Correcting dispersion corrections with density-corrected DFT

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Almost all empirical parameterizations of dispersion corrections in DFT use only energy errors, thereby mixing functional and density-driven errors. We introduce density and dispersion corrected DFT (D^2C -DFT), a dual-calibration approach that accounts for density delocalization errors when parametrizing dispersion interactions. We simply exclude density-sensitive reactions from the training data. We find a significant reduction in both errors and variation among several semilocal functionals and their global hybrids when tailored dispersion corrections are employed with Hartree-Fock densities.

dual-calibrated DFT | density functional theory | density-corrected DFT | dispersion corrections | delocalization errors

Significance

Density functional theory (DFT) is widely used in chemistry. Still, even the most popular approximations suffer from two fundamental limitations - they often predict that electrons are more delocalized than they should be and fail to capture weak long-range attractions between molecules. These issues are usually addressed separately by density correction schemes and empirical dispersion corrections. Here, we show that a "dual-calibration" approach, fitting dispersion corrections on density-corrected DFT energies, provides a more balanced and robust solution. The resulting density and dispersion-corrected DFT method improves accuracy across various functionals and chemical systems, particularly those susceptible to density delocalization errors, by simultaneously addressing two major error sources in DFT.

Introduction

Kohn-Sham density functional theory (KS-DFT) 1,2 is a standard method in computational chemistry, striking a balance between computational efficiency and accuracy. However, the quality of results depends on the quality of the approximate exchange-correlation functional used. Density-corrected DFT (DC-DFT) is a generic methodology for separating errors in the functional from errors in the self-consistent density. 5,6 The principles of DC-DFT follow directly from simple (functional) analysis and can be applied (in principle) to every approximate KS-DFT calculation ever run. In practice, for DC-DFT to yield improved energetics in (a subset of) KS calculations, several conditions must be met. In some cases, these conditions can be checked with highly accurate quantum chemical calculations [8,9], while

in others, improved energetics resulting from the procedure suggest the conditions have been met. In either case, DC-DFT has had an impact in a variety of challenging areas for standard DFT calculations. [10]-[24]

Semilocal density functional approximations (DFAs) often fail to capture long-range dispersion interactions accurately due to their local nature. [25] To remedy this, many dispersion corrections have been developed as empirical pair-wise additions to DFT energies, with parameters fitted to databases dominated by weak interactions. 26-31 These corrections are computed for each XC functional and then tested on other databases where dispersion effects vary in magnitude. However, in all these databases, the energy error of the DFA mixes both functional and density-driven errors, which in turn affects the choice of optimum dispersion parameters. This effect is particularly acute for reactions that are density-sensitive, i.e., those whose density-driven errors are large. In some cases, density-driven errors are much larger than dispersion corrections and can hopelessly distort the choice of dispersion parameters. 32

The remedy is very simple. Drop any density-sensitive reactions from the training data, find the optimal dispersion parameters, and then test on all the data using Hartree-Fock (HF) densities. It is known that HF densities typically greatly reduce density-driven errors for density-sensitive reactions and do little (if any) harm for insensitive cases. [5]

To fully assess the impact of both dispersion and density corrections, we apply them to a wide range of XC approximations, aiming to demonstrate the generality of the improvements and the reduction in the variance of predictions across different functionals. Figure I illustrates our results, using the extensive GMTKN55 database for general-main group thermochemistry, kinetics, and non-covalent interactions 33. The left panel takes several popular semilocal XC approximations, showing their improvements upon first

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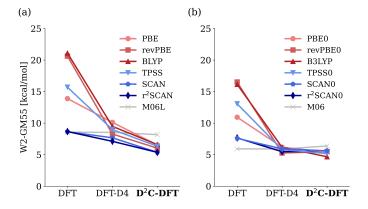


Figure 1: Performance of DFT, DFT-D4, and density and dispersion corrected DFT (D²C-DFT) for several semilocal functionals (a) and their corresponding hybrids (b), measured by the weighted total mean absolute deviation on the GMTKN55 database 33 (W2-GM55). D²C-DFT significantly enhances accuracy and reduces variability across a wide range of functionals. M06 and M06L show no such improvement.

adding dispersion corrections without DC-DFT analysis (the standard approach, DFT-D4) and then adding density corrections using the principles of DC-DFT (D²C-DFT). All but one of the functionals improve significantly with each correction. Moreover, their errors become comparable once both corrections are employed. The glaring exception is M06L. Because of its highly empirical construction, which likely mixes functional and density-driven errors and already includes dispersion, it shows no improvement whatsoever, moving it from among the best performers without correction to the worst after correction. We also point to the right panel, which are global hybrids made from those on the left, where all the same effects occur, but less dramatically, because a fraction of the HF density is already being used.

The current paper presents an example of what we call dual-calibration, a general approach in which empirical parameters in a DFT correction are chosen under DC-DFT principles. Dual-calibration aims to resolve density-driven and functional errors simultaneously, with the specific combination of density correction and dispersion correction being one such example. It was first used to create the BL1p functional [32] and later to create the HF-r²SCAN-DC4 method [21]. The latter, which employs an explicit dispersion formula, has shown remarkable accuracy across various applications, from water complexes to biomolecules, and has been further refined. [34] By studying many different DFAs, the present work demonstrates that D²C-DFT is a general procedure that should be adopted whenever empirical parameters are optimized in DFT calculations.

Background

Density-Corrected DFT

DC-DFT provides a framework for analyzing the energy error in any self-consistent KS-DFT calculation. [5,11] The energy error can be separated into the contribution due to errors in the functional (for a given density, n) and those in the self-consistent density (\tilde{n} , introduced by the use of an approximate XC potential). Using tildes for approximate quantities, we write:

$$\Delta E = \tilde{E}[\tilde{n}] - E[n]$$

$$= \underbrace{\tilde{E}[\tilde{n}] - \tilde{E}[n]}_{\text{density-driven error}(\Delta E_D)} + \underbrace{\tilde{E}[n] - E[n]}_{\text{functional error}(\Delta E_F)}$$
(1)

The errors in most KS-DFT calculations are dominated by the functional contribution, but even very small densitydriven errors can have an impact on overall performance, as we show below.

For theoretical studies of small molecules, accurate densities are sometimes available from quantum chemical methods, such as Coupled Cluster Singles and Doubles (CCSD). However, most practical applications use the HF density as a proxy for the exact one, as its computational cost is often comparable to the corresponding KS-DFT calculation. For cases where density-driven errors are significant, the HF density is usually a good proxy. In such cases, a useful indicator is the 'density sensitivity,' defined as,

$$\tilde{S} = |\tilde{E}[n^{\text{HF}}] - \tilde{E}[n^{\text{LDA}}]|. \tag{2}$$

This density sensitivity, which makes use of HF and LDA densities, provides a crude measure of how sensitive the density in a system is likely to be to the choice of XC functional. If the sensitivity is low, the density-driven error is small in the system. Conversely, if the sensitivity is high, there is a higher probability of significant density-driven error in the system. 35 We have shown that significant density sensitivity suggests that a system is suffering substantial density-driven error, which can be mitigated by employing HF density. In cases where the HF density is a good proxy for the exact density, we expect

$$D[n^{\rm HF}] = \tilde{E}[n^{\rm HF}] - \tilde{E}[n] \tag{3}$$

to be smaller than the density-driven error of the usual calculation. While it can make sense to vary the cut-off for different-sized systems or weaker bonds or define in a new way, 36,37 here, we simply apply 2 kcal/mol to all cases.

Likewise, it has been shown that ROHF (restricted openshell) can perform better than UHF (unrestricted) for DC-DFT calculations. [38] Here, for ease of implementation and to keep computational costs down, we use UHF throughout whenever HF breaks symmetry.

Dispersion-Corrected DFT

Since it is widely acknowledged that (semi)local DFT cannot correctly capture long-range dispersion interaction, [25] various methods to correct dispersion interaction have been developed. These include explicit nonlocal density functionals [39] [41], empirical functionals [42], [43], and additive correction methods [26] [31]. Nonlocal methods are often computationally expensive in codes using atom-centered basis sets, and empirical functionals are known to miss capturing accurate asymptotic behavior. [44]

Considering the cost-accuracy balance, which is also the philosophy of DFT, the practical choice is additive correction. In Grimme and coworkers' additive dispersion model, DFT-D4[26, [27], dispersion energy is defined as,

$$E_{\text{disp}}^{\text{D4}} = \underbrace{-\sum_{AB} \sum_{n}^{6,8} s_n \frac{C_n^{AB}}{R_{AB}^n} f_{\text{damp}}^n(R_{AB})}_{E^{(2)}} + E^{(3)}, \quad (4)$$

where the sums are over all atom pairs A and B separated by R_{AB} , and the D4 model swiftly evaluates two-body $(E^{(2)})$ and three-body $(E^{(3)})$, contributions to the energy based on dispersion coefficients (C_n^{AB}) between atoms using fractional coordination numbers and atomic charges from a structure.

