



To blow-up or not to blow-up for a granular kinetic equation

José A. Carrillo^{a,*}, Ruiwen Shu^b, Li Wang^c, Wuzhe Xu^d

^a Mathematical Institute, University of Oxford, Oxford OX2 6GG, UK

^b Department of Mathematics, University of Georgia, Athens, GA 30602, USA

^c School of Mathematics, University of Minnesota, Twin cities, MN 55455, USA

^d Department of Mathematics and Statistics, University of Massachusetts Amherst, Amherst, MA 01003, USA

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ABSTRACT

A simplified kinetic description of rapid granular media leads to a nonlocal Vlasov-type equation with a convolution integral operator that is of the same form as the continuity equations for aggregation-diffusion macroscopic dynamics. While the singular behavior of these nonlinear continuity equations is well studied in the literature, the extension to the corresponding granular kinetic equation is highly nontrivial. The main question is whether the singularity formed in velocity direction will be enhanced or mitigated by the shear in phase space due to free transport. We present a preliminary study through a meticulous numerical investigation and heuristic arguments. We have numerically developed a structure-preserving method with adaptive mesh refinement that can effectively capture potential blow-up behavior in the solution for granular kinetic equations. We have analytically constructed a finite-time blow-up infinite mass solution and discussed how this can provide insights into the finite mass scenario.

1. Introduction

Granular flows are omnipresent in nature, from large scale mudslides to small scale table top experiments [1,2]. Rapid granular flows or granular gases consist of a large number of small discrete grains, which interact by instantaneous hard-sphere like collisions [3]. The physical modeling of granular gases has been revolutionized in the past decades due to the maturity of kinetic theory. Unlike the ideal gas particles, the distinct feature of granular particles is the intrinsic inelasticity of the collisions between grains. As a result, there is a dissipation of energy, which brings a perturbed system quickly to rest.

Deriving kinetic equations from statistical mechanics of particle systems undergoing inelastic collisions faces important issues such as the inelastic collapse [4], i.e. infinite many collisions in finite time. However, the kinetic description of granular gases [5–7] has been successful in computing transport coefficients for hydrodynamic descriptions used in experiments far from their supposed limits of validity [8–14]. We refer to the seminal book of Brilliantov and Pöschel [3] and the recent review [15] for further physical and mathematical details and a comprehensive list of references in the subject.

Let us consider particles moving in a one dimensional setting. Denoting by $v, v_* \in \mathbb{R}$ the velocities of two particles before their collision, and v' and v'_* the velocities after collision, we can write the postcollisional velocities in terms of the precollisional ones as

$$v' = v - \varepsilon(v - v_*), \quad v'_* = v_* + \varepsilon(v - v_*)$$

for $0 \leq \varepsilon = \varepsilon(|v - v_*|, \theta) \leq 1$. This immediately implies the energy dissipation:

$$|v'_*|^2 + |v'|^2 - |v|^2 - |v_*|^2 = -\varepsilon(1 - \varepsilon)(v - v_*)^2 \leq 0,$$

and therefore, ε controls the loss of kinetic energy and is referred as the restitution coefficient while θ is a parameter controlling the strength of inelasticity. Typical restitution coefficients are of the form $\varepsilon(|v - v_*|, \theta) = \varepsilon_0(1 + \theta|v - v_*|^\beta)^{-1}$, with $\beta \geq -2$ and $0 \leq \varepsilon_0 \leq 1$. Then following the same formal derivation of the Boltzmann equation from the particle dynamics, the kinetic description of the above inelastic collision takes the form [16–19]

$$\partial_t f + v \cdot \partial_x f = \int_{\mathbb{R}} \int_{\mathbb{R}_+} B(\theta) \left\{ |v' - v'_*|^2 \frac{1}{J} f(v') f(v'_*) - |v - v_*|^2 f(v) f(v_*) \right\} \times dv_* d\theta, \quad (1.1)$$

where $f(t, x, v) : (0, \infty) \times \mathbb{R} \times \mathbb{R} \mapsto \mathbb{R}^+$ denotes the probability density function of grains and J is the Jacobian of the transformation from (v, v_*) to (v', v'_*) , and B is the rate function.

As opposed to the Boltzmann equation whose equilibrium is the Maxwellian, the equilibrium of (1.1) is a Dirac mass located at the mean velocity of particles, which means that all particles are at rest in the comoving frame. This is a direct consequence of energy dissipation. If in addition to the friction, the granular material is put in interaction

* Corresponding author.

E-mail addresses: carrillo@maths.ox.ac.uk (J.A. Carrillo), ruiwen.shu@uga.edu (R. Shu), liwang@umn.edu (L. Wang), wuzhexu@umass.edu (W. Xu).

with a thermal bath, then a combined effect of friction and diffusion may lead to a non-Gaussian, non-Dirac equilibrium [20–22].

In the quasi-elastic regime, i.e., $\varepsilon \simeq 1$, (1.1) can be formally approximated by [16,17,19]

$$\begin{aligned} \partial_t f + v \partial_x f &= \frac{\lambda}{2} \partial_v \left(f \int_{\mathbb{R}} |v-w|^{\beta+1} (v-w) f(w) dw \right) \\ &=: \frac{\lambda}{2} \partial_v ((\partial_v W *_{\nu} f) f), \end{aligned} \quad (1.2)$$

where the collision now is described by a nonlocal interaction with kernel

$$W(|v|) = \frac{|v|^\gamma}{\gamma}, \quad (1.3)$$

where $\gamma = \beta + 3$. This is also a reminiscent of the *grazing collision limit* of Boltzmann operators.

The spatially homogeneous version of (1.2), that bares the name of aggregation equation (upon replacing v by x), has been well studied in the literature [23,24]. In particular, starting from a bounded initial data, there is a sharp threshold that distinguishes the finite time blow-up versus global in time solution in terms of the Osgood condition. However, by allowing a spatial dependence with a free transport dynamics, it generates substantial difficulties in understanding its solution behavior [25]. In essence, the lingering question is:

Will the singularity formed in v -direction enhanced or mitigated by the shear?

Here, we tackle this problem through a series of numerical experiments followed by a heuristic argument. As with many other physical systems, such as the Navier–Stokes equations for fluids and the Boltzmann equation for rarefied gas, a naive discretization would easily lead to unstable, physics-violated numerical solutions. In principle, this issue is addressed with the aid of numerical analysis, which provide a theoretical guarantee and practical insight to the numerical schemes. Despite that numerical analysis has its own theory and tools, it is undoubtedly that PDE analysis is the stepping stone for the development of the numerical analysis of PDEs. Said differently, one needs to have at least some well-posedness result or apriori estimate of solution for a PDE before one can simulate it numerically with confidence. Otherwise it is ambiguous to assert that the computed numerical solutions are physically relevant.

However, when it comes to possible singular solutions, the theory is often lacking. One would instead rely on numerical solutions to give some guidance on the possible solutions [26–28]. This is the approach we are taking here. To this end, we will follow the splitting approach in [29] which gives rise to a local in time weak solution, and our design principle of the numerical scheme is to preserve as much physical quantities as possible, including conservation and dissipation. Our main challenge to numerically solve (1.2) is the possible singularity formation. For this problem, we will integrate a mesh refinement technique [27] that rearranges a certain amount of grid points dynamically, and therefore ensures the high resolution within the possible blow-up region. A slight modification that guarantees mass conservation will also be added.

Theoretically, for the case $\gamma = 2$, we construct a family of self-similar analytic solutions with infinite total mass. The maximal time T of existence for such solutions is always finite, and there is a critical threshold which determines whether the velocity support shrinks (in the super-critical case) or expands (in the sub-critical case) as $t \rightarrow T^-$. Heuristic arguments show that in the super-critical case such infinite-mass solution could be approximated by a finite-mass solution, leading to the finite-time blow-up of a finite-mass solution, while in the subcritical case a similar finite-mass solution will not have finite-time blow-up. This threshold for finite-mass solution is verified by numerical simulations. Further heuristic analysis suggests that this mechanism will not lead to a finite-time blow-up if $\gamma > 2$. The case $\gamma = 2$ has a special structure and global-in-time finite mass measure solutions have

been constructed by optimal transport theory [30], where finite mass concentration is neither proved or disproved.

In Section 2, we first summarize the solution properties of the spatially homogeneous version of (1.2), as well as the physical quantities that we intend to preserve for the spatially inhomogeneous version. In the next subsection, we introduce the mesh refinement technique and then build it into a variational formulation—the Jordan–Kinderlehrer–Otto scheme. Section 3 is devoted to the numerical tests. We first verify the performance of our scheme validating our numerical solution with theoretical predictions for the spatially homogeneous case. We then carry out a series of systematic experiments to conjecture the solution behavior in the spatially inhomogeneous case.¹ We finally offer theoretical evidence in Section 4 to substantiate our numerical findings.

2. Structure-preserving numerical schemes

We first discuss the main properties of solutions to the equations of interest. We focus next on the development of suitable numerical schemes keeping these properties at the discrete level.

2.1. Solution properties

Consider spatially homogeneous case of (1.2), i.e.,

$$\partial_t f = \partial_v ((\partial_v W *_{\nu} f) f) \quad (2.1)$$

with radial kernel $W(v) = w(|v|)$. We first cite the following theorem [24] regarding its sharp threshold dynamics.

Theorem 1. *Starting from a bounded, compactly supported, nonnegative initial data,*

- if $\frac{w'(r)}{w(r)}$ is monotone decreasing, $w''(r) > 0$ and $\int_0^1 \frac{1}{w'(r)} dr < +\infty$,

then the solution to (2.1) blows up in finite time. The blow-up time depends only on the initial data through its radius of support and total mass;

- if $\frac{w'(r)}{w(r)}$ is monotone decreasing, $w''(r) > 0$ and $\int_0^1 \frac{1}{w'(r)} dr = +\infty$,

then the solution to (2.1) stays bounded for all time and converges as $t \rightarrow \infty$ to a Dirac centered at the center of mass of the initial data.

The conditions (2.2) and (2.3) stem from the Osgood condition for well-posedness of the ODE $\frac{dR}{dt} = -w'(R)$, with the former violating the condition and the latter satisfying it. This theorem, applied to the kernel (1.3), implies that: infinite time blow-up for $\gamma > 2$ and finite time blow-up for $-1 < \gamma < 2$, see [23] for further properties. The case of $\gamma = 2$ leads to infinite time blow-up as it can be directly checked since it is a linear equation once conservation of mass and momentum are taken into account.

More particularly, when $\gamma = 3$ in (1.3), with a smooth enough initial condition centered at origin with a compact support, the solution as $t \rightarrow \infty$ takes the asymptotic form $f(t, v) \sim \frac{1}{2} \delta_{\frac{1}{t}} + \frac{1}{2} \delta_{-\frac{1}{t}}$. Therefore, when $t \rightarrow \infty$, the solution tends to form two symmetric Dirac delta functions, whose centers are converging to $v = 0$ at rate $\frac{1}{t}$. More details can be found [19,23,31,32].

When $\gamma = 1$, we have explicit prediction for the blow-up time since (2.1) can be viewed as the derivative of the Burgers equation [33].

¹ Matlab codes will be available on Github <https://github.com/woodssss/Granular-kinetic-equation>.

Indeed, $W = |v|$ and $\partial_v W = \text{sign}(v)$, $\partial_v^2 W = 2\delta(v)$. By setting $u(t, v) = -\int_{-\infty}^v f(t, y)dy$, we then have

$$\partial_t u + 2u\partial_v u = 0$$

with initial condition

$$u_0(v) = -\int_{-\infty}^v f_0(y)dy.$$

Then the blow-up time is equal to the shock formation time of Burgers equation, i.e.

$$T = -\frac{1}{2 \min_v u'_0(v)}.$$

Next we cite a few apriori properties of the solutions to the spatially inhomogeneous case proposed in [25,29].

Proposition 1. For $\gamma > 1$, the solution to (1.2) satisfies

- *Mass conservation:* $\frac{d}{dt} \int_{\mathbb{R} \times \mathbb{R}} f(t, x, v) dx dv = 0$.
- *Momentum conservation:* $\frac{d}{dt} \int_{\mathbb{R} \times \mathbb{R}} v f(t, x, v) dx dv = 0$.
- *Decrease of moments for $p \geq 2$:* $\frac{d}{dt} \int_{\mathbb{R} \times \mathbb{R}} |v|^p f(t, x, v) dx dv \leq 0$.
- *Increase of internal energy:* for a C^1 convex function U with $U(0) = 0$, we have $\frac{d}{dt} \int_{\mathbb{R} \times \mathbb{R}} U(f(t, x, v)) dx dv \geq 0$.

In this proposition, the third part, when $p = 2$, signifies the dissipation of energy, as expected in an inelastic collision. When $p = \infty$, it implies that for compactly supported initial data $f(0, x, \cdot)$, the support of $f(t, x, \cdot)$ remains compact for all $t > 0$. Moving to the fourth part, a notable selection is $U(f) = f \ln f$, representing the negative of the entropy.

2.2. Numerical method

In this subsection, we present a variational semi-Lagrangian scheme for the granular kinetic Eq. (1.2). As demonstrated in the theoretical paper [29], we divide the equation into two stages: a transport step that is explicitly addressed, and a collision step that is treated implicitly. More precisely, denote $f^n(x, v) \approx f(t_n, x, v)$, where $t_n = n\Delta t$ with Δt being the time step, we have:

$$\frac{f^{n+\frac{1}{2}} - f^n}{\Delta t} + v\partial_x f^n = 0, \quad (2.4)$$

and

$$\frac{f^{n+1} - f^{n+\frac{1}{2}}}{\Delta t} = \partial_v((\partial_v W *_{\nu} f^{n+1})f^{n+1}). \quad (2.5)$$

Periodic boundary condition in x and no flux boundary condition in v will be used throughout the paper.

2.3. Semi-Lagrangian scheme for the transport step

For the transport step (2.4), we will use the semi-Lagrangian scheme, which combines the method of characteristic and interpolation. Specifically, let $[-L_x, L_x]$ be the computational domain and define

$$x_i^0 = -L_x + i\Delta x/2, \quad 1 \leq i \leq N_x, \quad \Delta x = 2L_x/N_x,$$

as the initial grid points. As time progresses, the grids $\{x_i^n\}$ will be updated and become nonuniform due to mesh refinement (details of which will be provided in Section 2.4.3). Given $\{f^n(x_i^n)\}_{i=1}^{N_x}$, $\{f^{n+1}(x_i^{n+1})\}$ can be acquired by tracing back the characteristics:

$$f^{n+1}(x_i^{n+1}, v) = f^n(x_i^{n+1} - v\Delta t, v),$$

where $f^n(x_i^{n+1} - v\Delta t, v)$ is evaluated using piecewise cubic Hermite interpolation. To guarantee mass conservation, we re-weight the interpolated value by the total mass, namely,

$$f^{n+1}(x_i^{n+1}, v) := \max\{f^n(x_i^{n+1} - v\Delta t, v), 10^{-10}\} \frac{M_n}{M_{n+1}}, \quad (2.6)$$

where

$$M_n = \sum_{i=2}^{N_x-1} f^n(x_i^n, v) \frac{x_{i+1}^n - x_{i-1}^n}{2} + f^n(x_1^n, v)(L_x + \frac{x_2^n + x_1^n}{2}) + f^n(x_{N_x}^n, v)(L_x - \frac{x_{N_x}^n + x_{N_x-1}^n}{2}),$$

and let $y_i := x_i^{n+1} - v\Delta t$

$$M_{n+1} = \sum_{i=2}^{N_x-1} f^{n+1}(y_i, v) \frac{y_{i+1} - y_{i-1}}{2} + f^{n+1}(y_1^n, v)(L_x + \frac{y_2 + y_1}{2}) + f^{n+1}(y_{N_x}, v)(L_x - \frac{y_{N_x} + y_{N_x-1}}{2}),$$

where $f^{n+1}(y_i, v)$ is defined in (2.6).

2.4. Regularized JKO scheme with adaptive mesh refinement for collision step

This section centers around solving the collision step (2.5), for which we will adopt a regularized JKO scheme along with a mesh refinement technique. In the following, we first introduce the Fisher information regularized JKO scheme that was proposed in [34], and extend it to the non-uniform mesh case here. Secondly, in order to increase the resolution at the blow-up region, we use a mesh refinement technique that was originally proposed in [27], which adaptively redistribute a large portion of the grid points to the blow-up region.

