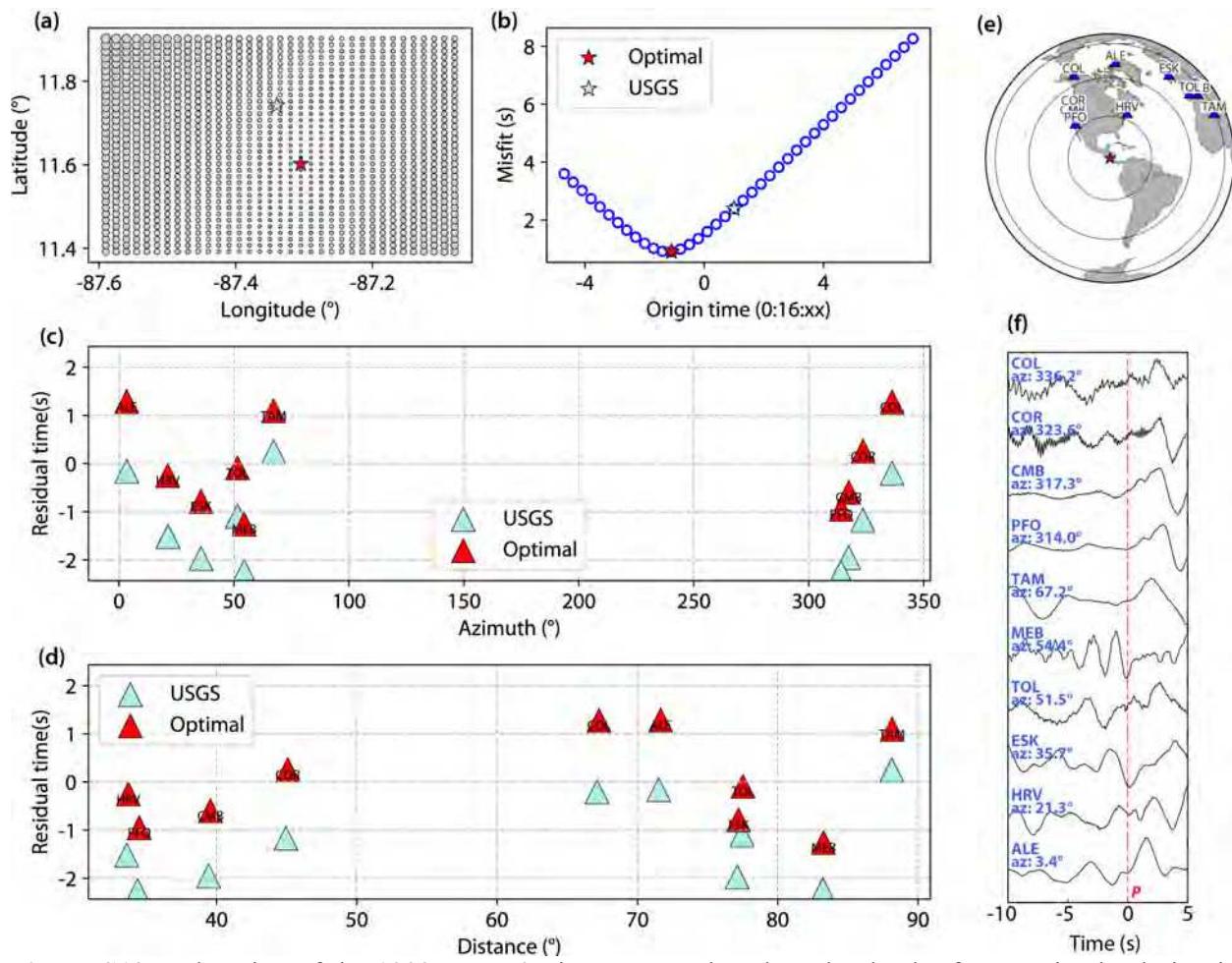
**Figure S10.** Relocation of the 2006  $M_W$  7.7 Java earthquake. The depth of our optimal solution is fixed at 12.7 km and 14.2 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



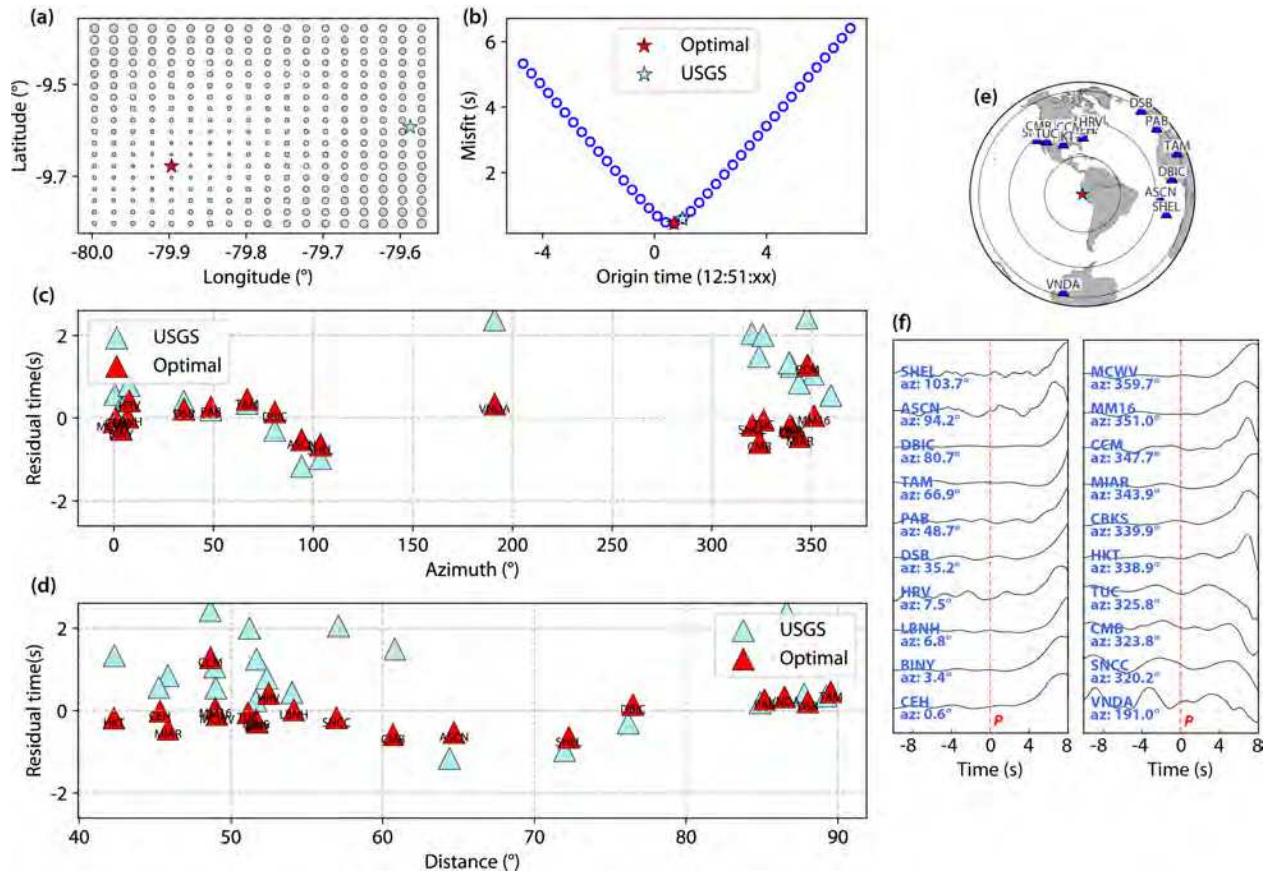