Additive corrections require appropriately fitted global parameters for each functional since the base energy from DFA already contains varying degrees of intermediate-range dispersion depending on the functional used. 45 In its major two-body term, there are global parameters s_6 and s_8 , which directly scale the multipolar contributions (dipole-dipole and dipole-quadrupole interactions, respectively), with $s_6 = 1$ unless it is paired with double-hybrid functionals. In contrast, s_8 is specifically tuned to the functional in use, adjusting the contribution of medium-range dispersion. 46

Further, the D4 model employs the Becke-Johnson choice of damping function 47 to prevent the dispersion energy from diverging at small interatomic separations:

$$f_{\text{damp}}^n(R_{AB}) = \frac{1}{1 + (R_1^{AB}/R_{AB})^n},$$
 (5)

where $R_1^{AB}=a_1R_0^{AB}+a_2$ is a scaled and offset Van der Waals radius R_0^{AB} . The parameters a_1 and a_2 are globally chosen for a given functional.

The fitting of these parameters s_8 , a_1 , and a_2 is a critical process in dispersion-corrected DFT. They need to be optimized to reflect the unique dispersion characteristics of different functionals without double-counting the dispersion effects already included in the base functional. Typically, the optimization process minimizes the discrepancy between calculated and reference interaction energies, leveraging databases rich in non-covalent interactions for effective parameterization. 48 Details on the DFT-D4 can be found in Ref. 26.

Functionals, Datasets, Basis sets, and Methods Used

Our study involves many standard (and some non-standard) choices of DFA, chemical databases, basis sets, and methodologies. In this section, we explain these choices.

Functionals: We include 7 popular semilocal functionals; 6 with few or no parameters, plus one heavily parametrized example. The 6 account for most DFT calculations at present and are computationally efficient for both molecules and materials. There are 3 generalized gradient approximations (GGAs) and three meta-GGAs (mGGAs) (including dependence on the kinetic energy density or Laplacian). They are PBE 49, a widely used non-empirical functional; revPBE 50, a modified version of PBE that performs well with dispersion corrections in GMTKN55; and BLYP 51, 52, an empirical functional combining the Becke88 exchange [51] and Lee-Yang-Parr correlation 52. We include 4 mGGA functionals: TPSS 53,54, a non-empirical functional designed for broad applicability: SCAN 55, Strongly-Constrained and Appropriately-Normed functional which satisfies all 19 known exact constraints [56]; r²SCAN [57], a numerically stable regularized version of SCAN; and M06L[58], a highly parameterized functional that implicitly captures some dispersion effects.

We also study global hybrids related to these seven. Hybrid functionals incorporate a fraction of exact exchange from HF theory, which helps reduce self-interaction error and improves band gaps. The optimal fraction of exact exchange is around 20-25% for thermochemistry and kinetics. 59 We consider 7 hybrid functionals: PBE0 59,60 (25% exact exchange), B3LYP 61 (20% exact exchange), and hybrid versions of the mGGA functionals with 25% exact exchange (TPSS0 62, SCAN0 63, r2SCAN0 64), as well as the M06 42 functional with 27% exact exchange.

Databases: To assess the performance and transferability of the D²C-DFT approach, we employ diverse datasets covering a wide range of chemical systems and properties.

Our primary dataset is GMTKN55[33], a comprehensive benchmark suite for general main group thermochemistry, kinetics, and non-covalent interactions. It consists of 55 distinct sub-datasets, totaling 1505 relative energies. We use the weighted total mean absolute deviation (WTMAD-2) metric to evaluate the performance of functionals on this dataset. Each dataset has a weight determined by its average energy divided by the average of all average energies of the 55 databases. Thus, strong bonds have low weights, and non-covalent interactions have high weights. Without this weighting, weak bonds have little effect on any averages. We use this weighting in all analyses of distributions in the database in this paper and refer to W2-GM55 as shorthand for WTMAD-2 on the GMTKN55 database.

The DIET dataset [65] is a smaller version of GMTKN55, designed to accelerate benchmarking while maintaining the diversity of the complete set. It contains 150 relative energies selected using a genetic algorithm, such that the ranking of functionals based on the DIET set closely mimics the ranking based on the complete GMTKN55 set, saving com-

putational costs when finding parameters. For D²C-DFT training, we used the mean absolute error (MAE) of the density-insensitive subset of DIET, considering the varying number of density-insensitive reactions (100-130) for different functionals. However, for the basis set comparison involving the complete DIET set, we employed the WTMAD-2 metric with the original weights from each reaction's home subset, referred to as W2-DIET in this paper. The P30 'poison' dataset 66 is also a subset of GMTKN55, representing the 30 most difficult reactions in it. We used the P30-5 subset, composed of systems up to 5 atoms, to check the impact of the dispersion correction fitting on the most challenging cases.

A specific dataset, the WATER27 [67], will play a vital role, consisting of 27 neutral and charged water clusters with up to 20 molecules. It tests a functional's ability to describe hydrogen bonding and other non-covalent interactions in aqueous systems. Extreme accuracy is vital in this subset for condensed phase simulations of water. [17], [21]

Beyond the GMTKN55 dataset, we use four others for validation and testing. The $S66\times8$ dataset [68] is an extended version of the S66 dataset [69] from GMTKN55, providing non-covalent binding energies for 66 small molecular dimers of biological importance at eight different intermolecular distances. This extension allows for a more comprehensive assessment of a functional's ability to describe non-covalent interactions across a range of distances. We used CCSD(T)/CBS energy from Ref. [70].

The Bauzá dataset [71] contains 30 complexes featuring halogen, chalcogen, and pnictogen bonds, which are essential non-covalent interactions. This dataset tests a functional's ability to describe σ -hole interactions and the impact of dispersion corrections on these systems. We used geometries and revised CCSD(T)/CBS values from Ref. [72].

Finally, the S6L [73] and L7 [74] datasets contain large supramolecular complexes, with S6L focusing on host-guest systems and L7 on extended molecular complexes. These datasets test the ability of functionals to capture dispersion interactions in large systems. These are vital as small errors in dispersion corrections are magnified in these larger systems. The S6L is the half of original S12L (2a, 2b, 4a, 5a, 6a, 7b), and we used reference DLPNO-CCSD (T_0) /CBS energy from Ref. [75] for both S6L and L7.

Basis sets: We employed the widely-used Karlsruhe def2 family [76], [77] of basis sets throughout this work, with the default being def2-QZVPPD, a large and accurate basis set that provides high-quality results. Polarization functions (P) allow for a better description of the distortion of atomic orbitals in molecules, while diffuse functions (D) improve the description of loosely bound electrons, which are essential for both anions and long-range interactions.

Methods: We compare five DFT methodologies. Self-consistent DFA calculations without any corrections (DFT); self-consistent DFA calculations with D4 model dispersion corrections (DFT-D4) 26,27; DFA calculations on HF densities (HF-DFT) 78,81; HF-DFT with dispersion correction

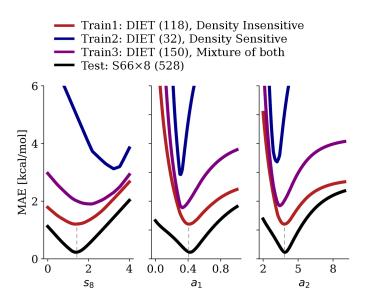


Figure 2: Parameter dependence of various training sets for D²C-B3LYP, showing the effects of individually adjusting each parameter while fixing the others at their optimal values. The dataset size for each set is indicated in parentheses. (See SI Appendix Fig. S1 for r²SCAN.)

using the original D4 parameters for the corresponding self-consistent DFA (HF-DFT-D4 $_{\rm orig}$); and our proposed method, HF-DFT with tailored dispersion correction optimized within the DC-DFT framework (D²C-DFT).