2.4.1. Regularized JKO scheme

To start, rewrite (2.1) as

$$\partial_t f = \partial_v(f \partial_v \frac{\delta E(f)}{\delta f}), \quad (2.7)$$

where the Energy functional is defined as

$$E(f) = \frac{1}{2} \int_{\mathbb{R}} \int_{\mathbb{R}} W(v-w) f(v) f(w) dv dw,$$

and $\frac{\delta E}{\delta f}$ is the functional derivative of E . Then (2.7) can be interpreted as a gradient flow of the E with respect to the Wasserstein metric, and therefore admits the following minimizing movement scheme, also called JKO scheme:

$$f^{n+1} \in \underset{f \in \mathcal{P}_{ac}(\Omega_v)}{\text{argmin}} \left\{ \frac{1}{2} d_W(f, f^{n+\frac{1}{2}})^2 + \Delta t E(f) \right\}. \quad (2.8)$$

Here $d_W(f, f^n)$ is the Wasserstein distance between f and f^n and $\mathcal{P}_{ac}(\Omega_v)$ is the set of probability measures on Ω_v that are absolutely continuous with respect to Lebesgue measure. In this paper, we use the Benamou-Brenier's dynamic formulation to numerically compute the Wasserstein distance. In particular, the distance between two measures $f_0 dx$ and $f_1 dx$ is

$$d_W(f_0, f_1)^2 = \min_{(f, m) \in C_1} \int_0^1 \int_{\Omega_v} \Phi(f(t, v), \|m(t, v)\|) dv dt, \quad (2.9)$$

where

$$\Phi(f, \|m\|) = \begin{cases} \frac{\|m\|^2}{f} & \text{if } f > 0, \\ 0 & \text{if } (f, m) = (0, 0), \\ +\infty & \text{otherwise,} \end{cases}$$

and the constraint set C_1 consists of

$$\partial_t f + \nabla_v \cdot m = 0 \text{ on } \Omega_v \times [0, 1], \quad m \cdot v = 0 \text{ on } \partial\Omega_v \times [0, 1], \\ f(\cdot, v) = f_0, \quad f(\cdot, 1) = f_1 \text{ on } \Omega_v.$$

Here v is the outer normal direction of Ω_v . Substituting (2.9) into (2.8), and write $m = fv$, we end up with a convex optimization with a linear constraint:

$$\begin{cases} (f, m) = \arg \min_{f, m} \frac{1}{2} \int_0^1 \int_{\Omega_v} \Phi(f, \|m\|) dv dt + \Delta t E(f(1, \cdot)) \\ \text{s.t.} \quad \partial_t f + \nabla_v \cdot m = 0, f(0, v) = f^n(v), \quad m \cdot v = 0. \end{cases} \quad (2.10)$$

In practice, to further increase the convexity of the objective function and preserve the positivity of the desired minimizer f^* , we add a Fisher information regularization term

$$\beta^{-2} \Delta t^2 \int_{\Omega_v} \|\nabla \log f\|^2 f dv$$

to the objective function and reach the following regularized problem, originally proposed in [34]:

$$f^{n+1}(v) \in \arg \inf_{m, f} \int_{\Omega_v} \frac{\|m(t)\|^2}{f(t)} + \beta^{-2} \Delta t^2 \|\nabla \log f(v)\|^2 f(v) dv + 2 \Delta t E(f), \quad (2.11)$$

such that

$$f - f^{n+\frac{1}{2}} + \nabla_v \cdot m = 0, \quad m \cdot v = 0. \quad (2.12)$$

Here the artificial time in the dynamic formulation (2.10) is replaced by one step finite difference in time. This is shown in [34] that significantly reduces the computational cost without violating the first order accuracy in time.

2.4.2. Non-uniform velocity discretization

Different from [34], we will use a non-uniform discretization in v so as to facilitate the later mesh refinement proposal. To be more specific, denote the computational domain in v to be $v \in [-L, L]$ and the number of grid points as N_v . Let $\{v_i\}_{i=1}^{N_v}$ be the grid points, and define

$$\Delta v_i := v_i - v_{i-1}, \quad h_i = \frac{1}{2}(\Delta v_i + \Delta v_{i+1}).$$

To discretize (2.11) and (2.12), we evaluate f at the grid points and m at the half grid points, i.e., i.e. $m_{i-\frac{1}{2}} = m(\frac{v_i + v_{i-1}}{2})$. Then the fully discretized version of (2.11) and (2.12) reads

$$f^{n+1} \in \arg \min F(f, m) + \Delta t E(f, m), \quad (2.13)$$

where

$$F(f, m) = \sum_{j=2}^{N_v} \left[\frac{2m_{j-\frac{1}{2}}^2}{f_j + f_{j-1}} + \frac{\beta^{-2} \Delta t^2}{\Delta v_i^2} (\log f_j - \log f_{j-1})^2 \frac{f_j + f_{j-1}}{2} \right] \Delta v_j,$$

and

$$E(f) = \sum_{i,l=1}^{N_v} W_{i,l} f_i f_l h_i h_l, \quad W_{i,l} = W(v_i - v_l).$$

The constraint function reads

$$f_j^{n+1} - f_j^n + \frac{m_{j+\frac{1}{2}} - m_{j-\frac{1}{2}}}{h_j} = 0. \quad (2.14)$$

Denote $u = [f, m]$, then the constraint (2.14) can be reformulated as a linear system $Au = b$. By introducing the indicator function, (2.13) together with (2.14) can be reformulated into an unconstrained optimization problem:

$$\min_u J(u) + \chi(u), \quad \chi(u) = \begin{cases} 0 & Au = b, \\ +\infty & \text{otherwise} \end{cases}, \quad (2.15)$$

where

$$J(u) := F(u) + \Delta t E(u). \quad (2.16)$$

As discussed in [34], thanks to the convexity of the $J(u)$, we can invoke the sequential quadratic programming to solve (2.15):

$$\begin{cases} z^{(l+1)} \in \arg \min_z \frac{1}{2} (z - u^{(l)})^T H^{(l)} (z - u^{(l)}) \\ \quad + \nabla J(u^{(l)})^T (z - u^{(l)}) + \chi(z), \\ u^{(l+1)} = u^{(l)} + t_l (z^{(l+1)} - u^{(l)}). \end{cases}$$

Here $H^{(l)}$ is an approximation of $\nabla^2 J(u^{(l)})$, for our particular form of J in (2.16), it has the form

$$(H^{(l)})_{i,j} = \begin{cases} (\nabla^2 F(u^{(l)}))_{i,i}, & \text{if } i = j, \\ 0, & \text{else.} \end{cases}$$

We now summarize our one step regularized JKO scheme in Algorithm 1, which is essentially the same as in [34]. Note that when the exit flag as defined in Algorithm 1 becomes $e_f = 1$, it signifies that the minimizer cannot be reached even after a sufficiently large number of iterations. This condition can be regarded as a numerical indicator of blow-up, and we will demonstrate its significance in the forthcoming numerical examples section.

Algorithm 1: One step regularized JKO scheme

Input : Grid points $\{v_j\}_j$, function value $\{f^n(v_j) =: f_j^n\}_j$, time step Δt , optimization step λ , max iteration number I_{max} and a stopping criteria ϵ

Output: $\{f_j^{n+1}\}_j$ and exit flag e_f

```

1 1. Compute  $\Delta v_i := v_i - v_{i-1}$  and  $h_i = \frac{\Delta v_i + \Delta v_{i+1}}{2}$ .
2 2. Let  $l = 0, \zeta = 1, e_f = 0$ ;
3 while  $l < I_{max}$  do
4   if  $\zeta > \epsilon$  then
5      $z^{(l+1)} \in$ 
6        $\arg \min_z \frac{1}{2} (z - u^{(l)})^T H^{(l)} (z - u^{(l)}) + \nabla J(u^{(l)})^T (z - u^{(l)}) + \chi(z)$ 
7       ;
8      $u^{(l+1)} = u^{(l)} + \lambda (z^{(l+1)} - u^{(l)})$ ;
9      $\zeta = \frac{\|u^{(l+1)} - u^{(l)}\|}{\|u^{(l)}\|}$ ;
10   else
11     Break;
12   end
13 3.  $e_f = 1$  if  $l = I_{max}$ ;

```

2.4.3. Adaptive mesh refinement

To investigate the possible blow-up in the solution, we need to keep increasing the resolution in regions where the solution is concentrated. Here we follow the approach developed in [27]. The main concept is to create a dynamic mapping between the original domain and the interval $[0, 1]^d$. This mapping ensures that the grid points in $[0, 1]^d$ remain uniformly spaced and finite. However, when transferring the grid back to the original domain, the points are concentrated more heavily in the vicinity of singularities.

To better illustrate the idea, let us consider (2.1) with $v \in [-L, L]$ and initial condition $f(0, v) = f_0(v)$. Choose two parameters δ_0 and δ , both in $(0, 1)$, and let

$$S_0 := \{v : |f^0(v)| \geq \delta_0 \|f^0\|_\infty\}, \quad (2.17)$$

which indicates the blow-up region. Now we intend to find a one-to-one map $v = \mu(s)$ between $v \in [-L, L]$ and an auxiliary variable $s \in [-1, 1]$. Specifically, we seek a mapping such that if we distribute N uniform grid points along s , then the resulting grid in v contains δN points within S_0 . By doing so, we can densely pack grid points in the neighbor of the concentration region, i.e., S_0 , while maintaining a sparser grid elsewhere. Note that the choice of the mapping function μ

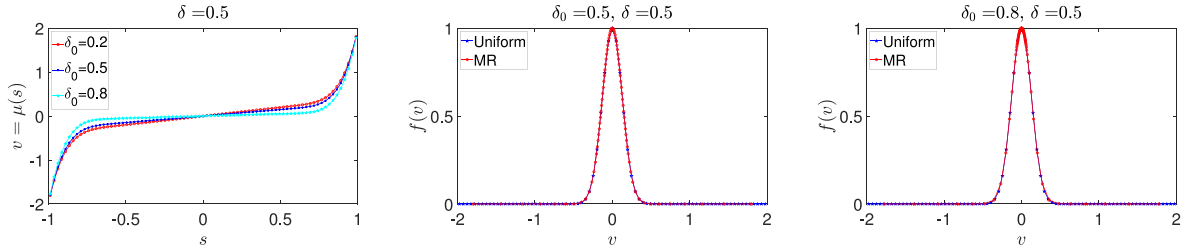


Fig. 1. Demonstration of Mesh refinement method for $f(v) = e^{-30v^2}$ for $v \in [-2, 2]$. The left figure shows the mapping function $\mu(s)$ with fixed $\delta = 0.5$ and various $\delta_0 = 0.2, 0.5, 0.8$. The middle and right one compare function $f(v)$ in uniform grid and after mesh refinement (MR).

highly depends on the shape of S_0 in (2.17), and we will focus on two specific cases.

The first case is when the concentration is symmetrically centered at a single point. Without loss of generality, let us assume it is centered at $v = 0$. In this case, we only need to find $\mu(s) : [0, 1] \rightarrow [0, L]$ and subsequently perform an odd extension to extend the mapping to the entire domain $s \in [-1, 1]$. In particular, define

$$S = \{v_j : f(v_j) > \delta_0(\max_i f(v_i)), v_j > 0\}, \quad r = \min\{v_j : v_j \in S\}, \quad (2.18)$$

then we seek a continuous and monotone increasing function $v = \mu(s)$, such that

$$\mu(0) = 0, \quad \mu(\delta) = r, \quad \mu(1) = L.$$

To this end, we use a straight line to connect $s = 0$ to $s = \delta$, and a concave up curve to connect $s = \delta$ to $s = 1$ if $\frac{r}{\delta} < L$; or a concave down curve otherwise. More specifically, μ takes the form

$$\mu(s) = \begin{cases} \frac{r}{\delta}s & s \in [0, \delta]; \\ a_1 s^5 + b_1 s & s \in [\delta, 1], \quad \text{if } r/\delta < L; \\ \frac{a_2}{1+e^{-b_2 s}} & s \in [\delta, 1], \quad \text{if } r/\delta \geq L, \end{cases} \quad (2.19)$$

where a_i and b_i , $i = 1, 2$ are two constants depending on r and δ to ensure continuity of $\mu(s)$ at the interface $s = \delta$.

In the second scenario, the concentrations are symmetrically positioned at two points. Once more, without loss of generality, we assume that the centers of these two concentrations are at $v = 0$. Similar to the previous case, we focus on the mapping in the positive part of the domain and perform an odd extension to cover the entire domain. Define

$$S = \{v_j : f(v_j) > \delta_0(\max_i f(v_i)), v_j > 0\}, \quad r_1 = \min\{v_j : v_j \in S\},$$

$$r_2 = \max\{v_j : v_j \in S\}.$$

Then we aim to find a continuous and monotone increasing function $\mu(s)$ such that

$$\mu(0) = 0, \quad \mu\left(\frac{1}{2} - \frac{\delta}{2}\right) = r_1, \quad \mu\left(\frac{1}{2} + \frac{\delta}{2}\right) = r_2, \quad \mu(1) = L. \quad (2.20)$$

As with the previous case, we use a straight line to connect $s = \frac{1}{2} - \frac{\delta}{2}$ to $s = \frac{1}{2} + \frac{\delta}{2}$, and use either $\mu(s) = as^5 + bs$ or $\mu(s) = \frac{a}{1+e^{-bs}}$ to connect $s = 0$ to $s = \frac{1}{2} - \frac{\delta}{2}$, and $s = \frac{1}{2} + \frac{\delta}{2}$ to $s = 1$, depending on the concavity. To be more precise, the expression for $\mu(s)$ takes the form:

$$\mu(s) = \begin{cases} \frac{r}{\delta}s & s \in [\frac{1}{2} - \frac{\delta}{2}, \frac{1}{2} + \frac{\delta}{2}], \\ a_1 s^5 + b_1 s & s \in [0, \frac{1}{2} - \frac{\delta}{2}] \quad \text{if } \frac{r_2 - r_1}{\delta} \geq L, \\ \frac{a_2}{1+e^{-b_2 s}} & s \in [0, \frac{1}{2} - \frac{\delta}{2}] \quad \text{if } \frac{r_2 - r_1}{\delta} < L, \\ a_3 s^5 + b_3 s & s \in [\frac{1}{2} - \frac{\delta}{2}, 1] \quad \text{if } \frac{r_2 - r_1}{\delta} < L, \\ \frac{a_4}{1+e^{-b_4 s}} & s \in [\frac{1}{2} + \frac{\delta}{2}, 1] \quad \text{if } \frac{r_2 - r_1}{\delta} \geq L, \end{cases} \quad (2.21)$$

where a_i and b_i , $i = 1, 2, 3, 4$ are constants depending on r_1 , r_2 and δ to ensure continuity of $\mu(s)$ at the interfaces.

To visualize the mapping, we examine a simple function $f(v) = e^{-20v^2}$ with $v \in [-2, 2]$. We plot this function using both a uniform mesh

and a non-uniform mesh. As illustrated in Fig. 1, the mesh refinement approach results in higher resolution near the function's peak.

Similarly, for the second case, we plot the function $f(v) = e^{-50(v-2)^2 - 50(v+2)^2}$ for $v \in [-4, 4]$ in Fig. 2. Once more, the mesh refinement approach provides higher resolution near the function's peaks.

In practice, given a mesh v_j^n and corresponding function values $f^n(v_j^n)$, we can construct an appropriate mapping $\mu^n(s)$ as described above to generate a new set of mesh grid v_j^{n+1} with higher resolution in the desired region. Note that our mesh refinement method differs somewhat from the approach proposed in [28], where the author employed a Gaussian function as the mapping function, along with three adjustable parameters. Our proposed mesh refinement technique is more general and comprises an arbitrary number of components. This enables easy manipulation of the mapping's shape, including its monotonicity and concavity. Our approach is not only more flexible but also capable of handling multiple blow-ups concurrently.

When interpolation is applied, a rescaling is imposed to ensure exact mass conservation. Specifically, after obtaining $f(v_j^{n+1})$, the following procedure is carried out:

$$f(v_j^{n+1}) := \max\{f(v_j^{n+1}), 10^{-10}\} \frac{M_n}{M_{n+1}}, \quad (2.22)$$

where

$$\begin{aligned} M_n &= \sum_{j=2}^{N_v-1} f(v_j^n) \frac{v_{j+1}^n - v_{j-1}^n}{2} + f(v_j^n) \left(L_v + \frac{v_2^n + v_1^n}{2} \right) \\ &\quad + f(v_{N_v}^n) \left(L_v - \frac{v_{N_v}^n + v_{N_v-1}^n}{2} \right), \\ M_{n+1} &= \sum_{j=2}^{N_v-1} f(v_j^{n+1}) \frac{v_{j+1}^{n+1} - v_{j-1}^{n+1}}{2} + f(v_j^{n+1}) \left(L_v + \frac{v_2^{n+1} + v_1^{n+1}}{2} \right) \\ &\quad + f(v_{N_v}^{n+1}) \left(L_v - \frac{v_{N_v}^{n+1} + v_{N_v-1}^{n+1}}{2} \right). \end{aligned}$$

2.4.4. Mesh refinement JKO scheme

Combining both the mesh refinement mechanism and the nonuniform JKO scheme results in a mesh refinement JKO scheme. First, we outline the following algorithm for one step of the JKO scheme:

For the evolution problem, we execute Algorithm 2 with two time stepping strategies. One is to use fixed time step and the other one is the adaptive time stepping. The latter choice is more expensive but useful in predicting the more accurate finite blow-up time. The algorithm is summarized in Algorithm 3.

Remark 1. In Algorithm 3, there are two indicators for blow-up. One is when the mesh size Δv_{\min} falls below a threshold ϵ_v , and the other is when optimization fails to converge within the maximum number of iterations. The latter condition is particularly valuable in the finite time blow-up case when a more accurate prediction of the blow-up time is desired. In such instances, as time approaches the blow-up time, a large time step may easily surpass the blow-up time, causing the optimization

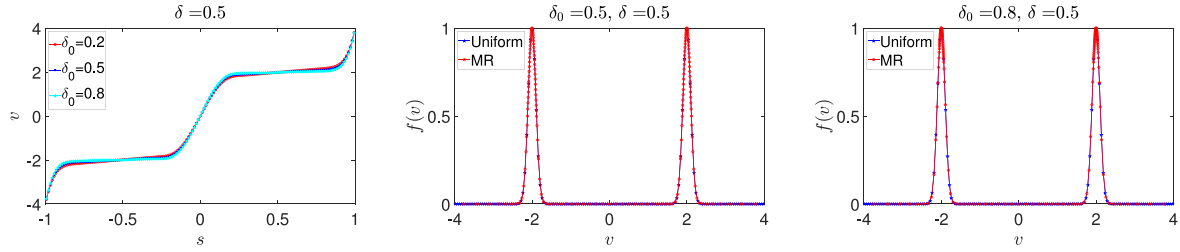


Fig. 2. Demonstration of Mesh refinement method for $f(v) = e^{-50(v-2)^2-50(v+2)^2}$ for $v \in [-4, 4]$. The left figure shows the mapping function $\mu(s)$ with fixed $\delta = 0.5$ and various $\delta_0 = 0.2, 0.5, 0.8$. The middle and right one compare function $f(v)$ in uniform grid and after mesh refinement(MR).

