

Centrality dependence of Lévy-stable two-pion Bose-Einstein correlations in $\sqrt{s_{NN}} = 200$ GeV Au + Au collisions

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The PHENIX experiment measured the centrality dependence of two-pion Bose-Einstein correlation functions in $\sqrt{s_{NN}} = 200$ GeV Au + Au collisions at the Relativistic Heavy Ion Collider at Brookhaven National Laboratory. The data are well represented by Lévy-stable source distributions. The extracted source parameters are the correlation-strength parameter λ , the Lévy index of stability α , and the Lévy-scale parameter R as a function of transverse mass m_T and centrality. The $\lambda(m_T)$ parameter is constant at larger values of m_T , but decreases as m_T decreases. The Lévy-scale parameter $R(m_T)$ decreases with m_T and exhibits proportionality to the length scale of the nuclear overlap region. The Lévy exponent $\alpha(m_T)$ is independent of m_T within uncertainties in each investigated centrality bin, but shows a clear centrality dependence. At all centralities, the Lévy exponent α is significantly different from that of Gaussian ($\alpha = 2$) or Cauchy ($\alpha = 1$) source distributions. Comparisons to the predictions of Monte-Carlo simulations of resonance-decay chains show that, in all but the most peripheral centrality class (50%–60%), the obtained results are inconsistent with the measurements, unless a significant reduction of the in-medium mass of the η' meson is included. In each centrality class, the best value of the in-medium η' mass is compared to the mass of the η meson, as well as to several theoretical predictions that consider restoration of $U_A(1)$ symmetry in hot hadronic matter.

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I. INTRODUCTION

In a previous paper [1] on Bose-Einstein Correlations (BECs)—also known as Hanbury Brown and Twiss (HBT) Correlations—the PHENIX Collaboration found that for 0%–30% centrality Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV, the two-particle BECs are well described by a Lévy-stable source distribution. However, the traditional description of the same dataset, using a Gaussian source distribution, was found to be inadequate [1]. A strong preference for the Lévy description had also been seen in $e^+ + e^-$ collisions at the Large Electron-Positron Collider [2] and in $p + p$, $p + \text{Pb}$, and $\text{Pb} + \text{Pb}$ collisions at the Large Hadron Collider [3–6], in Be+Be and Ar+Sc collisions at the Super Proton Synchrotron [7,8], and in Au + Au collisions at the Relativistic Heavy Ion Collider (RHIC) [9].

Presented here is a precise measurement of the centrality and transverse-mass dependence of the two-pion BEC

function in Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV by the PHENIX experiment at RHIC. This data sample, recorded in 2010, allows a fine transverse-mass binning and inference of the shape of the correlation function more precisely than was possible with earlier data sets. For the first time, these results are presented as a function of centrality for six centrality classes in the range 0%–60%. As was done for the 0%–30% centrality class in Ref. [1], the source parameters of the Lévy distribution (λ , R , α) are measured.¹ The centrality and the transverse-mass dependence of the Lévy-fit parameters are characterized with simple, theoretically

¹Note that followed are the conventions introduced in Refs. [10] and [11–13], which is based on a book by Nolan [14] on univariate Lévy-stable source distributions, including also multivariate but symmetric Lévy-stable distributions.

and empirically motivated fit functions. The centrality dependence of the parameters of these functions is investigated in detail.

The structure of this paper is as follows: Sections II and III present the PHENIX experimental setup and the selection of the data sample, respectively. Section IV explores the procedure of measurement and fitting of the two-pion correlation function. Section V discusses the systematic uncertainties. Section VI presents the extracted Lévy parameters of the source as a function of centrality. Section VII discusses Monte Carlo simulations and some of the possible physics interpretations of these results, which have a strong exclusion power due to their high precision. Section VIII summarizes and concludes. Finally an Appendix details our Monte Carlo simulations to interpret the PHENIX data.