Results and Discussion

Density and Dispersion-Corrected DFT (D²C-DFT)

D²C-DFT employs DFA calculations on HF densities of density-insensitive systems with new dispersion parameters accounting for DC-DFT. An earlier version applied it to the r²SCAN functional (the HF-r²SCAN-DC4 method [21]), but recent work finds inaccuracies when applied to large molecular systems. [34] Given that up to 60 % of interaction energy in large systems can stem solely from dispersion contributions, careful consideration in parameterization becomes crucial. This work aims to refine and expand HF-r²SCAN-DC4, focusing on its parameterization process.

For efficient parameterization, we use the DIET subset. For any specific DFA, we first identify any reactions that are density-sensitive and remove them to ensure density-driven errors have negligible impact on dispersion corrections. The density-sensitive reactions depend on the choice of DFA. Through global optimization using the density-insensitive training set from DIET, we find roughly 30 local minima. We next ranked these local minima based on their W2-GM55 and selected the top 20 parameter sets. These were then re-ranked based on their MAE performance on WATER27 to find the best 10, which ensures very high accuracy for WATER27. In the final stage, each was evaluated on the L7 dataset, which contains large organic complexes and ensures

Table 1: Errors in kcal/mol of DFT, DFT-D4, and D²C-DFT across 14 functionals, using five datasets: GMTKN55, WATER27, Bauzá, L7, and S6L. For the last three, "orig" indicates HF-DFT with dispersion correction using the original D4 parameters for the corresponding self-consistent DFA, where boldface denotes the smaller value (if better by more than 20%).

		$\mathrm{GMTKN55}^a$			$\mathrm{WATER}27^b$			$\mathrm{Bauz\acute{a}}^b$		$L7^{b,c}$		$\mathrm{S6L}^{b,c}$	
		DFT	DFT-D4	$\mathrm{D}^2\mathrm{C}$	DFT	DFT-D4	$\mathrm{D}^{2}\mathrm{C}$	$\overline{{ m D}^2{ m C}}$	orig	$\overline{\mathrm{D^2C}}$	orig	$\overline{\mathrm{D^2C}}$	orig
GGA	PBE	13.89	10.12	6.53	2.34	7.58	2.01	0.73	0.65	2.53	2.10	1.72	6.07
	revPBE	20.62	8.27	6.03	17.51	3.11	8.07	0.47	0.40	1.26	1.40	4.74	5.33
	BLYP	21.10	9.48	6.60	9.79	2.30	0.99	0.54	0.73	0.81	2.29	3.45	3.29
mGGA	TPSS	15.69	8.98	6.34	5.11	3.66	1.33	0.95	0.57	2.17	1.12	3.75	3.89
	SCAN	8.64	7.67	5.37	7.07	8.46	1.16	1.64	0.93	1.14	2.65	1.90	4.74
	r^2SCAN	8.66	7.11	5.36	4.24	6.30	0.95	1.11	0.82	1.09	2.10	2.26	4.16
	M06L	8.56	8.53	8.17	1.94	1.24	6.55	0.73	0.72	1.19	1.34	3.06	3.26
Hybrid GGA	PBE0	10.94	6.18	5.42	2.22	4.91	1.24	0.89	0.68	0.62	1.12	2.30	3.12
	revPBE0	16.51	5.26	5.47	14.24	3.77	8.67	0.56	0.41	1.24	1.16	3.57	3.52
	B3LYP	16.15	6.15	4.67	5.99	3.02	0.80	0.43	0.33	0.71	1.49	2.64	1.69
Hybrid mGGA	TPSS0	13.06	5.80	5.18	5.59	1.86	2.32	0.90	0.63	0.69	0.85	2.35	2.67
	SCAN0	7.59	5.91	5.61	4.76	6.36	1.65	1.57	0.90	1.62	0.87	1.98	1.05
	r^2SCAN0	7.65	5.47	5.53	2.76	5.05	1.07	1.23	0.84	0.64	1.72	0.77	2.74
	M06	5.92	5.89	6.37	2.65	1.63	5.03	0.83	0.80	2.49	1.17	5.78	1.75

 $[^]a$ W2-GM55

our method's ability to precisely depict dispersion forces even when dominant. The best-performing parameter set on L7, in terms of MAE, was ultimately chosen as the optimal parameter set for D^2C -DFT. The final optimized parameter sets for D^2C -DFT are available next to the standard choices for DFT-D4 from the literature. (SI Appendix, Table S1)

Our study, as illustrated in Fig. 2, emphasizes the critical importance of careful training set selection in parameterization. By comparing the outcomes of parameterization using density-insensitive, density-sensitive cases, and a mixture of both from the DIET dataset, we consistently find that exclusively training with density-insensitive cases leads to more precise parameter minima for the unseen test set. The figure demonstrates that calibrating the empirical parameters solely based on density-insensitive reactions (red) consistently leads to the most accurate parameters, as indicated by the minimum MAE (gray dashed line). In contrast, calibrating parameters using density-sensitive reactions (blue) or a combination of density-sensitive and density-insensitive reactions (purple) results in less accurate parameters due to density-driven errors in the loss function. The conventional parameterization process produces sub-optimal parameters. Despite density-sensitive cases being fewer in number, their higher error rates significantly distort the loss function. This analysis extends to the 30 most challenging reactions in the GMTKN55 dataset, the P30 set. Contrary to expectations, training on these difficult cases did not enhance parameter accuracy, reinforcing the necessity of excluding density-sensitive cases for optimal parameter fitting. (SI)Appendix, Figs. S2 and S3) A fundamental principle in our DC-DFT approach is that accurate parameterization requires training on only functional error. By adhering to this

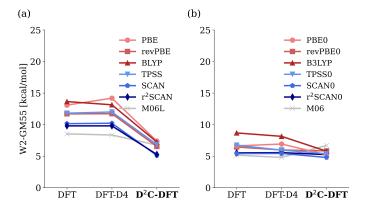


Figure 3: Same as Fig. 1, but including only density-sensitive cases, highlighting the importance of the HF density for semilocal approximations, but less so for hybrids.

strategy, D²C-DFT effectively separates both density-driven errors and dispersion portions from functional errors.

Overall performance of D2C-DFT

In Table $\boxed{\mathbb{I}}$ we report errors for many functionals across several databases. We analyze these fully, reading from the leftmost results to the right. Results for W2-GM55 were already illustrated in Fig. $\boxed{\mathbb{I}}$ demonstrating a significant error reduction relative to DFT and DFT-D4 for semilocal functionals and the same (but lesser) effect for their global hybrids. If we analyze the individual functionals, we first ignore the Minnesota functionals. Then D²C improves all semilocal functionals over either DFT or DFT-D4, sometimes by a factor of 3. When hybrids

 $^{^{\}it b}$ Mean absolute error

^c Calculated with dual-basis HF-DFT [34] using def2-SVPD/def2-TZVPPD basis set with counterpoise correction [82]

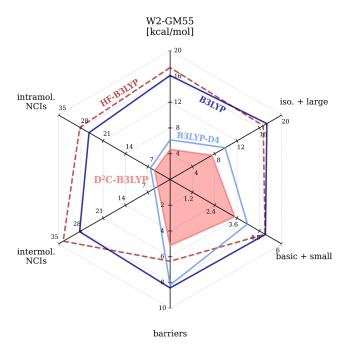


Figure 4: Performance of B3LYP in GMTKN55, divided into standard categories (*SI Appendix*, Tables S2 and S3), for four different methodologies. (See *SI Appendix*, Figs. S4, S5, S6, and S7 for all 14 functionals.)

are considered, this trend is much weaker. In two cases, results are (very slightly) worsened. We attribute this to the mixing of HF exchange, which likely mixes functional and density-driven errors, which are unaccounted for when choosing the amount of mixing. Figure 3 highlights the performance of D²C-DFT on just density-sensitive reactions, where density-driven errors are significant. The results for density-insensitive reactions, which constitute the majority of the GMTKN55 dataset, exhibit trends similar to those observed in Fig. 1 as they are major in GMTKN55. SI Appendix, Fig. S8

Finally, we note that, for both Minnesota functionals, almost no effect is seen either in D4 or D²C. We attribute this to their highly parameterized construction, which presumably thoroughly mixes both dispersion and density corrections, leading to no improvement when such effects are systematically accounted for. A critical paper claimed DFT was straying from the exact path due to deficiencies in densities, 83, but DC-DFT found no evidence of errors in chemically significant differences. However, the current analysis (Fig. 1) demonstrates the problems with such approaches.