Algorithm 2: One step mesh refinement JKO scheme

Input : Parameters $\delta_0, \delta \in (0, 1)$ and a uniform mesh $\{s_j\}_j$, grid points $\{v_j^n\}_j$, function value $\{f^n(v_j^n)\}_j$, time step size Δt and threshold value ϵ_v .

Output: updated grid $\{v_j^{n+1}\}$, function value $f^{n+1}(v_j^{n+1})$ and blow-up indicator e_f .

1. Acquire $f^{n+1}(v_j^n)$ and e_f from Algorithm 1.
2. Determine r according to (2.18) and construct μ via (2.19); or determine r_1 and r_2 according to (2.20) and construction μ via (2.21). Update mesh by setting $v_j^{n+1} = \mu(s_j)$.
3. Compute $f^{n+1}(v_j^{n+1})$ using the interpolation and rescaling (2.22).

Algorithm 3: Mesh refinement JKO scheme with/without adaptive time stepping

Input : Parameters $\delta_0, \delta \in (0, 1)$ and a uniform mesh $\{s_j\}_j$, initial grid points $\{v_j^0\}$, initial condition $\{f^0(v_j^0)\}_j$, initial time step Δt_0 , final time T and threshold values ϵ_t and ϵ_v .

Output: Numerical solution of equation (2.1) with high resolution at singularity and numerical blow-up time T_b .

1. Let $t = 0, n = 0, \Delta t = \Delta t_0, \Delta v_{\min} = \min_j \{v_{j+1}^0 - v_j^0\}, e_f = 0$;
2. **while** $t < T$ and $\begin{cases} \Delta t > \epsilon_t & \text{with adapt time step,} \\ \min \Delta v > \epsilon_v & \text{with fixed time step} \end{cases}$ **do**
3. 1. Run Algorithm 2 to get $\{v_j^{n+1}\}$ and $\{f^{n+1}(v_j^{n+1})\}$ and update e_f ;
4. 2. If using adapt time step
5. **while** $e_f = 1$ **do**
6. | $\Delta t = \Delta t/2$;
7. | Run Algorithm 2 to get $\{v_j^{n+1}\}$ and $\{f^{n+1}(v_j^{n+1})\}$ and update e_f ;
8. **end**
9. 3. $t = t + \Delta t, n = n + 1$.
10. **end**
11. $T_b(\epsilon) = t$.

step to diverge. In response, the time step can be decreased, and this step can be recomputed iteratively, progressively reducing the time step until the optimization converges. See Fig. 5 and the discussion in Section 3.1 for more details.

2.5. Combining the two steps

We now summarize the final algorithm in Algorithm 4 for calculating the spatially inhomogeneous granular kinetic equation. Note that for the time step, we can use either a fixed time step Δt_0 or adapt the

time step based on the mesh size $\Delta t = \min\{\Delta t_0, \min_j 0.9 \frac{\Delta v_{\min}}{L_v}, \min_j 0.9 \frac{\Delta x_{\min}}{L_x}\}$. In most examples, a fixed time step will be employed unless the accurate prediction of the finite blow-up time necessitates the use of an adaptive time step.

Algorithm 4: Semi-Lagrangian JKO scheme

Input : Parameters $\delta_0, \delta \in (0, 1)$ and a uniform mesh $\{s_j\}_j$, initial time step Δt_0 , time stepping strategy, final time T , stopping threshold $\epsilon_{x,v}$, initial mesh grid $\{x_i^0\}, \{v_j^0\}$ and initial condition $f(0, x, v) = f_0(x, v)$.

Output: Numerical solution of equation (1.2) with high resolution at singularity, numerical blow-up time T_b and blow-up indicators B_x, B_v .

1. Compute $\Delta v_{\min} = \min_j \{v_{j+1}^0 - v_j^0\}, \Delta x_{\min} = \min_i \{x_{i+1}^0 - x_i^0\}$, and set $B_x = 0$ and $B_v = 0, \Delta t = \Delta t_0$;
2. **while** $t < T, \Delta v_{\min} > \epsilon_{x,v}$ and $\Delta x_{\min} > \epsilon_{x,v}, e_f = 0$ **do**
3. 1. Either use fixed time step $\Delta t = \Delta t_0$ or adjust the time step according to: $\Delta t = \min\{\Delta t_0, \min\{0.9 \frac{\Delta v_{\min}}{L_v}, 0.9 \frac{\Delta x_{\min}}{L_x}\}\}$;
4. 2. Transport step
5. **for** $j = 1 \dots N_v$ **do**
6. | Compute $f^*(x_i^n, v_j^n)$ for $i = 1 \dots N_x$ by (2.6) ;
7. **end**
8. 3. Collision step
9. **for** $i = 1 \dots N_x$ **do**
10. | Compute $f^{n+1}(x_i^n, v_j^n)$ by Algorithm 1 given $f^*(x_i^n, v_j^n)$ and update e_f ;
11. **end**
12. 4. Mesh refinement
13. For one bump case, determine
14. $S = \{(x_i, v_j) : f^*(x_i, v_j) > \delta_0(\max_{i,j} f^*(x_i, v_j)), x_i > 0, v_j > 0\}$, then $r_x = \min\{x_i : (x_i, \cdot) \in S\}$ and $r_v = \min\{v_j : (\cdot, v_j) \in S\}$, then find $\mu_x(s)$ via r_x and (2.19), find $\mu_v(s)$ via r_v and (2.19).
15. For two bumps case, determine $S = \{(x_i, v_j) : f^*(x_i, v_j) > \delta_0(\max_{i,j} f^*(x_i, v_j)), x_i > 0, v_j > 0\}$, then $r_{x1} = \min\{x_i : (x_i, \cdot) \in S\}, r_{x2} = \max\{x_i : (x_i, \cdot) \in S\}, r_{v1} = \min\{v_j : (\cdot, v_j) \in S\}$ and $r_{v2} = \max\{v_j : (\cdot, v_j) \in S\}$. Then find $\mu_x(s)$ via r_{x1}, r_{x2} and (2.21), find $\mu_v(s)$ via r_{v1}, r_{v2} and (2.21).
16. 5. Interpolation
17. Use ‘pchip’ interpolation to get function value at the new grid points $\{f^{n+1}(x_i^{n+1}, v_j^{n+1})\}$ and rescale the total mass dimensionwise according to (2.6) and (2.22) respectively;
18. 5. Compute $\Delta x_{\min} = \min_i \{x_{i+1}^{n+1} - x_i^{n+1}\}, \Delta v_{\min} = \min_j \{v_{j+1}^{n+1} - v_j^{n+1}\}$;
19. 6. $t = t + \Delta t$.
20. **end**
21. $T_b = t$; set $B_x = 1$ if $\Delta x_{\min} < \epsilon_{x,v}$; set $B_v = 1$ if $\Delta v_{\min} < \epsilon_{x,v}$ or $e_f = 1$.

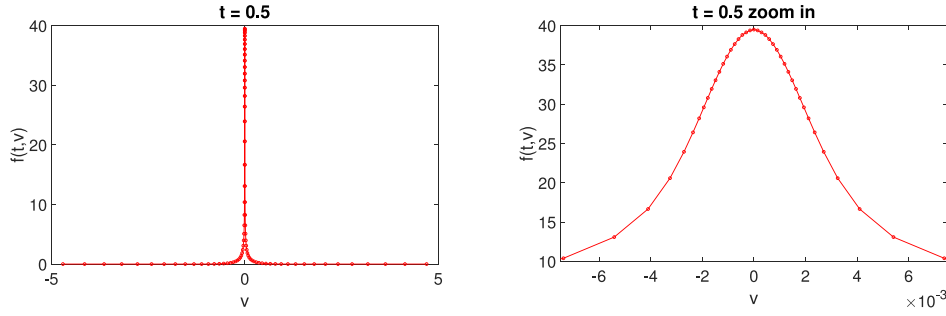


Fig. 3. Numerical solution with initial condition $g_1(v) = e^{-2v^2}$, $\delta = 0.2$, $\delta_0 = 0.5$ and fixed time step $\Delta t = 0.01$ with the kernel $W(v) = |v|$.

Table 1

Numerical blow-up time with various Δt , ϵ_v and initial conditions with the kernel $W(v) = |v|$. Here we use $N_v = 121$, $\delta = 0.5$ and $\delta_0 = 0.5$.

		$\Delta t = 0.02$	$\Delta t = 0.01$	$\Delta t = 0.005$	$\Delta t = 0.0025$	$\Delta t = 0.00125$
$f(0, v) = g_1(v)$	$\epsilon_v = 1e-3$	0.48	0.48	0.47	0.4675	0.4512
	$\epsilon_v = 1e-3/2$	0.5	0.49	0.49	0.4875	0.4725
	$\epsilon_v = 1e-3/4$	0.52	0.51	0.5	0.5	0.4875
	$\epsilon_v = 1e-3/8$	0.54	0.52	0.51	0.5075	0.4975
	$\epsilon_v = 1e-3/16$	0.58	0.53	0.515	0.5075	0.50375
$f(0, v) = g_2(v)$	$\epsilon_v = 1e-3$	0.24	0.24	0.235	0.235	0.2275
	$\epsilon_v = 1e-3/2$	0.26	0.25	0.245	0.2425	0.2375
	$\epsilon_v = 1e-3/4$	0.26	0.26	0.25	0.2475	0.245
	$\epsilon_v = 1e-3/8$	0.28	0.26	0.255	0.2525	0.25
	$\epsilon_v = 1e-3/16$	0.28	0.26	0.26	0.255	0.2525
$f(0, v) = g_3(v)$	$\epsilon_v = 1e-3$	0.12	0.12	0.12	0.1175	0.1137
	$\epsilon_v = 1e-3/2$	0.12	0.13	0.125	0.1225	0.1187
	$\epsilon_v = 1e-3/4$	0.14	0.13	0.125	0.125	0.1225
	$\epsilon_v = 1e-3/8$	0.14	0.13	0.13	0.1275	0.1265
	$\epsilon_v = 1e-3/16$	0.14	0.13	0.13	0.1275	0.12625

Table 2

Numerical blow-up time with various Δt and ϵ_v for initial condition $g_4(v)$ with $N_v = 121$, $\delta = 0.5$ and $\delta_0 = 0.5$.

	$\Delta t = 0.01$	$\Delta t = 0.005$	$\Delta t = 0.0025$	$\Delta t = 0.00125$
$\epsilon_v = 1e-3$	0.49	0.48	0.4775	0.4775
$\epsilon_v = 1e-3/2$	0.51	0.505	0.50125	0.50125
$\epsilon_v = 1e-3/4$	0.53	0.52	0.52	0.51875
$\epsilon_v = 1e-3/8$	0.56	0.54	0.5325	0.52875

3. Numerical examples

This section is dedicated to presenting the numerical findings regarding the (non)blow-up behavior of the granular kinetic equation. Specifically, there are several solution behaviors:

- (1) no blow-up occurs;
- (2) blow-up initially forms in the spatial x direction;
- (3) blow-up initially forms in the velocity v direction;
- (4) blow-up occurs simultaneously in both the spatial and velocity directions.

The last case is theoretically possible but exceptionally rare and challenging to construct. Therefore, our focus will be solely on the first three cases.

We first validate the capability of our proposed numerical solver (Algorithm 3) for the spatially homogeneous case, where analytical results are well-established. We then disclose the numerical outcomes derived from the implementation of our Algorithm 4. Based on these results, we formulate a conjecture regarding the blowup behavior for the spatially inhomogeneous granular kinetic equation.

3.1. Finite time blow-up verification: homogeneous problem

We begin by examining the spatially homogeneous case and establishing an appropriate blow-up criterion ($\epsilon_{x,v}$ in Algorithm 4). To

this end, we consider (1.2) with the kernel $W(v) = |v|$ and an initial condition $f(0, v) = g(v)$. The analytical blow-up time for this scenario is given by:

$$T = \frac{1}{2 \max_v g(v)}.$$

To evaluate our numerical approach, we consider three one-bump initial conditions: $g_1(v) = e^{-2v^2}$, $g_2(v) = 2e^{-2v^2}$, and $g_3(v) = 4e^{-2v^2}$ and one two-bumps initial condition $g_4(v) = e^{-10(v-1.5)^2} + e^{-10(v+1.5)^2}$, with analytic blow-up times $T_1 = 0.5$, $T_2 = 0.25$, $T_3 = 0.125$ and $T_4 = 0.5$, respectively. To fully investigate the blow-up behavior, we employ two time-stepping strategies. One is with a fixed time step Δt . The numerical results with the initial condition g_1 at $t = 0.5$ are depicted in Fig. 3, where $\Delta t = 0.01$ and the solution $f(t = 0.5, v)$ nearly converges to a Dirac delta function centered at $v = 0$.

More detailed results are provided in the Tables 1, 2 and Fig. 4. From Tables 1, 2, we observe that as we reduce the stopping criterion (as indicated by the columns), the numerical blow-up time increases. Conversely, when we decrease the time step size, the numerical blow-up time decreases. Notably, as we reduce both the stopping criterion and the time step size simultaneously, the numerical blow-up time approaches convergence with the analytical blow-up time. In Fig. 4, we observed that the minimum Δv decreased very rapidly near the analytic numerical blow-up time. In this particular test, we have chosen to omit the exit flag as a stopping criterion, allowing the algorithm to continue running even if the minimizer is not reached, and this is why we observe oscillations at the tail of the results when the simulation time extends beyond the analytical blow-up time.

The second strategy is to employ adaptive step sizes, with the initial step size set at $\Delta t_0 = 0.01$, and subsequently halving it until the minimization process converges. As depicted in Fig. 5, one can observe that the time step size Δt decays significantly as it approaches the analytic blow-up time. Based on these observations, we believe that our proposed numerical solver is capable of reproducing the blow-up solution and accurately capturing the analytic blow-up time.

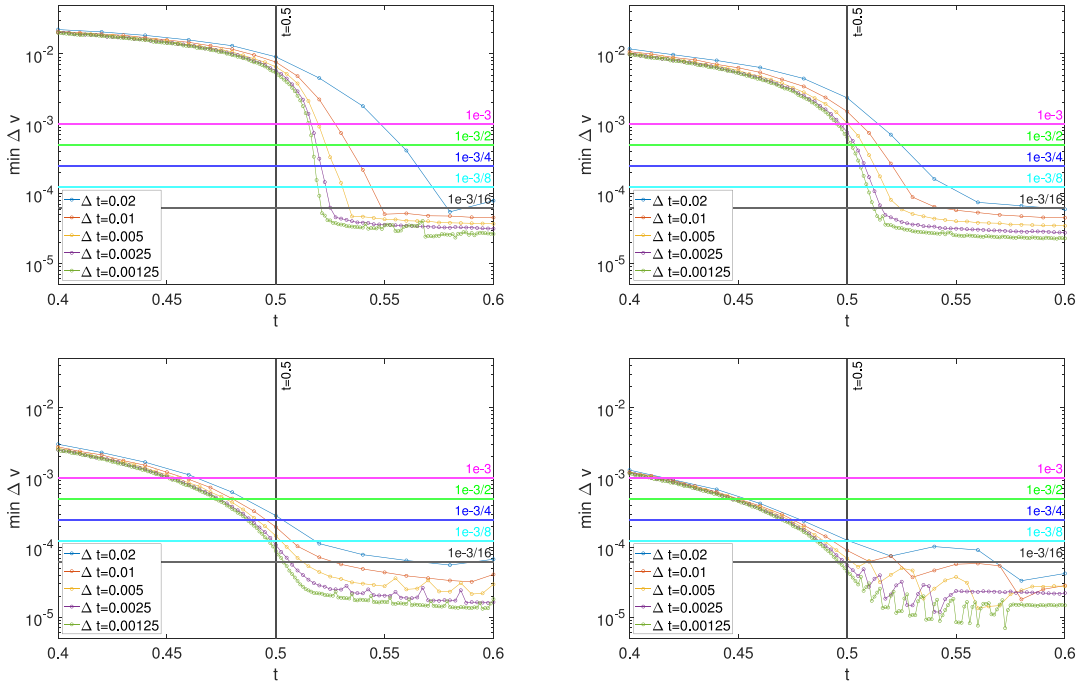


Fig. 4. Time versus minimum Δv for initial condition $g_1(v)$ with fixed Δt with the kernel $W(v) = |v|$. Here $N_v = 121$ and $\delta = 0.5$, the images from top left to bottom right correspond to $\delta_0 = 0.05, 0.1, 0.5, 0.8$ respectively.

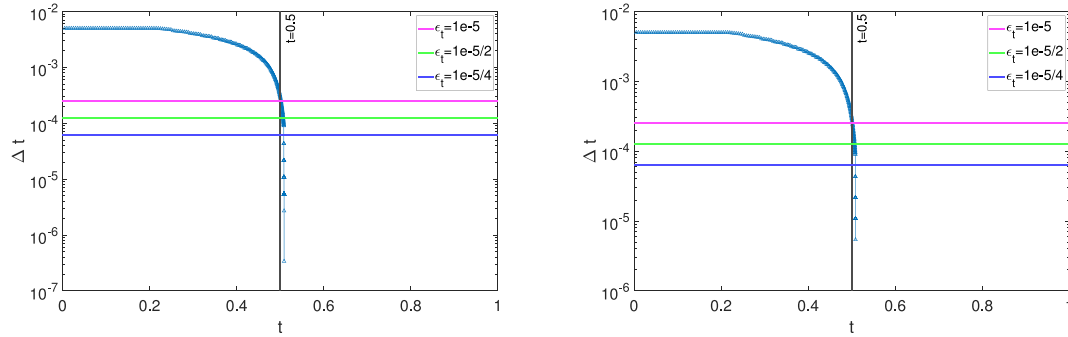


Fig. 5. Time versus adaptive time step size Δt with the kernel $W(v) = |v|$. Here, we consider the initial condition $g_1(v)$ with $N_v = 121$ and $\delta = 0.5$. On the left is $\delta_0 = 0.05$, and on the right is $\delta_0 = 0.5$.