**Figure S11.** Relocation of the 1994  $M_w$  7.8 Java earthquake. The depth of our optimal solution is fixed at 11.2 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



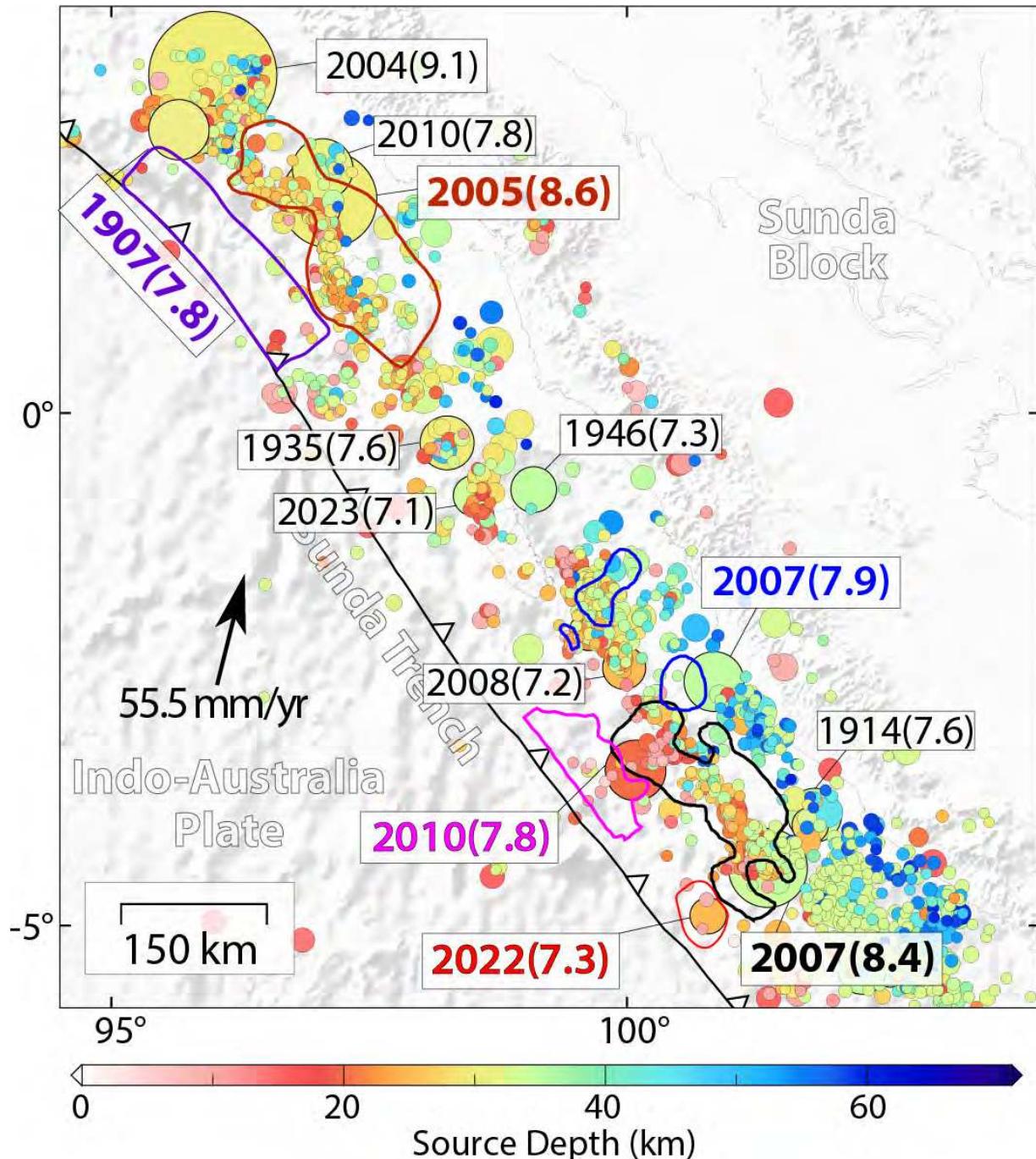
**Figure S12.** Relocation of the 2012  $M_w$  7.3 El Salvador earthquake. The depth of our