



Letter

Upper bound on the speed of sound in nuclear matter from transport

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ABSTRACT

We point out that there is an upper bound on the speed of sound squared given by $c_s^2 \leq 0.781$ valid for all known systems described by relativistic transient hydrodynamics where calculations of certain ratios of hydrodynamic transport coefficients can be performed from first principles. Assuming this bound is valid for ultradense matter implies that the maximum mass of isolated (non-rotating) neutron stars cannot be larger than 2.7 solar masses.

1. Introduction

The measurements of neutron stars with masses $\gtrsim 2 M_\odot$ (where $M_\odot$ is the mass of the sun) [1–3] suggest that the cold neutron star equation of state (EoS) is stiff enough such that its speed of sound, c_s , surpasses its conformal limit, $c_s \leq 1/\sqrt{3}$, at a given density [4]. This is also supported by the binary neutron-star merger gravitational-wave event GW170817 [5–7], which further constrained the masses and effective tidal deformability of the inspiraling neutron stars [8–16]. Further support to a supraconformal speed of sound may come from the observation of gravitational waves from the merger of a black hole with a compact object of mass $2.6 M_\odot$ [17], if the latter is a neutron star [18–23].

Due to asymptotic freedom [24,25], the speed of sound in QCD is expected to approach the conformal limit from below at sufficiently large densities [26,27]. This fact, combined with the knowledge about the nuclear physics equation of state at low densities [8,28–30] and the recent observation of large neutron star masses, implies that c_s in QCD is not monotonic as this quantity must display at least one peak as a function of density (when taking into account the region of asymptotically large density). Assuming there is only a single peak, essential questions are: (i) How large is c_s at the peak? (ii) At what density does c_s peak? (iii) What is the correct effective theory of ultradense matter that explains the answers to (i) and (ii)? Ultimately, future observations [31–34] will constrain the answer to (i) and (ii) [35–41] and provide useful guidance towards answering (iii).

Relativistic causality and covariant stability [42] impose that the speed of sound cannot surpass the speed of light, i.e., $c_s \leq 1$. At high temperatures and zero baryon chemical potential, lattice QCD calculations [43,44] and holographic models [45–48] find that the conformal limit for the speed of sound is respected, $c_s \leq 1/\sqrt{3}$. At nonzero baryon density, however, discussions concerning the (subluminal) upper bound to the speed of sound go back several decades [49–52]. It is now known that the speed of sound can surpass the conformal result in a variety of systems such as QCD at large isospin density [53], two-color QCD [54], holographic models [55–59], resummed perturbation theory [60,61], quarkyonic matter [62–66], and other models at high density [67–82]. However, despite recent progress [83,84], very little is known about the transport properties of ultradense matter when the speed of sound nears the speed of light.

In this work, we point out that ultradense matter with nearly luminal speed of sound must have very unusual transport properties. We discuss a scenario where fundamental properties of thermodynamics and transport in relativity can constrain the equation of state of ultradense matter. Assuming the validity of relativistic transient hydrodynamics in which dissipative stresses obey additional relaxation equations [85], we argue that causality and stability of matter imply that there is an upper bound on the speed of sound squared given by $c_s^2 \leq 0.781$ in all known systems¹ where calculations of certain ratios of first-order hydrodynamic

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¹ We only consider systems where the equilibrium state is unique and the correlation length is finite.

transport coefficients and their corresponding relaxation times can be performed from first principles. In this case, a violation of the bound implies that such systems either exhibit very unusual transport coefficients or cannot be consistently described by transient fluid dynamics,² signaling the presence of exotic transport behavior that simple relaxation dynamics cannot describe.³ Given current astrophysical constraints, we show that imposing this bound on c_s implies that the maximum mass of an isolated neutron star is bound by $M < 2.7 M_\odot$, assuming chiral effective theory is valid up to two times saturation density. Therefore, the new transport bound provides further support to the claim that the $2.6 M_\odot$ compact object observed to merge with a $23 M_\odot$ black hole causing the GW190814 event [17] can be a neutron star.