3.2. Infinite time blow-up verification: homogeneous problem

We now consider the spatially homogeneous case of (1.2) with kernel $W(v) = |v|^3$, whose solution converges asymptotic approaches $f(t, v) = \frac{M}{2} \delta_{\frac{1}{t}} + \frac{M}{2} \delta_{-\frac{1}{t}}$ (M is the total mass) as time approaches infinity. To demonstrate that our method can capture the desired blow-up behavior, we consider an initial condition with two symmetric bumps:

$$f(0, v) = e^{-10(x-1.5)^2} + e^{-10(x+1.5)^2}. \quad (3.1)$$

Running algorithm 3 with $\delta_0 = 0.8$, $\delta = 0.5$, $\Delta t = 0.01$ and $N_v = 101$, we observed in Fig. 7 that the two initially symmetric bumps shrinking to two symmetric Dirac delta bumps and approaching $v = 0$ from two sides. Additionally, we examined the rate at which the two peaks converge to the origin, and as shown in Fig. 8, they converge with a rate of $O(\frac{1}{t})$, which is consistent with the analytical result. In contrast to the finite time blow-up cases (see Fig. 4), the minimum of Δv decreases very slowly.

Comparing Fig. 4 with Fig. 8, it is important to point out that the minimum of Δv exhibits distinct patterns. In cases with finite-time blow-up, the minimum of Δv experiences a rapid decline as time approaches or slightly surpasses the analytic blow-up time. Conversely,

in scenarios without finite-time blow-up, the minimum of Δv decreases at a considerably slower rate. This observation suggests that one can confirm the presence of finite-time blow-up by monitoring the rate of decay of the minimum of Δv , serving as an additional criterion for detecting finite-time blow-up behavior. In practice, when employing a fixed time stepping strategy as in Algorithm 3, we suggest setting the threshold for detecting blow-up in the single bump scenario at $\epsilon_v = 1e-3/16$ (as shown in Fig. 4, Table 1) and setting $\epsilon_v = 1e-3/8$ for two bumps case (see Fig. 6, Table 2). Furthermore, in situations aiming at accurately capturing the analytical blow-up time, we implement the adaptive time step strategy with a threshold set as in Algorithm 3 at $\epsilon_t = 5e-6$ (See Fig. 5).

3.3. Spatially inhomogeneous case with $\gamma = 3$

As previously discussed in Section 1 and numerically verified in Section 3.2, the spatially homogeneous case of (1.2) with kernel $W(v) = |v|^3$ exhibits an infinite-time blow-up solution. In this section, we aim to numerically investigate how the spatial dependence will affect the behavior of the solution solving (1.2).

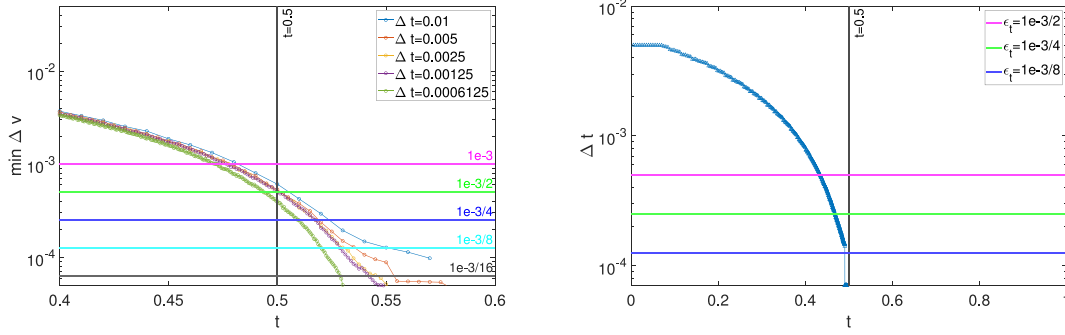


Fig. 6. Plot with initial condition $g_4(v)$ and kernel $W(v) = |v|$, $N_v = 121$, $\delta = 0.5$ and $\delta_0 = 0.5$. On the left is time versus minimum Δv and the right is time versus adaptive time step size Δt .

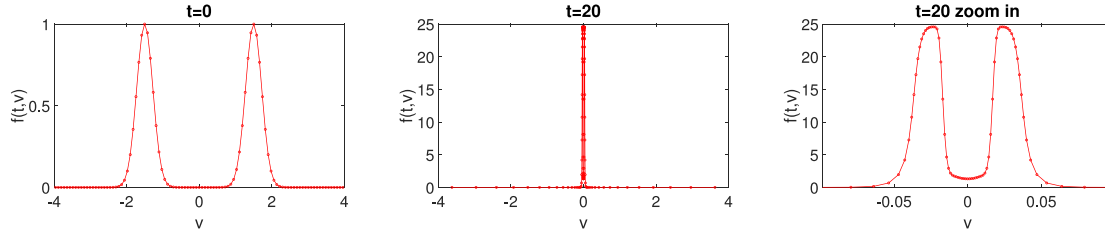


Fig. 7. Numerical solution with $\delta_0 = 0.8$, $\delta = 0.5$, $\Delta t = 0.01$ and $N_v = 101$ with the kernel $W(v) = |v|^3$. The left plots the initial condition, the middle depicts $f(20, v)$, and the right plot provides a zoom-in of $f(20, v)$ at $v = 0$.

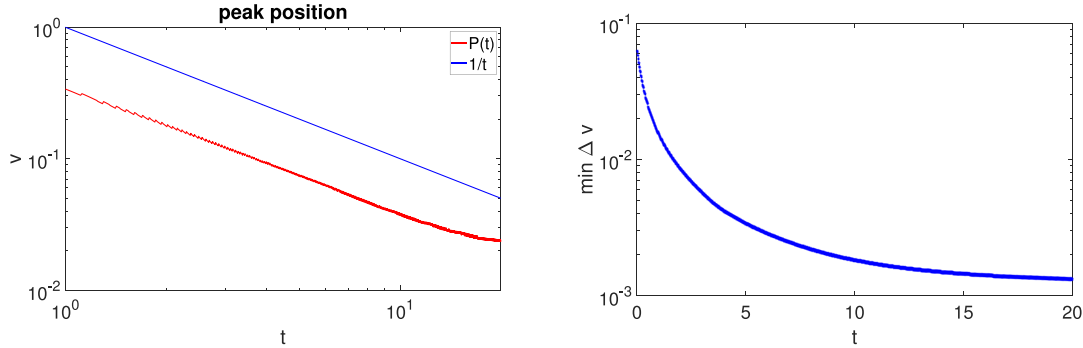


Fig. 8. Numerical solution with $\delta_0 = 0.8$, $\delta = 0.5$, $\Delta t = 0.01$ and $N_v = 101$ with the kernel $W(v) = |v|^3$. The left plot is the position of maximum value of $f(t, v)$, i.e., $P(t) = \max_v f(t, v)$ and the right plot is the minimum Δv along time.

To achieve this, we consider the following initial condition, which consists of two symmetrically located bumps with respect to the origin:

$$f_0(x, v) = e^{-a(x+c)^2} e^{-b(v-d)^2} + e^{-a(x-c)^2} e^{-b(v+d)^2}. \quad (3.2)$$

Note that when the spatial dependence is disregarded, (3.2) reduces to the two bump case discussed in Section 3.2. Consequently, the two bumps, situated with opposite spatial locations, will both move towards the origin while aggregating their velocity, as described in Section 3.2. It becomes intriguing to examine which direction exhibits blow-up first, if indeed there is a blow-up at all.

In practice, we set $L_x = 4$, $L_v = 4$, $a = b = 6$, $c = 1.5$, $d = 2$ and then execute Algorithm 4 with $\delta_0 = 0.02$, $\delta = 0.5$, and various values of the strength of inelasticity λ solving (1.2). Note that the values of δ_0 and δ in this case differ from those used in the spatially homogeneous scenario. This difference is necessary because the small value of δ_0 helps prevent oscillations caused by low resolution in regions with small values of f . A fixed time step of $\Delta t = 0.05$ will be used throughout the evolution.

Here, we present two representative numerical results for $\lambda = 4$ and $\lambda = 10$. Additional results for other values of λ (specifically, $\lambda = 2, 6, 8$)

can be found in Section 3.3. For the case of $\lambda = 4$, the evolution of the solution unfolds in three distinct stages. At first, the two initial bumps move towards the origin. Subsequently, they converge at the origin, forming a concentration. Finally, due to the smallness of the concentration coefficient, the transport effect prevails over the concentration effect, leading to a separation of the two bumps after they pass through the origin. Such an evolution is illustrated in Fig. 9. A more detailed plot at $t = 3$ is presented in Fig. 10, explicitly showing that the two bumps have bypassed each other. Most notably, the bottom right plot in Fig. 10 depicts the decreasing and then increasing changes in Δx and Δv , revealing the concentration and separation motion of the two bumps. Additionally, the peak value of $f(t, x, v)$ undergoes a gradual increase and then decrease, indicating no blow-up.

For $\lambda = 10$, the evolution of the solution also undergoes three stages. The first two stages are similar to those observed for $\lambda = 4$, but with a more pronounced concentration effect. In this case, the solution rapidly forms a concentration that resembles a Dirac delta at the origin, as depicted in Fig. 11.

While a stronger concentration effect might suggest a higher likelihood of a blow-up solution, we still observe that after an extended

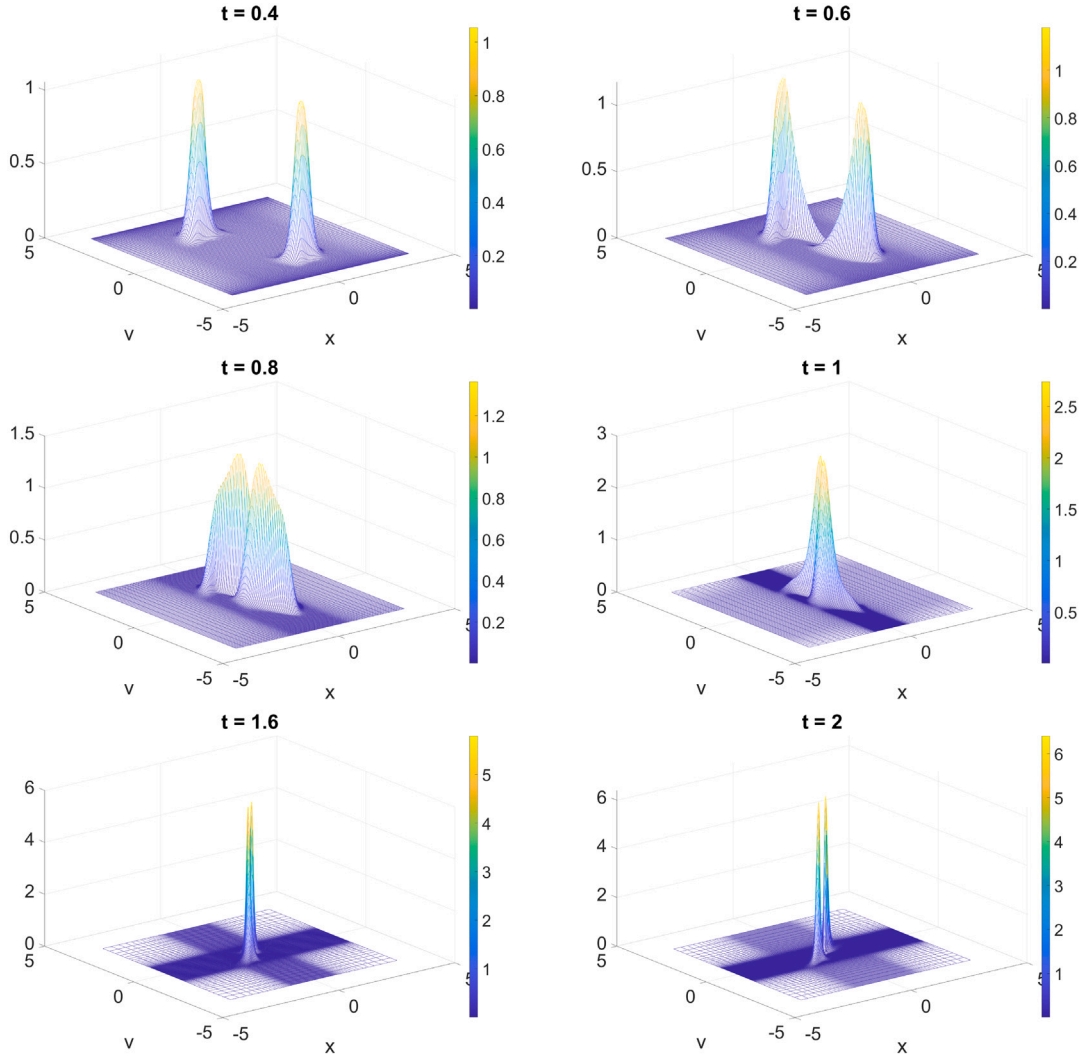


Fig. 9. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, $\lambda = 4$. From top left to bottom right are time snapshots at $t = 0.4, 0.6, 0.8, 1, 1.6, 2$.

mixing time at the origin, the solution eventually spreads out along the x direction due to the transport effect. This can be seen in the zoom-in plot and contour plot on the top right and bottom left of Fig. 12. This observation is further supported by the bottom right plot in Fig. 12, where the minimum of Δx initially decreases and then increases, implying the concentration on x at first and finally spreading out by the transport effect.

We numerically verify the persistence of properties introduced in Proposition 1 over time for $\gamma = 3$ and $\lambda = 4$ and $\lambda = 10$, with the results plotted in the first two rows of Fig. 15 for the settings of Figs. 9 and 11 respectively. In Fig. 15, the total mass and momentum are preserved. The energy (second order moment), third order moment, fourth order moment and the entropy are decreasing over time. Here the entropy is defined by $U(f) = -f \log(f)$.

3.4. Spatially inhomogeneous case with $\gamma = 3/2$

When $\gamma < 2$, a blow-up solution can emerge within finite time in spatially homogeneous cases. When considering transport, examples with finite-time blow-up in velocity space can be readily constructed. However, it is nontrivial to devise an appropriate initial condition that leads to blow-up in both the x and v simultaneously or blow-up in the x direction before the v direction. For illustrative purposes, we provide examples of finite-time blow-up in the v direction for $\gamma = 3/2$.

Typically, when two symmetric bumps initially locate at second quadrant and fourth quadrant respectively, the blow-up forms in v direction before and after two bumps meet at $x = 0$. More particularly, for initial condition (3.2) with $a = b = 6$, $c = d = 2.5$ and $\lambda = 10$, the blow-up arises on v in finite time before two bumps meet at $x = 0$, see result in Fig. 13. When $a = 6$, $b = 3$, $c = d = 2.5$ and $\lambda = 5$ are selected, the blow-up arises on v in finite time after two bumps have passed the $x = 0$, see Fig. 14.

Finally, we conduct a numerical verification to ensure that the conservation or dissipation properties outlined in Proposition 1 persist over time up to the blow-up time. Conservations of mass and momentum, dissipation of moments and entropy are all well captured. The results are collected in the last two rows of Fig. 15, with the third row corresponding to the settings in Fig. 13 and the bottom row corresponding to the settings in Fig. 14, respectively.

3.5. To blow-up or not to blow-up: Numerical conclusions

In the last two subsections, we have analyzed, based on our time adaptive algorithm, a possible scenario of blow-up in the inhomogeneous problem (1.2) with initial data (3.2). The numerical experiments firmly supports the evidence of no blow-up for $\gamma = 3$ (see Figs. 9–12) while numerical blow-up is detected for $\gamma = 3/2$ (see Figs. 13–14). Although theoretical results corroborating these numerical findings are lacking at present, we have conducted all possible theoretical checks on

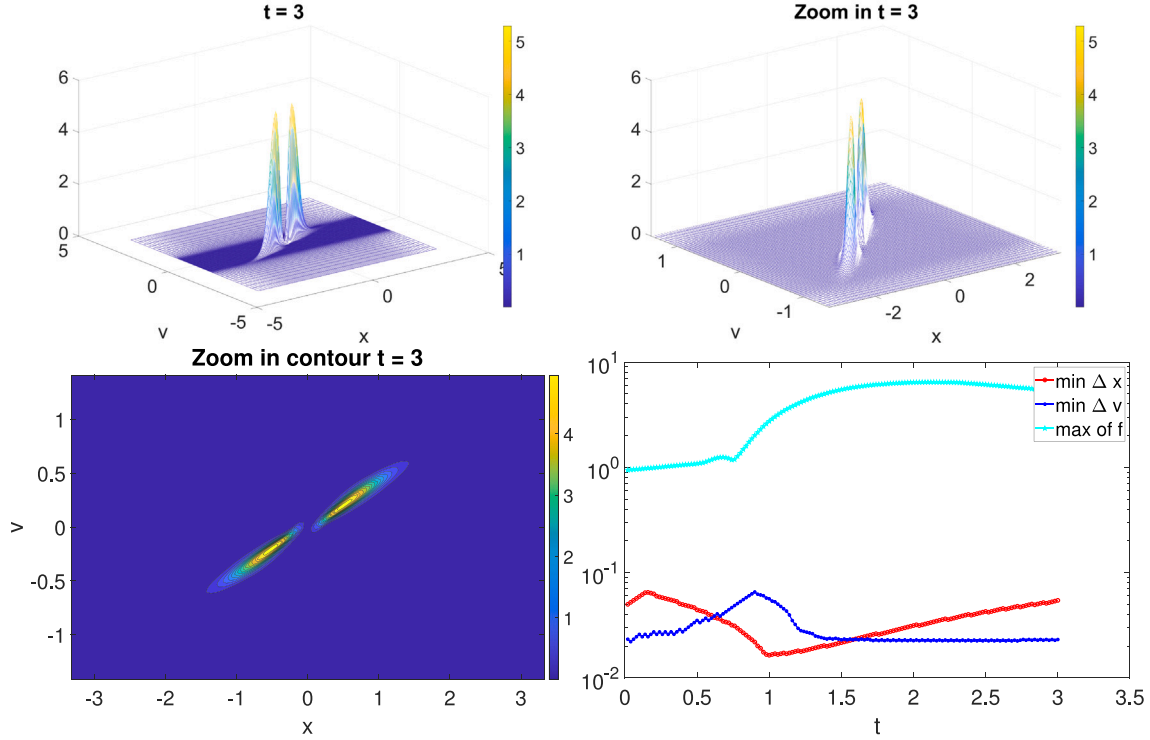


Fig. 10. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, $\lambda = 4$. Top left: numerical solution $f(t = 4, x, v)$. Top right: zoom-in plot of $f(t = 4, x, v)$. Bottom left: contour plot of $f(t = 4, x, v)$. Bottom right: record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

the numerical solutions to ensure they satisfy the correct conservations and dissipations, see Fig. 15.