In particular, PHENIX data on the transverse-mass and centrality dependence of the Lévy intercept parameter $\lambda(m_T)$ are compared to centrality-dependent Monte Carlo simulations of resonance decay chains. In all but the most peripheral centrality class (50%–60%), the Monte Carlo simulations are found to be inconsistent with the measurements unless a significant reduction of the in-medium mass of the η' meson is included. In each centrality class, the best value of the in-medium η' mass is determined from χ^2 and confidence level (CL or p value) maps, based on a comparison of PHENIX data and Monte Carlo simulations. The resulting values of in-medium modified η' masses are compared to the mass of the η meson as well as to several theoretical predictions that consider restoration of $U_A(1)$ symmetry in hot hadronic matter as discussed in Secs. VII and VIII and then further detailed in the Appendix. Throughout this paper, units are used such that $\hbar = c = 1$. Also, the utilized fits represent the fitted data with CL in the statistically allowed $0.1\% \leq \text{CL} \leq 99.9\%$ interval.

II. THE PHENIX EXPERIMENT

The data used in this analysis are the same, apart from the centrality selection, as in the previous PHENIX Lévy HBT analysis with a 0%–30% centrality selection [1]. The PHENIX experimental apparatus relevant to this analysis is thus also the same. Briefly, the PHENIX detector is subdivided into the central-arm spectrometer (covering $2 \times 90^\circ$ azimuthal and $|\eta| \leq 0.35$ pseudorapidity acceptance), which is used here to focus mainly on hadron, electron, and photon identification and measurement. In the forward direction for each beam, two muon-arm spectrometers are used to focus mainly on identification and measurement of muons. There are also various event-characterization and triggering detectors in place. Of particular benefit here is good identification and measurement of charged pions. Ref. [1] provides further details.

III. DATA SAMPLE

The data sample used in this analysis comprises Au + Au collisions recorded by the PHENIX detector at $\sqrt{s_{NN}} = 200$ GeV in 2010. The minimum-bias data sample contains ≈ 7.3 billion events which is reduced to ≈ 4.4 billion with the 0%–60% centrality selection. The centrality dependence of

the transverse-mass trends is explored here in term of the numbers of participants (N_{part}), which was determined via Glauber-model calculations by the PHENIX experiment based on Ref. [15].

The present analysis shares almost all details with the previous analysis of Ref. [1], including using the word “cuts” to refer to selection criteria. The similarities and differences between the current and previous analyses are detailed below. In the present analysis, the Lévy fit parameters are determined in 23 bins of transverse mass from 0.248 to 0.876 GeV and in six, 10% wide centrality bins in the range of 0%–60%. Well-measured tracks are selected using the same single-track cuts as in Ref. [1]. The event-selection criteria (except the centrality selection) and the particle-identification techniques are also the same as in Ref. [1]. The single-track cuts and their variations are considered as sources of systematic uncertainties, as described in Sec. V.

The particle identification (PID) of pions is based on time-of-flight information and the path length information given by the track model. As in Ref. [1], a general cut on transverse momentum, $p_T > 0.16$ GeV, is applied to all pions. The cuts used in the PID are also considered as sources of systematic uncertainties.

In addition to the cuts on single tracks, pair cuts are imposed to minimize two-track effects: track merging and track splitting. Merging occurs when two tracks are so close to each other that the reconstruction algorithm considers them to be one track. Splitting is the opposite of the merging effect: one track is falsely reconstructed as two. These ambiguous pairs can be removed from the sample by geometrical cuts on their $\Delta\varphi$ - Δz plane, where $\Delta\varphi$ denotes the azimuthal angle difference of the hit positions and Δz is the difference of the z coordinates of the pair, as determined by the drift chambers (DCs), lead-scintillator electromagnetic calorimeter (PbSc), and time of flight (TOF) in the east and west arms of the PHENIX spectrometer.

These pair cuts were carefully investigated in the previous Lévy analysis [1]. However, due to the different centrality selections, the pair-cut settings here are slightly modified. The pair cuts are defined in the $\Delta\varphi$ - Δz plane as

$$\Delta\varphi > \Delta\varphi_0 \left(1 - \frac{\Delta z}{\Delta z_0}\right) \quad \text{and} \quad \Delta\varphi > \Delta\varphi_1 \quad (\text{DC and PbSc}), \quad (1)$$

$$\Delta\varphi > \Delta\varphi_0 \left(1 - \frac{\Delta z}{\Delta z_0}\right) \quad (\text{TOF east}), \quad (2)$$

$$\Delta\varphi > \Delta\varphi_0 \quad \text{and} \quad \Delta z > \Delta z_0 \quad (\text{TOF west}). \quad (3)$$

The default values of the $\Delta\varphi_0$, $\Delta\varphi_1$, and Δz_0 can be found in Table I, where also listed are the alternative values, which are used in Sec. V to determine the systematic uncertainties.