2. Speed of sound bound from transport

Fluid dynamics describes the long-time, long-wavelength dynamics of conserved quantities in many-body systems [91]. Ideal fluid dynamics describes the evolution of the energy density ϵ , number density n , and flow velocity in the absence of dissipation, with the information about the system's microscopic properties encoded in the equation of state, e.g., $P = P(\epsilon, n)$, where P is the pressure computed in thermodynamic equilibrium. In this limit, sound waves propagate with frequency $\omega = \pm c_s k$, where $c_s = \sqrt{dP/d\epsilon}$ denotes the isentropic speed of sound in the fluid, and $k = |\mathbf{k}|$ is the wavenumber.

Small dissipative corrections can be taken into account by expressing the dissipative fluxes in a systematic derivative expansion of the hydrodynamic fields [91], which introduces the well-known transport coefficients: the shear viscosity η , bulk viscosity ζ , and charge conductivity σ .

Constraining any transport coefficients in nuclear matter is challenging since it involves extracting near-equilibrium information at strong coupling from QCD. However, imposing causality can provide useful information about transport coefficients when the speed of sound nears the speed of light.

In causal and stable first-order theories [92–96], the long-wavelength behavior of the sound wave dispersion relation is given by⁴ $\omega = \pm c_s k - i\Gamma k^2/2 + \mathcal{O}(k^3)$, where the damping coefficient is

$$\Gamma = \frac{\zeta + \frac{4}{3}\eta}{\epsilon + P} + \frac{\sigma}{(\epsilon + P)c_s^2} \left(\frac{\partial P}{\partial n} \right)_\epsilon^2. \quad (1)$$

Causality [97] and the stability of the equilibrium state [42] impose that dispersion relations must obey the fundamental inequality $\text{Im } \omega \leq |\text{Im } k|$, $\omega, k \in \mathbb{C}$. Since the second law of thermodynamics imposes that η, ζ, σ are non-negative, $\Gamma \geq 0$, one can show⁵ that $\Gamma \rightarrow 0$ when $c_s \rightarrow 1$ [42]. This statement was generalized in Ref. [98], where it was shown that the contribution from all the other terms also vanishes, making luminal sound waves exactly non-dispersive and described by the trivial dispersion relation, $\omega = k$. This implies that at densities where the speed of sound of ultradense matter approaches the speed of light, sound waves propagate without damping or dispersion.

In relativistic transient fluid dynamics [85], transport is further characterized by relaxation time coefficients that determine when the dissipative fluxes approach their first-order values. For instance, the shear

relaxation time coefficient τ_π has been calculated for many theories (see, e.g., Tab. 2.1 in Ref. [99]). In all of these systems, it was found to be tightly constrained as

$$\tau_\pi = C_\pi \times \frac{\eta}{\epsilon + P}, \quad C_\pi \in [2(2 - \ln 2), 6.1], \quad (2)$$

with the upper limit of C_π originating from theories evaluated at *weak coupling*⁶ [101]. We, therefore, find the following bound

$$\frac{4\eta}{3\tau_\pi(\epsilon + P)} = \frac{4}{3C_\pi} \geq \frac{4}{3 \times 6.1}, \quad (3)$$

from relativistic transport in theories where calculations can be done from first principles.

By contrast, while the bulk relaxation time τ_Π has also been calculated in many theories, the corresponding bound on the ratio of $\frac{\zeta}{\tau_\Pi(\epsilon + P)}$, where ζ is the bulk viscosity, is much less constrained,

$$\frac{\zeta}{\tau_\Pi(\epsilon + P)} = \frac{1}{C_\Pi} > 0, \quad (4)$$

because there are examples where $\zeta = 0$, $\tau_\Pi \neq 0$ [102]. Finally, not much is known from first principles about similar ratios involving the charge conductivity and the relaxation times, and thus, we shall not consider it further in this work.