Focusing on D²C-B3LYP, its superior performance can be attributed to the synergy between the overly repulsive Becke88 exchange [51] and purely attractive dispersion corrections. Radar plots separate the relative performance of methods on different types of GMTKN55 databases. [SI] Appendix Tables [S2] and [S3] Figure [4] showcases this balanced improvement across multiple GMTKN55 subgroups, indicating a harmonious error correction in various chemi-

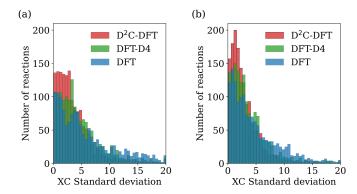


Figure 5: The standard deviation of predictions from (a) six semilocal and (b) six hybrid functionals (no M functionals) for every reaction in GMTKN55. The width of the distribution shrinks from DFT to D^2C -DFT, showing improved agreement. For semilocals, the averages are 7.8, 5.4, and 3.8 respectively, with 50%, 40%, and 25% above 5 kcal/mol. The corresponding numbers for hybrids are 5.5, 4, and 3, and 40%, 25%, and 16%, respectively.

cal systems. Here, we see that HF-B3LYP largely worsens performance in many categories, while D^2C -B3LYP outperforms B3LYP-D4 in all categories. Radar plots for all other functionals are available. (SI Appendix, Figs. S4, S5, S6, and $\overline{S7}$)

Possibly more importantly, we see decreased variation in DFT results, both across different functionals and, for a given functional, across the GMTKN55 database. Figure 5 showcases the first phenomenon, showing the standard deviation among six functionals. Its distribution in GMTKN55 becomes significantly narrower with D²C-DFT, illustrating a notable improvement in precision. Also, by examining the radar plots for all functionals grouped by rung (SI Appendix, Figs. S4, S5, S6, and S7), we observe that the D²C variants exhibit pink hexagons of similar size and shape, indicating similar behavior of functionals in the D²C framework.

The second narrowing is for the error distribution within a given functional, as shown in <u>SI Appendix</u>, Figs. <u>S9</u> and <u>S10</u>. This reduced variability among DFA answers can be crucial in real-world situations where a benchmark answer is unavailable.

In <u>SI Appendix</u>, we provide all 100k calculations used to generate the W2-GM55 data in this paper. We emphasize that it contains an extraordinary wealth of specific information that can interest specific communities. Some communities care only about one particular functional (such as users of PBE 49) or users of B3LYP 61, largely an orthogonal set). Others might care deeply about performance on a specific type of system, such as non-covalent interactions in general 33 or water in particular 67. Meanwhile, developers might be most interested in self-interaction errors 33. The current paper provides only a broad overview and some interesting slices.

To emphasize the importance of looking within data sets, Table 2 gives results for each separate dataset with B3LYP,

Table 2: Performance of B3LYP across the GMTKN55 subdatasets, with mean absolute errors reported in kcal/mol. The first column (Dens.) indicates whether half or more of the systems in the dataset are density-sensitive by the threshold $\tilde{S} > 2$ (see Eq. 2). The second column (Disp.) indicates whether half or more of the systems in the dataset are dominated by dispersion effects. A system is considered dispersion-dominated if the dispersion-corrected (D4) energy differs from the self-consistent DFT energy by more than 50%. In the majority of cases, D²C outperforms other methods. (See *SI Appendix*, Table S3 for a description of each dataset.)

Basic properties and reaction energies for small systems		Dens.	Disp.	DFT	DFT-D4	HF-DFT	D ² C-DFT
W4-11	Basic	propertie	es and r	eaction	energies for	small syste	ems
Real Column Section Section							
DIPCS10	G21EA	NO	NO	2.28	2.31	2.73	2.69
NO	G21IP	NO	NO	3.77	3.74	4.14	4.11
No	DIPCS10	NO	NO	4.56	4.57	4.19	4.33
ALKBDE10							
No							
AL2X6							
HEAVYSB11							
NBPRC							
ALK8							
RC21							
BH76RC NO NO 2.36							
BH76RC							
TH51							
TAUT15							
NO							
Reaction large systems somerisation reactions							
MB16-43							
DARC							
RSE43							
BSR36							
CDIE20							
ISO34							
ISOL24							
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INV24	BHPERI	NO	YES	4.35	1.13	4.75	2.42
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BUT14DIOL NO NO 0.41 0.46 0.68 0.26	BUT14DIOL	NO	NO	0.41	0.46	0.68	0.26

highlighting in bold whenever one of the methods is more than 20% better than its nearest rival. We have also listed which datasets are density-sensitive (more than half the

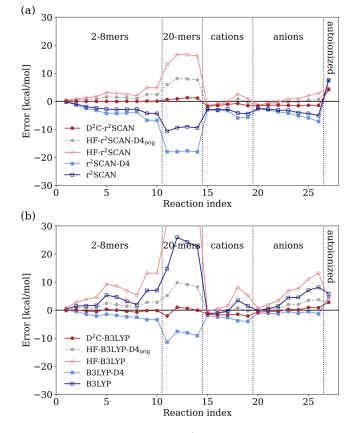


Figure 6: Performance of (a) r²SCAN and (b) B3LYP functionals on the WATER27 dataset. The D²C approach outperforms other methods, particularly for larger water clusters. This improvement is crucial for accurately modeling aqueous environments. (See SI Appendix, Figs. S11 S12 S13 and S14 for the other functionals.)

reactions are) and which are dispersion dominated (DFT and DFT-D4 errors differ by more than half the DFT error). Clearly, D²C does best most often, but it is interesting to compare it with DFT-D4, i.e., the standard dispersion methodology without density correction. It is not intuitive why using HF densities does not improve some of these datasets. Doubtless, there are accidental error cancellations between the underlying functional error, the error in the dispersion estimate, and the density-driven error, none of which are driven to zero with our methodology.

The middle columns of Table $\boxed{1}$ highlight the outperformance of D^2C on the WATER27 dataset. We regard achieving high accuracy for water clusters as a crucial component of any D^2C methodology. The principle of using HF densities to fix issues in semilocal functionals for water models was first noted by Dasgupta et al. $\boxed{17}$, $\boxed{18}$, and further refinement led to the HF-r²SCAN-DC4 method $\boxed{21}$. Here, the improvement for the six semilocal functionals is even more considerable than in W2-GM55 and is unpacked for r²SCAN in Fig. $\boxed{6}$ (a). Starting from the pure functional, D4 corrections make things worse. Insertion of HF densities significantly overcorrects. Using HF and D4 (original

parameters) gives a better performance, but much better results for the crucial 20-mers occur with D²C. Such a small error is needed for models of water, which is why we use WATER27 as a validation set.

Figure (b) shows the corresponding figure for B3LYP, showing the same overall trends and comparable results. A unique observation is the underperformance of D²C-revPBE and D²C-revPBE0 in WATER27. We interpret this as a result of the resolution of charge transfer errors inherent in revPBE family functionals through the use of HF density, eliminating the compensatory error cancellation. [84] WATER27 plots for the remaining functionals are available. [SI Appendix] Figs. [S11] [S12] [S13] and [S14] This accuracy can also be seen in water hexamer and 20-mers, where D²C-r²SCAN provides the most accurate predictions compared to other dispersion-corrected HF-r²SCAN variants, and getting (almost) all relative energies correct. [SI Appendix] Figs. [S15] and [S16]

The rightmost columns of Table 1 detail D^2C -DFT's evaluation across the Bauzá, L7, and S6L datasets. The Bauzá set is analogous to the WATER27 but is not in the validations set. It is discussed in detail in the next section. The rest test the capability of D^2C with large molecular structures, as exemplified by S6L and L7.

In summary, our study generalizes the D²C-DFT method from HF-r²SCAN-DC4 across various functionals, achieving a balanced performance in diverse chemical situations. This success stems from the effective use of HF density in mitigating density-driven errors and precise dispersion parameterization within the DC-DFT framework, thus validating the dual-calibration approach's principles and effectiveness.