As in Ref. [1], in addition to these cuts, if multiple tracks are found that are associated with hits in the same tower of the PbSc slat of the TOF east, or strip of the TOF west detector, then all but one (randomly chosen) are removed. This ensures that no ghost tracks remain in the sample after the above-mentioned pair cuts.

TABLE I. The values of the coordinates for the pair-selection (cuts) criteria and the alternative values used to determine systematic uncertainties.

Pair cuts	DC			TOF east		TOF west		PbSc		
	$\Delta\varphi_0$ (rad)	Δz_0 (cm)	$\Delta\varphi_1$ (rad)	$\Delta\varphi_0$ (rad)	Δz_0 (cm)	$\Delta\varphi_0$ (rad)	Δz_0 (cm)	$\Delta\varphi_0$ (rad)	Δz_0 (cm)	$\Delta\varphi_1$ (rad)
Default cut settings	0.12	8.0	0.017	0.12	12	0.075	14.0	0.12	16	0.015
Loose drift chamber cut	0.11	7.0	0.016	0.12	12	0.075	14.0	0.12	16	0.015
Strict drift chamber cut	0.13	9.0	0.018	0.12	12	0.075	14.0	0.12	16	0.015
Loose ID detector cuts	0.12	8.0	0.017	0.11	11	0.070	13.0	0.11	15	0.013
Strict ID detector cuts	0.12	8.0	0.017	0.13	13	0.080	15.0	0.13	17	0.017

Within statistical uncertainties of the fit parameters, using only positive pions gives results consistent with using only negative pions. Consequently, only the results from combined fits to both $(++)$ and $(--)$ charge combinations of identified pion pairs are presented here.

IV. MEASURING AND FITTING THE TWO-PARTICLE CORRELATION FUNCTION

A. Measuring the correlation function

In principle, a detailed shape analysis of the two-particle correlation functions could require three-dimensional measurements, but the lack of statistical precision could make such measurements impractical. The goal here is to obtain precise results in several transverse-mass and centrality bins. Thus, the correlation functions use a single variable Q . PHENIX preliminary results on a multivariate Lévy analysis are also available [16], but go beyond the scope of the present paper.

Our relative momentum variable Q is chosen as the modulus of the three-momentum difference in the longitudinal comoving system (LCMS) [17]. With the Bertsch-Pratt decomposition of the relative momentum to the “side, out, and long” components [18,19] in the LCMS, the variable can be written as

$$Q \equiv |\mathbf{q}_{\text{LCMS}}| = \sqrt{q_{\text{out,LCMS}}^2 + q_{\text{side,LCMS}}^2 + q_{\text{long,LCMS}}^2}. \quad (4)$$

The motivation for this variable comes from the experimental observation [10] that in a Gaussian three-dimensional BEC analysis, a quadratic sum appearing in the BEC functions provided nearly equal BEC radius parameters, i.e., $R_i \approx R$ for $i = \text{side, out, long}$. Thus, this quadratic sum can be simplified

as

$$\sum_{i=\text{side, out, long}} R_i^2 q_i^2 \approx R^2 \left(\sum_{i=\text{side, out, long}} q_i^2 \right) = R^2 Q^2, \quad (5)$$

which depends on the relative momentum only through the one-dimensional variable Q , the magnitude of the relative momentum of the pair in the LCMS. It was shown in Ref. [10] that the same quadratic sum appears in three-dimensional symmetric Lévy-stable distributions. So our choice for the relative momentum variable Q actually includes, as a special case, the Lévy analysis of three-dimensional Bose-Einstein correlation functions under the condition that the HBT radii in all three spatial dimensions are equal within experimental uncertainties. An approximate equality of the transverse-mass and centrality-dependent BEC radii was found in several experiments, e.g., S + Pb collisions at the CERN Super Proton Synchrotron energies by the NA44 Collaboration [20], and in $\sqrt{s_{NN}} = 200$ GeV Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV by the PHENIX and STAR Collaborations in the Gaussian approximation in Refs. [21,22]. PHENIX preliminary multivariate Lévy fits in Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV are also used in Ref. [16]. Especially note that the PHENIX preliminary three-dimensional Lévy-analysis for Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV finds that the three BEC radii in the LCMS are approximately equal (except perhaps at small transverse mass).