Further general statements about transport can also be made in the context of transient fluid dynamics. By requiring that the entropy measured by all inertial observers is maximized in equilibrium in the presence of constraints (such as energy and/or particle number), one can derive a number of thermodynamic inequalities that ensure that the stability of the equilibrium state holds in all Lorentz reference frames [103]. These inequalities also imply causality in the linear regime around equilibrium [104]. In transient fluid dynamics, one can obtain such thermodynamic inequalities by computing the so-called information current E^μ [104], which must be a time-like future-directed 4-vector with non-positive divergence, $\nabla_\mu E^\mu \leq 0$, to satisfy the covariant version of the second law of thermodynamics. In the case of Israel-Stewart theory with shear and bulk viscosity, this implies the following thermodynamic inequality [105,106]

$$c_s^2 \leq 1 - \frac{4}{3C_\pi} - \frac{1}{C_\Pi}, \quad (5)$$

which constitutes an upper bound on the speed of sound⁷ from relativistic thermodynamic principles⁸. Furthermore, we see that matter in a state where the equilibrium speed of sound $c_s \rightarrow 1$ must have unusual transport behavior as not only η and ζ but also their corresponding ratios involving the relaxation times must vanish.

It should be stressed that the bound (5) relies on the assumption that transient fluid dynamics, where additional degrees of freedom (such as the shear and bulk stresses) are included as independent variables characterizing an extended quasi-equilibrium thermodynamic state [109], faithfully represent the thermodynamic properties of nuclear matter. This nontrivial assumption implies that the near-equilibrium properties of ultradense matter display not only typical fluid behavior but also visco-elastic properties, which are commonly found in liquids [110–112].

Using the additional information (3, 4) on C_π, C_Π , one can make the bound (5) sharp:

$$c_s^2 \leq 0.781. \quad (6)$$

² This is not connected to the non-analytic behavior induced by the backreaction of sound waves [86].

³ While unusual transport coefficients can be expected from superfluid phases, the onset of superfluidity leads to multiple sound modes [87], which makes causality and stability analyses considerably more convoluted [88]. Therefore, we leave the exploration of superfluid states to future work. However, we note that our bound holds even in the case of a pion condensate, where a global $U(1)$ symmetry is spontaneously broken [53,89,90] (see Fig. 1).

⁴ The $\mathcal{O}(k^3)$ terms are determined by the UV regulators parameterizing the choice of hydrodynamic frame.

⁵ Take $k = i\alpha$ (with $0 < \alpha \ll 1$) and impose the inequality $\text{Im } \omega \leq |\text{Im } k|$.

⁶ A strict bound $C_\pi \geq 5$ has been proved in Ref. [100] for all weakly coupled theories, along with a pioneering calculation of τ_π to NLO in perturbative QCD.

⁷ We only consider the equilibrium speed of sound; for a definition of a far-from-equilibrium speed of sound, see [107].

⁸ We note that (5) also follows from applying the inequality $\text{Im } \omega \leq |\text{Im } k|$, or by simply computing the maximum propagation speeds in linearized Israel-Stewart-like theories [108, Eq. (49)].

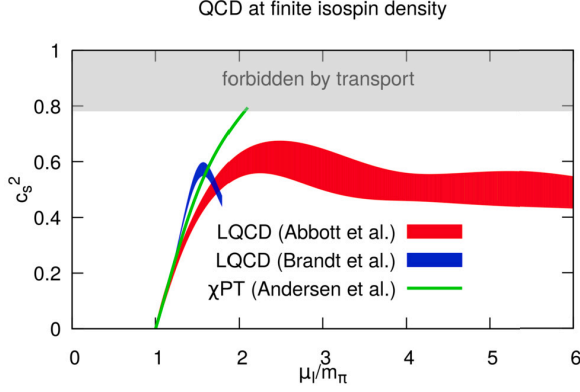


Fig. 1. Speed of sound squared at $T = 0$ in QCD as a function of isospin chemical potential μ_I . Shown are results from lattice QCD (LQCD) for $N_f = 2 + 1$ from Refs. [114,115] with $m_\pi \simeq 170$ MeV; LQCD for $N_f = 2 + 1$ from Refs. [116,117] with $m_\pi \simeq 135$ MeV; and chiral perturbation theory (χ PT) results for $N_f = 2$ to NLO with $m_\pi \simeq 135$ MeV [118]. The transport bound (6) is indicated by the grey region marked as “forbidden by transport”.

