

Site Heterogeneity and Broad Surface-Binding Isotherms in Modern Catalysis: Building Intuition Beyond the Sabatier Principle

James M. Mayer

Department of Chemistry, Yale University, 225 Prospect St., New Haven CT 06520
james.mayer@yale.edu

ABSTRACT: Learning the science of heterogeneous catalysis and electrocatalysis always starts with the simple case of a flat, uniform surface with an ideal adsorbate. It has of course been recognized for a century that real catalysts are more complicated. For the increasingly complex catalysts of the 21st century, this Perspective argues that surface heterogeneity and non-ideal binding isotherms are central features, and their implications need to be incorporated in current thinking. A variety of systems are described herein where catalyst complexity leads to broad, non-Langmuirian surface isotherms for the binding of hydrogen atoms – and this occurs even for ideal, flat Pt(111) surfaces. Modern catalysis employs nanoscale materials whose surfaces have substantial step, edge, corner, impurity, and other defect sites, and they increasingly have both metallic and non-metallic elements M_nX_m , including metal oxides, chalcogenides, pnictides, carbides, doped carbons, etc. The surfaces of such catalysts are often not crystal facets of the bulk phase underneath, and they typically have a variety of potential active sites. Catalytic surfaces *in operando* are often non-stoichiometric, amorphous, dynamic, and impure, and often vary from one part of the surface to another. Understanding of the issues that arise at such nanoscale, multi-element catalysts is just beginning to emerge. Yet these catalysts are widely discussed using Brønsted/Bell-Evans-Polanyi (BEP) relations, volcano plots, Tafel slopes, the Butler-Volmer equation, and other linear free energy relations (LFERs), which all depend on the implicit assumption that the active sites are all “similar” and that surface adsorption is close to ideal. These assumptions underly the ubiquitous intuition based on the Sabatier Principle, that the fastest catalysis will occur when key intermediates have free energies of adsorption that are not too strong nor too weak. Current catalysis research often aims to minimize the complexity of non-ideal isotherms through experimental and computational design (e.g., the use of single crystal surfaces), and these studies are the foundation of the field. *In contrast, this Perspective argues that the heterogeneity of binding sites and binding energies is an inherent strength of these catalysts.* This diversity makes many nanoscale catalysts inherently a high-throughput screen wrapped in a tiny package. Only by making the heterogeneity part of the foundation of catalysis models, sorting the types of active sites and dissecting non-ideal binding isotherms, will modern catalysis learn to harness the inherent diversity of real catalysts. Controlling rather than avoiding diversity is needed to optimize complex modern catalysts and catalytic conditions.

I. Introduction

Further development of heterogenous catalysis and electrocatalysis will require understanding of the binding and reactivity of surface-bound intermediates. While commercially viable catalysts are rarely ‘designed,’ there are many examples of understanding mechanisms and structure/function relationships being the foundation for remarkable advances. This Perspective modern catalysts becoming more nanoscale and more compositionally complex prompts a re-evaluation of common intuition and approaches.

Understanding the binding and the reactivity of surface intermediates has been a central challenge in catalysis for a hundred years. With the growth of surface science, it became possible to simplify the problem by using clean single crystal metal surfaces as catalysts. This has led to great

advances, such as atomic resolution of the individual steps in the catalytic oxidation of CO on a platinum(111) surface.¹ The results from single crystal surface science can, in some situations, be directly connected to catalytic processes.³

However, even with well-ordered, effectively defect-free single-crystal metal surfaces, kinetic treatments of surface reactions often make simplifying approximations. Kinetic models frequently—though certainly not always (*cf.* ^{4,5})—use the assumption that adsorption is *ideal*, that the surface coverage of each adsorbate X (θ_X) follows the Langmuir isotherm.⁶ Ideal behavior allows rate and equilibrium expressions to simply use surface coverages, which means that binding energies and reaction barriers are well described by single free energy values (see below).

For traditional catalysts, assuming adsorbate ideality is often sufficient, especially for trends in catalysis by similar surfaces.⁶ This has led to a common intuition that non-

ideality does not have a large effect on catalysis. Boudart and Djega-Mariadassou's classic 1984 text *Kinetics of Heterogeneous Catalytic Reactions* derives this conclusion for modestly non-ideal adsorbates whose binding energies vary in a regular way (see below).⁷ Somorjai and Li's 700+ page *Introduction to Surface Science and Catalysis* mentions only the Langmuir isotherm and concludes:⁸

Although the model is physically unrealistic for describing the adsorption of gases on real surfaces, its successes ... [are] due to the relative insensitivity of macroscopic adsorption measurements to the atomic details of the adsorption process.⁸

Somorjai and Li's statement is less accurate today because of the complexity of modern catalysts. Modern catalysts are often small nanoparticles or irregular surfaces that have a large proportion of edge, corner, and defect sites. These are more reactive than sites on low-index crystal facets,⁹⁻¹¹ and can make the surface more prone to rearrangement, over time and even within a catalytic cycle.¹²⁻¹⁷ Traditional intuition has often been based on noble metal catalysts, while current research is increasingly exploring earth-abundant materials for reasons of cost, scarcity, and sustainability.¹⁸ Many catalysts of current interest are binary and ternary materials, including oxides, chalcogenides, pnictides, carbides, etc.¹⁹⁻²² The surfaces of such materials often have stoichiometries and structures that differ from the bulk, and their most catalytically active regions can be hydrous and/or amorphous.²³⁻²⁶ The surfaces of binary materials M_nX_m are often quite reactive, including to air, water, and the sometimes harsh conditions of thermal and electrochemical catalysis. These surfaces are therefore likely to have strongly-bound impurities. All of these properties lead to relatively high levels of surface diversity, in composition, structure, and electronic structure.

Today, the central intuition for heterogeneous catalysis is the century-old Sabatier Principle, that surface intermediates should not bind too strongly or too weakly (Figure 1²⁷). However, high surface diversity and non-ideality prompt a reconsideration of this intuition. Figure 1 is almost self-evident for a surface with a single, ideal site. For more complex catalysts that bind substrates at various sites, a single surface bond strength is usually not a sufficient predictor.

The importance of surface diversity and non-ideality was well-known long ago, but this topic often takes a back seat in discussions of modern catalysis. Still, there have always been studies in this area, and the last few years have seen a resurgence (cf.,²⁸⁻³⁵). Increasingly powerful analytical tools have shown experimentally that many catalysts are non-ideal and dynamic. A few representative examples could include nanoscale gold catalysts,³⁶ Fischer-Tropsch catalysts,³⁷ electrosynthesis,³⁸ colloidal chalcogenide and pnictide nanoparticles,³⁹⁻⁴⁶ H on metal phosphide catalysts,⁴⁷ and the RuO₂ and IrO₂ studies described below. Theoretical and computational studies are re-emphasizing non-ideality and heterogeneity, and the leading studies have much to teach us.^{28-31,48-52} It must be emphasized that the statements in this Perspective are not new or unique. The goal here is to raise awareness of surface heterogeneity and non-ideality and their many implications for catalysis.^{15,53-57}

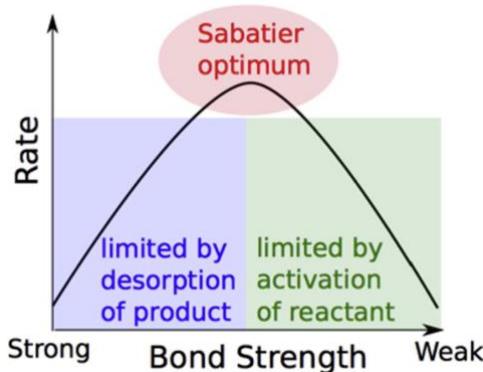


Figure 1. Illustration of the Sabatier Principle.²⁷ Reprinted from *Journal of Catalysis*, Vol 328, Andrew J. Medford, Aleksandra Vojvodic, Jens S. Hummelshøj, Johannes Voss, Frank Abild-Pedersen, Felix Studt, Thomas Bligaard, Anders Nilsson, Jens K. Nørskov, From the Sabatier principle to a predictive theory of transition-metal heterogeneous catalysis, Page 37, 2015, with permission from Elsevier.

II. Overview of this Perspective

Non-ideal behavior should not be a correction or an afterthought—this is a typical property of catalytic surfaces. Catalysts with multiple kinds of active sites can be valuable, operating under a range of conditions and binding different surface intermediates in various ways. They could potentially be viewed as an all-in-one high-throughput experiment. Single-site catalysts are exciting and attractive for many reasons, but there are challenges in knowing what single structure to aim for. Even for proteins with clearer structure-function relationships and refined computational approaches, a recent deep-learning paper identified 7,648 designs for individual experimental testing.⁵⁸ Machine learning approaches are becoming increasingly important in catalysis, using computational and/or experimental datasets.⁵⁹⁻⁶⁶ In addition to testing many different catalysts, it would be valuable to also examine and harness the inherent diversity in multi-element, nanoscale catalysts.

The focus here is on oxidation/reduction (redox) reactions, and the binding of molecular fragments such as H and OH. In the parlance of inorganic chemistry, these are X-type rather than L- or Z-type ligands.⁶⁷⁻⁶⁹ The revised intuition advocated here may be less dramatic for the binding of stable molecules such as CO. The examples below largely describe stoichiometric reactions involving surface hydrogen, the X-type adsorbate that has been most studied. Studies of individual reaction steps are a valuable approach to understanding catalytic systems, following the model of homogeneous catalysis.⁷⁰ Still, as emphasized by a reviewer, “only true probe of active site chemistry is the kinetic profile of the reaction itself” and the *ex situ* measurements described here need to be considered in that light.

The next few sections establish a common background for the Perspective, introducing various kinds of Linear Free Energy Relationships (LFERs) and surface adsorption isotherms. Connections are made between treatments for surfaces *vs.* homogeneous solutions, and between thermal *vs.* electrochemical reactivity. Systems with non-ideal adsorption are presented in Sections VI, VII, and X. Sections VIII and IX explain how substantial non-ideality

complicates the foundational LFER treatments in catalysis. The Perspective closes with ideas for understanding and harnessing surface diversity and revised LFERs for improved catalysis.

III. Historical View of Heterogeneity and Non-Ideality in Catalysis

Non-ideal surface adsorption—a diversity of surface sites and adsorbate interactions—has long been recognized as an important feature of catalysis and electrocatalysis. Mathematical forms of non-ideal isotherms were introduced in the early 20th century, and they have been extensively used.⁷¹

A 1958 paper by Parsons derived the shape of volcano plots first assuming ideal Langmuir isotherms, and immediately followed that with a discussion on non-ideality (his Section 4), the start of which is quoted here.⁷² Parsons' conclusions are discussed in Section IX.