Furthermore, the radii are approximately equal to the radius found in a one-dimensional analysis [16], i.e., the Lévy source is indeed approximately spherical in the LCMS. Hence, this one-dimensional Lévy analysis, which focuses on the centrality and transverse-mass dependence, can be considered also as a reasonable approximation of a three-dimensional Lévy analysis. A detailed justification of this choice of Q can also be found in Ref. [1], where the same variable is used.

Our choice of Q can be written using the measured momenta of identified pions, $p_i^\mu = (E_i, p_{i,x}, p_{i,y}, p_{i,z})$. For $i = 1, 2$,

$$Q \equiv |\mathbf{q}_{\text{LCMS}}| = \sqrt{(p_{1,x} - p_{2,x})^2 + (p_{1,y} - p_{2,y})^2 + \frac{4(p_{1,z}E_2 - p_{2,z}E_1)^2}{(E_1 + E_2)^2 - (p_{1,z} + p_{2,z})^2}}, \quad (6)$$

where the z axis coincides with the beam axis.

The correlation function is measured as

$$C_2(Q) = \frac{A(Q) \int_{Q_{\text{int,min}}}^{Q_{\text{int,max}}} B(m_T, Q)}{B(Q) \int_{Q_{\text{int,min}}}^{Q_{\text{int,max}}} A(m_T, Q)}, \quad (7)$$

where $A(Q)$ is the actual Q distribution of pairs of identical pions coming from the same event, $B(Q)$ is the corresponding distribution of pairs of identical pions from different events, and $Q_{\text{int,min}}$ and $Q_{\text{int,max}}$ denote the lower and upper bound of the integrals. Note that $A(Q)$ may contain correlations in addition to the quantum-statistical correlation due to the indistinguishability of the measurement of identical bosons. Other correlations could be the consequence of conservation laws, resonance decays, kinematics, detector acceptance effects, etc.

Ideally, the background $B(Q)$ distribution, or reference sample, is identical to the sample of like-charged pion pairs in all respects, except for the Bose-Einstein interference effect itself [23]. As such a $B(Q)$ distribution does not exist in nature, it has to be generated using approximation schemes. The choice here is explained below. Careful testing showed stability over the choice of the fit range and systematic variations of acceptance, PID, and other cuts detailed among the systematic uncertainties.

In choosing the background distribution $B(Q)$, each member of the pairs is selected from different events. Hence, $B(Q)$ contains only trivial kinematic correlations between independent particles that are distributed with the same single-particle spectra and in the same kinematic range. The background distribution is affected by the centrality-selection cuts and detector-acceptance effects, rapidity, and transverse-mass cuts. In $A(Q)$ the pairs are correlated not only due to Bose-Einstein correlations, but also by other initial- or final-state interactions. Such effects from interactions include branching processes of jets, hadronization effects, resonance decays, energy, and momentum conservation laws, possible other kinematic effects, such as elastic scattering, Coulomb, and strong final-state interactions.

In general, $A(Q|MC)$ and $B(Q|MC)$ should be generated from Monte-Carlo simulations that describe these single-particle data and all other non-BEC correlation effects. These Monte-Carlo distributions are then used to correct $B(Q)$ for its lack of non-BEC correlations, as was done, for example, in $e^+ + e^-$ collisions by the L3 Collaboration [2]. However, in high-energy heavy-ion physics there is a good reason why this complicated and Monte Carlo dependent procedure is not necessary. Namely, the presence of high-multiplicity events, with $\langle n \rangle \gg 1$, where $\langle n \rangle$ is the mean charged-particle multiplicity at midrapidity. The usual kinematic correlations, which are due to resonance decays or conservation laws, are proportional to the mean multiplicity, $\langle n \rangle$, while Bose-Einstein correlations grow with the mean number of pairs $\langle n(n-1) \rangle$. Thus in high-multiplicity heavy ion collisions, Bose-Einstein correlations (together with Coulomb and strong final-state interactions that also grow proportionally with the number of pairs) outnumber all the other correlations. Hence, the MC approach can be safely abandoned. In the expression $A(Q)/B(Q)$ the Bose-Einstein, Coulomb, and strong final-state correla-

tions dominate, and the other correlations are suppressed as $\langle n \rangle / \langle n(n-1) \rangle \propto 1/\langle n \rangle \ll 1$.