optimal solution is fixed at 19.9 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



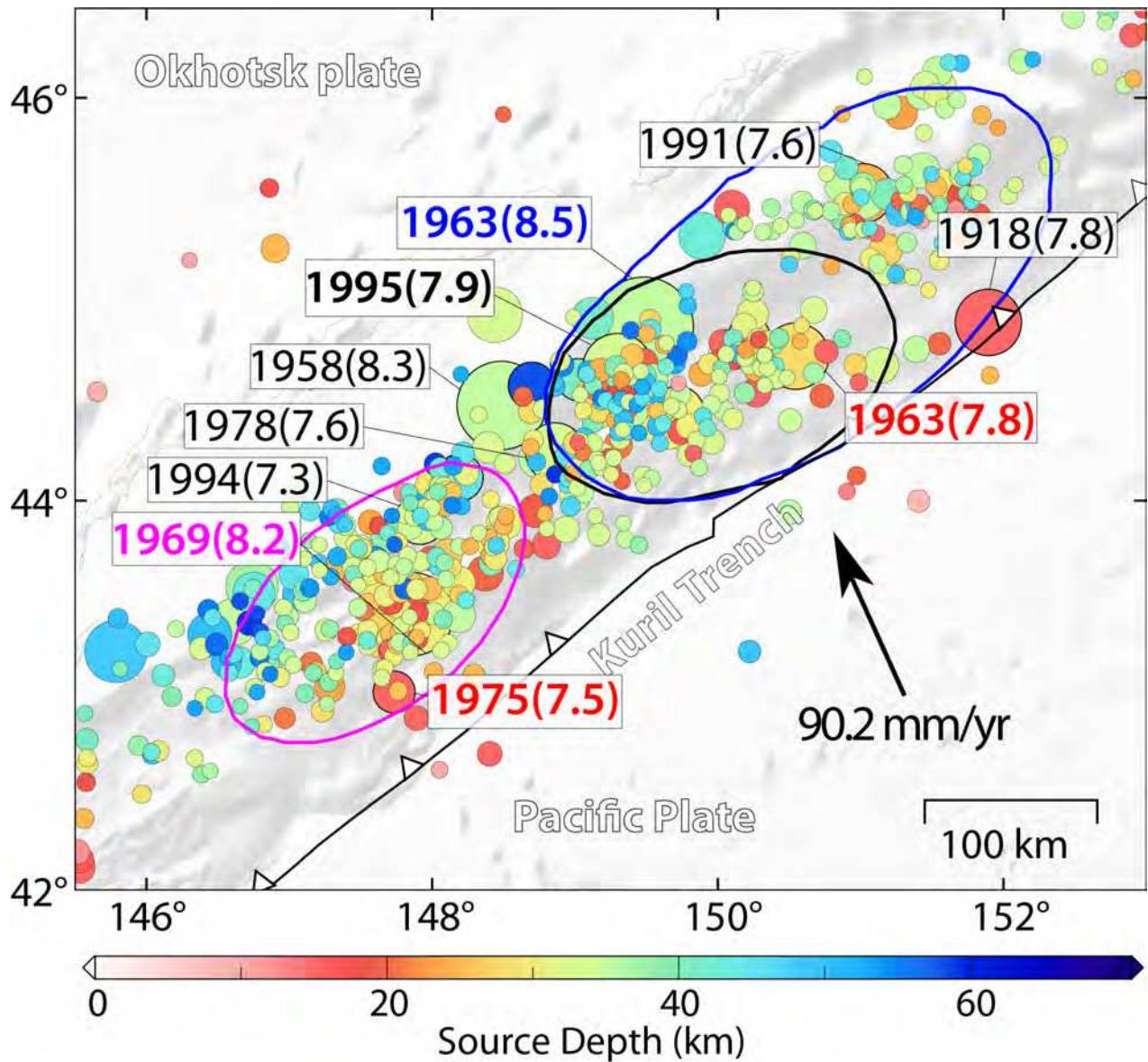
**Figure S13.** Relocation of the 1992  $M_w$  7.6 Nicaragua earthquake. The depth of our optimal solution is fixed at 19.9 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