If the bound above is violated, this indicates either that the shear viscous coefficients evade Eq. (3) somehow, or that transport cannot be studied using the framework of transient hydrodynamics without violating fundamental thermodynamic principles. This, per se, indicates the presence of nontrivial transport dynamics in ultradense matter that cannot be captured within well-known relativistic hydrodynamic formulations [113] where simple relaxation equations govern the evolution of the dissipative fluxes.

In the next section, we compare this bound to results in theories where c_s^2 can be calculated (or estimated) from first principles.

3. Speed of sound in nuclear matter

Results for the speed of sound squared for cold nuclear matter are available from first-principles QCD calculations such as lattice QCD, effective field theory, and perturbative QCD. In particular, results for c_s^2 are shown in Figs. 1, 2, 3 for QCD at finite isospin chemical potential, QCD at finite baryon chemical potential, and two-color QCD at finite chemical potential, respectively. In addition to the results shown, results on thermodynamics at finite-density for QCD in the strong coupling limit from Ref. [119] and various model calculation results for c_s^2 are available [120,121].

Comparing the results for c_s^2 from first principles calculations shown in Figs. 1, 2, and 3 to the transport bound (6), one finds that the transport bound on the speed of sound is obeyed in all cases.

It is worth noticing that the bound in Eq. (6) is respected even for low temperature where superfluidity occurs, [53,89,90,122], as can be seen in Figs. 1, 2.

4. Implications for the neutron star mass – radius relation

The speed of sound of ultradense matter plays a crucial role in the structure of neutron stars. A larger c_s^2 allows for larger values of the stellar mass — as larger pressure gradients can counterbalance a larger gravitational force — but also gives rise to larger stars — as matter becomes harder to compress. Observations of “small” neutron star radii of $R \sim 13$ km [127–130], as well as tidal deformability constraints from GW170817 [5–7], require speeds of sound well below the conformal value of $c_s^2 \leq 1/3$ at low densities, while very massive stars of $M \sim 2 M_\odot$ [1–3] require much larger values of c_s^2 at high density, most likely above $c_s^2 = 1/3$ [4,8,131–134].

We thus investigate the implications of the transport bound $c_s^2 \leq 0.781$ for the neutron star mass–radius relation. To find the most massive configurations allowed by this bound, we employ EoSs that saturate

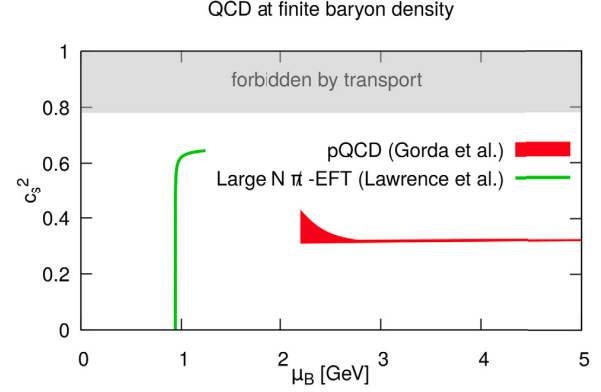


Fig. 2. Speed of sound squared at $T = 0$ in QCD as a function of baryon chemical potential μ_B . Shown are results from perturbative QCD (pQCD) for $N_f = 3$ to partial N3LO Refs. [123–125] and large N_f EFT results to LO from Ref. [122]. The transport bound (6) is indicated by the grey region marked as “forbidden by transport”.

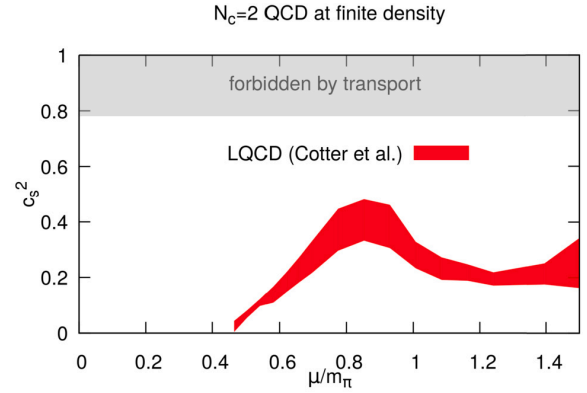


Fig. 3. Speed of sound squared at $T \simeq 50$ MeV in two-color QCD as a function of chemical potential μ . Shown are results for c_s^2 extracted from data provided in Ref. [126] with $m_\pi \simeq 700$ MeV. The transport bound (6) is indicated by the grey region marked as “forbidden by transport”.