In BEC measurements using charged particles, Coulomb repulsion modifies the correlation function at low Q values creating the ‘‘Coulomb hole’’ [17]. To account for the final-state Coulomb interaction, the Coulomb wave function is integrated over the source of pions. Strong final-state interactions could also be taken into account using phase shifts that modify the Coulomb wave function. However, the strong final-state interaction of the pion pairs is small compared to the experimental precision [17,24]. All the Lévy fits here (without corrections for the strong final-state interactions) are of good quality with $CL \gg 0.1\%$. The strong final-state interactions are found not to affect the quality of Lévy fits, nor to change the parameters significantly [17,24]. Therefore, such corrections are not considered in this analysis. In contrast, the effect of Coulomb interactions is clearly visible as a Coulomb hole at small values of Q . The fitting function is appropriately modified, as detailed Sec. IV B 2.

B. The fitting function

1. The Lévy shape

For the shape of the correlation function, it is not possible to *a priori* assume or know, what is or what should be the shape-model of the Bose-Einstein correlation functions. Measurements of this quantity test the hypothesis of being consistent with a Lévy shape. Using the plane-wave approximation and assuming a spherically symmetric, Lévy-type source, the two-particle BEC function has the simple form

$$C_2(Q) = 1 + \lambda \exp[-Q^\alpha R^\alpha], \quad (8)$$

where λ is the strength of the correlation, R is the scale parameter in physical units, and α is the Lévy index of stability [10]. This hypothesis is tested not only by successful fits to the PHENIX data with $CL \gg 0.1\%$, which are stable and robust with systematic variations of particle identification, acceptance, and pair cuts, as detailed in Section V. The validity of the Lévy source shape is also checked by employing a Lévy expansion technique [12,13,25]. This method utilizes a complete set of polynomials that are orthonormal with respect to a Lévy weight function [12], so this method is able to characterize, model independently, any deviation from a Lévy-stable source shape. To first order of the expansion, no significant deviation from the Lévy shape is found in any of the centrality and transverse-mass ranges investigated here.

The relevance of stable distributions to the analysis of (Coulomb-corrected) Bose-Einstein correlations was studied in Ref. [10], following the general mathematical ideas summarized in Ref. [14]. Univariate stable distributions are usually characterized by the Fourier transform of their density distributions that are called the characteristic functions. Following the convention of the previous analysis [1] and the theoretical paper where the idea first appeared [10], the $S(\alpha, \beta, \gamma, \delta; 1)$ notation is used. The parameter α is the Lévy index of stability (or characteristic exponent) that is limited to the domain $0 < \alpha \leq 2$. The asymmetry parameter β is limited to the domain $-1 \leq \beta \leq 1$. The scale parameter γ is nonnegative and the location parameter δ can be any real

number with $-\infty < \delta < \infty$. The book by Nolan [14] details exhaustively the $\alpha = 1$ special case and the ubiquitous nature of Lévy-stable source distributions. It provides illustrations of the stable densities in the $S(\alpha, \beta, \gamma, \delta; 1)$ parametrization.

The asymmetric τ -model [26] was utilized successfully by the L3 Collaboration [2] to interpret two-jet data. Two of the experimentally testable predictions that are directly related to asymmetrical ($\beta \neq 0$) source distributions were investigated. The first indication of such an asymmetry was related to a dip (an anticorrelated region) in the two-particle Bose-Einstein correlation function. Observance of such a dip in PHENIX data would be a strong experimental indication of the presence of an asymmetric source. However, such a dip is not observed in any of the investigated transverse-mass or centrality bins. Secondly, if the asymmetric τ -model has relevance in the present analysis, the BEC function would depend on any relative-momentum component only through the invariant momentum q_{inv} . In this case, the BEC functions increase with decreasing values of q_{inv} , even if Q is kept constant. Such behavior is not observed here; on the contrary, the BEC function is seen to increase with decreasing Q , even if q_{inv} is kept at constant values. Therefore, a vanishing asymmetry parameter, $\beta = 0$, is assumed.