4. Effect of Heterogeneity of Surface or of Interactions Between Adsorbed Hydrogen Atoms: Temkin Model

The feature of the above discussion which is farthest from plausibility is the use of a Langmuir adsorption isotherm since it is well known that the heats of adsorption of hydrogen on metals fall with coverage....⁷²

Conway's *Accounts of Chemical Research* article from 1981 emphasizes the multiplicity of sites for an adsorbate A with surface coverage θ_A :⁷³

... results obtained by thermal desorption in gas/solid experiments and especially in electrochemical measurements show that [the equilibrium constant for adsorption] is usually neither independent of θ_A nor a continuous function of θ_A ; it often has a series of discrete values over distinguishable small ranges of θ_A as θ_A is changed from 0 to 1. This constitutes the phenomenon of "multiple state" adsorption in monolayers.⁷³

For gas-phase hydrogen adsorption on Pt, recent measurements and calculations by García-Diézquez, Hibbitts, and Iglesia found that ΔH_{ads} decreased with increasing surface coverage, by ~30 kJ mol⁻¹ (~7 kcal mol⁻¹).⁷⁴ Pt nanoparticles (1.6 - 9.1 nm diameter) had ca. 40 kJ mol⁻¹ (~10 kcal mol⁻¹) stronger binding than a flat Pt(111) surface. This is:

because their surfaces expose a larger fraction of low-coordination atoms. The nonuniformity of Pt surfaces, repulsion among co-adsorbates, and high adsorbate mobility starkly contrast with the requirements for Langmuirian descriptions of binding and reactions at surfaces, despite the ubiquitous use and success in practice of the resulting equations in describing adsorption isotherms and reaction rates on surfaces.⁷⁴

Such 'structure sensitivity' is well known for Pt and other catalysts, and has long been recognized to contribute to site heterogeneity in metal nanoparticle surfaces.^{4,75-82}

IV. Background 1: Kinetic-Thermodynamic Linear Free Energy Relationships (LFERs), Volcano plots

Surface binding energies have been an important part of catalysis since the Sabatier principle a century ago (Figure 1).⁸³⁻⁸⁵ Quantitatively, the formation and reactions of

adsorbates are described by free energies. Kinetic-thermodynamic linear free energy relationships (LFERs) correlate rate parameters, $\ln(k)$ or ΔG^\ddagger , with driving force, $\ln(K_{\text{eq}})$ or $\Delta G^\circ_{\text{reaction}}$. In solution chemistry, LFERs have long been a central tool for understanding reactions,⁸⁶⁻⁹² starting from the Brønsted 'catalysis law' and Evans and Polanyi's generalization to free energies (eqs 1 and 2). Eq 2 is now called a BEP relation, for Brønsted (or Bell)-Evans-Polanyi. BEP relations are widely used in surface reactivity, as is illustrated by the 2010 *Chemical Review* "Reactivity Theory of Transition-Metal Surfaces: A Brønsted–Evans–Polanyi Linear Activation Energy–Free-Energy Analysis."⁸⁹ The common volcano plots can be viewed as two LFERs that intersect at the Sabatier optimum (Figure 2, see also Section IX).

$$\ln(k) = \alpha \ln(K_{\text{eq}}) + \ln(k_0) \quad (1)$$

$$\Delta G^\ddagger = \alpha \Delta G^\circ + \Delta G_0^\ddagger \quad (2)$$

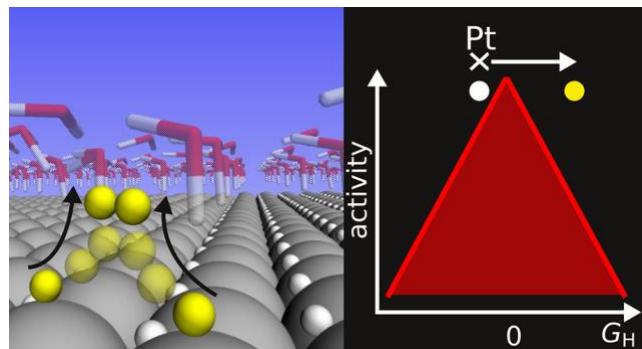


Figure 2. Left: Illustration of the Tafel mechanism for H_2 formation at Pt(111), where two OPD-H atoms (yellow balls) come together to form H_2 . The yellow atop OPD-H are distinguished from the inert, silvery UPD-H in hollow sites (see Section VI). Right: Volcano plot correlating catalytic activity with their free energy to bind $\frac{1}{2}\text{H}_2$ (G_H). Reprinted (adapted) with permission from *ACS Catal.* **2020**, *10*, 121–128.⁹³ Copyright 2019 American Chemical Society.

In parallel with LFERs for solution reactions, similar equations were advanced for electrochemical catalysis: the Tafel equation and then the Butler–Volmer equation (eqs 3,4). These relate the current density $\log(j)$, a measure of rate, to the overpotential η that describes the driving force. η is the difference between the applied potential and the equilibrium potential under the catalytic conditions (eq 5), and j_0 is the formal current density in each direction at $\eta = 0$. These are very similar to eqs 1 and 2 since $\ln(K_{\text{eq}})$, ΔG° , and nE° are conceptually the same in this context (eq 6).

$$\log(j) = \log(j_0) \pm \alpha \eta \frac{2.303F}{RT} \quad (3)$$

$$\log(j) = \log(j_0) \left\{ \beta \eta \frac{2.303F}{RT} - (1 - \beta) \eta \frac{2.303F}{RT} \right\} \quad (4)$$

$$\eta = E_{\text{applied}} - E_{\text{equilibrium}} \quad (5)$$

$$\Delta G^\circ = -RT \ln(K_{\text{eq}}) \quad \Delta G^\circ = -nFE^\circ, \quad \Delta G = -nF\eta \quad (6)$$

The central assumption of LFERs, a volcano plots, and the Sabatier Principle is that the reactions being compared are 'similar.' Often 'similar' is defined intuitively (all metal phosphide surfaces) or in a circular fashion, by taking reactions that follow the same LFER as similar. The question of similarity is central to the application of LFERs and the

Sabatier Principle, and to this Perspective, as discussed in Section V.ii below.

Eqs 1-4 all relate rate to driving force. When the k , $\Delta G^\ddagger, j$, K_{eq} , ΔG° and η all refer to a single step, then all the proportionality constants α and β have the same value. Across a series of reactions, α and β report what fraction of the change in driving force appears in the barrier, with values often taken to be $\sim\frac{1}{2}$.

However, LFERs are often extended beyond their original definition to correlations that involve *overall* rates for a multistep catalytic reactions [$\log(j)$, $\log(k)$, turnover frequency (TOF)].⁹⁴⁻⁹⁶ These rates are correlated with thermochemical parameters for one step, or for the entire catalytic cycle (for electrochemical overpotentials). This jump from a single step to a multistep process requires linear *scaling relationships*, which relate the ΔG° and ΔG^\ddagger_i of each of the i kinetically significant steps. Scaling with the overall ΔG° or η can also be needed.⁹⁷ Such relationships have been established computationally for a variety of systems, predominantly on relatively simple catalytic surfaces.^{72,94,98-102} However, such scaling is less likely to hold for multistep catalysis on complex surfaces, that can have different intermediates on different sites. For instance, stepped surfaces have different scaling than flat ones.^{90,103} In general, pre-equilibrium steps $\ln(K_{eq,i})$ scale with 100% of the ΔG^\ddagger_i for that step (eq 6), while the *barrier* for the rate limiting step scales with only a fraction of the driving force for that step (eq 2).⁹⁶

LFERs are at the center of most semi-quantitative analyses of surface catalysis, and they are a foundation of the long history of volcano plots (e.g., Figure 2 for the HER).^{72,91,93-106} Both the rising and falling lines are essentially LFERs that correlate a measure of a catalytic rate (e.g., $\ln(k)$ or $\ln(j)$) with a thermochemical descriptor. In Figure 2, the descriptor is the hydrogen binding energy ΔG_H (referenced to H_2).⁸⁹ Volcano plots are simplified models^{93,99,107} that are widely used as a quantitative implementation of the Sabatier principle. Many sets of catalytic reactions show a volcano relationship and these plots have provided very valuable rationales and predictions.^{27,90} Still, the simplifications inherent in these plots are harder to justify for complex multi-element, nanoscale catalysts, as described below.

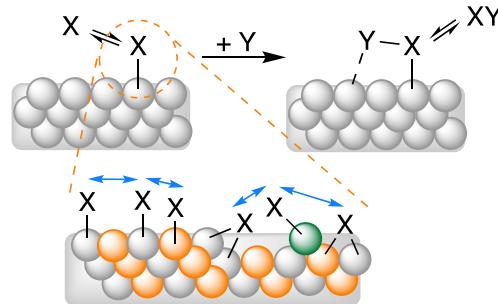
V. Background 2: Non-ideality and Isotherms

Ideal behavior for an adsorbate X means that its activity a_X parallels its surface coverage θ_X . The definition of an ideal *solution* of X is similar, that the activity of X is given by its concentration $[X]$. Chemists and chemistry textbooks almost always deal with solutions that are close to ideal. Deviations are notable at high concentrations, for ions in lower-polarity solvents, with strongly-interacting mixed solvents, and in some other situations. These are often the conditions of practical catalysis and electrocatalysis, so catalysis scientists and engineers are in general much more familiar with non-idealities¹⁰⁸⁻¹¹⁰ (and they can skip this section or object to its simplifications). In this Perspective, a_X is used for *surface activities* and concentrations will be used for solutions (assumed ideal).

i. Site heterogeneity and adsorbate interactions

Adsorption of a species X can be non-ideal for different reasons. Site heterogeneity means that adsorbates can occupy different kinds of sites. These can be different sites on a single-crystal surface, e.g., bridge vs. atop, surface defects such as vacancies or steps, or different elements on the surface (Scheme 1). Non-ideality also results from interactions between adsorbates (blue arrows in Scheme 1).

Scheme 1. Simplified mechanism for $X + Y \rightarrow XY$ at a single crystal metal surface (top) versus the complexities of X binding to an irregular, impure binary material (bottom) with inter-adsorbate interactions (blue arrows).



Generally, non-ideality is more common on surfaces than in solution. Gileadi has a particularly clear explanation of this difference:¹¹¹

[On a pure metal,] there is an average distance of 0.28 nm (2.8 Å) between atoms. The equivalent bulk concentration is 78 M, which is evidently much above anything that can be observed experimentally. Thus, we find that the molecules are very crowded on the surface. A relatively low fractional coverage of $\theta = 0.1$ is equivalent to a bulk concentration of about 2.5 M, in which interactions between solute molecules are expected to be the rule, rather than the exception....¹¹¹

He also has a colorful analogy for the presence of multiple surface sites:¹¹¹

Solid surfaces are rarely homogeneous.... This is much like people entering a movie theater with unmarked seats: the best seats [most negative ΔG_{ads}] are taken first, and the worst remain for late-comers....¹¹¹

One could add that, when possible, most moviegoers prefer to leave open seats between their group and other patrons—a rough analogy with repulsive site-site surface interactions.