In this section, we have not shown numerical results for the case of $\gamma = 2$. This case is quite particular since it is the borderline case between infinite-time blow-up and finite-time blow-up in the homogeneous case [24]. Moreover, the homogeneous equation for $\gamma = 2$ is equivalent to a linear equation when conservation of mass and momentum are taken into account for which it is easy to show that infinite time blow-up happens. Moreover, global-in-time measure solutions have been constructed for the inhomogeneous problem (1.2) with $\gamma = 2$ in [30] allowing for possible concentrations of mass. Next section is devoted to show that such concentration of mass happens for infinite mass initial data. We will provide a self-similar infinite mass blowing-up solution. Moreover, we will give heuristic arguments and provide further numerical evidence that such a finite-time blow-up of solutions happens for finite-mass initial data for $\gamma = 2$.

4. Explicit solutions for $\gamma = 2$

In this section we give some explicit solutions to (1.2) for $\gamma = 2$ and its modifications, which provide intuition about whether (1.2) has global wellposedness of L^1 solutions, see [30] for global in time measure solutions.

4.1. Infinite-mass self-similar solutions for $\gamma = 2$

We consider (1.2) with $\gamma = 2$. It can be rewritten as

$$\begin{aligned} \partial_t f + v \partial_x f &= \frac{\lambda}{2} \rho \partial_v ((v - u)f), \quad \rho(t, x) = \int_{\mathbb{R}} f(t, x, v) dv, \\ \rho(t, x) u(t, x) &= \int_{\mathbb{R}} v f(t, x, v) dv, \end{aligned} \quad (4.1)$$

using the macroscopic quantities ρ and u . We will construct a class of infinite-mass self-similar solutions for (4.1). Consider an ansatz of the form

$$f(t, x, v) = m(t) \phi(a(t)(v - b(t)x)^2), \quad (4.2)$$

where $m(t) > 0$, $a(t) > 0$, $b(t)$ are to be determined. Here ϕ is a smooth nonnegative function with $\int_{\mathbb{R}} \phi(x^2) dx = 1$. A typical choice is $\phi(x) = \frac{1}{\sqrt{2\pi}} e^{-x}$. It has density and bulk velocity

$$\rho(t, x) = \rho(t) = \frac{m(t)}{\sqrt{a(t)}}, \quad u(t, x) = b(t)x.$$

Since ρ is constant in x , (4.2) clearly has infinite total mass. See Fig. 16 for illustration.

Substituting (4.2) into (4.1) and dividing by m on both sides, we get

$$\begin{aligned} \frac{m'}{m} \phi - \left(-a'(v - bx)^2 + 2a(v - bx)b'x + 2ab(v - bx)v \right) \phi' \\ = \frac{\lambda m}{2\sqrt{a}} (\phi + 2a(v - bx)^2 \phi'), \end{aligned}$$

i.e.,

$$\begin{aligned} \left(\frac{m'}{m} - \frac{\lambda m}{2\sqrt{a}} \right) \phi - (v - bx) \\ \times \left(-a'(v - bx) + 2ab'x + 2abv + \lambda m \sqrt{a}(v - bx) \right) \phi' = 0. \end{aligned}$$

Setting the underlined parts equal to zero (collecting v terms and x terms in the second part), we get

$$\begin{cases} m' = \frac{\lambda m^2}{2\sqrt{a}} \\ a' = 2ab + \lambda m \sqrt{a} \\ a'b + 2ab' - \lambda m \sqrt{ab} = 0 \end{cases}$$

This shows that (4.2) solves (4.1) as long as $(m(t), a(t), b(t))$ solves the above ODE system.

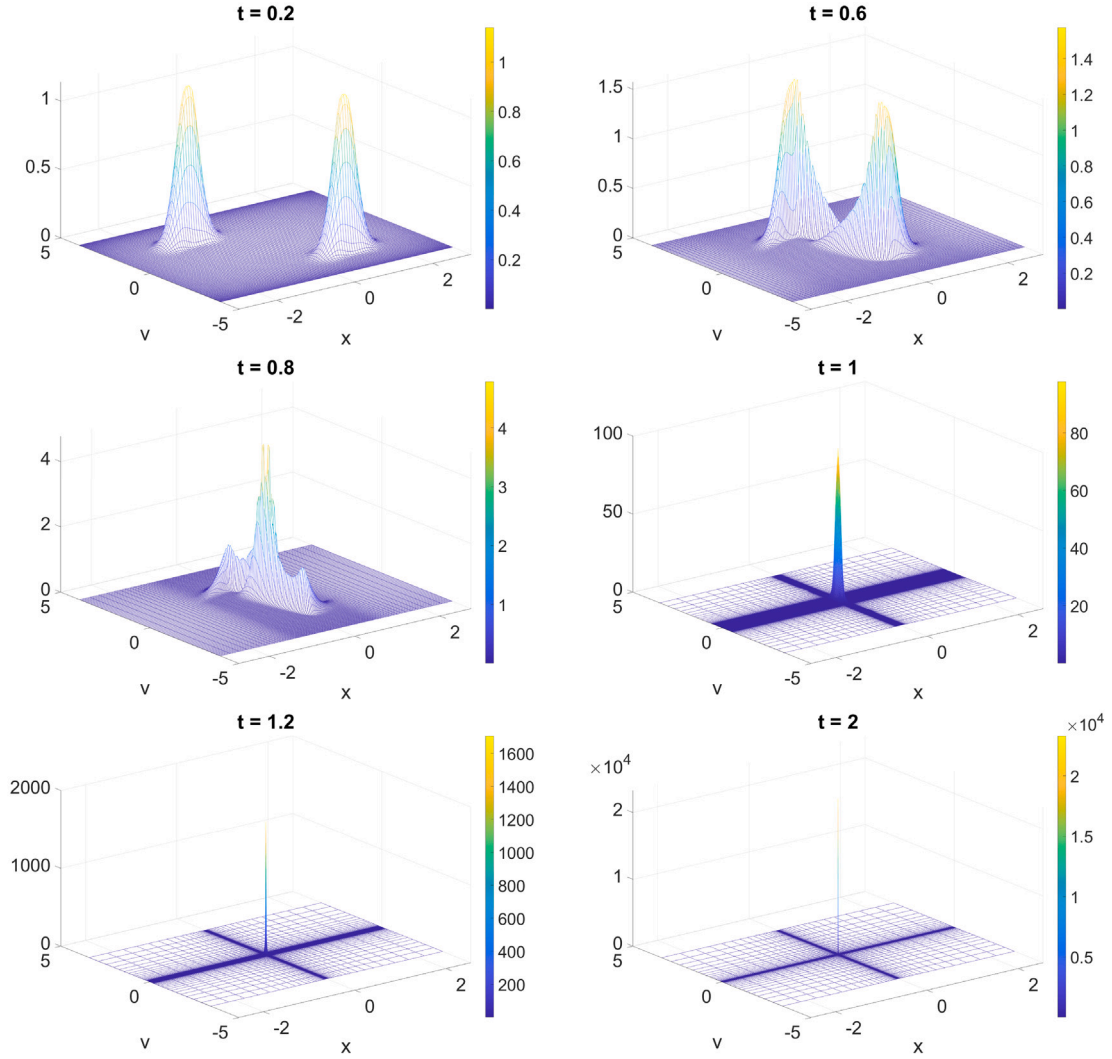


Fig. 11. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, $\lambda = 10$. From top left to bottom right are numerical solution at $t = 0.2, 0.6, 0.8, 1, 1.2, 2$.

4.1.1. Solving the ODE system

We substitute the second equation into the third and get

$$\begin{cases} m' = \frac{\lambda m^2}{2\sqrt{a}} \\ a' = 2ab + \lambda m\sqrt{a} \\ b' = -b^2 \end{cases}$$

Using the density variable $\rho = m/\sqrt{a}$, we then have

$$\begin{cases} \rho' = -b\rho \\ m' = \frac{\lambda}{2}\rho m \\ b' = -b^2 \end{cases} \quad (4.3)$$

Denote the initial condition of (ρ, m, b) as (ρ_0, m_0, b_0) with $\rho_0 > 0$, $m_0 > 0$, $b_0 < 0$. First notice that the b equation has explicit solution

$$b(t) = -\frac{1}{T-t}, \quad T = \frac{1}{-b_0} > 0,$$

which blows up to $-\infty$ at time T . Substituting into the ρ equation, we get

$$\rho(t) = \rho_0 \exp\left(-\int_0^t b(s)ds\right) = \frac{\rho_0 T}{T-t}. \quad (4.4)$$

Then, substituting into the m equation, we get

$$\begin{aligned} m(t) &= m_0 \exp\left(\frac{\lambda}{2} \int_0^t \rho(s)ds\right) \\ &= m_0 \exp\left(\frac{\lambda}{2} \rho_0 T \ln \frac{T}{T-t}\right) = m_0 \left(\frac{T}{T-t}\right)^{\frac{\lambda}{2} \rho_0 T}. \end{aligned} \quad (4.5)$$

4.1.2. Analysis of the blow-up behavior

Although the solution to (4.3) always blows up at T , we need to analyze whether the corresponding blow-up behavior of the infinite-mass solution (4.2) can be approximated by a finite-mass solution. This will be done by calculating the typical width of the velocity support for (4.2), which is

$$\frac{1}{\sqrt{a(t)}} = \frac{\rho(t)}{m(t)} = \frac{\rho_0}{m_0} \left(\frac{T}{T-t}\right)^{1-\frac{\lambda}{2}\rho_0 T}.$$

Denote $\epsilon(t) = T-t$ as the time before blow-up. Then we get the relations

$$b \sim -\epsilon^{-1}, \quad \rho \sim \epsilon^{-1}, \quad m \sim \epsilon^{-\frac{\lambda}{2}\rho_0 T}, \quad \frac{1}{\sqrt{a}} \sim \epsilon^{\frac{\lambda}{2}\rho_0 T-1}, \quad (4.6)$$

as $t \rightarrow T^-$ (where the notation \sim means that the two quantities are related by a constant multiple). Take a box $[-\epsilon, \epsilon] \times [-1, 1]$ in the (x, v) -plane (up to constant multiples of its size). Then the mean velocities

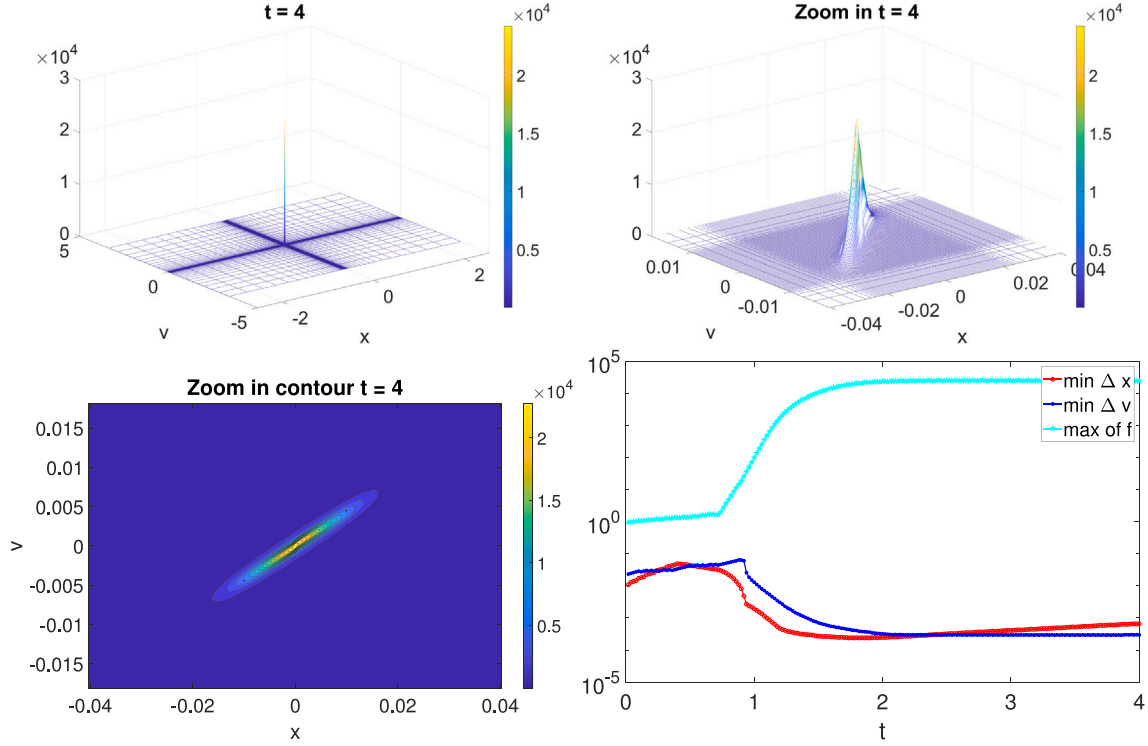


Fig. 12. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, $\lambda = 10$. Top left: numerical solution $f(t = 4, x, v)$. Top right: zoom-in plot of $f(t = 4, x, v)$. Bottom left: contour plot of $f(t = 4, x, v)$. Bottom right: record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

$(x, b(t)x)$, $x \in [-\epsilon, \epsilon]$ lies in the box, and the total mass $\iint_{x \in [-\epsilon, \epsilon]} f dv dx \sim 1$ since $\rho \sim \epsilon^{-1}$. We propose the threshold condition

$$\lambda \rho_0 T > 2. \quad (4.7)$$

- If (4.7) holds, then the width of the velocity support $\frac{1}{\sqrt{a}} \rightarrow 0$ as $t \rightarrow T^-$. This means that most mass of f with $x \in [-\epsilon, \epsilon]$ lies inside the box $[-\epsilon, \epsilon] \times [-1, 1]$ with fixed height. It is likely that such a infinite-mass solution can be approximated by a finite-mass solution, leading to finite-time blow-up for a finite-mass solution.
- If (4.7) fails with $\lambda \rho_0 T < 2$, then the velocity support $\frac{1}{\sqrt{a}} \rightarrow \infty$ as $t \rightarrow T^-$. Such an infinite-mass solution cannot be well-approximated by a finite-mass solution for t close to T because $\sup_{(x,v) \in \text{supp } f(t, \cdot, \cdot)} |v|$ for the latter cannot grow in time. Therefore we do not expect finite-time blow-up for a finite-mass solution if its initial data is close to (4.2) in this case.

4.2. A conditional blow-up result

In this subsection we give a conditional blow-up result for finite-mass solutions to (4.1) based on an analysis of the characteristic flow.

The characteristic flow $(X(t; x, v), V(t; x, v))$ of a C^1 solution to (4.1) satisfies the ODE system

$$\dot{X} = V, \quad \dot{V} = \frac{\lambda}{2} \rho(t, X)(u(t, X, V) - V); \quad X(0; x, v) = x, \quad V(0; x, v) = v.$$

For the infinite-mass solution (4.2), we have

$$\rho(t, X) = \rho(t), \quad u(t, X, V) = b(t)X,$$

and thus

$$\dot{X} = V, \quad \dot{V} = \frac{\lambda}{2} \rho(t)(b(t)X - V).$$

Along the characteristic originating from (x, v) , the quantity $\alpha = \alpha(x, v) := \sqrt{a}(V - bX) = \frac{m_0}{\rho_0}(v + \frac{x}{T})$ is conserved in time (by calculating

its derivative explicitly). Therefore we get

$$\dot{X} = bX + \frac{\alpha}{\sqrt{a}} = -\frac{1}{T-t}X + \frac{\alpha \rho_0}{m_0} \left(\frac{T}{T-t} \right)^{1-\frac{\lambda}{2}\rho_0 T}$$

and thus its explicit solution

$$\frac{X(t; x, v)}{T-t} = \frac{x}{T} + \frac{\alpha \rho_0}{m_0} \frac{1}{1-\frac{\lambda}{2}\rho_0 T} \left(\left(\frac{T}{T-t} \right)^{1-\frac{\lambda}{2}\rho_0 T} - 1 \right).$$

Now we assume the super-critical case (4.7), and consider a characteristic starting from (x, v) with $x < 0$ and $\alpha(x, v) > 0$ (i.e., $v > -\frac{x}{T}$). Then we have $1 - \frac{\lambda}{2}\rho_0 T < 0$ and thus

$$\frac{X(t; x, v)}{T-t} \leq \frac{x}{T} - (v + \frac{x}{T}) \frac{1}{1-\frac{\lambda}{2}\rho_0 T}.$$

Therefore, for the nonempty set of velocities such that

$$-\frac{x}{T} < v \leq -\frac{x}{T} \cdot \frac{\lambda}{2} \rho_0 T$$

then we have $X(t; x, v) < 0$ for any $t \in (0, T)$, see Fig. 17 Left for illustration. If the profile ϕ in the ansatz is compactly supported, then by taking x sufficiently negative it is always possible to choose $(x, v) \notin \text{supp } f(0, \cdot, \cdot)$ satisfying

$$-\frac{x}{T} < v < -\frac{x}{T} \cdot \frac{\lambda}{2} \rho_0 T. \quad (4.8)$$

Such a characteristic enables us to give a conditional finite-time blow-up result for finite-mass solutions.