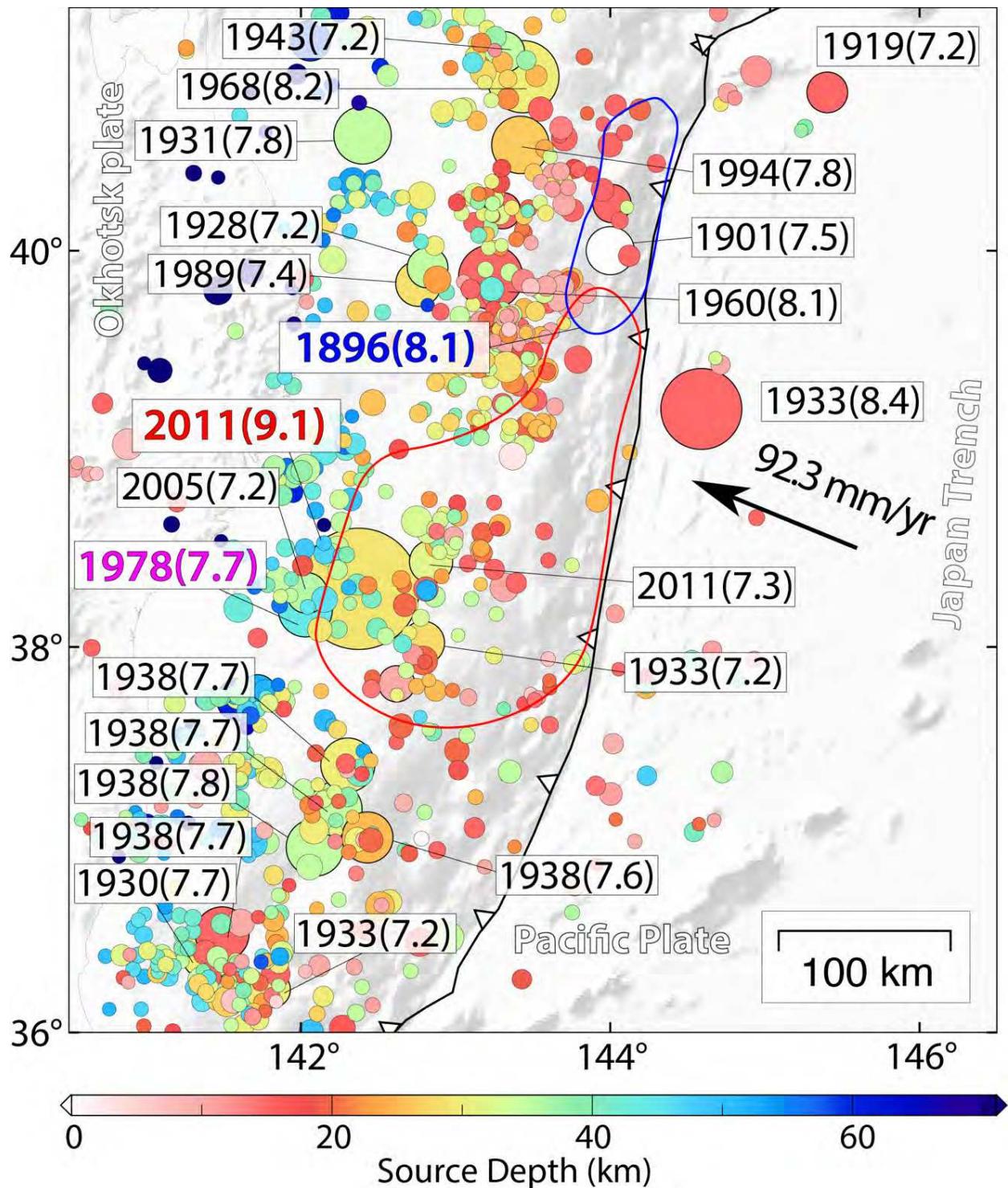
**Figure S14.** Relocation of the 1996  $M_w$  7.6 Peru earthquake. The depth of our optimal solution is fixed at 16.9 km by projecting the corresponding epicenter onto the Slab 2.0 model. Other symbols are same as Figure S3.