The quotes from Gileadi, Parsons, and Conway here and above predate the current movement toward nanoparticle and multi-element catalysts, which makes non-ideality even more important.

ii. Sorting different surface sites

In traditional surface science, analyzing adsorption typically starts with the identification of the different types of surface sites.¹¹² Often flat crystal and stepped surfaces are investigated separately, as these typically fall on different LFERs.^{103,113} Interactions among adsorbates are considered as a correction term, as needed. When distinct sites behave differently, the reaction is said to be ‘structure sensitive.’¹¹² *Sorting into different sites is central to this Perspective, and to any treatment of heterogeneity.*

When comparing different reactions, surface sites, or catalysts, these must be “similar” in order for them to follow

a single LFER, volcano plot, or the Sabatier Principle. The mathematical definition of similarity for this purpose is that eq 2 can be simplified to use ΔG° as the *only* parameter for the series of reactions (eq 7).

$$\Delta G^\ddagger = \alpha \Delta G^\circ + \Delta G_0^\ddagger \quad (2) \quad \Rightarrow \quad \Delta G^\ddagger = f(\Delta G^\circ) \quad (7)$$

This requires that the α and ΔG_0^\ddagger parameters are constant or vary in a way that is uniquely defined by ΔG° . This simplification has been justified for simple surfaces based on computed scaling relationships. For instance, DFT calculations of energy barriers for A–B bond cleavage reactions (E_{ts}) on flat single crystal metal surfaces scale very well for different metals and different bonds, with eq 7 becoming $E_{ts} = 0.84E_{diss}(A-B) + 1.92$.⁹⁰ Still, the scaling is different for stepped surfaces.^{90,103}

The simplest and most common sorting, first proposed by Boudart, is to use a find Most Abundant Reactive Intermediate (MARI) and perhaps a Most Abundant Surface Intermediate (MASI), when these can be identified and quantified.^{114,115} This is often a sufficient first approximation. When comparing different catalysts, it adds the assumptions of common MARI and MASI species and coverages.

An additional complication is that eqs 2 and 7 implicitly assume Langmuir isotherms (unless all the parameters become functions of θ). Non-ideal adsorption adds additional features that must follow requirements for similarity such as scaling relationships (see below).

It is the hypothesis of this Perspective that the simplification in eq 7 cannot be done for complex, nanoscale surfaces composed of multiple elements, and that scaling relationships will not hold. The α and ΔG_0^\ddagger parameters depend on many things beyond ΔG° , such as whether the reactive H is bound to one element or another, or in an atop vs. a 3-fold vs. a step site. It is my hope that this hypothesis will be tested by experiments and computations that include the full complexity of nanoscale M_nX_m and single entity¹¹⁶ catalysts under operating conditions.

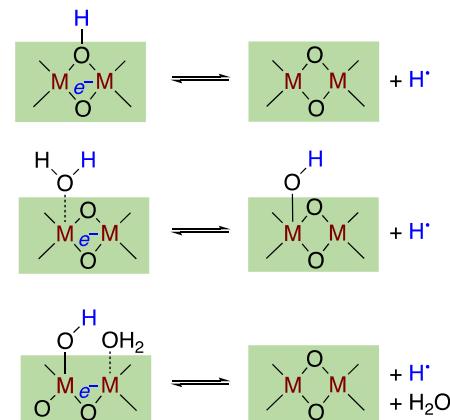
The challenge is that there is usually little or no understanding of how to do this sorting for complex surfaces. The number and nature of the various sites, their energetics, and their reactivity are challenging to explore. For instance, in the Fischer-Tropsch process, the stoichiometry, surfaces, and binding sites for H, CO, and $C_xH_yO_z$ intermediates on the nanoscale iron-carbide particles are essentially unknown. The use of complex surfaces with various sites brings the sorting problem to the fore.

Therefore, it is valuable to measure a binding isotherm for a surface in its entirety, without any sorting. The isotherm will likely cover multiple kinds of sites and could provide insight into the sorting process, as discussed below.

A brief foray into metal oxide surfaces illustrates the sorting challenges, and it sets the stage for the case studies below. Redox reactions of oxides in contact with protic media generally occur by proton-coupled electron transfer (PCET), the transfer of $H^+ + e^-$ or, equivalently, a hydrogen atom (H^\cdot).^{117,118} Scheme 2 sketches a few possible pathways for oxidation of a protonated or hydrated surface. The proton

can be transferred from an oxygen that part of the lattice or from a surface hydroxide or water molecule. The oxide or hydroxide formed upon H^+ loss could become part of the lattice. More generally, the less crystalline the surface, the more hydrated, porous, and amorphous, the less meaningful is any distinction between lattice O and adatom sites, and the more muddled are the definitions of surface sites and surface coverages.

Scheme 2: H-atom loss from a hydrated metal oxide surface could occur from a lattice oxygen, a surface hydroxide, or a weakly bound water molecule. The green boxes indicate the oxide lattice, but this distinction is likely less clear in practice.



The challenges in sorting blur the traditional distinction between surface heterogeneity and adsorbate-adsorbate interactions. These effects are also often comingled in experimental measurements. Therefore, adsorption will be described here just as ‘non-ideal,’ in most cases without specifying an origin.

iii. Characteristics of ideal and non-ideal adsorption

Ideal adsorption is characterized by a single free energy of adsorption, ΔG_{ads}° . We primarily use here $\theta = 1/2$ as the standard state for ΔG_{ads}° . This allows direct analogies with the typical solution standard state where all dissolved species have equal concentrations (1 M for ideal solutions). The analogous situation on a surface is when half the surface sites are occupied by the adsorbate, $\theta = (1 - \theta) = 1/2$.

A common alternative surface standard state uses the properties of the adsorbate at low coverage, $\theta \rightarrow 0$, extrapolated to $\theta = 1$ (though see¹¹⁹⁻¹²³). The $\theta \rightarrow 0$ extrapolation avoids interactions between adsorbates but is less intuitive.

When adsorption is non-ideal, the *apparent* free energy of adsorption is a function of coverage: $\overline{\Delta G}_{ads}^\circ(\theta)$, and the activity a_X need not be proportional to θ_X . Non-ideality makes kinetic and equilibrium expressions complex, so it is helpful to use a qualitative metric such as the width of the free energy isotherm described by $\overline{\Delta G}_{ads}^\circ(\theta)$. This width is roughly equivalent to the range of surface–X bond dissociation free energies (BDFEs) and to the width of an anodic or cathodic peak in a cyclic voltammogram (CV).

Moderately wide isotherms have only a modest effect on kinetic parameters. Adsorbates with $\overline{\Delta G}_{ads}^\circ$ 2.7 kcal mol⁻¹ (120 meV) weaker than the $\theta = 1/2$ standard state ΔG_{ads}° should undergo reactions roughly ten times faster than

adsorbates at $\theta = \frac{1}{2}$ (at 298 K using eqs 1,2 with $\alpha = \frac{1}{2}$). However, this increase in k is often offset by the lower population of more weakly bound sites. This balance is the origin of the traditional intuition that moderate non-ideality of a single active site is not typically an important factor in catalytic reactions (though see ³⁴).

iv. Ideal behavior: the Langmuir isotherm

Ideal adsorption of X is described mathematically by the Langmuir isotherm, shown in one form in eq 8. For a molecular adsorbate like CO, the activity is taken as its partial pressure or its solution concentration. When the X is an unstable fragment, such as a hydrogen atom, a_X is often referenced indirectly, for instance H^* to H_2 gas.

$$\frac{\theta}{1-\theta} = K_{\text{eq}}[X] = (e^{-\Delta G^\circ/RT})[X] \quad (8)$$

$$A + X \rightleftharpoons AX \quad \frac{[AX]}{[A]} = K_{\text{eq}}[X] \quad (9)$$

$$\text{Ox} + e^- \rightleftharpoons \text{Red} \quad E = E^\circ - \frac{RT}{nF} \ln \left(\frac{[\text{Red}]}{[\text{Ox}]} \right) \quad (10)$$

$$\overline{\Delta G}_{\text{ads}}(\theta) = \Delta G^\circ_{\text{ads}} + \frac{RT}{F} \ln \left(\frac{\theta}{(1-\theta)} \right) \quad (11)$$

The Langmuir isotherm for X binding to a surface is mathematically identical to the equilibrium expression for X binding to a solution species (eq 9), and to the Nernst equation (eq 10). The $\theta/(1-\theta)$ term is analogous to the ratio of bound-to-unbound ($[AX]/[A]$) or reduced-to-oxidized ($[\text{Red}]/[\text{Ox}]$). The Langmuir isotherm can also be written in parallel to the Nernst equation, as a coverage-dependent free energy $\overline{\Delta G}_{\text{ads}}(\theta)$ (eq 11). Just as the effective acidity of a buffer varies from the K_a when $[\text{Acid}] \neq [\text{Base}]$, and the potential varies from E° when $[\text{Red}] \neq [\text{Ox}]$, $\overline{\Delta G}_{\text{ads}}(\theta)$ has a Gaussian distribution around $\Delta G^\circ_{\text{ads}}$. All of these have a full-width at half-max (FWHM) of 90 mV (2 kcal mol⁻¹, or a 30-fold change in concentration).

v. Non-ideal behavior: Frumkin and Temkin isotherms

Adsorption is said to be non-ideal when the coverage-dependent free energy $\overline{\Delta G}_{\text{ads}}(\theta)$ does not follow a Langmuir isotherm. A common way to treat this is to simply add a term that is linear with θ . Eq 12 shows one version of the Frumkin isotherm as a Langmuir isotherm with an added $C(\theta - 0.5)$ term. The Temkin isotherm is simpler, omitting the $\theta/(1-\theta)$ term. These are two of the many isotherms, which can each be expressed in different ways and used for various purposes.^{71, 86,124} The parameter C of the added term has units of free energy, often given as unitless $f = C/RT$. Since the $C(\theta - 0.5)$ term is zero at the $\theta = 0.5$ standard state, all of these isotherms have the same standard state $\Delta G^\circ_{\text{ads}}$, but the Frumkin and Temkin isotherms are wider ($C > 0$) or narrower ($C < 0$) than the ideal Langmuir.

$$\overline{\Delta G}_{\text{ads}}(\theta) = \Delta G^\circ_{\text{ads}} + \frac{RT}{nF} \ln \left(\frac{\theta}{(1-\theta)} \right) + C(\theta - 0.5) \quad (12)$$

(Frumkin)

$$\overline{\Delta G}_{\text{ads}}(\theta) = \Delta G^\circ_{\text{ads}} + C(\theta - 0.5) \quad (\text{Temkin}) \quad (13)$$

The Frumkin and Temkin isotherms were originally introduced to model inter-adsorbate interactions, though they are also useful when the origin of the non-ideality is not known. When the C term is small (and at low θ with a $\theta = 0$ standard state), eq 12 reduces to the Langmuir isotherm (eq 11). Alternatively, when C is large, or when θ is close to

$\frac{1}{2}$ and $\theta/(1-\theta)$ is close to 1, then the isotherm is dominated by the new linear term. Then the Frumkin and Temkin isotherms are the same (eq 12 \Rightarrow eq 13), and the dependence of the coverage on free energy is *linear*. This contrasts with ideal behavior, where concentrations, rate constants, and equilibrium constants for solutes and adsorbates vary *exponentially* with energies, not linearly.

The isotherms plotted in Figure 3⁸⁶ show the difference between the Langmuir isotherm ($f = 0$) versus the Frumkin isotherms with linear sections that vary with f [$C = f(0.59 \text{ kcal mol}^{-1})$ at 298 K]. The horizontal axis is the driving force, using the $\theta \rightarrow 0$ standard state, plotted as the applied electrochemical potential minus E° (bottom axis) and as $\overline{\Delta G}_{\text{ads}}(\theta) - \Delta G^\circ_{\text{ads}}$ (top axis). The equivalent plot using the $\theta = \frac{1}{2}$ standard state has the same curves, just shifted so that the isotherms all intersect at $\theta = 0.5$, where $(E - E^\circ) = (\overline{\Delta G}_{\text{ads}}(\theta) - \Delta G^\circ_{\text{ads}}) = 0$.