Theorem 2. Assume (4.7) and denote $f(t, x, v)$ as the infinite-mass solution (4.2) with ϕ compactly supported. Let $(x_1, v_1) \notin \text{supp } f(0, \cdot, \cdot)$ satisfy $x_1 < 0$ and (4.8), and denote $(X_1(t), V_1(t))$ as its corresponding characteristic.

Let $\tilde{f}(t, x, v)$ be a C^1 solution to (4.1), whose initial data satisfies

$$\tilde{f}(0, -x, -v) = \tilde{f}(0, x, v)$$

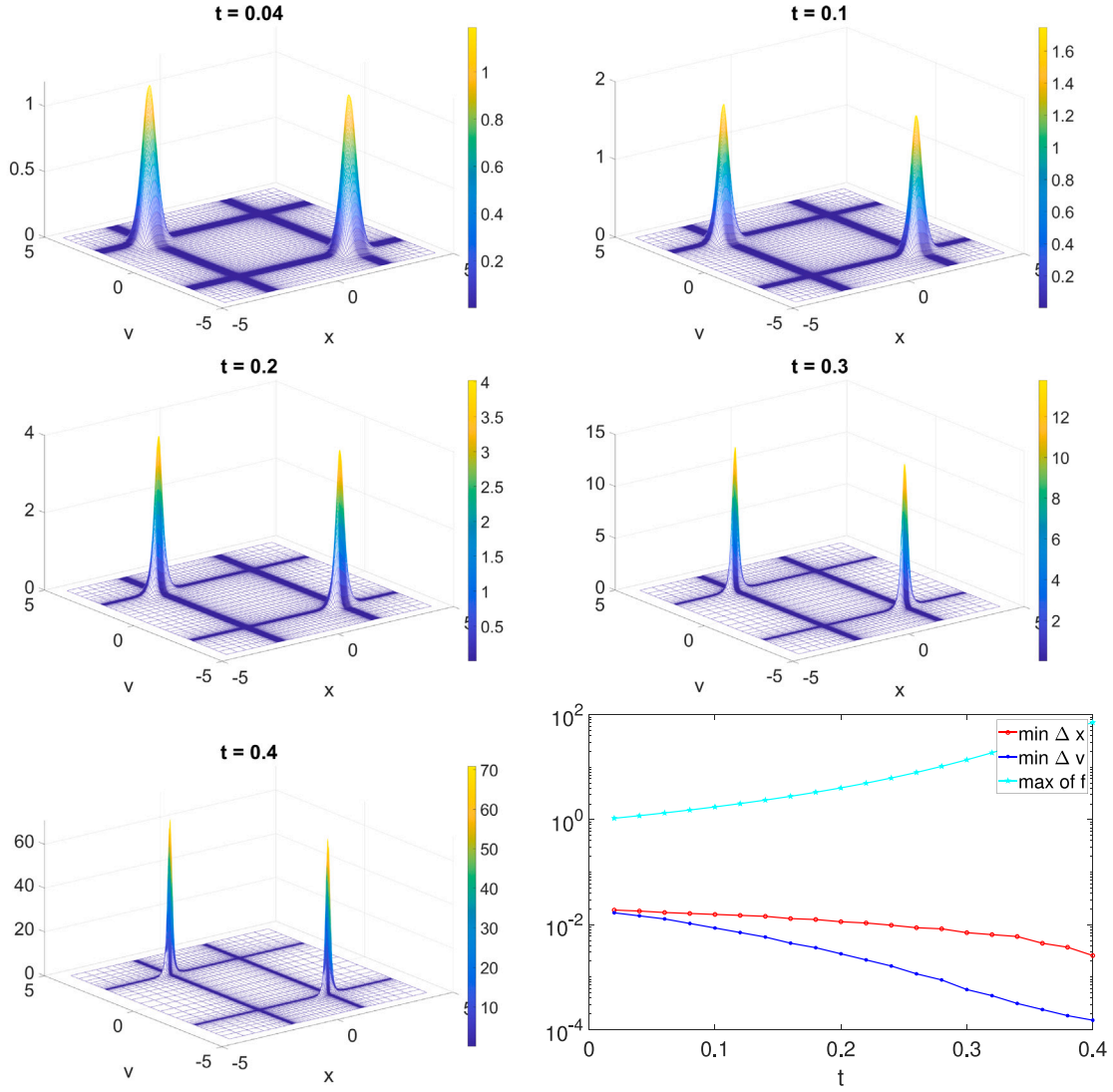


Fig. 13. Numerical solution of (1.2) with kernel $W(v) = |v|^{3/2}$ for $N_x = 121$, $N_v = 121$, $a = b = 6$, $c = d = 2.5$ and $L_x = L_v = 5$. Top left to bottom left are the numerical solution at $t = 0.04, 0.1, 0.2, 0.3, 0.4$ and the bottom right is the record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

and

$$\tilde{f}(0, x, v) = f(0, x, v), \quad \forall |x| \leq |x_1|, \forall v.$$

Denote \tilde{T} as the maximal existence time for the solution \tilde{f} .

If there holds

$$(X_1(t), v) \notin \text{supp } \tilde{f}(t, \cdot, \cdot), \quad \forall t < \max\{T, \tilde{T}\}, v \geq V_1(t), \quad (4.9)$$

then we have $\tilde{T} \geq T$, i.e., the solution \tilde{f} has a finite-time blow-up.

The extra condition (4.9) can be checked numerically as the following (see Fig. 17 Right for illustration). Suppose (4.9) breaks down at time T_* , then the proof below shows that \tilde{f} and f agrees for any $t < T_*$ and $|x| < |X_1(t)|$. As a consequence, $(X_1(T_*), V_1(T_*)) \notin \text{supp } \tilde{f}(t, \cdot, \cdot)$ and $(X_1(T_*), b(T_*)X_1(T_*)) \in \text{supp } \tilde{f}(t, \cdot, \cdot)$. Since (4.9) breaks down at time T_* , there exists some $v_* > V_1(t)$ such that $(X_1(T_*), v_*) \in \text{supp } \tilde{f}(T_*, \cdot, \cdot)$. Since $V_1(T_*)$ is between $b(T_*)X_1(T_*)$ and v_* , we see that $\text{supp } \tilde{f}(T_*, X_1(T_*), \cdot)$ has at least two connected components. Therefore, if one can check numerically that the level sets of $\tilde{f}(t, \cdot, \cdot)$ intersecting with any vertical line always have one connected component, this would indicate that (4.9) holds.

Proof. Let (X, V) denote the characteristic flow of f , and (\tilde{X}, \tilde{V}) denote the characteristic flow of \tilde{f} . We claim that for any $(x, v) \in \text{supp } \tilde{f}(0, \cdot, \cdot)$ with $x < x_1$, we have

$$\tilde{X}(t; x, v) < X_1(t), \quad \forall t < \max\{T, \tilde{T}\}.$$

Similar holds if $x > -x_1$ since \tilde{f} is odd in (x, v) .

The paragraphs before Theorem 2 show that $X_1(t) < 0$ for any $t < T$. Therefore the claim would imply that

$$\tilde{f}(t, x, v) = f(t, x, v), \quad \forall |x| \leq |X_1(t)|, \forall v$$

by tracing back the characteristics in $\text{supp } \tilde{f}(t, \cdot, \cdot)$. Therefore the finite-time blow-up of f at time T implies that of \tilde{f} at T provided that $\tilde{T} \geq T$, and the conclusion is obtained.

To prove the claim, we assume the contrary that $t_* > 0$ is the earliest time such that $\tilde{X}(t_*; x, v) = X_1(t_*)$ is reached. Therefore we necessarily have $\tilde{V}(t_*; x, v) \geq V_1(t_*)$. Then, since $(x, v) \in \text{supp } \tilde{f}(0, \cdot, \cdot)$, we may propagate to time t_* and get $(\tilde{X}(t_*; x, v), \tilde{V}(t_*; x, v)) \in \text{supp } \tilde{f}(t_*, \cdot, \cdot)$. Combining with $\tilde{X}(t_*; x, v) = X_1(t_*)$ and $\tilde{V}(t_*; x, v) \geq V_1(t_*)$, we get a contradiction with (4.9). \square

4.3. Approximation of $\gamma > 2$ solutions by $\gamma = 2$ solutions

We consider (1.2) with $\gamma > 2$, and provide a heuristic argument by approximating it with $\gamma = 2$ solutions. In the spirit of the previous subsection, we assume that the typical width of the support of f in the

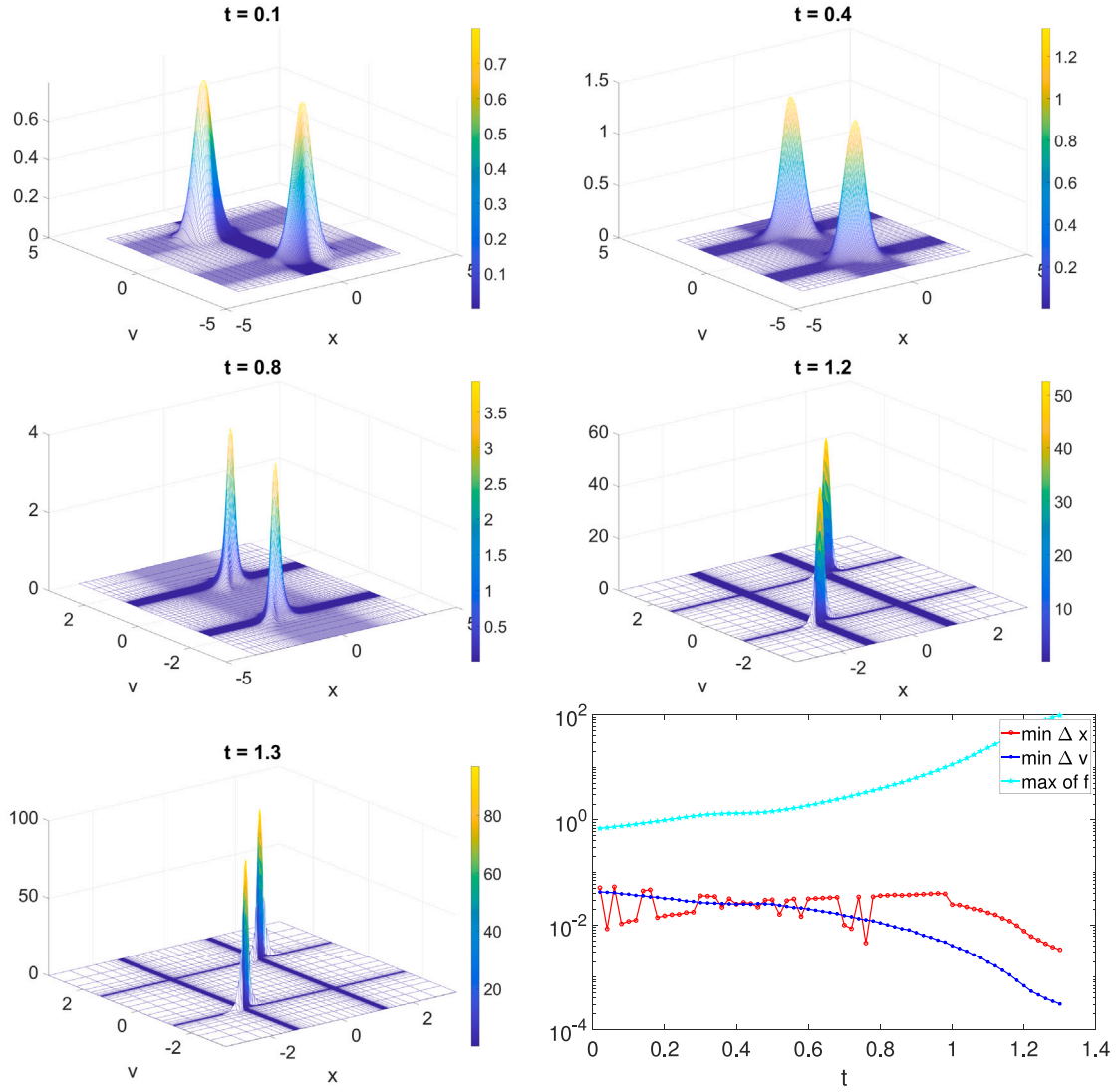


Fig. 14. Numerical solution of (1.2) with kernel $W(v) = |v|^{3/2}$ for $N_x = 121$, $N_v = 121$, $a = 6$, $b = 3$, $c = d = 2.5$ and $L_x = L_v = 5$. Top left to bottom left are the numerical solution at $t = 0.1, 0.4, 0.8, 1.2, 1.3$ and the bottom right is the record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

v -direction is $R_v > 0$ which depends on t but does not depend on x . Then we approximate (1.2) by replacing $|v - w|^{\gamma-2}$ with $R_v^{\gamma-2}$, and get

$$\partial_t f + v \partial_x f = \frac{\lambda}{2} R_v^{\gamma-2} \rho \partial_v ((v - u) f).$$

For the infinite-mass solution (4.2), we may take $R_v = \frac{1}{\sqrt{a}}$. This gives

$$\partial_t f + v \partial_x f = \frac{\lambda}{2} \left(\frac{1}{\sqrt{a}} \right)^\beta \rho \partial_v ((v - u) f), \quad \beta := \gamma - 2 > 0 \quad (4.10)$$

which is analogous to (4.1), with an extra coefficient on the RHS. Notice that (4.2) is still a solution to (4.10) if (ρ, m, b) satisfies (4.3) with $\frac{\lambda}{2}$ replaced by $\frac{\lambda}{2} \left(\frac{1}{\sqrt{a}} \right)^\beta$:

$$\begin{cases} \rho' = -b\rho \\ m' = \frac{\lambda}{2} \rho^{1+\beta} m^{1-\beta} \\ b' = -b^2 \end{cases}$$

The b, ρ solution is the same as before, given by (4.4) and (4.5). For the m solution, we integrate the m equation and get

$$(m^\beta)' \sim \frac{1}{(T-t)^{1+\beta}}, \quad m(t) \sim \frac{1}{T-t}.$$

Therefore we get the width of velocity support

$$\frac{1}{\sqrt{a}} \sim 1$$

as $t \rightarrow T^-$.

To discuss the physical meaning of the last result, (following notations of the previous subsection) we notice that the total mass in the box $[-\epsilon, \epsilon] \times [-1, 1]$ is bounded as $t \rightarrow T^-$. However, for any characteristic curve $(X(t), V(t))$ starting from $X_0 < 0$, $V_0 = b_0 X_0 + \frac{\alpha}{\sqrt{a_0}}$, $\alpha \in \mathbb{R}$, we have

$$\dot{X} = V, \quad \dot{V} = -\frac{\lambda}{2} \left(\frac{1}{\sqrt{a}} \right)^\beta \rho (V - bX),$$

and therefore $\sqrt{a}(V - bX) = \alpha$ is conserved (by calculating its time derivative explicitly). This gives

$$\dot{X} = bX + \alpha \frac{\rho}{m} = -\frac{1}{T-t} X + \alpha C,$$

for some constant $C > 0$ (since both ρ, m behave like $1/(T-t)$). Then we integrate to get

$$X(t) = \frac{T-t}{T} \left(X_0 + \alpha C T \ln \frac{T}{T-t} \right).$$

Recall that $X_0 < 0$. If $\alpha > 0$, then any characteristic $X(t)$ reaches 0 before time T . See Fig. 18 for illustration.

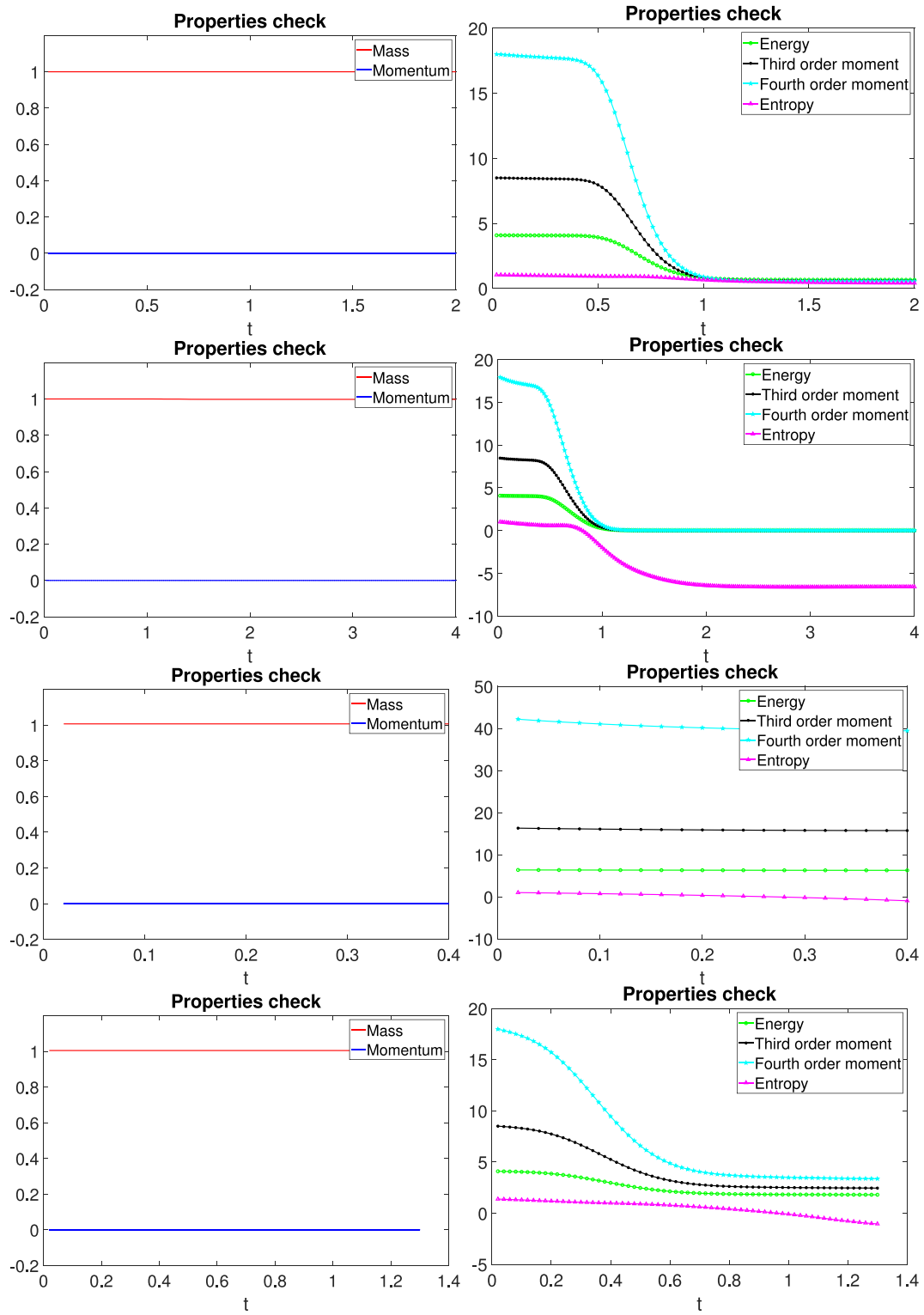


Fig. 15. Numerical conservation and dissipation properties. First and second row correspond to $\gamma = 3$ and the numerical solutions in Figs. 9 and 11 respectively. Third and fourth row correspond to $\gamma = 3/2$ and the numerical solutions in Figs. 13 and 14 respectively.