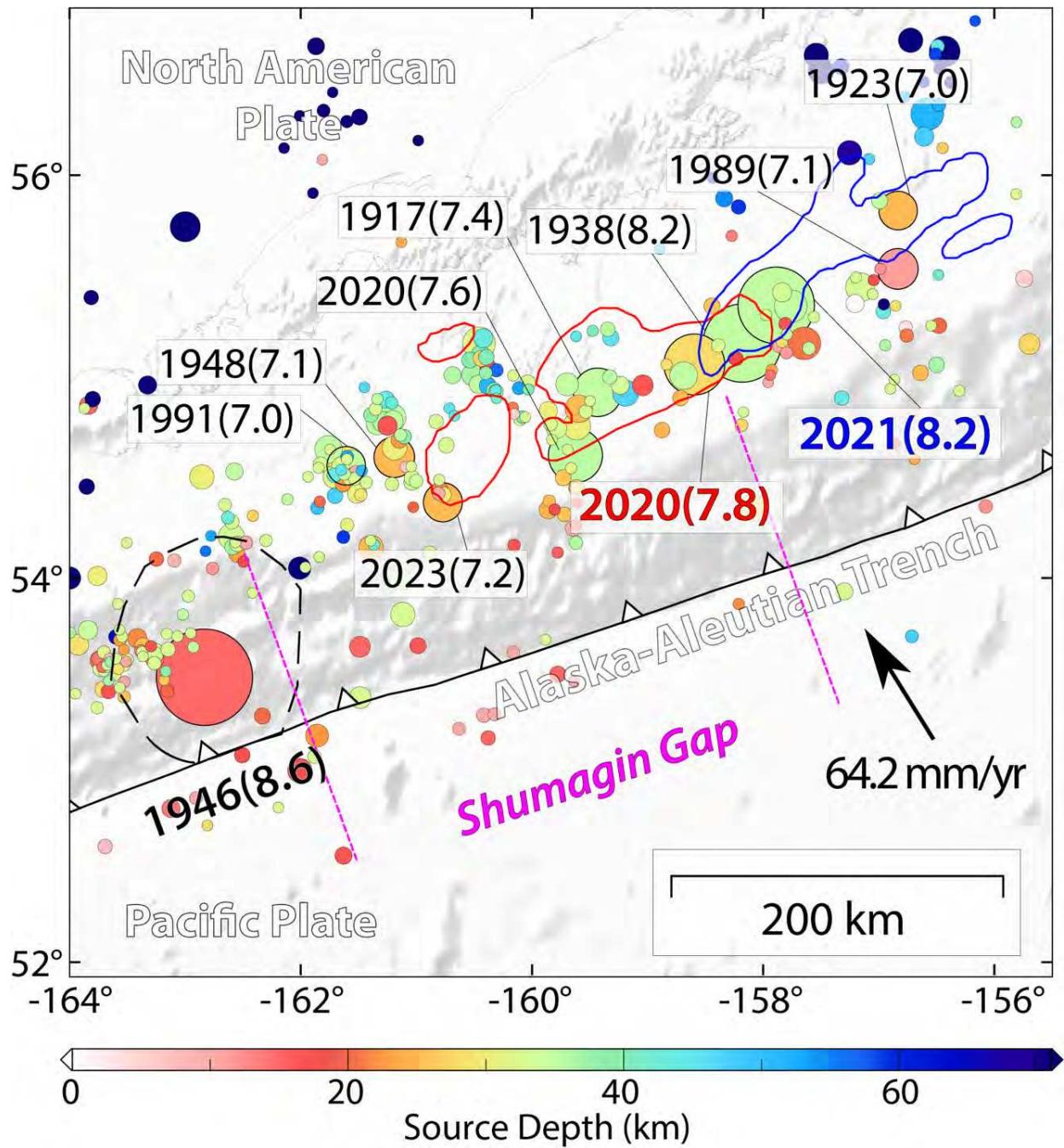
**Figure S15.** Seismicity in the southern Sumatra subduction zone. Circles show earthquake epicenters from the USGS- NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 1907  $M_w$  7.8 Sumatra (Kanamori et al., 2010; Martin et al., 2019; deep-purple curve), 2005  $M_w$  8.6 Nias (Hsu et al., 2006; brown curve), 2007  $M_w$  8.4 Bengkulu (Konca et al., 2008; black curve),  $M_w$  7.9 Pagai (Konca et al., 2008; blue curves), 2010  $M_w$  7.8 Mentawai (Yue et al., 2014; magenta curve), 2022  $M_w$  7.3 Sumatra (this study; red curve) earthquakes. Black arrow represents plate motion direction and rate of the Indo-Australia plate relative to the fixed Sunda block computed using model MORVEL (DeMets et al., 2010).