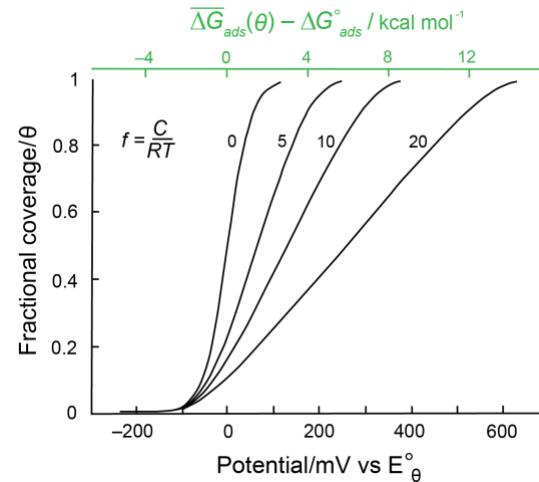


Figure 3. The dependence of θ on applied potential for different values of the unitless parameter $f = r/RT$. $f = 0$ corresponds to the Langmuir isotherm. The potentials are vs. the standard potential, using the $\theta \rightarrow 0$ standard state [and X at standard state, eq 8]. $e(E_{\text{applied}} - E^\circ)$ and $\overline{\Delta G}_{\text{ads}}(\theta) = \Delta G^\circ_{\text{ads}}$ are the relative free energy of adsorption vs. that standard state. The equivalent plot with the $\theta = 0.5$ standard state used in this Perspective has the same curves, just shifted so that they all intersect at $\theta = 0.5$ (zero driving force). Reproduced and adapted from *Electrode Kinetics for Chemists, Chemical Engineers and Materials Scientists*: Gileadi, E. 1996, figure 21, p. 268.⁸⁶ Reproduced with permission of John Wiley and Sons and Copyright Clearance Center.

One implication of the linear dependence of θ on energy for the Frumkin isotherm is that the population of adsorbates in each energy range ($d\theta/d\Delta G$) is roughly constant. In this case, the faster rates for more weakly-bound adsorbates are *not* offset by a lower population. This balance was described above as an origin of the intuition that broad isotherms do not strongly affect rate constants. A flat region of a $d\theta/d\Delta G$ plot indicates a Frumkin isotherm, as in the roughly rectangular section for UPD-hydrogen on Pt(111) at the left of Figure 4 in the next section. The challenges of developing LFERs for non-ideal adsorbates are raised below (Section VIII).

As a final general comment, LFERs give barriers and rate constants but catalysis depends on *rates* rather than rate constants. This is emphasized in the degree of rate control analysis,^{54,99-102,114,115} and is why Tafel slopes for multistep reactions depend on the sequence of elementary steps, pre-

equilibria, barriers, and surface coverages.^{72,94,98-101} One example is the oxidation of ammonia over noble metal surfaces, for which the selectivity between N₂, NO, and N₂O is calculated to be due to relative surface coverages on Pt and Pd rather than differences in barriers (as for Rh).¹⁰²

A recent *Chemical Review* provides a detailed mathematical analysis of complex reversible reactions, including non-idealities, concluding with:⁵⁴

There is, therefore, a pressing need for the development of analytical diagnostics which enable identification of geometric and/or electronic effects outside of goodness-of-fit to conventional metrics such as reaction order or activation energy. Is there, for example, a kinetic or thermodynamic “signature” of site-ensemble requirements or lateral adsorbate interactions which can be elicited from steady-state or transient rate measurements either near or far from equilibrium?⁵⁴

VI. Hydrogen on Pt(111)

This Perspective emphasizes isotherms for hydrogen adsorption (*H) because H transfers are involved in many kinds of catalysis, and because its isotherms are the most studied.¹²⁵⁻¹³⁸ While H might be thought of as a simple adatom, its isotherms are often not simple.

The cyclic voltammogram (CV) for a Pt(111) surface in aqueous 0.1 M HClO₄ is remarkably complex (Figure 4), given the simplicity of the system. The single-crystal surface is flat with identical surface Pt atoms, and the only relevant solution species are H₂O and H₃O⁺ (ClO₄⁻ is chemically inert). The CV is truncated just above the H₂ thermodynamic potential (0 V vs. the reversible hydrogen electrode, RHE, eq 14) because Pt is an exceptional HER catalyst and the cathodic current becomes large approaching 0 V.

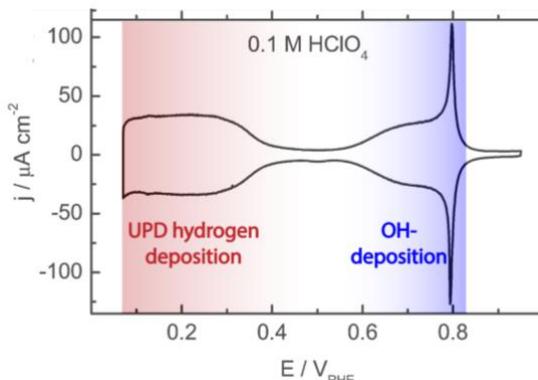
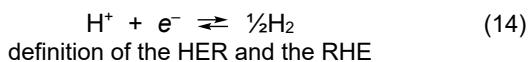
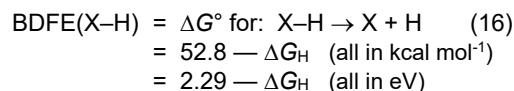
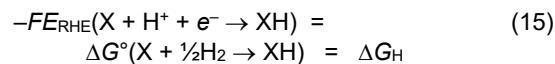


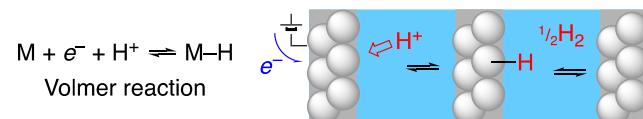
Figure 4. Cyclic voltammogram of Pt(111) in 0.1 M HClO₄, with labels for the UPD-H and OH adsorption regions added.¹³⁹ Adapted from García, G.; Koper, M. T. M. *Phys. Chem. Chem. Phys.* **2008**, *10*, 3802¹³⁹ with permission from the Royal Society of Chemistry.

The mirror relationship between the anodic and cathodic currents shows that the CV reports on thermodynamic properties. When potentials are referenced to RHE, then FE° is the same as the traditional surface-science measure of H affinity, the free energy for addition of $\frac{1}{2}\text{H}_2$ (ΔG_H , eq 15). E° vs. RHE, ΔG_H , and the surface–H bond dissociation free energies (BDFEs, eq 16) are all conceptually the same.¹¹⁸



The feature at the left of the CV in Figure 4 is due to the underpotential deposition of hydrogen (UPD-H), so named because H is added to the Pt surface *under* the potential needed to form H₂ ($V_{\text{RHE}} > 0$). The electrochemical process is called the Volmer reaction, and it forms the same UPD-H as from Pt(111) and H₂ gas in the same aqueous media (Scheme 3). The CV measures the current ($d[\#e^-]/dt$), so the shape of the CV wave is the derivative of the coverage vs. applied potential (Figure 3 vs. Figure 4). The UPD-H feature is at least 300 mV wide, much wider than the Langmuir isotherm's Gaussian 90 mV. The flat top of this feature implicates a Frumkin isotherm with a repulsive interaction between surface H, with $C \cong +250 \text{ mV} (+6 \text{ kcal mol}^{-1})$.

Scheme 3: Formation of surface H by proton reduction (the Volmer reaction, left and center), and by addition of H₂ (right).



At higher potentials, the complex CV shape for Pt(111) is due to the deposition of hydroxide. The broad feature starting at ~0.6 V is the initial accumulation of surface OH ($C > 0$), then the sharp feature (~0.78 V) indicates cooperative adsorption to form an ordered monolayer ($C < 0$). Catalysis on Pt(111) at high potentials is strongly influenced by this isotherm.^{79,140-142} It should also be noted that the isotherms are different on different Pt crystal faces.

There is general consensus that Pt(111) has a second, more important kind of adsorbed H: *overpotential* deposited H (OPD-H).^{129-131,143-146} OPD-H is very active in catalysis, forming H₂ and hydrogenating substrates, while UPD-H is quite unreactive. The isotherm for OPD-H is more challenging to examine because of the rapid H₂ formation, but an atop site is indicated by the v(PtH) IR adsorption at $2050 \pm 50 \text{ cm}^{-1}$ (depending on the conditions) and its appropriate shifting in D₂O.^{129-131,143-147} UPD-H, however, lies in a three-fold hollow site. Computational studies in general support this picture.^{93,148}

The message of H on Pt(111) is that there are two kinds of H that are very different in structure and reactivity. Kinetic treatments of H-transfer reactions on Pt(111) have only considered the OPD-H because of the unreactive nature of UPD-H. The *complete* isotherm for H on Pt(111), however, would include both UPD-H and OPD-H. This illustrates the importance of sorting into distinct sites (Section V.ii above). However, such sorting is not yet possible for the more complex catalysts being currently developed.

Pt(111) is one of the simplest, most ordered, and most studied catalytic surfaces. Even on this surface, the isotherm is not simple. *Adsorbates and intermediates on most catalytic surfaces have complex isotherms that impact catalysis.*

VII. Diversity of Hydrogen on Metal Oxides

Metal oxides can be active catalysts, co-catalysts, and catalyst supports. Examples include IrO_2 and others for the OER^{25,149-151} and cerium oxide in automotive ‘three-way catalysts’.^{152,153} Oxide surfaces have received great attention, from single crystals in ultra-high vacuum (UHV) to minerals in natural waters,¹⁵⁴⁻¹⁵⁶ including the thermochemistry of bulk oxides and hydroxides.^{157,158} Still, the surface binding or loss of molecular fragments such as H or O is only beginning to be documented.¹⁵⁵

Addition of an H atom is equivalent to addition of a proton and an electron ($\text{H}^\bullet = e^- + \text{H}^+$) or $\frac{1}{2} \text{H}_2$. H can be added in a chemical or electrochemical PCET reaction, as in Schemes 2 and 3 above.^{117,118} Addition of H atoms to a metal oxide generally forms a surface hydroxide and adds an electron to metal-based orbitals, so H is an *n*-dopant.¹⁵⁹ The electron can be added to an ion with a localized valence (Ce^{3+} in CeO_2 , Figure 5A), to an *n*-type multi-metallic trap state (TiO_2), or to a metallic band structure (IrO_2). Protons are ubiquitous on oxide surfaces exposed to air or water, as hydroxides or bound H_2O .¹⁶⁰

i. H on colloidal ceria nanocrystals

Colloidal cerium oxide nanoparticles (NPs) are a relatively simple system that has been found to be a highly non-ideal. Suspended in THF, with oleate capping ligands, these ~ 1.8 nm NPs react reversibly with soluble H-atom transfer (HAT) reagents, such as 2,6-dimethylhydroquinone (Figure 5B).¹⁶¹⁻¹⁶³ The stoichiometry of this reaction was experimentally established, that it only exchanges H atoms. Therefore, at equilibrium, the affinity of the NPs for H is the same as that of the reagent: the O–H surface BDFE of the NPs is equal to the average O–H BDFE of the reagent. Measuring the equilibrium stoichiometry with reagents of different BDFEs mapped out the isotherm for H-binding to the ceria NPs (Figure 5C).¹⁶²

The measured CeO-H BDFEs span 13 kcal mol⁻¹ (0.56 eV) with changes in the average redox state of the surface of the nanoceria colloid (Figure 5A,C).¹⁶² At higher % Ce^{3+} , adding H^\bullet is more difficult implying weaker O–H BDFEs. The BDFE variation corresponds to a change of 10^9 in equilibrium constant for HAT. The wide range is likely due to increasing lattice strain as more Ce^{4+} is converted to larger Ce^{3+} ions. The strain makes it more difficult to accommodate more Ce^{3+} ions, and therefore weakens the O–H bonds formed. The mostly linear shape of Fig. 5C indicates that these data can be well fit to a Frumkin isotherm.