this bound above a threshold $n > n_{\text{stiff}}$ for the baryon density. At densities $n < n_{\text{stiff}}$ below this threshold, we generate ~ 200 low-density EoSs consistent with uncertainty bands from chiral effective theory (χ EFT) results [8], combined with a GPPVA(TM1e) crust EoS [135–139] (we refer the reader to the supplementary material for details). Most recent constraints from multiple sources on the symmetry energy in Ref. [140] suggest that the χ EFT uncertainty does not cover the full range of posterior values reported in [140]. In particular, the authors of Ref. [140] find that considerably larger symmetry energy values are allowed, which would lead to higher neutron star masses than when considering χ EFT. Reconciling the uncertainty bands between χ EFT and Ref. [140] is beyond the scope of this work, so we choose to use the χ EFT results in the following. We match each of these EoSs $\epsilon_{\text{soft}}^{(i)}(P)$, where superscript i specifies which EoS, to a high-density EoS $\epsilon_{\text{stiff}}^{(i)}(P)$ with constant speed of sound $c_s^2 = c_{s(\text{bound})}^2 = 0.781$, while keeping ϵ a continuous function of P :

$$\epsilon_{\text{stiff}}^{(i)}(P) = \epsilon_{\text{soft}}^{(i)}(P_{\text{stiff}}) + \frac{P - P_{\text{stiff}}}{c_{s(\text{bound})}^2}, \quad (7)$$

where P_{stiff} is the value of the pressure as $n \rightarrow n_{\text{stiff}}^-$ from the left.

We then employ the Tolman–Oppenheimer–Volkoff (TOV) equation to obtain mass–radius relations for each of these piecewise equations of state. Similar procedures have been extensively employed in the liter-

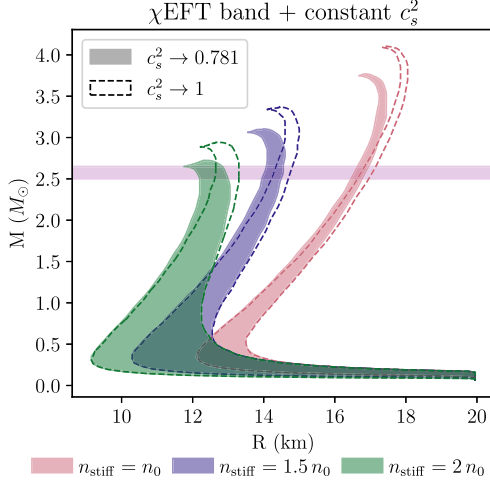


Fig. 4. Mass-radius relations for EoS that are in agreement with χ EFT results at baryon densities $n < n_{\text{stiff}}$ and saturate the new bound $c_s^2 \leq 0.781$ at densities $n > n_{\text{stiff}}$. Different shaded bands represent results for switching densities n_{stiff} of 1 (pink), 1.5 (purple), and 2 (green) times the baryon density at nuclear saturation, with the band thickness reflecting uncertainties in χ EFT results. Dashed lines represent results in the case where $c_s^2 = 1$ saturates the causal bound for $n > n_{\text{stiff}}$, with the same color scheme. The horizontal magenta band indicates the uncertainty in the mass of the compact object of unknown nature observed in GW190814 [17].

ature to investigate the evidence in favor of supraconformal speeds of sound [4,8,10,13,141]. The region in the mass-radius diagram obtained via this process is shown in Fig. 4 as shaded bands. We show results for different values of n_{stiff} , of 1, 1.5 and 2 times the baryon density at nuclear saturation $n_0 \approx 0.16 \text{ fm}^{-3}$.

Note that our constructed EoS are not meant to be realistic models of the nuclear EoS behavior, but instead are meant to check if χ EFT plus transport bounds are relevant for astrophysics.