Origins of density delocalization errors

In this section, we delve into mitigating density-driven errors, which are intrinsically linked to delocalization errors in DFT. Delocalization error remains a major challenge in DFT 85-87, impacting dissociation limits 88, band gaps 89, 90, and charge transfer predictions 91. Some (but not all) can be alleviated with DC-DFT methods. We use technology developed in Refs. 12 and 15.

A classic (and extreme) case of density delocalization error is stretching an ionic bond to extremes. 92 Figure 7(a) shows conventional DFAs struggling with severe delocalization error in the NaCl dissociation curve, leading to incorrect energies and fractional charges as they approach the dissociation limit. In contrast, D²C-DFT, using HF density, accurately captures the correct charge and densities and thus produces much more accurate energies at the dissociation limit. To demonstrate the underlying reason, we conduct a fractional analysis for infinitely separated Na and Cl atoms, as depicted in Fig. 7(b). While the DFAs exhibit a concave upward curve leading to fractional electron convergence and inaccuracies in dissociation limits (typically about 0.4 of an electron remains transferred, even as $R \to \infty$), HF's concave downward curve has sharp downward cusps exactly at integers, just like the exact curve

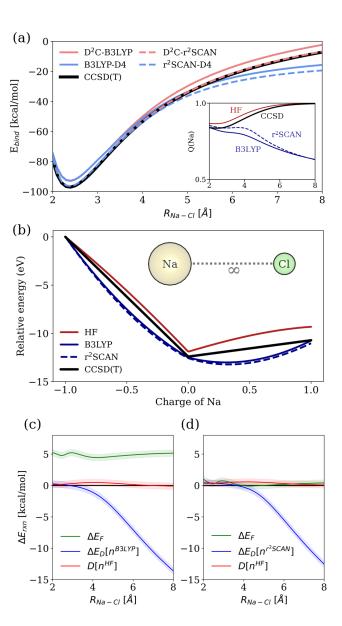


Figure 7: Origin of errors in DFA bond dissociation curves of NaCl: (a) displays the NaCl dissociation curve, with an inset displaying the calculated partial charge on the Na atom; (b) examines fractional electron analysis for isolated Na and Cl with the x-axis indicating charge on Na; (c) and (d) utilize Kohn-Sham inversion to dissect errors within B3LYP and $\rm r^2SCAN$, respectively. The $\pm 0.5~\rm kcal/mol$ bands illustrate the inversion's uncertainty. 15

(which consists of linear segments [93]). Note that the HF energies are hopelessly inaccurate, while in the vicinity of the integer, the DFA energies are much more accurate. Hence, the value of HF-DFT, as the HF density and charges remain very accurate, and the functional error of the DFA is very small on the right density.

To gain insights into the remaining error of D^2C -B3LYP, we conducted Kohn-Sham inversion 15, 94 as shown in Fig. 7(c), using a reference CCSD density 8, to strictly separate functional and density-driven errors, as in Eq. 1.

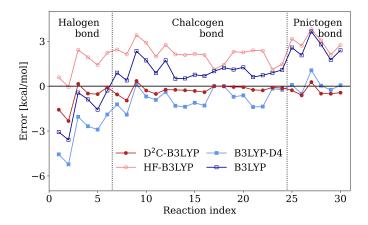


Figure 8: Evaluation of B3LYP behavior using the Bauzá dataset, testing the accuracy of σ -hole interactions. The use of the HF density in D²C significantly reduces delocalization errors. (See SI Appendix, Fig. S17 for r²SCAN.)

Notably, B3LYP exhibits a constant functional error of approximately 5 kcal/mol, while its density-driven error varies from zero to a significant amount as stretching occurs. In contrast, the density difference of HF remains uniformly close to zero due to its correct behavior. So, the errors in D²C-B3LYP's dissociation curve in (a) stem from B3LYP's uncorrected functional errors, as D²C-DFT primarily addresses dispersion errors—a minor component in the NaCl system (\sim 0.1 kcal/mol). On the other hand, in the case of r²SCAN illustrated in Fig. 7(d), the density difference of HF and further functional error of r²SCAN are also nearly zero, which results in D²C-r²SCAN's near perfect dissociation curve in (a). Here, r²SCAN is much better than B3LYP. In either case, the HF density is an excellent proxy for the exact density.

Moving from the illustrative case of NaCl dissociation to the Bauzá dataset, we expand our analysis of delocalization errors. The Bauzá dataset, a collection of Halogen, Chalcogen, and Pnictogen-involved systems, offers an ideal ground to assess the accuracy of σ -hole interaction representation, which is known to pose significant delocalization error, [72] and was attributed to large density-driven errors in our prior work. [14]

As demonstrated in Fig. 8, when dispersion correction is applied to the self-consistent B3LYP, B3LYP-D4 often results in overbinding, worsening results. However, D²C-B3LYP leads to a remarkable decrease in errors in almost every reaction. Applying the HF density means accurate dispersion is restored in dispersionless situations, 45, 95 substantially improves results. This can also be thought of in terms of the absence of self-interaction error in HF. 96 Averages across types are reported in Table 3.

Like the NaCl dissociation case, the effects of delocalization in the self-consistent density in the dissociation of halogen complexes were studied in Ref. [14]. The conventional DFT tends to overbind due to artificial charge transfer. In stark contrast, HF-DFT precisely mirrors the

Table 3: Mean absolute errors (in kcal/mol) for the Bauzá dataset, categorized by interaction type and overall performance (All).

	Halogen	Chalcogen	Pnictogen	All
B3LYP	1.64	1.02	2.54	1.45
B3LYP-D4	3.22	0.83	0.33	1.21
HF-B3LYP	1.44	2.14	2.93	2.16
D^2C -B3LYP	0.86	0.28	0.43	0.43

reference curve, adeptly tackling a major element of delocalization error - stretched geometries. This is consistent with other studies that highlight the impressive capability of HF-DFT in mitigating other aspects of delocalization error, including ions-in-solution 11, spin gaps 13, and torsional barriers 16.

Transferability of Parameters across basis sets

For D²C-DFT applications, we recommend the def2-QZVPPD basis set, quadruple-zeta level with diffusion functions, chosen for its minimal basis set errors. Our parameters were optimized using this set, but this may become unaffordable for large systems. Here, we test the transferability of our dispersion parameters to smaller basis sets, using W2-DIET for efficiency. Figure 9 shows that D²C-DFT's parameters are transferable effectively down to the triple zeta level with a diffused basis set (def2-TZVPPD) or at the quadruple-zeta level without diffusion functions (def2-QZVP). Below this, performance significantly diminishes. Despite ongoing discussions about the necessity for basis set-specific parameter fitting in dispersion correction, [97,98] we recommend using def2-QZVPPD as the first choice, with feasible applicability down to def2-TZVPPD.

For larger systems, we suggest the dual-basis method [99] for HF-DFT discussed in Ref. [34]. This approach, starting with a double-zeta level calculation, allows a triple- or quadruple-zeta level computation at a much lower computational cost, striking a balance between efficiency and accuracy, which is particularly apt for larger molecules where computational efficiency is paramount.

Conclusion

Our study illustrates the general principle of dual calibration: Separating distinct sources of error when parameters are being fitted in empirical density functionals. Here, the two sources are dispersion corrections and density corrections. By separating these two, we improve results compared to either individually (DFT-D4 for dispersion and HF-DFT for density-correction) or even compared to both if the original D4 parameters are used.

To demonstrate its generality, we applied our procedure to a broad range of non-empirical and only slightly empirical functionals that enjoy considerable popularity in their

			Primitives				
basis	ВЗІХР	B3LYP-D4	diff.	diff. HF-B3LYP D ² C-B3LYP diff.		diff.	Primitives
QZVPPD	16.11	6.42	-5%	16.85	4.49	-	1098
QZVP	15.94	6.78	-	16.64	4.66	4%	996
TZVPPD	15.89	6.53	-4%	16.74	4.62	3%	602
TZVP	15.60	8.55	26%	15.73	6.02	34%	356
SVPD	12.54	11.34	67%	13.25	8.80	96%	304
SVP	19.64	18.92	179%	17.19	15.51	245%	202

Figure 9: Transferability of dispersion parameters in B3LYP-D4 and D²C-B3LYP, optimized repectively for the def2-QZVP and def2-QZVPPD, across various basis set sizes within the Karlsruhe (def2-) family. For context, we note the number of primitive Gaussians for the C_8H_{18} molecule. Rows are color-coded based on the difference relative to optimized basis sets (Positive values indicate degraded performance.): Green for <10%, yellow for 10-50%, and red for >50%. This analysis leverages the DIET, a subset of GMTKN55. (See SI Appendix Fig. S18 for r²SCAN.)

respective fields of application. By eliminating density-sensitive cases from our training data, our dispersion corrections differ somewhat from the standard choices, as they correct only functional errors. On the WTMAD-2 measure of the GMTKN55 database, our results are indistinguishable from HF-DFT-D4 using the original parameters. However, because of our DC-DFT training and validation, we do significantly better in crucial cases like WATER27 and large dispersion-dominated databases. Our results here can be considered a generalization of (and improvement over) that of Ref. [21], which applied these principles to the r²SCAN semilocal functional.