A subsequent study examined the kinetics of H-atom transfer from $\text{CeO}_{2-x}\text{H}_y$ to the hydrazyl radical DPPH as a function of θ_{H} .¹⁶³ By using only a small amount of DPPH, the θ_{H} , CeO-H BDFE, and $\Delta G^\circ_{\text{HAT}}$ did not change significantly over the reaction. Rate constants for reactions initially poised at different θ_{H} gave a LFER connecting $\ln(k_{\text{HAT}})$ with $\ln(K_{\text{eq}})$, the latter from the $\Delta G^\circ_{\text{HAT}}$ of each reaction. The linear relationship showed that all of the CeO-H species belonged to the same type of surface H, though the BDFE was changing substantially. Surprisingly, the Brønsted α for these reactions was only 0.2, much less than that often assumed 0.5.

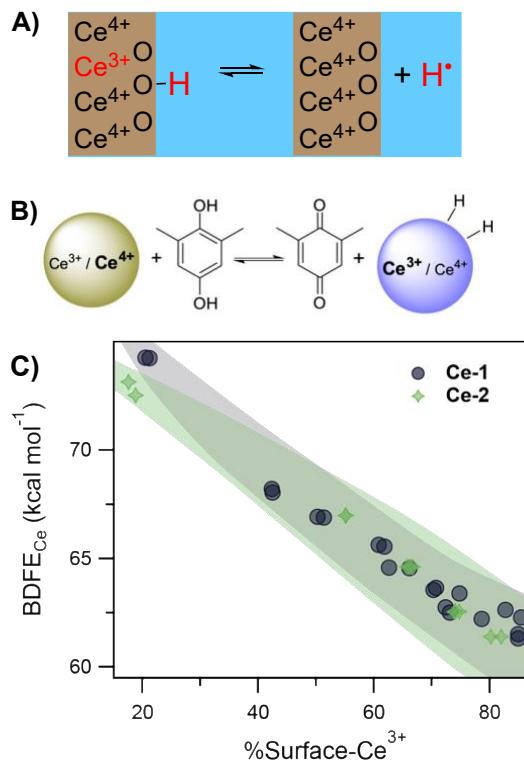
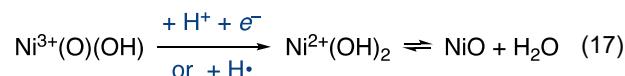


Figure 5. **A)** Schematic of H-atom addition and removal from a cerium oxide (ceria) surface. The free energy of the forward reaction is the surface O–H BDFE. **B)** At equilibrium between colloidal $\text{CeO}_{2-x}\text{H}_y$ NPs and a hydroquinone/quinone couple, the CeO-H BDFE (BDFFe_{ce}) is that same as the average BDFE in the H_2Q . **C)** Plot of BDFFe vs. % surface Ce^{3+} of two batches of $\text{CeO}_{2-x}\text{H}_y$ NPs. The shaded regions show the estimated uncertainties. **B** and **C** Reprinted (adapted) with permission from *J. Am. Chem. Soc.* 2021, 143, 2896.¹⁶² Copyright 2021 American Chemical Society.

ii. H on nickel oxide

NiO has been widely studied as an earth abundant p-type semiconducting oxides, and—with some iron—a good water oxidation electrocatalysts. In this Perspective, electrochemical studies of NiO thin films are introduced as a compliment to the thermal reactions of ceria NPs in the last section. NiO has long been known to typically show $\text{Ni}^{3+}/\text{Ni}^{2+}$ and $\text{Ni}^{4+}/\text{Ni}^{3+}$ PCET redox couples (e.g., eq 17).^{164,165} Eq 17 also shows the $\text{Ni}(\text{OH})_2/\text{NiO} + \text{H}_2\text{O}$ equilibrium to indicate the complexity of hydrous metal oxides.



The NiO-H isotherm was measured by spectroelectrochemistry, at each of the potentials shown by dotted lines over the CV in 65A.¹⁶⁵ For the $\text{NiOOH}/\text{Ni}(\text{OH})_2$ couple, the % Ni^{2+} was taken as equal to the θ_{H} , following eq 17. The data followed a Frumkin isotherm (solid line in Figure 6B) rather than ideal behavior (dashed line). The Frumkin C parameter was 6 kcal mol⁻¹ (0.25 eV).¹⁶⁵

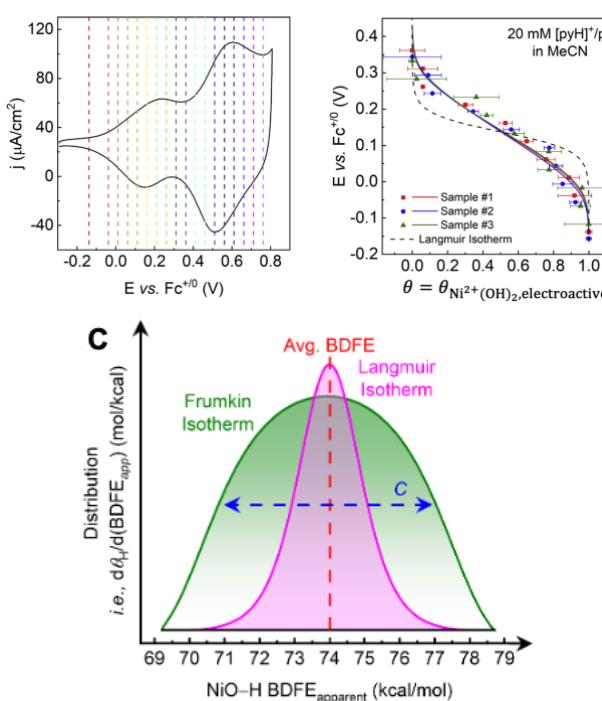


Figure 6. **A)** CV of a NiO thin film in pyH^+/py buffered MeCN. **B)** $\theta_{\text{H}} = \theta_{\text{Ni}^{2+}}$ at each of the potentials shown in A) by the vertical dashed lines. The solid black line is the fit of three datasets to a single Frumkin isotherm; the dashed line is the best Langmuir isotherm fit. **C)** Representation of the NiO isotherm as a range of BDFEs (green) that is much wider than the ideal Langmuir isotherm. *Permission requested.* [Reprinted with permission from *Chem.* **2022**, *8*, 3324–3345. Copyright 2022 Cell Press.]

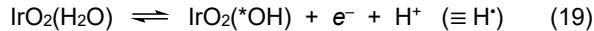
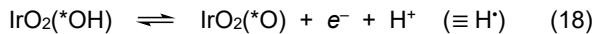
Similar measurements were done in aqueous solutions and ‘dry’ organic solvents (MeCN and DMF), and with a wide range of different buffers, over 10^{16} in proton activity. Remarkably, *all of the isotherms were the same* when referenced to RHE. The center of each isotherm, the average NiO–H BDFE, was $74 \pm 2 \text{ kcal mol}^{-1}$ under all conditions, and the Frumkin C parameters were all the same. This was not expected for electrochemical measurements under quite different conditions. However, such solvent and medium independence is typical of molecular hydrogen atom transfer reactions.¹¹⁸ Thus even this electrochemical reaction has the character of H^+ transfer (eq 17 under the arrow).

iii. H on IrO_2 and RuO_2

IrO_2 and RuO_2 are among the most active metal oxides for redox catalysis such as the OER, with the most active forms believed to be porous and in part amorphous.^{24,166–175} RuO_2 is used commercially in the dimensionally stable anodes (DSA) to produce Cl_2 as part of the large-volume chlor-alkali process.^{176–178} These oxides have been studied in great detail and *in operando* experiments are generating many new insights.^{25,26,170,173,182} Two recent sets of studies have emphasized isotherms and are interesting to compare. Suntivich *et al.* have examined highly ordered epitaxial thin films of $\text{IrO}_2(110)$ and $\text{RuO}_2(110)$ on TiO_2 ,^{171–172} while the groups of Stephens, Durrant, and Rao (SDR) have collaborated to examine hydrous, amorphous iridium oxide films electrodeposited on FTO.¹⁷³

Both studies reported two reversible redox features in electrochemical experiments (SDR observed a third). Most

of the metal ions in the porous film were electroactive, in contrast to only the undercoordinated sites (ucs) at the surface of the crystalline sample. Suntivich *et al.* described the electrochemical process as the binding energies for $^*\text{OH}$ and $^*\text{O}$ to Ir ucs.^{171–172} However, the PCET description^{117,118} used by SDR seems more reasonable to us, as the DFT calculated structures of $\text{IrO}_2(110)(^*\text{OH})$ and $\text{IrO}_2(110)(^*\text{O})$ differ by an H atom (Figure 7B and C).¹⁷¹ One redox couple was assigned as PCET from a water perhaps weakly bound to an Ir ucs (eq 19; cf. Scheme 2 above).



The isotherms reported in these studies were non-ideal, and they needed to be fit with Frumkin rather than Langmuir isotherms. The Frumkin C parameters for the $\text{MO}_2^{4+/5+}$ couples were similar (3–7 kcal mol^{−1}), both for highly ordered crystalline films and for the disordered, porous, amorphous films. This rough independence of the non-idealities on the form of the catalyst is surprising. It suggests that the apparent repulsion between the adsorbates ($C > 0$) is not due to a direct through space interaction, or due to surface inhomogeneity. The observed similarity supports the interpretation by both teams that the apparent repulsion is mediated through the aqueous double-layer.