To gauge the significance of the new c_s^2 bound, we compare the maximum allowed mass for $c_s^2(\text{bound}) = 0.781$ to the maximum mass when the less restrictive bound $c_s^2 \leq 1$ — from relativistic causality [49] — is saturated, which can be found by replacing $c_s^2(\text{bound}) = 1$ in Eq. (7). Contours for the obtained regions in the mass-radius diagram, for each different value of n_{stiff} , are represented as dashed lines in Fig. 4.

We find that the transport bound $c_s^2 \leq 0.781$ decreases the maximum mass by $\sim 8\%$ with respect to the maximum value allowed by the causal bound. This is the case for all the different soft EoS, regardless of the density n_{stiff} at which the EoS becomes stiff, with an uncertainty of 0.2%. The new bound also tends to decrease the radius of the maximally massive configuration by $3.6 \pm 0.3\%$. Effects of the same bound become small for lighter stellar masses, below $\sim 2 M_\odot$.

5. Discussion

Relativistic causality and covariant stability impose strong constraints on transport quantities and the speed of sound in relativistic fluids. Combining these constraints with bounds on transport coefficients observed in a wide range of systems, we have found that, for all of these systems, assuming the validity of transient relativistic fluid dynamics implies that the speed of sound squared is bounded from above by $c_s^2 \leq 0.781$. We investigated the applicability of this bound in first-principles calculations of the speed of sound in QCD and two-color QCD. All of the known first-principles calculations obey the new bound. This bound can be sharpened by further constraining the ratios C_π, C_Π from shear and bulk viscosity (and performing a similar analysis including effects from charge conductivity), especially in systems at finite baryon density.

Not much is known about transport properties beyond first order at low temperature and finite density, the exception being calculations of so-called thermodynamic transport coefficients [122,142–144]. Nonetheless, in Sec. 3, we have checked that our new bound on the speed of sound is respected in first-principles results both for QCD at zero temperature and for two-color QCD at a temperature around 50 MeV. The results discussed in this work can lead to a renewed interest in calculating transport coefficients, particularly the shear and bulk relaxation times, in order to assess to what extent the bounds in Eq. (3) and (4) apply in the cold and dense regime.

It would be interesting to extend our analysis to the case where superfluids may be present, which we leave for future work. While our analysis is not directly applicable in this case, we have seen that our bound on c_s^2 is not violated in first-principles results for QCD at high isospin density, where a superfluid pion condensate is expected [53,89,90], and non-relativistic chiral QCD where the analogues of Cooper pairs form [122].⁹

We also investigated the consequences of this bound for the neutron star mass-radius relation. Our results show that, if χ EFT is taken to be a good description of nuclear matter up to two times saturation density ($n_{\text{stiff}} = 2 n_0$), then a single neutron star mass above $2.7 M_\odot$ (and below $\sim 2.9 M_\odot$) would provide strong evidence for the breaking of the $c_s^2 \leq 0.781$ bound, indicating that ultradense matter must exhibit exotic transport properties. On the other hand, our new bound is fully compatible with the possibility that the $\sim 2.6 M_\odot$ compact object observed via gravitational wave event GW190814 could be a neutron star.

It is worth noting that our proposed bound to the speed of sound seems to be compatible with observation-driven analyses of the neutron-star equations of state, e.g., [41,146–148], which suggest that values of $c_s^2 \gtrsim 0.781$ are statistically disfavored.

A potential follow-up direction to this work is provided by extending our modeling analysis for neutron stars in light of the constraints on the nuclear matter EoS up to three times n_0 given in Ref. [140].

Future Bayesian analyses of neutron star and gravitational wave observations could be employed to more systematically assess the implications of our new bound on c_s^2 for the cold and dense nuclear matter EoS and for the neutron star mass-radius relation. Such analyses could also estimate the probability that this new bound is violated inside neutron stars and, therefore, help establish how likely it is that the matter in the core of massive neutron stars displays unusual transport properties.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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⁹ A superfluid state may also be present in two-color QCD [145], where the bound $c_s^2 \leq 0.781$ also seems to hold.

Appendix A. Supplementary material

Supplementary material related to this article can be found online at <https://doi.org/10.1016/j.physletb.2024.139184>.

Data availability

Data will be made available on request.

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