While double-hybrid functionals have recently achieved impressive accuracy on GMTKN55, their reliance on costly MP2 calculations limits their applicability to larger systems. [100] [101]. Such functionals suffer much less from dispersion- and density-driven errors, so our dual calibration approach would have a much smaller impact. Nonetheless, given the impressive accuracy that these functionals already achieve, even a much smaller improvement might be significant.

We also point out that the dual calibration approach, perhaps surprisingly, does not rely on the HF density being 'better' (i.e., yielding more accurate energies in approximate functionals) than self-consistent densities, even for density-sensitive problems. The critical step is to find global dispersion parameters (for each approximate functional) on only density-insensitive cases, thereby minimizing any arbitrary bias produced by density-driven errors. The fact that density-sensitive cases can be significantly improved when HF densities are used with dispersion corrections suggests that, at least in those cases, the HF density is better, but only in the sense defined above. Finding a more accurate density than HF density remains a pressing and challenging issue, as it would unveil the true functional errors without making assumptions. Additionally, identifying the various types of functional errors beyond dispersion errors would

be required for a clear resolution of them.

In summary, the dual-calibration method, embodied by the $\mathrm{D}^2\mathrm{C}\text{-}\mathrm{DFT}$, stands as a promising method with capabilities to improve DFT systematically. The results shown here suggest that some version of this would be helpful in almost all searches for empirical parameters. Its continued development and refinement hold great potential for advancing the DFT functionals and their applications.

Computational Methods

All DFT calculations were performed at Pyscf 102 103. Dispersion corrections were facilitated by employing the dftd4 program for D4 corrections 26,27. Dispersion parameters of D²C-DFT were optimized using the SciPy 104 shgo optimizer 105. Results for datasets L7 and S6L were calculated with dual-basis HF-DFT 34,99 targeting def2-TZVPPD with counterpoise correction 82. A fractional electron picture was drawn with the method of Refs. 106 and 107. A Mulliken population analysis 108 based on meta-Lowdin atomic orbitals 109 is used to determine partial charges of NaCl. For the error decomposition, Kohn-Sham inversion was conducted using KS-pies 110, in the Wu-Yang scheme 94, targeting CCSD density 8 at aug-cc-pVTZ basis set 111.

SI Appendix

Dispersion parameters for 14 $\rm D^2C$ -DFT and DFT-D4, Abbreviation used for GMTKN55, Description of subsets within GMTKN55, Simplified view on parameterization strategy, Performance on GMTKN55 for 14 functionals, Distribution of weighted errors for 14 functionals across GMTKN55, Performance on WATER27 for 14 functionals, Performance on water hexamers and 20-mers, Evaluation of $\rm r^2SCAN$ correction variants using the Báuza dataset, Transferability of dispersion parameters for $\rm r^2SCAN$ variants.

Raw data of calculations in GMTKN55 database

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The authors declare no competing interest.

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Supplementary Materials for

Correcting dispersion corrections with density-corrected DFT

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Table S1: DFT-D4 parameters used in our work. s_6 is set to 1. Refer to Subsection "D²C-DFT" in the main text.

		D ² C-DFI	-	DFT-D4					
	s_8	a_1	a_2	cf.	s_8	a_1	a_2	Ref	
PBE	1.78595387	0.88511469	2.32863362	a	0.95948085	0.38574991	4.80688534	[1]	
PBE0	1.53360351	0.78689150	3.25641582		1.20065498	0.40085597	5.02928789	<u>[1]</u>	
revPBE	1.80966761	0.58558155	2.71965468		1.74676530	0.53634900	3.07261485	<u>(1)</u>	
revPBE0	1.74590783	0.40385673	4.26818294		1.57185414	0.38705966	4.11028876	<u>[1]</u>	
BLYP	1.65244425	0.56438878	2.65715701		2.34076671	0.4448865	4.09330090	<u>(1)</u>	
B3LYP	1.35513689	0.41757850	3.84594813		2.02929367	0.40868035	4.53807137	<u>[1]</u>	
TPSS	1.54044984	0.69473318	2.51512802		1.76596355	0.42822303	4.54257102	<u>(1)</u>	
TPSS0	1.50843498	0.60162555	3.65500533		1.62438102	0.40329022	4.80537871	<u>(1)</u>	
SCAN	1.72616184	0.06450398	8.62911596		1.46126056	0.62930855	6.31284039	<u>(1)</u>	
SCAN0	3.69655894	0.16214976	8.90158495		6.1187	0.3750	8.1124	[2]	
r ² SCAN	0.02734375	0.74707031	3.34667969	b	0.6019	0.5156	5.7734	[3]	
r ² SCAN0	3.97877459	0.75987648	5.45977445		0.8992	0.4778	5.8779	[3]	
M06L	0.75781250	0.81445313	6.16992188		0.59493760	0.71422359	6.35314182	[1]	
M06	1.30522230	0.83568617	4.37780185		0.16366729	0.53456413	6.06192174	1	

^a Parameters revised from Ref. [4]

^b Parameters revised from Ref. [5]

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Table S2: Abbreviation, description, and the number of reactions included in the GMTKN55. classified further by 5 subgroups. In the main text, the Fig. 4 uses this classification.

abbreviation	description	#
GMTKN55	Full GMTKN55	1505
basic + small	Basic properties and reaction energies for small systems	473
iso. + large	Reaction energies for large systems and isomerization reactions	243
barriers	Reaction barrier heights	194
intermol. NCIs	Intermolecular noncovalent interactions	304
intramol. NCIs	Intramolecular noncovalent interactions	291

Table S3: Description of the subsets within the GMTKN55 database. This table is from Ref. [6]. GMTKN55 database is utilized throughout the main text.