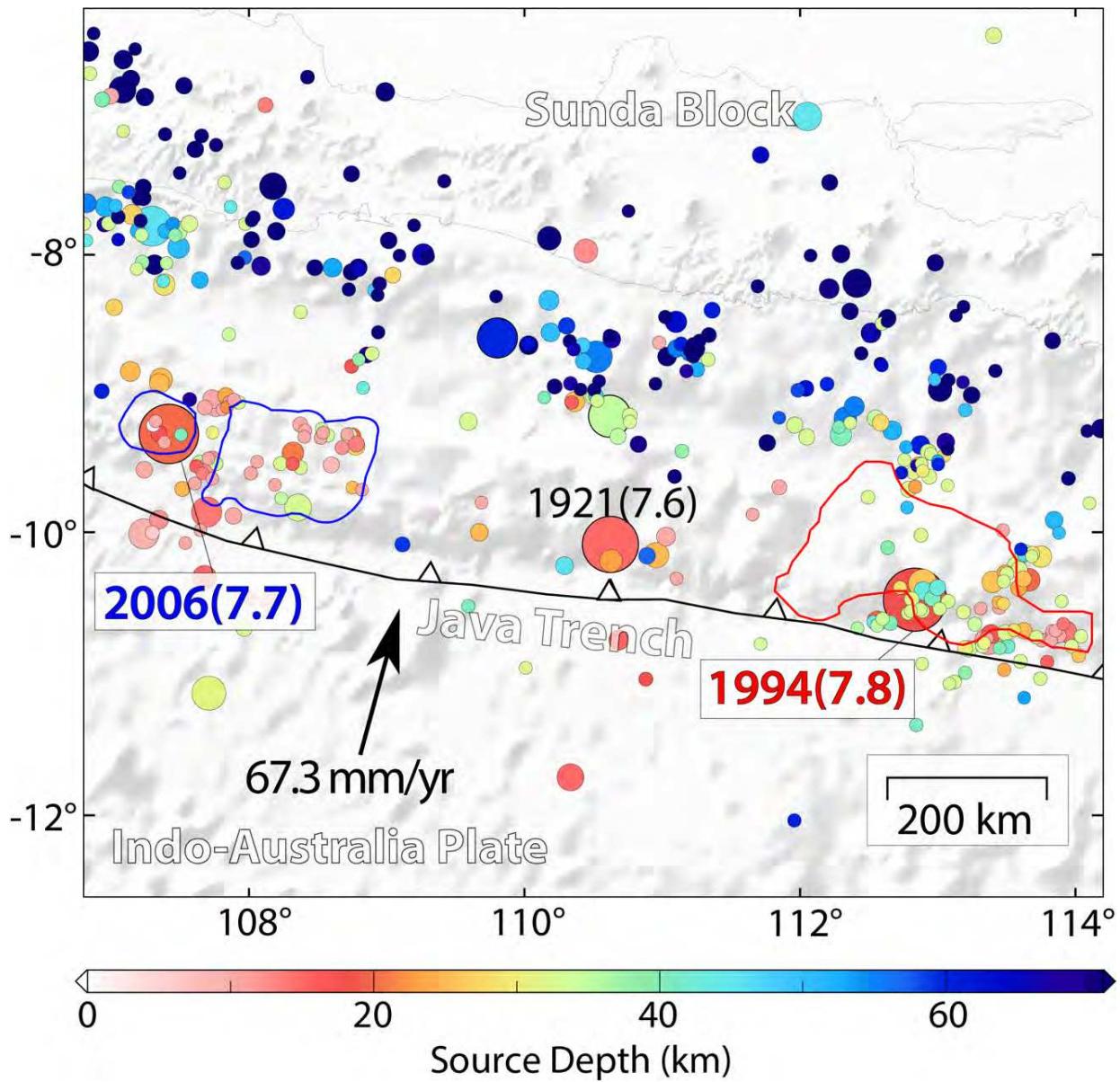
**Figure S16.** Seismicity in the Kuril subduction zone. Circles show earthquake epicenters from the USGS- NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 1963  $M_w 8.5$  (blue curve), 1969  $M_w 8.2$  (magenta curve), and the 1995  $M_w 7.9$  (black curve) Kuril events (Fukao. 1979; Pelayo & Wiens, 1992; Pérez, 2000). Black arrow represents plate motion direction and rate of the Pacific plate relative to the fixed Okhotsk plate computed using model MORVEL (DeMets et al., 2010).