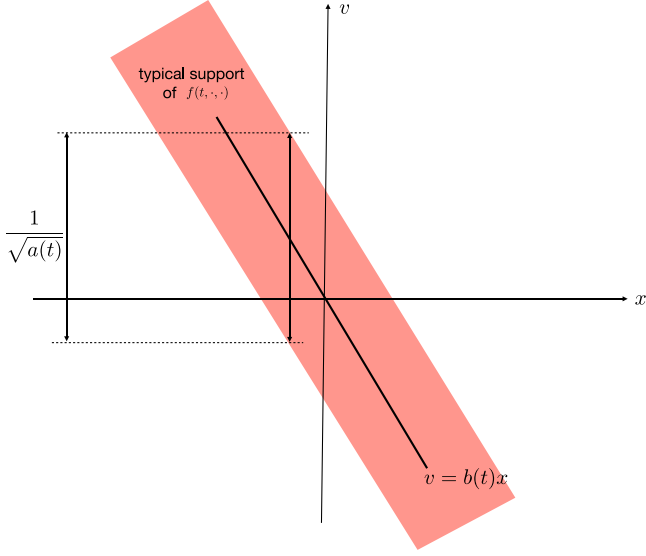


Fig. 16. Illustration of the infinite-mass solution to (4.1) or (4.10) with initial data (4.2).

Therefore, in the infinite-mass solution, for any particle arbitrarily far from 0 initially, it always contribute to the mass in $[-\epsilon, \epsilon] \times [-1, 1]$ when t is close to T . Such contribution, for particles far enough, will not be possible in a finite-mass solution. Therefore we expect that this infinite-mass solution cannot be well-approximated by a finite-mass solution, i.e., a finite-mass solution of (4.10) with a similar initial data would not have a finite-time blow-up with the same mechanism.

Next we analyze what this might imply for the global wellposedness for (1.2). Suppose we take the initial data (4.2) with ϕ supported in $[-1/2, 1/2]$. Then the solution $f(t, x, v)$ to (4.10), for fixed t, x , is always supported in $[bx - \frac{1}{2\sqrt{a}}, bx + \frac{1}{2\sqrt{a}}]$ in the v -direction. For such a function f , we always have $|v - w| \leq \frac{1}{\sqrt{a}}$ if $v \in \text{supp } f(t, x, \cdot)$. This suggests that the approximation from (1.2) to (4.10) is making the aggregation stronger. Since (4.10) is not likely to have finite-time blow-up for finite-mass solutions, we deduce heuristically that the same is true for (1.2) corroborating our numerical findings for $\gamma = 3$ in Section 3.3.

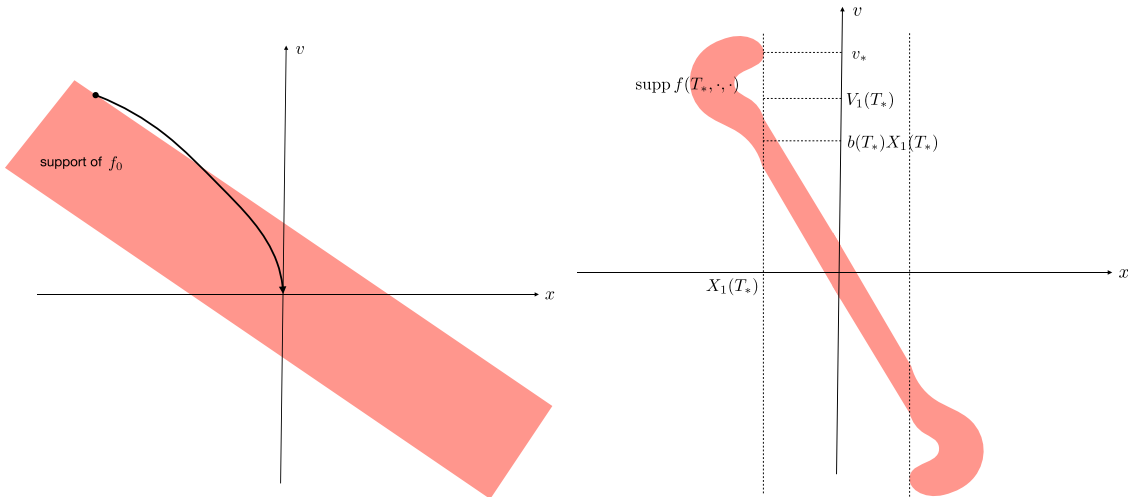


Fig. 17. Left: a particle trajectory for (4.1) which satisfies $X(t; x, v) < 0$ for any $t \in (0, T)$. Right: a possible scenario in which (4.9) breaks down.

4.4. Numerics for spatially inhomogeneous case with $\gamma = 2$

This section dedicates to provide further numerical evidence to our analytical findings for the case $\gamma = 2$. As discussed previously in Section 1, for the spatially homogeneous case, the solution is the critical case between finite-time and infinite time blow-up solution in v . As a result, it introduces the more intricate question of whether there will be blow-up in x before the blow-up occurs in v , and also whether the shear in x will accelerate the blow-up in v .

To examine the solution behavior, we will adhere to the theoretical findings in this section and validate the conjecture in Section 4.1.2. Specifically, we consider the following initial condition:

$$f_0(x, v) = \begin{cases} \frac{1}{\sqrt{2\pi}} e^{-a(bx+v)^2}, & -x_1 \leq x \leq x_1; \\ \frac{1}{\sqrt{2\pi}} e^{-a(bx+v)^2} e^{-1000(x-x_1)^2}, & \text{else,} \end{cases} \quad (4.11)$$

This initial condition is constructed to resemble the infinite mass case studied in Section 4.1. Here $\rho_0 = \int f_0(0, v) dv = \sqrt{\frac{1}{2a}}$. It is conjectured that when the relation (4.7) is satisfied, with $a = 120$ and $b = 10$, one expects a blow-up in the solution around $T = 0.1$ when $\lambda > \frac{2}{\rho_0 T} = 20\sqrt{240} \approx 310$. The blow-up is anticipated to behave similarly to the infinite mass scenario represented in (4.6). Therefore, our first task is to verify that the asymptotic behavior of the solution indeed follows (4.6). More particularly, as λ approaches the blow-up threshold of 310 (i.e., $\frac{\lambda}{2\rho_0 T} \rightarrow 1$), we expect that $\rho \sim \epsilon^{-1}$ and $m \sim \epsilon^{-\frac{\lambda}{2}\rho_0 T}$. However, the exact $\epsilon(t)$ is unknown since the analytic blow-up time is also unknown. Therefore, we numerically check if $\frac{\log \rho}{\log m} \rightarrow 1$ instead, which is verified in Fig. 19.

Next we numerically check the blow-up conditions stated in Theorem 2. Practically, this involves verifying that the level sets of the numerical solution intersect any vertical line with just one connected component, the numerical results are demonstrated in Fig. 20. Additional results with more values of λ can be found in Appendix A.2.

Furthermore, by experimenting with various aggregation constants λ , we are able to numerically validate both the blow-up criterion and the corresponding blow-up times $T = 0.1$, see results in Table 3.

In addition, we present the evolution for $\lambda = 310$ in Fig. 21. In contrast to the slow decay and the eventual increase of the minimum of Δx in Fig. 12, the minimum of Δx decreases dramatically and rapidly falls below the blow-up criterion, which provides additional evidence of blow-up on x direction.

Table 3

Numerical solution of (1.2) with kernel $W(v) = |v|^2$ for $N_x = 121$, $N_v = 201$, $\delta = 0.5$, $\delta_0 = 0.5$ and adaptive time stepping strategy with $\epsilon = 5e-6$. The table shows the numerical blow-up time T_b for various λ .

λ	220	260	280	290	300	304	306	308	310
T_b	0.10349	0.10154	0.10066	0.10047	0.10027	0.10018	0.10016	0.10012	0.10009

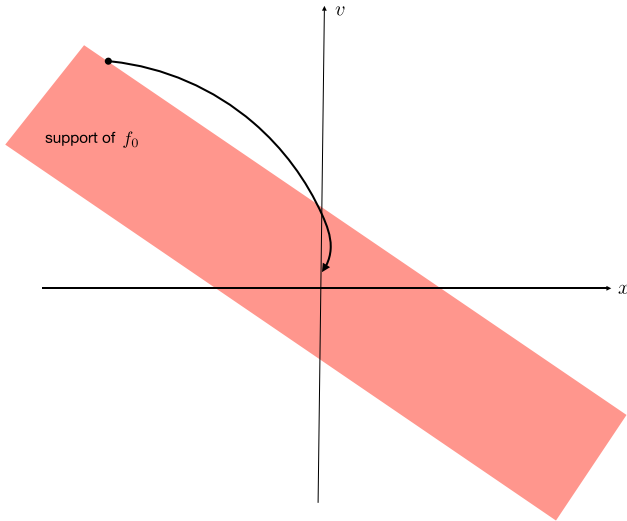


Fig. 18. A particle trajectory for (4.10) which passes $x = 0$.

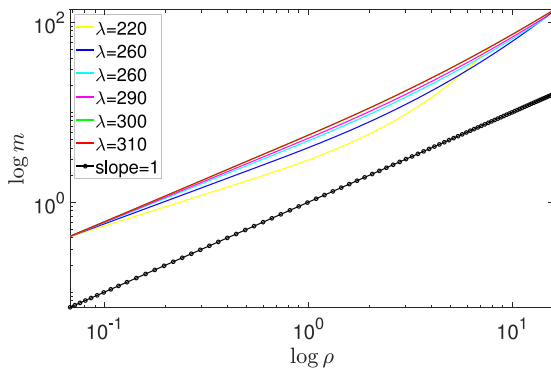


Fig. 19. Numerical solution of (1.2) with kernel $W(v) = |v|^2$ for $N_x = 121$, $N_v = 201$, $\delta = 0.5$, $\delta_0 = 0.5$ and adaptive time stepping strategy with $\epsilon = 5e-6$. The figure shows that $\frac{\log \rho}{\log m}$ approaches to 1 as $\lambda \rightarrow 310$.

5. Conclusion

In this paper, we investigate the well-posedness of the granular kinetic equation, a challenging problem characterized by the possibility of finite or infinite-time blow-up. To tackle this challenge, we propose a structure-preserving numerical scheme designed to mimic the underlying gradient flow of the associated energy functional. Consequently, our scheme upholds key physical properties, including conservation and entropy decay. Additionally, we integrate a mesh refinement mechanism that dynamically assigns a denser grid to regions of high concentration. This approach enhances resolution in concentrated regions, enabling the detection of potential blow-ups in the solution. To employ this numerical tool for investigating solution behavior, we meticulously choose a set of representative initial conditions for our experiments. Our numerical findings can be summarized in three aspects:

1. For $\gamma = 3$, despite the initial formation of a pronounced concentration in the x direction, the transport effect eventually takes over and disperses the concentration. This observation prompts us to

conjecture that for $\gamma > 2$, the solution will not exhibit finite-time blow-up in either the x or v direction.

2. In the critical case with $\gamma = 2$, through the construction of special solutions, we establish that there is always a finite-time blow-up for the infinite-mass solution. Subsequently, we perform a finite-mass approximation to the infinite-mass scenario, and numerical evidence suggests that the blow-up can occur in the x direction within a finite time.
3. For $\gamma < 2$, we presented two examples, both demonstrating finite-time blow-up in the v direction. It remains an open question whether a blow-up will also occur in the x direction. While we believe that blow-up in both x and v directions can occur simultaneously, constructing a suitable initial condition for this scenario poses a significant challenge and will be a topic for future investigation.

CRediT authorship contribution statement

José A. Carrillo: Conceptualization, Formal analysis, Funding acquisition, Investigation, Methodology, Project administration, Supervision, Writing – original draft, Writing – review & editing. **Ruiwen Shu:** Conceptualization, Formal analysis, Investigation, Writing – original draft, Writing – review & editing. **Li Wang:** Conceptualization, Formal analysis, Funding acquisition, Investigation, Methodology, Project administration, Supervision, Writing – original draft, Writing – review & editing. **Wuzhe Xu:** Data curation, Investigation, Methodology, Software, Validation, Visualization, Writing – original draft, Writing – review & editing.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix

Some Supplementary numerical examples are presented for completeness.

A.1. Supplementary results for spatially inhomogeneous cases with $\gamma = 3$

Additionally, numerical results for $\lambda = 2$, $\lambda = 6$ and $\lambda = 8$ are presented in Figs. 22, 23 and 24 respectively.

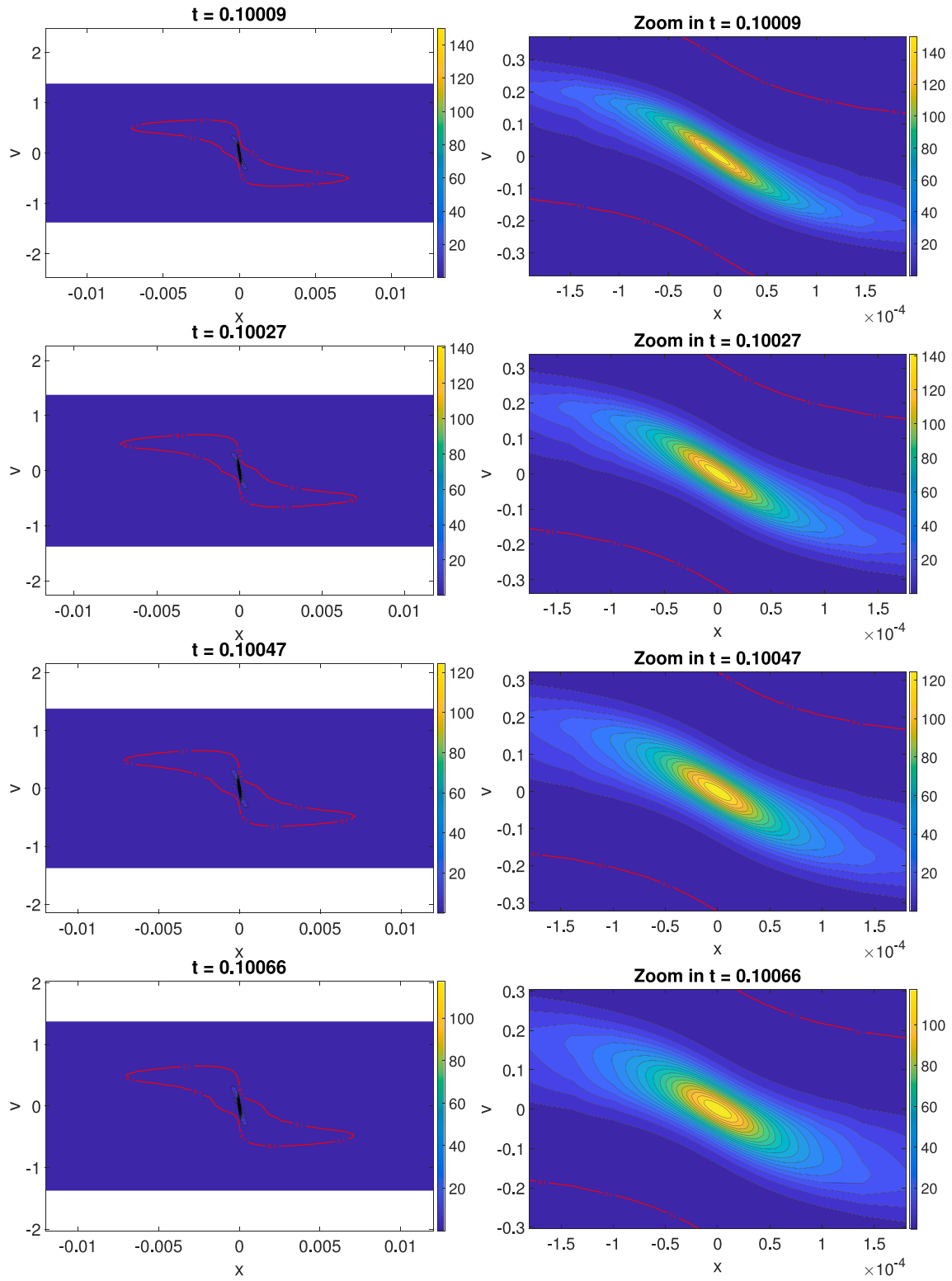


Fig. 20. Numerical solution of (1.2) with kernel $W(v) = |v|^2$ and initial condition (4.11) for $\lambda = 310, 300, 290$ and 280 with $N_x = 121$, $N_v = 201$, $\delta_0 = 0.5$ and $\delta = 0.5$ in both x and v . From top to bottom the plots show numerical solutions and their zoom-in at the numerical blowup time.