Set	Description	# ^a	Weight
	Basic properties and reaction energies for small systems		
W4-11	Total atomisation energies	140 (152)	0.18519
G21EA	Adiabatic electron affinities	25 (50)	1.69045
G21IP	Adiabatic ionisation potentials	36 (71)	0.22064
DIPCS10	Double-ionisation potentials of closed-shell systems	10(20)	0.08687
PA26	Adiabatic proton affinities (incl. of amino acids)	26 (52)	0.30065
SIE4x4	Self-interaction-error related problems	16 (23)	1.68539
ALKBDE10	Dissociation energies in group-1 and -2 diatomics	10 (20)	0.56450
YBDE18	Bond-dissociation energies in ylides	18 (29)	1.15351
AL2x6	Dimerisation energies of AlX ₃ compounds	6 (11)	1.58402
HEAVYSB11	Dissociation energies in heavy-element compounds	11 (22)	0.97961
NBPRC	Oligomerisations and H ₂ fragmentations of NH ₃ /BH ₃ systems; H2 activation reactions with PH ₃ /BH ₃ systems	12 (21)	2.05136
ALK8	Dissociation and other reactions of alkaline compounds	8 (17)	0.90796
RC21	Fragmentations and rearrangements in radical cations	21 (41)	1.59222
G2RC	Reaction energies of selected G2/97 systems	25 (47)	1.10878
BH76RC	Reaction energies of the BH76 set	30	2.65710
FH51	Reaction energies in various (in-)organic systems	51 (87)	1.83289
TAUT15	Relative energies in tautomers	15 (25)	18.6605
DC13	13 difficult cases for DFT methods	13 (30)	1.03385
DC15	Reaction energies for large systems and isomerisation reactions	15 (50)	1.05505
MB16-43	Decomposition energies of artificial molecules	12 (59)	0.12135
DARC	Reaction energies of Diels-Alder reactions	43 (58)	1.75046
RSE43	Radical-stabilisation energies	14 (22)	7.47666
BSR36		43 (88)	3.50930
CDIE20	Bond-separation reactions of saturated hydrocarbons	36 (38)	
	Double-bond isomerisation energies in cyclic systems	20 (36)	14.0172
ISO34	Isomerisation energies of small and medium-sized organic molecules	34 (63)	3.90116
ISOL24	Isomerisation energies of large organic molecules	24 (48)	2.59321
C60ISO	Relative energies between C ₆₀ isomers	9 (10)	0.57851
PArel	Relative energies in protonated isomers	20 (31)	12.2751
ВН76	Reaction barrier heights Barrier heights of hydrogen transfer, heavy atom transfer,	76 (86)	3.05353
	nucleophilic substitution, unimolecular and association reactions	*****	
BHPERI	Barrier heights of pericyclic reactions	26 (61)	2.72312
BHDIV10	Diverse reaction barrier heights	10 (20)	1.25383
INV24	Inversion/racemisation barrier heights	24 (48)	1.78484
BHROT27	Barrier heights for rotation around single bonds	27 (40)	9.06110
PX13	Proton-exchange barriers in H ₂ O, NH ₃ , and HF clusters	13 (29)	1.70375
WCPT18	Proton-transfer barriers in uncatalysed and water-catalysed reactions	18 (28)	1.62456
	Intermolecular noncovalent interactions		
RG18	Interaction energies in rare-gas complexes	18 (25)	98
ADIM6	Interaction energies of n-alkane dimers	6 (12)	16.9250
S22	Binding energies of noncovalently bound dimers	22 (57)	7.78378
S66	Binding energies of noncovalently bound dimers	66 (198)	10.3969
HEAVY28	Noncovalent interaction energies between heavy element hydrides	28 (38)	45.7859
WATER27	Binding energies in $(H_2O)_n$, $H^+(H_2O)_n$ and $OH^-(H_2O)_n$	27 (30)	0.70051
CARBHB12	$\label{eq:Hydrogen-bonded} \mbox{Hydrogen-bonded complexes between carbene analogues and H_2O, NH_3, or HCl}$	12 (36)	9.41787
PNICO23	Interaction energies in pnicogen-containing dimers	23 (69)	13.3006
HAL59	Binding energies in halogenated dimers (incl. halogen bonds)	59 (105)	12.3775
AHB21	Interaction energies in anion-neutral dimers	21 (63)	2.52777
CHB6	Interaction energies in cation-neutral dimers	6 (18)	2.12208
IL16	Interaction energies in anion-cation dimers	16 (48)	0.52125
	Intramolecular noncovalent interactions		
IDISP	Intramolecular dispersion interactions	6 (13)	3.99625
ICONF	Relative energies in conformers of inorganic systems	17 (27)	17.4010
ACONF	Relative energies of alkane conformers	15 (18)	30.9901
AMINO20x4	Relative energies in amino acid conformers	80 (100)	23.3076
PCONF21	Relative energies in tri- and tetrapeptide conformers	18 (21)	35.0503
- -	Relative energies in melatonin conformers	51 (52)	11.4343
MCONF		- ()	
	Relative energies of sugar conformers	17 (19)	12.3565
MCONF SCONF UPU23	Relative energies of sugar conformers Relative energies between RNA-backbone conformers	17 (19) 23 (24)	12.3565 9.93253

^a Relative energy counts and required single-point computations (in parenthesis)

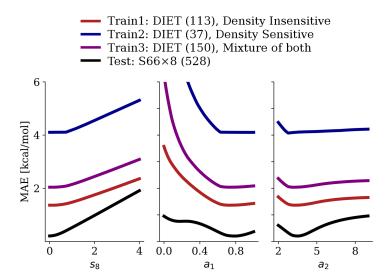


Figure S1: Simplified view on D^2C - r^2SCAN 's parameterization strategy, highlighting the distinction between training set types. Three graphs illustrate the effects of individually adjusting each parameter while fixing the others at optimal values. The datasets utilized in this analysis are from Refs. $\boxed{2}$ and $\boxed{8}$. Refer to Fig. $\boxed{2}$ in the main text for D^2C -B3LYP results.

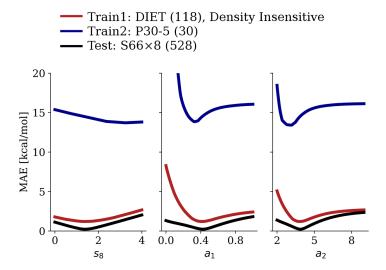


Figure S2: Simplified view on D^2C -B3LYP's parameterization strategy, using P30-5 dataset [9] instead, highlighting the distinction between training set types. Refer to Fig. [2] in the main text which utilizes a different training set.

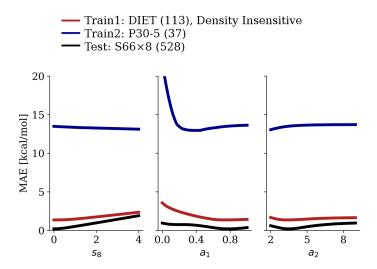


Figure S3: Simplified view on D²C-r²SCAN's parameterization strategy, using P30-5 dataset [9] instead. Refer to Fig. [S1] which uses a different training set.

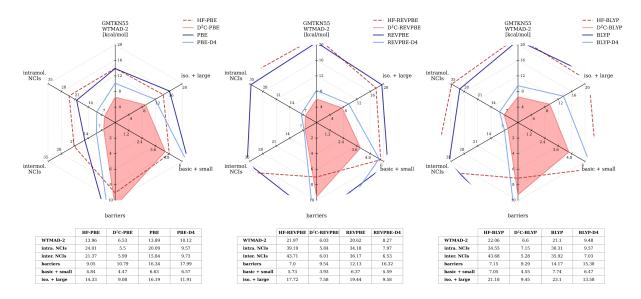


Figure S4: Performance of the GGA functionals variants across the GMTKN55 database. 6 Each section signifies a distinct category. basic + small: Basic properties and reaction energies for small systems, iso + large: Reaction energies for large systems and isomerization reactions, barriers: Reaction barrier heights, intermol NCIs: Intermolecular noncovalent interactions, intramol NCIs: Intramolecular noncovalent interactions, GMTKN55 WTMAD-2: Weighted Total Mean Absolute deviation. Refer to Fig. 4 in the main text for comparison with B3LYP variants.

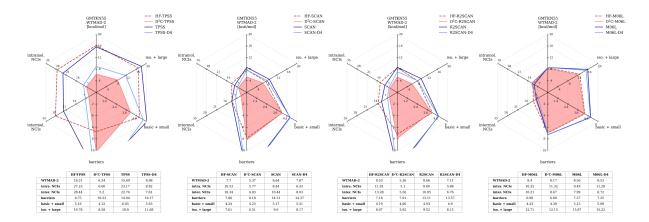


Figure S5: Performance of the mGGA functionals variants across the GMTKN55 database.

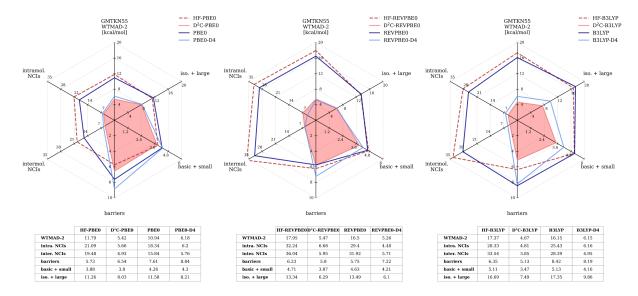


Figure S6: Performance of the Hybrid GGA functionals variants across the GMTKN55 database.

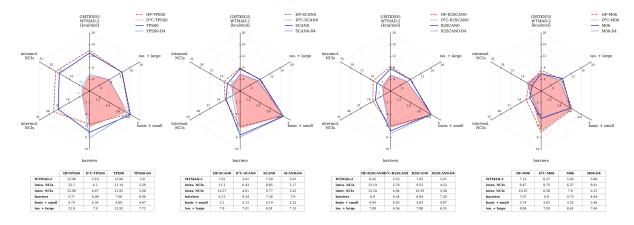


Figure S7: Performance of the Hybrid mGGA functionals variants across the GMTKN55 database.