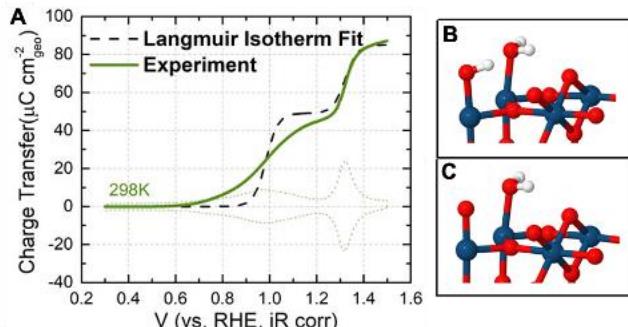


Figure 7. **A)** Plot of charge transferred (Q) vs. potential, which is the isotherm, for an $\text{RuO}_2(110)$ thin film in 0.1 M KOH. The fainted dotted trace is the CV (dQ/dV). The dashed best-fit line shows that using two Langmuir isotherms does not follow the data. **B)** and **C)** are portions of the DFT-computed structures for $\text{IrO}_2(100)(^*\text{OH})$ and $\text{IrO}_2(100)(^*\text{O})$, respectively, with the OH and O at the upper left (Ir in blue-grey). Part **A** reprinted with permission *requested* from Hu, B.; Kuo, D.-Y.; Paik, H.; Schlam, D. G.; Suntivich, J. *J. Chem. Phys.* **152**, 094704 (2020).³² Copyright 2020 AIP Publishing. Parts **B** and **C** reprinted with permission *Journal of the American Chemical Society* **2017**, *139*, 3473–3479.¹⁷¹

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The redox couples measured in these and other studies of IrO_2 shift with pH, usually more than the predicted Nernstian shift of 59 mV per pH expected for $1\text{H}^+/1\text{e}^-$ stoichiometries (19 and 20, E vs. SHE). Both teams interpreted this shift as being due to Ir–O bonds became stronger at higher pH, while still retaining the $1\text{H}^+/1\text{e}^-$ chemical processes. However, the same >59 mV per pH “super-Nernstian” shift for many hydrated oxides has been interpreted differently in the pseudo-capacitor literature, including for RuO_2 .^{183,184} In that field, super-Nernstian behavior indicates a stoichiometry with more than one H^+ per e^- . In one TiO_2 system, there is evidence for a class of trap states with a 2H^+ per 1e^-

stoichiometry.¹⁸⁵ These issues illustrate the complexity of oxide surfaces, as discussed in Section V.ii above. Most important, this discussion emphasizes the *critical role of surface stoichiometry*, even if that is difficult to determine. Understanding the OER by IrO_2 will require measuring $\text{H}^+:\text{e}^-$ and $\text{H}_2\text{O}:\text{Ir}$ surface stoichiometries through each of the intermediates.

iv. H on TiO_2

TiO_2 is one of the most studied oxides, in its many forms.¹⁸⁶⁻¹⁸⁹ Here only one computational study is described, the perhaps extreme case of surface $\text{TiO}-\text{H}$ BDFEs.¹⁹⁰ One set of $\text{TiO}-\text{H}$ was formed *in silico* by adding H^+ to the TiO_2 surface, yielding a $^*\text{OH}$ coupled to e^- in a Ti d trap state ($\text{TiO}_2\bullet\text{e}^-\text{H}^+$, roughly Ti^{3+}OH). This set was compared with surface OH groups formed by non-redox hydration of the TiO_2 surface: $(\text{Ti}^{4+})_2\text{O}^{2-} + \text{H}_2\text{O} \rightarrow 2(\text{Ti}^{4+}\text{OH}^-)$ (Figure 8A and B).

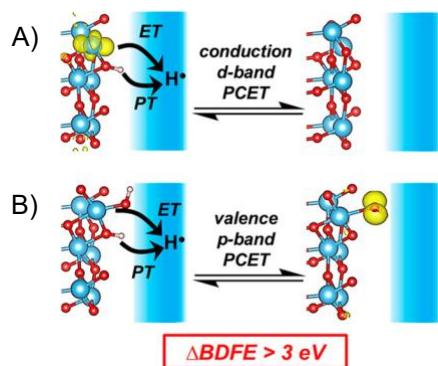


Figure 8. **A)** PCET from a Ti d-band state, formally $[\text{TiO}_2]^-[\text{H}]^+ \rightarrow [\text{TiO}_2]^0 + (\text{H}^+ + \text{e}^-)$, and PCET from **B)** an O p-band state, $[\text{TiO}_2]^0[\text{OH}]^-[\text{H}]^+ \rightarrow [\text{TiO}_2]^0[\text{O}]^+ + (\text{H}^+ + \text{e}^-)$. The converged geometries for each structure have a yellow isosurface representing the charge density associated with the Ti^{3+} polaron (top) or valence band hole (bottom).¹⁹⁰ Reprinted (adapted) with permission from *J. Phys. Chem. Lett.* **2021**, *12*, 40, 9761-9767. Copyright 2021 American Chemical Society.

The computed BDFE for the $\text{Ti}^{3+}\text{O}-\text{H}$ bonds in $\text{TiO}_2\bullet\text{e}^-\text{H}^+$ was ca. 40 kcal mol⁻¹. The $\text{Ti}^{4+}\text{O}-\text{H}$ bond, however, was dramatically stronger, ~120 kcal mol. (The examples within each of the sets were found to be similar.) The difference of ~80 kcal mol⁻¹ (~3.5 eV) is remarkable, especially for O-H bonds that are very similar in structure, bond length, and stretching frequency. The $\text{Ti}^{3+}\text{O}-\text{H}$ bonds are very weak (thermodynamically uphill of H_2) while the $\text{Ti}^{4+}\text{O}-\text{H}$ BDFE is close to that in H_2O . Previous experimental data for TiO_2 films, TiO_2 colloids, and molecular Ti^{3+}OH compounds, also have low BDFE.^{188,191}

The primary difference between these two O-H bonds is the source of the electron being removed to make H^+ . For “ $\text{Ti}^{3+}\text{O}-\text{H}$,” the e^- comes from a high-energy trap state near the conduction band. For $\text{Ti}^{4+}\text{O}-\text{H}$, however, the e^- is removed from a low-energy O lone pair near the valence band to leave an oxyl radical.¹⁹⁰

Thus, typical *n*-type TiO_2 surfaces in contact with water are covered with identical-looking hydroxyl groups, some of which have very strong O-H bonds (like water) and others have very weak O-H bonds that are better H-atom donors than H_2 . The different classes of O-H groups are distinguished by the presence or absence of a high-energy

trap-state electron that can be transferred with the proton. From a thermochemical perspective, the e^- and H^+ could be quite distant from each other, as long as their transfer doesn’t lead to an unfavorable buildup of charge. Molecular PCET reactions have been observed when the proton and electron were more than 1 nm apart in the reactants or products.¹⁹² The full isotherm for H on TiO_2 thus has structurally similar components highly separated in energy.

VIII. LFERs with broad isotherms

For the reasons developed above, LFERs are much more complicated for elementary reaction steps involving adsorbates with non-ideal binding isotherms, vs. the ideal cases in Section IV. Therefore, most treatments of LFERs and volcano plots assume ideality (though certainly not all treatments, *e.g.*,^{28,108,193-196}). In the context of the current Perspective, there are a number of issues to highlight.

First, LFERs require free energies, and determining ΔG° starts with defining a standard state. The $\theta = \frac{1}{2}$ standard state at first glance seems closer to the Sabatier principle and easier to apply. However, when the variety of sites and their activities, densities and populations are not known, “ $\theta = \frac{1}{2}$ ” is challenging.

Consider OPD-H and UPD-H on the ideal $\text{Pt}(111)$ surface (Section VI above). In this case, $\theta_{\text{H}} = \frac{1}{2}$ could refer to *i*) half of the UPD-H sites being occupied, *ii*) half of the OPD sites occupied, or the *iii*) half of all the sites occupied (most of the UPD and few of the OPD). Choices *i*) and *iii*) are unreasonable given the inertness of UPD-H and the high HER activity of $\text{Pt}(111)$, and *ii*) is ill-defined due to our limited understanding of OPD-H. The $\theta = 0$ standard state would not include any of the OPD active sites on $\text{Pt}(111)$. If very little OPD-H is formed at RHE ($\Delta G^\circ\text{H} = 0$), this free energy would unlikely be the point of highest activity—as concluded by Peterson *et al.* (Figure 2 above).⁹³

Broad overall isotherms also can complicate the “sorting challenge” (Section V.ii), that is, knowing which materials and reactions are similar enough to be treated with a single LFER or volcano plot. In some cases different sites on the same surface need to be treated with different BEP correlations, such as step sites vs. terraces on a metal surface.^{90,91,197} As one example of a more complex process, the lowest energy path for H_2 loss from hydrogenated Mo_2N was computed to be the heterolytic combination of a hydridic MoH with a protic NH.¹⁹⁸ In this case, using $^*\text{H}$ for a generic surface hydrogen is not appropriate—there are (at least) two kinds of $^*\text{H}$ with quite different properties. Such issues are expected to be common for multi-element, nanoscale catalysts that have a range of sites and mechanisms.

The recasting of LFERs for broad isotherms has been the subject of many studies over many years, including^{4,72,124,193-196,199-205}. This has been done algebraically, such as with the simplified Frumkin formula (eq 12), and with numerical treatments using experimental isotherms such as for OH on $\text{Pt}(111)$ (Figure 4). However, neither is a simple task. Kinetic and thermodynamic treatments of non-ideal situations are complex because rate ‘constants’ and apparent free energies vary with coverage.

The mathematical treatments often require the Brønsted law to connect energies to rate constants (eq 1), and typically assume $\alpha = 0.5$. However, this parameter can vary substantially. An early paper by Nørskov *et al.* found a “universal α ” of 0.84 for computed X–Y bond dissociation on flat transition metal surfaces.^{90,197} For the Volmer reaction at Hg and Au electrodes with various proton donors and in various solvents, Conway, and Jackson and Surendranath, reported α values from 0.3 to 0.7,^{206,207} which has recently been generalized.²⁰⁸ Our kinetic studies of the hydrogen atom transfer from the CeO_xH_y nanoparticles in (Section VII.i) gave an α of 0.2 (over a 0.6 eV range).¹⁶³ The origins of these different α values are not clear. Even for well-defined homogenous, molecular HAT/PCET reactions, the large variation of α is an ongoing topic of research.^{140,209-213}

IX. Volcano Plots with Broad Isotherms

Volcano plots have been discussed since Parsons and Gerischer’s independent predictions in 1958^{72,214} and early experimental tests by Conway and Bockris,²¹⁵ Ruetschi and Delahay,²¹⁶ and Trasatti.²¹⁷ In 2005, Nørskov and coworkers built a volcano plot from experimental rates and computed metal–H binding energies, using “a 2 × 2 surface cell [of] a three layer fcc(111) slab.”⁹⁸ Calculations at hydrogen coverages of $\theta = 1/4$ and $\theta = 1$ both gave clear volcano-shaped plots. That work has been extended in many directions.²⁷

Most modern volcano plots have assumed regular surfaces and Langmuir isotherms, sometimes suggesting ‘universal’ relationships.^{27,90,218-220} Still, the simplifications inherent in volcano plots have often been examined (*cf.*,^{93,107,221,222}). In the context of this perspective, the variations among step edges, defects, and other surface heterogeneities can be highlighted. Different sites often have often have different BEP relations (different α and $k_0/\Delta G_0^\ddagger$, eqs 1, 2).^{77,90,103,112,113,223,224} Thus, inhomogeneity and non-ideality should be relevant to volcano plots.