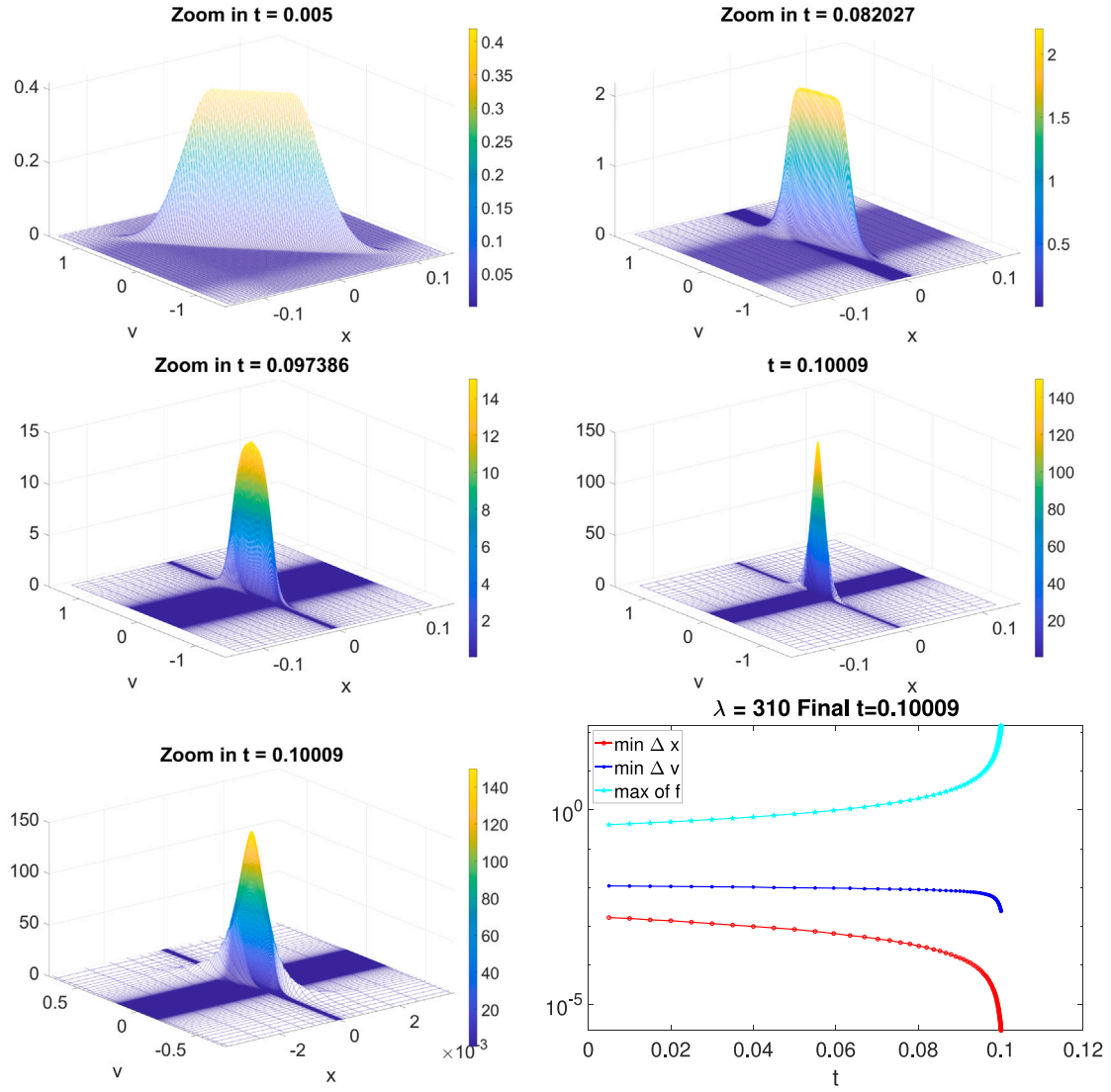


Fig. 21. Numerical solution of (1.2) with kernel $W(v) = |v|^2$ for $N_x = 121$, $N_v = 201$, $\lambda = 310$. Top two rows are time snapshots. Bottom left: the zoom-in plot at numerical blow-up time. Bottom right: record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

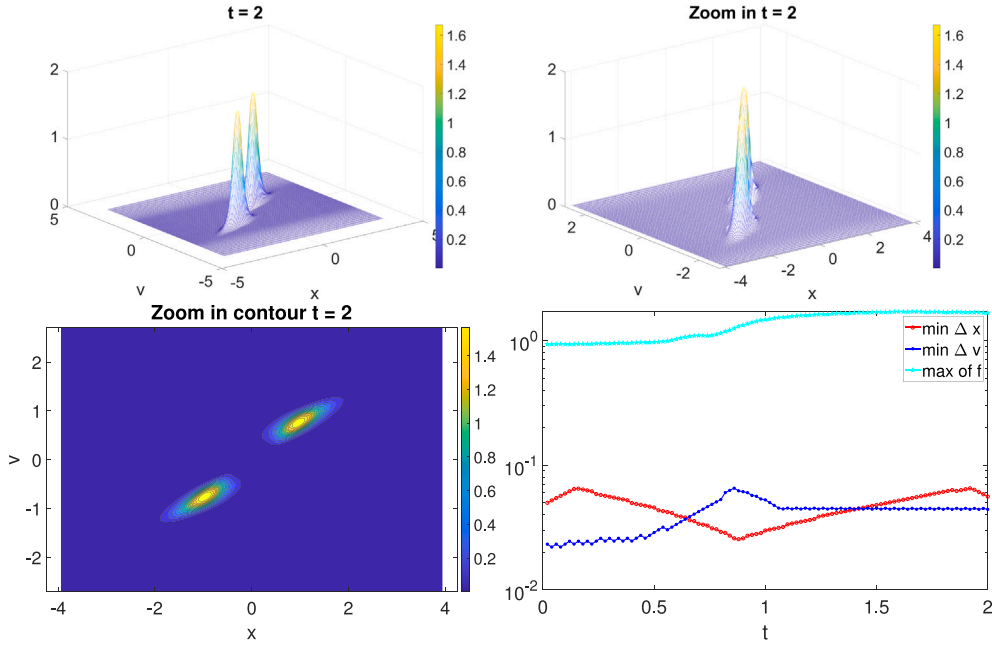


Fig. 22. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, and $\lambda = 2$. Top left is the $f(t = 4, x, v)$, top right is the zoom in of $f(t = 4, x, v)$, bottom left is the contour zoom in plot of $f(t = 4, x, v)$ and bottom right is the record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

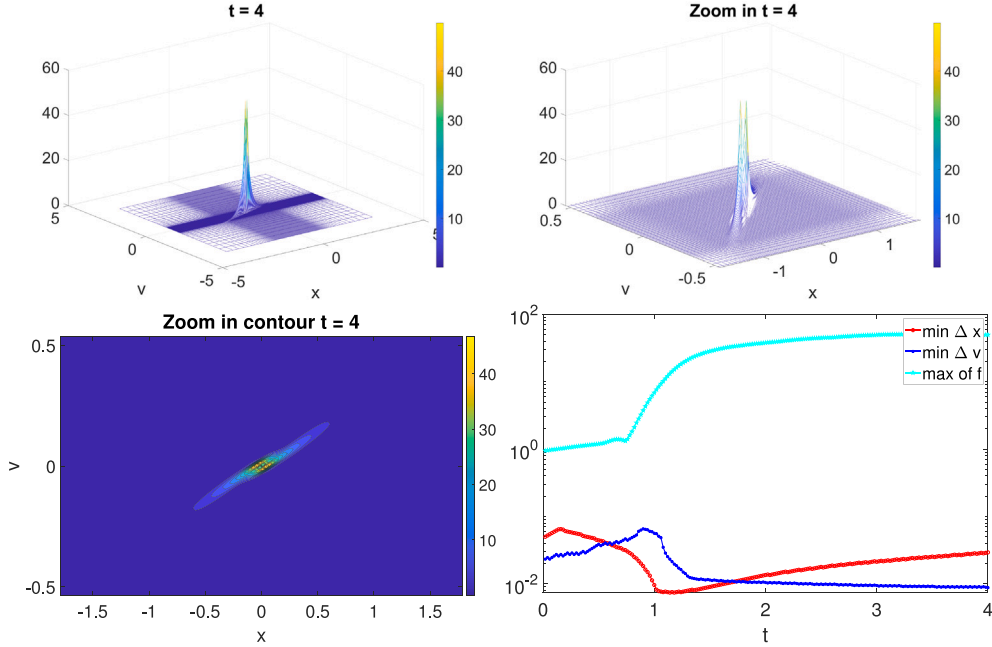


Fig. 23. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, and $\lambda = 4$. Top left is the $f(t = 4, x, v)$, top right is the zoom in of $f(t = 4, x, v)$, bottom left is the contour zoom in plot of $f(t = 4, x, v)$ and bottom right is the record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

A.2. Supplementary results for spatially inhomogeneous cases with $\gamma = 2$

In this part, we present supplementary examples to numerically check the blow-up condition described in Theorem 2. In Fig. 25, we present supplementary results for $\gamma = 2$ with initial condition (4.11)

using different values of λ , and the snapshots of evolution with $\lambda = 310$ are presented in Fig. 26. In these two figures, the level sets at $f(t, x, v) = 0.1$ intersect any vertical line with only one connected component, which verify the blow-up condition.

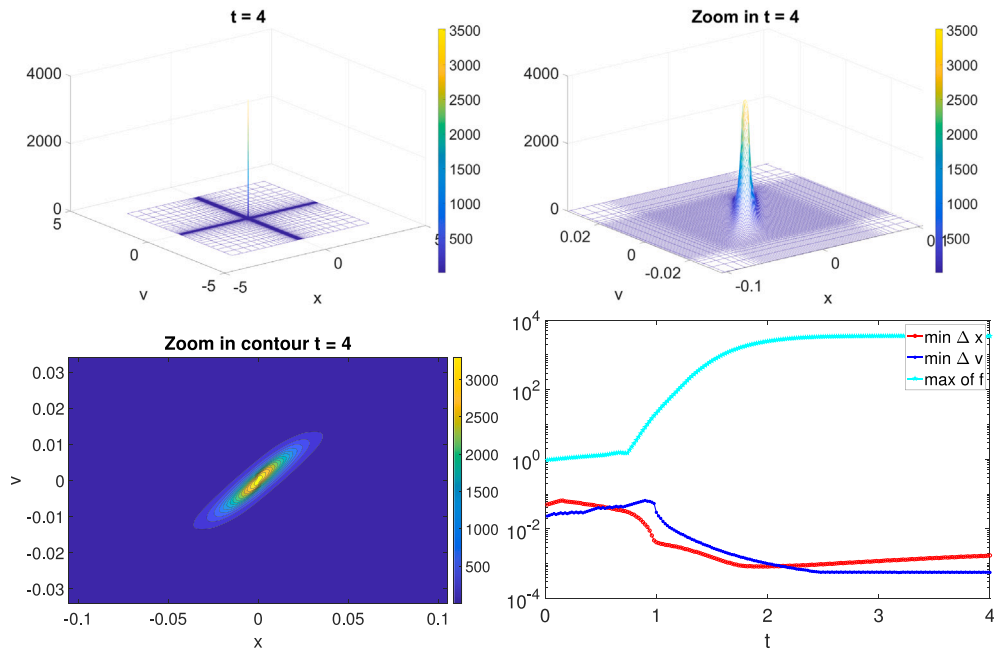


Fig. 24. Numerical solution of (1.2) with kernel $W(v) = |v|^3$ for $N_x = 121$, $N_v = 121$, and $\lambda = 6$. Top left is the $f(t = 4, x, v)$, top right is the zoom in of $f(t = 4, x, v)$, bottom left is the contour zoom in plot of $f(t = 4, x, v)$ and bottom right is the record of minimum of Δx , Δv and maximum of $f(t, x, v)$ at each time step.

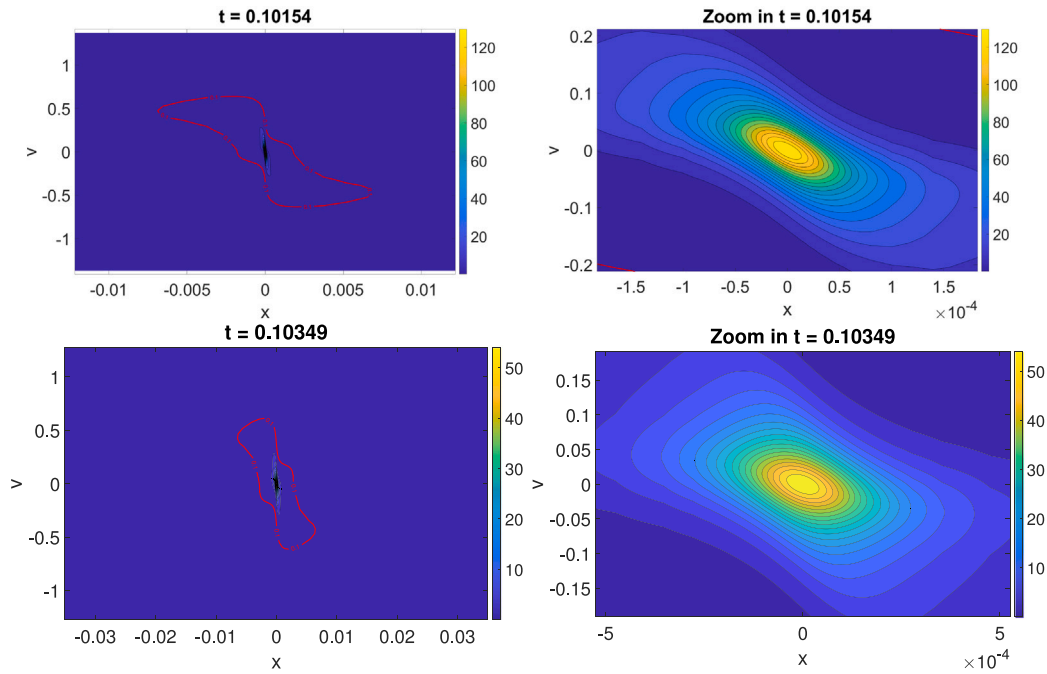


Fig. 25. From top to bottom are numerical solutions of (1.2) with kernel $W(v) = |v|^2$ and their zoom-in plots at the numerical blowup time for initial condition (4.11) with $\lambda = 260$ and 220 . Here we use $N_x = 121$, $N_v = 201$, $\delta_0 = 0.5$ and $\delta = 0.5$ for both the x and v .

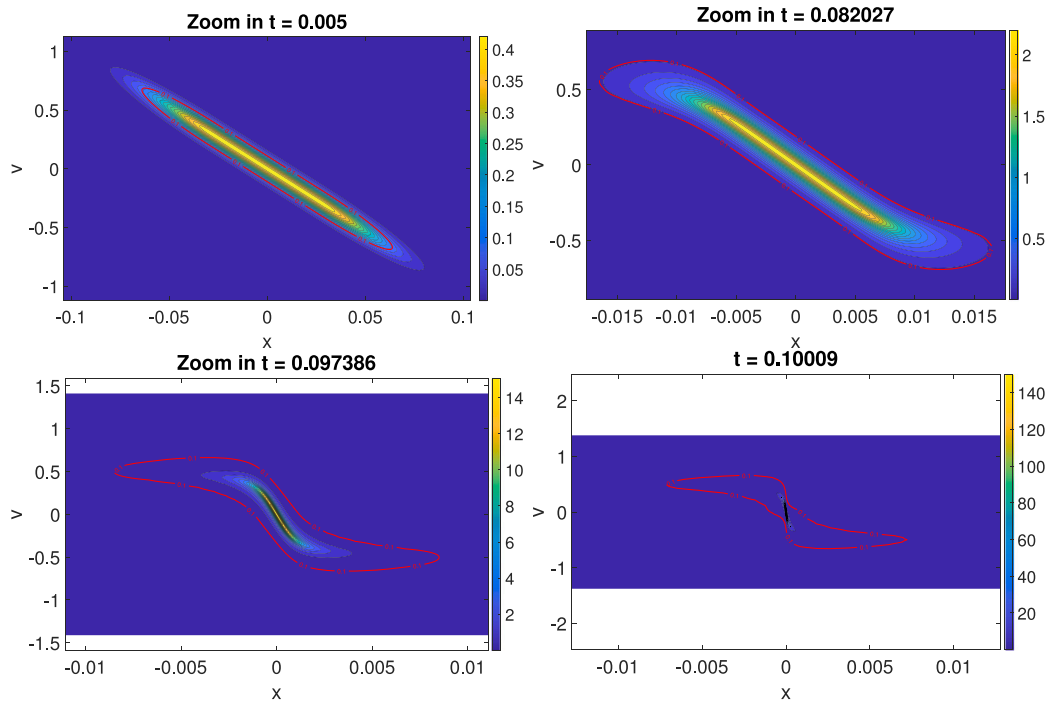


Fig. 26. Snapshots of the contour plots for the evolution of the numerical solution to (1.2) with kernel $W(v) = |v|^2$ for $N_x = 121$, $N_v = 201$, $\lambda = 310$.

Data availability

Data will be made available on request.

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