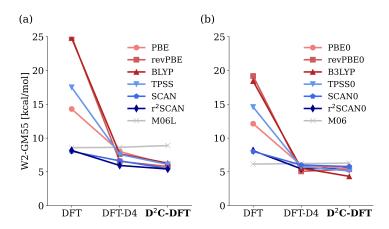


Figure S8: Performance of DFT, DFT-D4, and D^2C -DFT on the density-insensitive reactions in GMTKN55 database [6]. See Figs. [1] and [3] in the main text for all reactions in GMTKN55 and only density-sensitive reactions in GMTKN55.

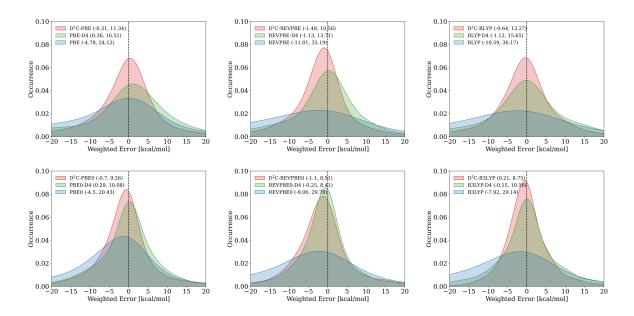


Figure S9: Distribution of Weighted Errors for each GGA-based functionals across 1505 reactions in GMTKN55, including mean and standard deviation (in parenthesis). The reduction in variability is also observed in individual functional with D^2C -DFT. Refer to Fig. 5 in the main text to see the standard deviations across functionals.

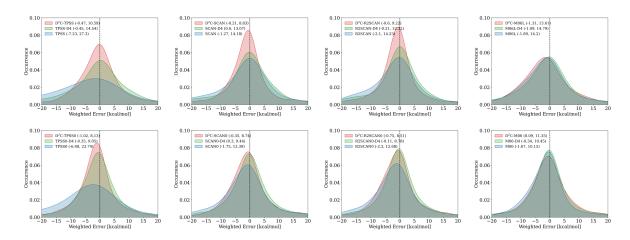


Figure S10: Distribution of Weighted Signed Errors for each mGGA-based functionals across 1505 reactions in GMTKN55, including mean and standard deviation (in parenthesis).

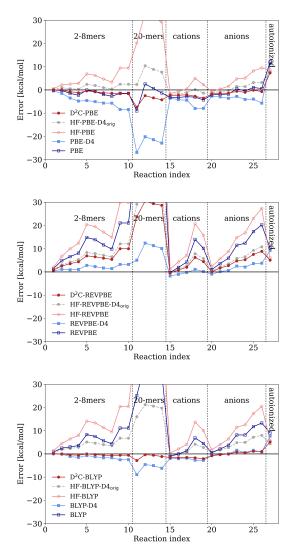


Figure S11: Performance of GGA functionals on WATER27 dataset [10]. To make a negative error implies an overbind effect, the signs of the errors are inverted from the original reference. See Fig. 6 in the main text for the performance on r²SCAN and B3LYP.

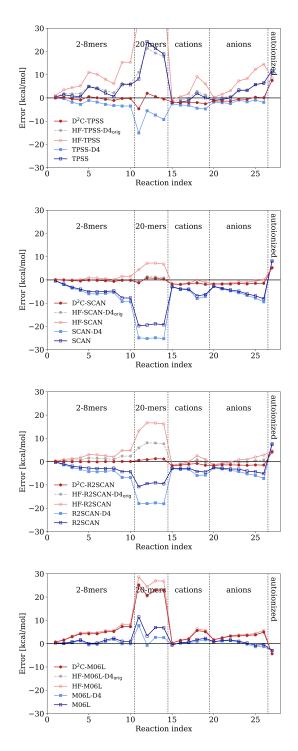


Figure S12: Performance of mGGA functionals on WATER27 dataset.

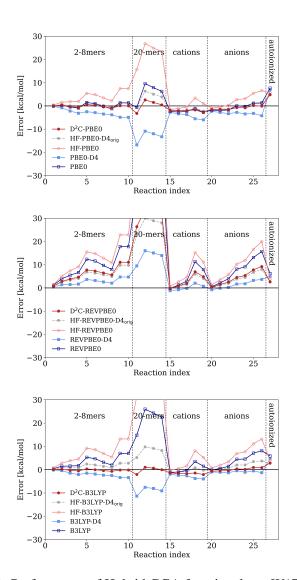


Figure S13: Performance of Hybrid GGA functionals on WATER27 dataset.

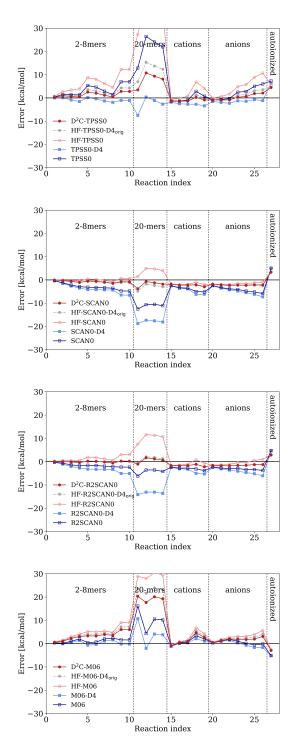


Figure S14: Performance of Hybrid mGGA functionals on WATER27 dataset.

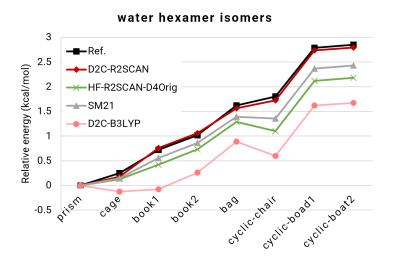


Figure S15: Water hexamers' relative energy compared to prism structure. Geometries are from Ref. [II] and reference CCSD(T)/CBS energy from Ref. [I2]. D²C-r²SCAN, HF-r²SCAN-D4Orig, and SM21 are the same HF-r²SCAN schemes coupled with different dispersion parameters. Dispersion parameter for D4Orig is from Ref. [I] and SM21 is from Ref. [2]. See Subsection 'Overall performance of D²C-DFT' in the main text.

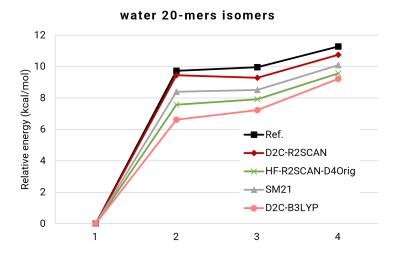


Figure S16: Water 20-mers' relative energies from WATER27[10].

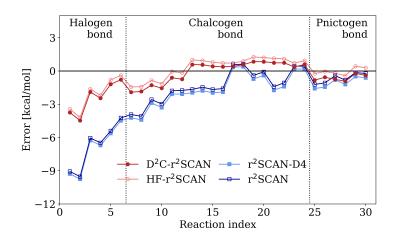


Figure S17: Evaluation of r^2SCAN correction variants using the Bauza dataset [13]. This dataset encompasses a wide range of reactions involving Halogen, Chalcogen, and Pnictogen elements, making it an ideal benchmark for testing the accuracy in modeling σ -hole interactions. See Fig. 8 in the main text for the B3LYP results.

			W2-DIET	[kcal/mol]			Primitives
basis	r ² SCAN	r ² SCAN-D4	diff.	HF-r ² SCAN	D ² C-r ² SCAN	diff.	C_8H_{18}
QZVPPD	8.88	7.24	-5%	7.97	4.88	-	1098
QZVP	9.13	7.60	-	7.94	4.98	2%	996
TZVPPD	8.80	7.35	-3%	7.90	4.99	2%	602
TZVP	10.05	8.83	16%	7.63	5.83	19%	356
SVPD	9.77	12.02	58%	6.63	9.16	88%	304
SVP	16.26	17.38	129%	12.52	14.35	194%	202

Figure S18: Transferability of dispersion parameters in r^2SCAN -D4 and D^2C - r^2SCAN , optimized for the def2-QZVP and def2-QZVPPD, across various basis set sizes within the Karlsruhe family [14, 15]. For context, we note the number of primitive Gaussians for the C_8H_{18} molecule. Difference relative to def2-QZVPPD (diff.) is color-coded: Green marks mean degradation below 10 %, yellow for 10-50 %, and red for over 50 %. This analysis leverages the DIET150[8], a subset of GMTKN55[6]. The parameter of D^2C - r^2SCAN can be safely used down to def2-TZVPPD. See Fig. 9 in the main text for the B3LYP result.

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