In his original paper, Parsons examined how the shape of a volcano plot for HER would change with a non-ideal isotherm for H binding.⁷² He solved algebraically for the exchange current i_0 as a function of the standard state binding energy ΔG° (written as Δg° in Figure 9). i_0 is the extrapolated current in each direction at zero driving force (similar to the j_0 above). The standard state was taken to be $\theta = 1/2$, and a Temkin isotherm was chosen for mathematical convenience. In this analysis, the different values of Δg° implicitly refer to different catalysts that are “similar” in the ways discussed above. The conclusions were summarized with the text below (slightly edited) and the drawings in Figure 9.

... the effect of assuming a Temkin isotherm on the curve of $\log(i_0)$ against ΔG° is to insert a horizontal region in the curve at its maximum as shown in fig. 8B.... Experimental studies of adsorption of gaseous hydrogen suggest that this horizontal section may be about 7 kcal mol⁻¹ (0.3 eV) long.⁷²

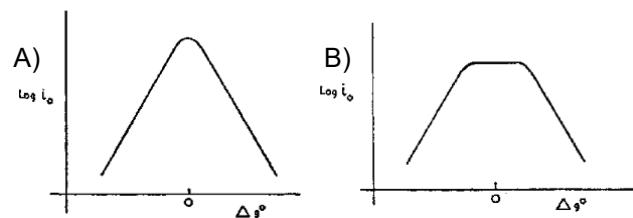


Figure 9. “Form of the relation between exchange current at a hydrogen electrode [$\log(i_0)$] and the standard free energy of adsorption of hydrogen on the electrode surface [ΔG^0 , labeled as Δg°], assuming that the adsorbed atoms (A) obey a Langmuir adsorption or (B) obey Temkin’s logarithmic adsorption isotherm.” Reproduced from *Trans. Faraday Soc.*, 1958, 54, 1053-1063; with permission from the Royal Society of Chemistry; permission conveyed through Copyright Clearance Center, Inc.

Parsons’ estimation of a 7 kcal mol⁻¹ (0.3 eV) width for the flat region in Figure 9B is a substantial fraction of the total width of a typical volcano plot. This conclusion is particularly significant when the flat section includes the expected Sabatier optimum. The width is roughly speaking the range of surface–H BDFEs, which is *ca.* 7 kcal mol⁻¹ (0.3 eV) for a flat metal surface in contact with a gas or liquid phase. For instance, the width of just UPD-H binding on Pt(111) is about 0.3 eV (Figure 4), similar to the width for supported Pt nanoparticles + H₂(g).⁷⁴ The measured isotherm for H on CeO₂ nanoparticles is almost twice as wide. The wider and more irregular isotherms for more complex catalysts should give larger changes to the volcano shape. This insight from 1958 should be part of discussions in modern catalysis.

X. Case Study: H on cobalt phosphide and its volcano plot

i. Experiments and theory for CoP

Transition metal phosphides (TMPs) are being explored as substitutes for noble metals in reductive catalysis such as the HER.^{225,226} TMPs are not simple materials, typically with multiple stoichiometries and structures. As stated by Cossairt and co-workers in their recent Perspective,⁴⁷

“The structural complexity of TMP surfaces introduces several H adsorption sites with varying [BDFEs].”⁴⁷

The stoichiometric reaction of H₂ with cobalt phosphide (CoP)²²⁵⁻²²⁷ has been explored.²²⁸ Acid-washed, high surface area CoP adds hydrogen to give a highly hydrogenated material (*ca.* CoP(H)_{0.3} *bulk* stoichiometry). Titrations with various molecular reagents showed the [CoP]–H UPD-H isotherm to be over 17 kcal mol⁻¹ (0.7 eV) wide.

These experiments confirmed the results of a prior computational study of H on CoP.²²⁹ That study found a variety of H adsorption sites, on various types of Co and P surface atoms, with occupancies and energies depending crystal face and coverage (Figure 10; panel e shows only the energies for H on the CoP(111) surface). In retrospect, the metal-oxide materials in Section VII are a simpler case because usually H binds only to oxygen.

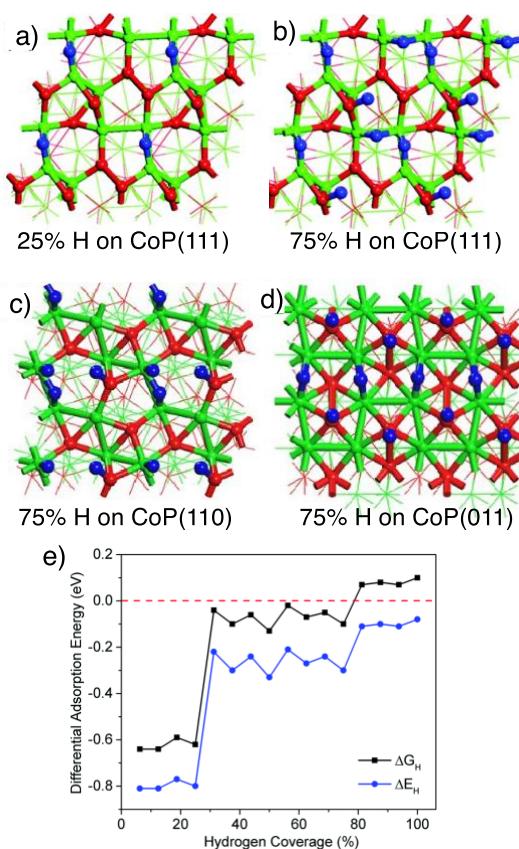


Figure 10. a-d) Sites for H (blue balls) on various facets of CoP and different coverages, as indicated. Co = green; P = red. e) Differential adsorption energy (ΔE_H) and adsorption free energy (ΔG_H) as a function of hydrogen coverage on CoP(111) [BDFE = $\Delta G_H + 2.23$ eV]. Reproduced from *Phys. Chem. Chem. Phys.*, 2008, 10, 3802-3811;²²⁹ with permission from the Royal Society of Chemistry; permission conveyed through Copyright Clearance Center, Inc.

ii. CoP in the volcano plot for HER by transition metal phosphides

Figure 11 places the experimental isotherm as a burgundy-colored bar over the volcano plot developed for the HER by TMP. Additionally, the experiments do not probe the $\Delta G_H > 0$ region where binding energies are unfavorable. The computed isotherm in Figure 10e for the CoP(111) surface is similar to the experimental one.²²⁹

Nørskov and co-workers were well aware of the complexity of TMPs when they assembled this volcano plot, and they took an interesting approach to address this complexity. As detailed in their SI (edited for brevity with italics added):²³⁰

[For] 2~5 surfaces of each TMP, the ΔG_H was calculated for hydrogen coverages from $\theta_H = 0$ ML to 1 ML. Here, θ_H is defined as the fraction of a monolayer with respect to the number of active sites on the surface. The number of active sites on the surface is defined as number of possible sites where hydrogen can adsorb at a reasonable binding energy ($\Delta G_H < 0.7$ eV). The coverage at which hydrogen desorption becomes more favorable than further adsorption is taken to be the coverage where HER occurs. *For each TMP, the surface with the most thermoneutral ΔG_H was considered to be the most active amongst the stable surfaces; the ΔG_H at this surface is thus taken to be the descriptor for that TMP.*²³⁰

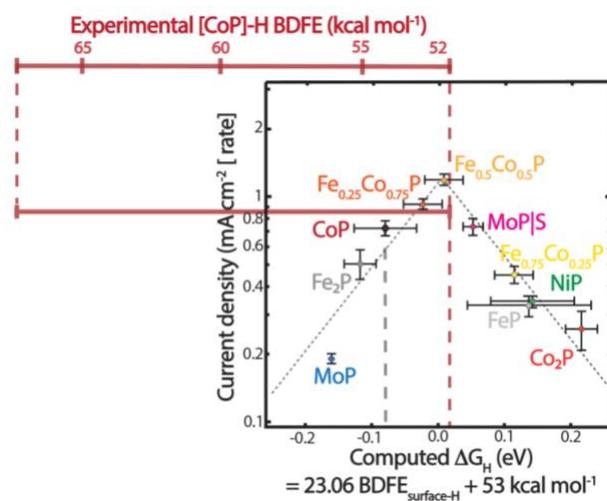


Figure 11. Composite image adapted and reprinted with permission from *J. Am. Chem. Soc.*, 2023, 145 (13), 7050-7064.¹¹⁷ Copyright 2023 American Chemical Society. Experimental data (burgundy horizontal bars)²²⁸ are plotted over a square plot of an activity volcano for the HER for TMPs, current density at $\eta = 100$ mV *vs.* computed ΔG_H (Reproduced from *Ener. Environ. Sci.*, 2015, 8 (10), 3022-3029,²³⁰ with permission from the Royal Society of Chemistry; permission conveyed through Copyright Clearance Center, Inc.). The single ΔG_H shown for each material derive from a large set of computations.

The italics were added to the quote above to emphasize that this approach used the Sabatier Principle as the sole arbiter of which sites were most active. Likely this was chosen in part for computational convenience, but it seems to reflect the intuition of the field as well: only the ΔG° is relevant. This assumption was made even though it encompassed different active site structures. The points in Figure 11 include four active sites with H in P-atop sites, three with H in M-H-M bridge sites, and one with H in a three-fold metal hollow site, M_3H .²³⁰ It goes against my inorganic chemist's intuition to correlate these sites with a single LFER, just because these sites have $\Delta G^\circ_{\text{H}_2}$ closest to 0 for their material.

As discussed in Section IV above, scaling relationships are required to simplify a traditional LFER with three parameters to an equation that depends only on ΔG° .²⁷ For TMPs, a common scaling relation requires that the various PH and M_nH have the same reactivity when their $\Delta G^\circ_{\text{H}}$ values are the same. I know of no precedent in molecular chemistry for such similarity between transition metal and non-metal compounds.

More generally, scaling relationships typically do not hold for more complex catalysts. An early computational study by Abild-Pedersen *et al.* showed that scaling relations for binding of EH_x fragments often differ by 0.5 eV or more between close-packed vs. stepped single crystal surfaces of the same element.¹⁰³ H adsorption on atop sites is ≥ 0.3 eV weaker than adsorption on the most stable site on the same transition-metal surface. This is one explanation for OPD and UPD-H on Pt(111) not falling on the same LFER. Greely later summarized the variability of scaling relationships.¹¹³

These results imply that careful tuning of nanoparticle shape [of ideal transition metal crystals] could, in fact, change the scaling relationships and thus alter the optimal metal

catalyst The structure sensitivity of the scaling relationships thus adds an extra dimension....¹¹³

While this manuscript was in revision, two papers were published that incorporated Frumkin isotherms into volcano plots for OER by cobalt oxide³¹ and iridium oxide.⁵² These studies built on prior experimental work (e.g., that in Section VII.iii¹⁷³) and concluded (italics added):

By combining predictions from density functional theory with parameters obtained from electroadsorption isotherms, we demonstrate that a destabilization of catalytic intermediates occurs with increasing coverage.³¹

[The Frumkin isotherm] (~ 0.19 eV/monolayer) may account for a significant fraction of the observed OER catalytic activity in both [cobalt] materials."³¹

... we present a modified volcano plot that elucidates how the intrinsic water oxidation kinetics can be increased by optimizing both the binding energy *and the interaction strength between the catalytically active states.*⁵²

This Perspective generalizes these quotes to argue that understanding catalysis on complex surfaces requires more than the ΔG° -based Sabatier optimum. Using LFERs requires at least some ability to sort the different sites, some measurements of the Brønsted α and ΔG_0^\ddagger parameters, and some knowledge of the isotherms for each site. When multiple distinct sites contribute to the catalysis, each site should probably have its own point on the volcano.