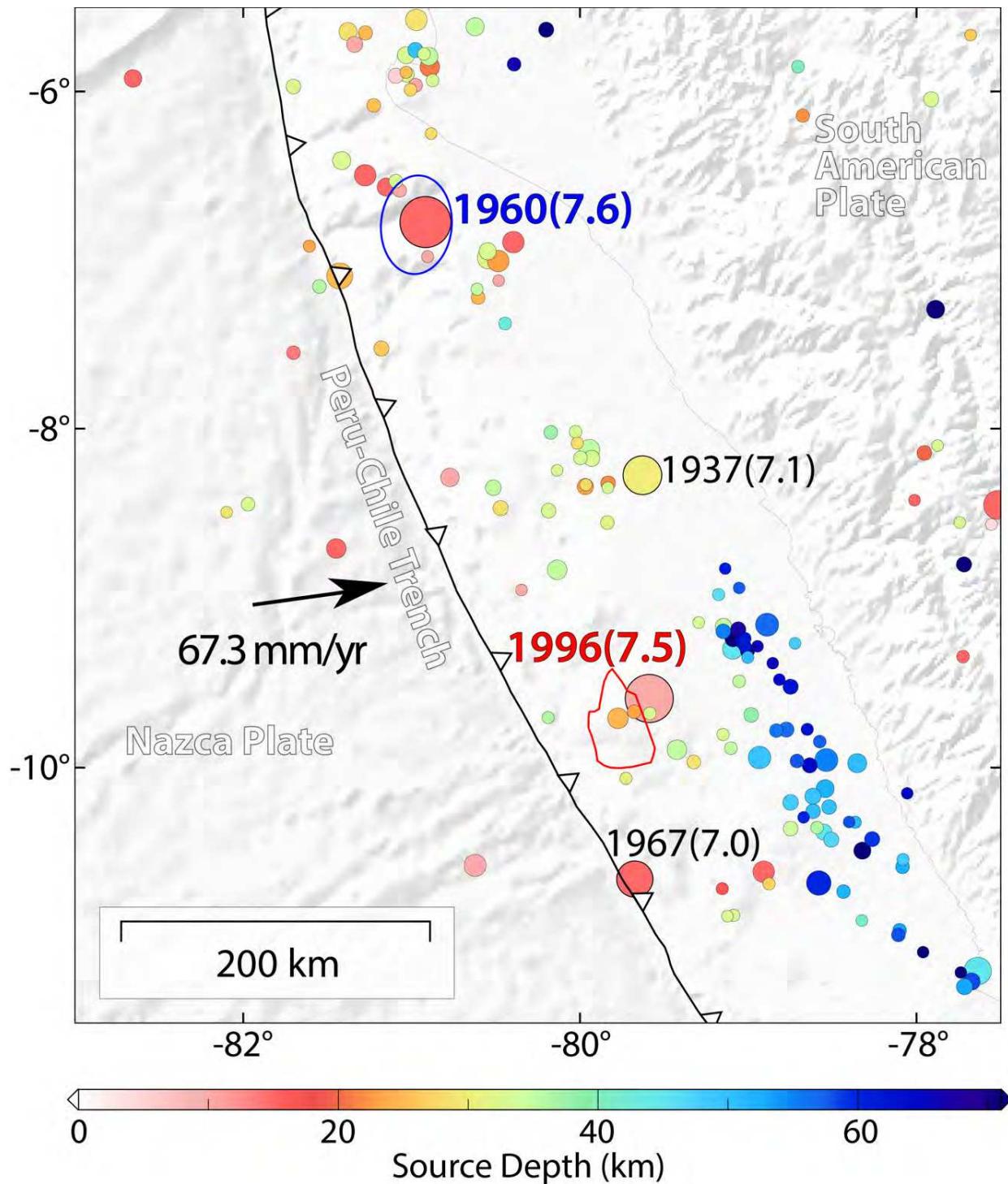
**Figure S17.** Seismicity in the Japan subduction zone along Honshu. Circles show earthquake epicenters from the USGS- NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 1896  $M_w 8.1$  Sanriku (blue curve), and the 2011  $M_w 9.1$  Tohoku (red curve) events (Ye et al., 2021). Black arrow represents plate motion direction and rate of the Pacific plate relative to the fixed Okhotsk plate computed using model MORVEL (DeMets et al., 2010).



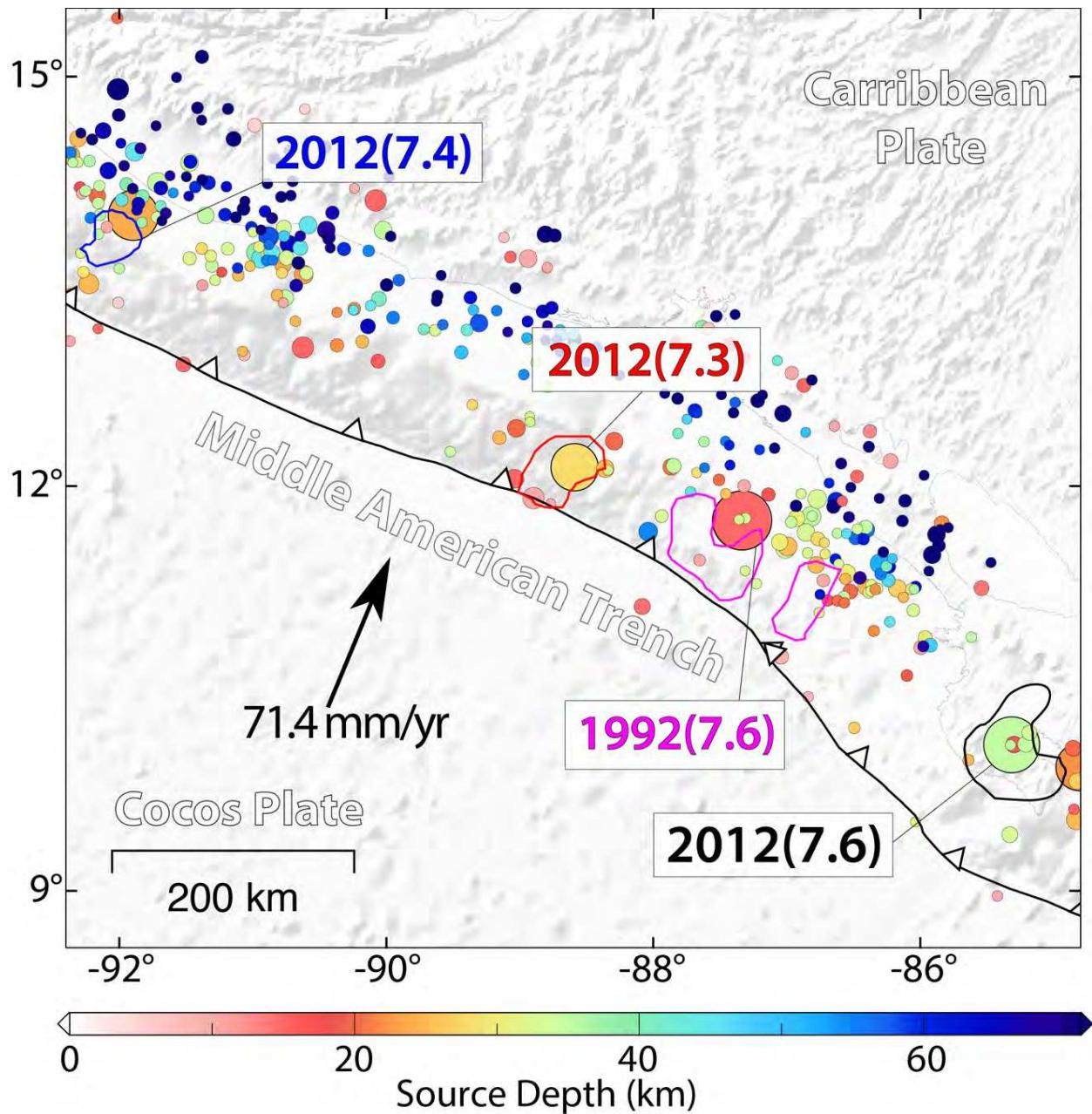
**Figure S18.** Seismicity along the Shumagin Gap Alaska-Aleutian subduction zone. Circles show earthquake epicenters from the USGS-NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 2020  $M_w$  7.8 Shumagin (Liu et al., 2023; red curves), the 2021  $M_w$  8.2 Chignik (Ye et al., 2022; blue curves) events. Black arrow represents plate motion direction and rate of the Pacific plate relative to the fixed North American plate computed using model MORVEL (DeMets et al., 2010).



**Figure S19.** Seismicity in the Java subduction zone. Circles show earthquake epicenters from the USGS-NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 1994  $M_w 7.8$  Java (Bilek and Engdahl, 2007; red curve), the 2006  $M_w 7.7$  Java events (Ammon et al., 2006; blue curves) events. Black arrow represents plate motion direction and rate of the Indo-Australia plate relative to the fixed Sunda block computed using model MORVEL (DeMets et al., 2010).



**Figure S20.** Seismicity in the Peru-Chile subduction zone. Circles show earthquake epicenters from the USGS-NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the slip distributions for the 1960  $M_w$  7.6 Peru (Pelayo and Wiens, 1992; Bilek, 2010; blue curve), 1996  $M_w$  7.5 Peru (Ihmlé et al., 1998; red curve) events. Black arrow represents plate motion direction and rate of the Nazca plate relative to the fixed South American plate computed using model MORVEL (DeMets et al., 2010).



**Figure S21.** Seismicity in the Middle American subduction zone. Circles show earthquake epicenters from the USGS-NEIC catalog since 1900, scaled with earthquake magnitude, and color-coded with source depth. Bold contours show the main slip distributions for the 1992  $M_w$  7.6 Nicaragua (Ye et al., 2013; magenta curves), 2012  $M_w$  7.3 El Salvador (Ye et al., 2013; red curve), 2012  $M_w$  7.4 Guatemala (Ye et al., 2013; black curve), and 2012  $M_w$  7.6 Costa Rica events (Ye et al., 2013; black curve). Black arrow represents plate motion direction and rate of the Cocos plate relative to the fixed Caribbean plate computed using model MORVEL (DeMets et al., 2010).