XI. Paths forward

The non-idealities of nanoscale interfaces presented here are not complexities to be avoided but rather opportunities for new advances in catalysis. This Perspective encourages the increasingly powerful experimental and computational tools in this field to be more frequently focused on the complexity of modern catalysts.

A central theme of this Perspective is that parameters other than the binding energy are important. Ideal LFER treatments and volcano plots using only ΔG° as a descriptor have been very valuable and impactful, but better models and intuition are needed for more complex catalysts.

That binding energies are not a sufficient descriptor is not a new conclusion. This has long been evident for single materials as “surface sensitivity.”^{112,231} Computational studies by Abild-Pedersen *et al.* showed the variety of scaling relationships for different types of single crystal metal surfaces almost 20 years ago.¹⁰³ Yet the Sabatier Principle is taken as the sole guide to catalytic activity in many current papers (though certainly not all). This Perspective aims to change that intuition.

i. Determining Isotherms and Kinetic Brønsted α 's.

The quantitative hypothesis advanced here is that different sites follow different scaling relationships because of their variety of Brønsted α s and intrinsic barriers ΔG_0^\ddagger (eq 2).

The first step in testing this hypothesis—and in learning how to take advantage of site differentiation and non-idealities—is to directly measure the binding isotherms, α s and ΔG_0^\ddagger s. This will require a focus on stoichiometric reactions of surfaces, in addition to overall catalytic metrics. There are some measurements and computations of this kind (*cf.*,^{163,130,206-232}) but not yet enough to derive general

principles. Studies of stoichiometric reactions of complex nanoparticles will in effect average over all of the variety sites on the surface. This will be a valuable complement to single-crystal studies.

Methods are increasingly available to measure isotherms, and this is made easier for nanoscale catalysts because of their high surface areas. Colloidal nanoparticles can often be studied by the classical equilibrium and kinetic methods of solution chemistry. Electrochemistry provides powerful and highly sensitive techniques to monitor adsorbates on electrodes (with care to distinguish faradaic from capacitive currents).^{233,234} My laboratory is examining stoichiometric reactions with molecules in solution, as described above and in²³⁵⁻²⁴⁰. Measuring isotherms and rate constants at different coverages generate LFERs. The shape(s) and slope(s) of the LFER(s) should indicate the number of kinetically distinct sites within the isotherm, taking “similar” sites as those that fall on the same LFER.

Operando, in situ, and ex situ studies of reactive intermediates are increasingly common and are very informative. *Operando* measurements predominantly examine species that accumulate on a catalytic surface,²⁴¹ so some intermediates will need to be studied *ex situ*. Surface intermediates can likely be prepared using organic chemical approaches, as illustrated by the use of CH₃I in surface science to form surface methyl groups (e.g.,²⁴²). As one possible example, organic methods to deliver a formyl group²⁴³ could perhaps form this important surface species and allow studies of its chemistry. Time-dependent applied potentials can also select different surface intermediates in favorable cases.^{244,245}

Theory, computation, and experiment need to go hand-in-hand in characterizing adsorbate structures, energetics, and reactivity as a function of coverage. Complex, dynamic, non-stoichiometric surface structures can be challenging for theory, even to know where to start. Machine learning and related approaches are promising approaches to bridge the “complexity gap,”⁵⁶ and they will be helped by experimental datasets.

ii. “Seeing” the Surface: Spectroscopy, Microscopy, and other tools.

The characterization and study of surface intermediates is not a new quest, but it is increasingly accessible with the advent of increasingly powerful tools.^{241,246} These include x-ray techniques; optical, vibrational, and other spectroscopies; electron microscopies and the expanding experiments that can be done in the microscope;^{16,17} scanning probe microscopies; and high spatial resolution reactivity mapping.^{38,247,248} Continuing challenges include connecting signal intensity with absolute surface coverages and identifying the kinetically competent vs. spectator species. Still, the future is exciting in this area.

Another issue is that most x-ray and electron tools are less sensitive to lighter atoms. The increased use of techniques sensitive to hydrogen and deuterium would be valuable, including TOF-SIMS, atom-probe tomography,²⁴⁹ Rutherford back scattering,²⁵⁰ and neutron experiments. Vibrational spectroscopy with isotopic substitution can be very powerful.²⁴⁴ Still, the spectra can be overlapping, high frequency X–H stretching modes are not easily seen by

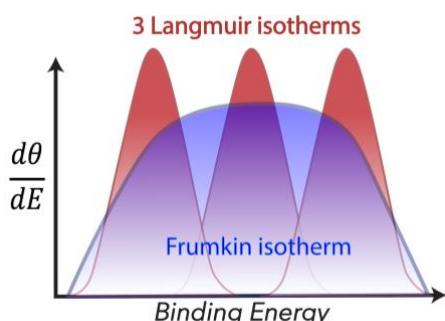
Raman and multidimensional techniques, and O–H and N–H modes are often broadened by hydrogen bonding. Still, great results are increasingly being obtained.²⁴¹ Imagine the insights that could be gained from an ability to characterize and quantitate $-\text{CO}_2\text{H}$, $-\text{CH}_2\text{O}$, $-\text{CHO}$, and/or $-\text{CH}_x$ groups on the surfaces of CO_2 reduction catalysts.

iii. Modeling Complex Isotherms

For surface diversity and non-ideality to become important parts of basic catalysis discussions, simple and accessible models are needed. Historically, there have been two limiting simplifications. One can assume that all the different sites behave similarly, can be treated as a Langmuir isotherm, and follow a single LFER. As noted above, this can be sufficient for traditional catalysts.¹⁹⁶ At the other extreme, the sites can be considered to be very different and only one type is catalytically active—Boudart's most abundant reactive intermediate (MARI).^{114,115} Since modern catalysts are not well modelled by either limit, an intermediate model seems appropriate: approximating a broad and/or irregular isotherm as the sum of Langmuir isotherms (Scheme 4). Various multi-Langmuir models have long been used, for different applications (cf.,^{201,251–255}). As Winterlin, Ertl, *et al.* wrote about NO dissociation on a Ru(0001) surface:¹⁰

For "real" catalysis ... the overall reactivity will be the result of weighted contributions from various surface structure elements, dominated by the active sites.¹⁰

Scheme 4. Modelling a Frumkin isotherm (blue) as the sum of three Langmuir isotherms (burgundy).



The multi-Langmuir model is intuitive when the broadness of the isotherm comes from distinct surface sites. Each component Langmuir isotherm is then treated as ideal, removing the complexities of non-ideal activities, rate 'constants', free energies, *etc.* Scheme 4 shows a Frumkin isotherm modeled as three ideal adsorbate states, each with its own coverage θ_i , saturation coverage $\theta_{i,\text{max}}$, standard free energy of adsorption, ΔG_i° , and rate constant k_i .

On complex, ill-defined surfaces of nanoscale, multi-element catalysts, the surface structure elements and active sites are not known. Flipping the perspective of the quote above, the combination of a measured isotherm and catalytic activity as a function of surface could be a powerful way to analyze the surface elements and their activities.

Applying Scheme 4 to H on ceria NPs (Section VII.i), the separation of the three Langmuir isotherms would be ~ 6 kcal mol⁻¹ (0.26 eV). With a Brønsted α of 0.5, this range of BDFEs would predict a variation in rate constants of $\sim 10^4$ between the outlying isotherms; with the experimental $\alpha =$

0.2,¹⁶³ the range of k 's among these groups is a factor of 15. As a more complex example, consider the Tafel reaction combining two surface H's ($^*\text{H}$) to form H_2 (Figure 2 left). With one kind of $^*\text{H}$ well described by a single Langmuir isotherm, the rate law is simple second order (eq 20, with rapid surface diffusion). When there are i different sites, H_2 can be formed from two $^*\text{H}_i$ or from different sites, $^*\text{H}_i$ plus $^*\text{H}_j$ (eq 21).^{256,257} The latter mechanism was suggested for H_2 desorption from Mo_2NH_x : a hydridic $\text{Mo}-\text{H}^{\delta-}$ plus a protic $\text{N}-\text{H}^{\delta+}$.¹⁹⁸

$$\frac{d[\text{H}_2]}{dt} = k(\theta_{^*\text{H}})^2 \quad (20)$$

$$\frac{d[\text{H}_2]}{dt} = \sum_i k_{ii} (\theta_{^*\text{H}_i})^2 + \sum_{i,j} k_{ij} (\theta_{^*\text{H}_i})(\theta_{^*\text{H}_j}) \quad (21)$$

Electrochemical reaction steps can be similarly modeled, with multiple E_i° analogous to the ΔG_i° above. In a multi-Langmuir model, the Butler-Volmer equation (BV) becomes a sum of terms, each exponential in their $(E_{\text{applied}} - E_i^\circ)$, and the current should not be simply exponential with potential. The presence of non-ideal isotherms is one likely reason why empirical Tafel slopes are typically not the integer multiples of $(30 \text{ mV/decade})^{-1}$ predicted by simple treatments.²⁵⁸

XII. Final thoughts

The increasing complexity of modern catalyst materials prompts a renewed emphasis on the non-ideal behavior of adsorbed catalytic intermediates. Compared to traditional metallic catalysts, nanoscale, multi-element, reactive catalysts often have a wide variety of surface sites, structures, stoichiometries, thermochemistry, adsorbate interactions, mechanisms, and reactivity. This diversity is a strength of these catalysts, a strength that should be harnessed to advance catalysis science and engineering.

Boudart's "three-step approach to surface kinetics" in *Kinetics on Ideal and Real Surfaces* (1956) included:⁶

[In the] third step The nonideality of the surface process, including heterogeneity and interactions, must then be taken into account either by postulating a surface distribution of energy sites or by introducing real adsorption isotherms determined separately. While the third step ... is hardly feasible in the vast majority of cases, there are a number of reasons for going beyond the first step whenever possible.

Now, 68 years later, it is time to tackle and take advantage of the non-ideality of complex catalysts.

ASSOCIATED CONTENT

There is no Supporting Information for this Perspective.

AUTHOR INFORMATION

Corresponding Author

* james.mayer@yale.edu

Department of Chemistry, Yale University, 225 Prospect St.,
New Haven CT 06520

Orchid ID: 0000-0002-3943-5250

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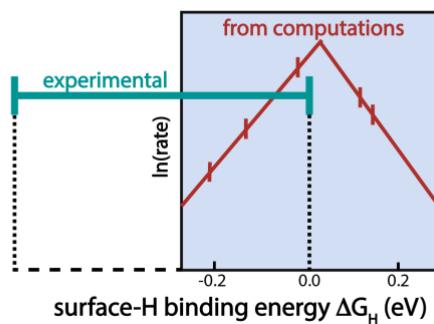
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TOC GRAPHIC

Volcano plot for HER by transition metal phosphides with experimental values for CoP
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