The correlation function is thus based on the assumption of the simplest case of univariate and symmetric ($\beta = 0$) stable distributions. The scale parameter γ is replaced by the physical parameter R [10]. In the standardized notation of Nolan [14], this corresponds to the $(\alpha, \beta = 0, \gamma = R/2^{\frac{1}{\alpha}}, \delta; 1)$ convention. The Fourier-transformed source-density distribution has a simple form, $f(q) = \exp(iq\delta - \frac{1}{2}|qR|^\alpha)$, and its modulus square leads to the simple form of the BEC function of Eq. (8). Note that the correlation function does not depend on δ .

The relationship with the Gaussian source distribution is also apparent in Eq. (8) as it corresponds to the $\alpha = 2$ special case. As detailed in Ref. [10], the physical-scale parameter R corresponds only in this $\alpha = 2$ special case to the

root-mean-square value of the source. For all $0 < \alpha < 2$, the root-mean-square value is divergent, as is well known for the Cauchy or Lorentzian special case ($\alpha = 1$). In the $\alpha = 1$ case the physical Lévy-scale parameter R corresponds to the half width at half maximum (HWHM) of the source distribution. In fact for all values of α , R is proportional to the HWHM, the constant of proportionality depending on α [27]. Another notable property of the Lévy-stable source distributions is that in the $\alpha < 1$ cases, even the first moment of the source distribution is divergent. In high energy particle and nuclear physics, the R values correspond to a few femtometers [1,2].

2. The Coulomb interaction

As in Ref. [1], the Coulomb final-state interaction is characterized using the Sinyukov-Bowler method [28,29]. This method corresponds to the integration of the two-particle Coulomb wave function [30] for a core-halo type of particle-emitting source [31,32]. However, such a Coulomb wave-function integration cannot be performed analytically in the case of a Lévy source. Hence, numerical approaches are needed. The previous Lévy BEC analysis used an iterative method based on a numerical table which contains the values of the integral for a range of values in the parameters and in the variable Q . The details can be found in Ref. [1]. In this paper a parametrization [33,34] is based on the aforementioned numerical table, which is considerably faster.

The momentum difference variable of the Coulomb correction is the invariant four-momentum difference q_{inv} rather than the variable Q , which is used in the present analysis. Neglecting the difference between the two variables could introduce a systematic uncertainty of $\approx 5\%$ [1]. In the present analysis, the difference is determined by measuring the actual pair distribution in both Q and q_{inv} , $A(Q, q_{\text{inv}})$. Using this two-dimensional distribution, the Coulomb correction is incorporated with a weighted average.

The final form of the fitting function is then

$$C_2(Q; \lambda, R, \alpha, N, \varepsilon) = 1 - \lambda + \lambda C_2^{(0)}(Q; R, \alpha, N, \varepsilon) w(Q; R, \alpha) \\ \text{with } C_2^{(0)}(Q; R, \alpha, N, \varepsilon) = [1 + \exp(-R^\alpha Q^\alpha)]N(1 + \varepsilon Q) \\ \text{and } w(Q; R, \alpha) = \frac{\sum_k A(Q, q_{\text{inv},k})K(q_{\text{inv},k}; R, \alpha)}{\sum_k A(Q, q_{\text{inv},k})}, \quad (9)$$

where $K(q_{\text{inv},k}; R, \alpha)$ is the Coulomb correction given by the parametrization [33,34]. In the definition of the weight function $w(Q; R, \alpha)$, the index of summation k runs over those bins in q_{inv} for a given value of Q , where the number of actual pairs after the two-track cuts is nonvanishing. This summation thus averages over the q_{inv} -dependent Coulomb correction for a Lévy-type source characterized by R and α , in q_{inv} bins only where $A(Q, q_{\text{inv},k})$ is nonzero. This method generalizes on Eqs. (2) and (3) of Ref. [21] for a Lévy-shaped Bose-Einstein correlation function and also accounts for the q_{inv} dependence of the Coulomb correction. The large Q behavior was found to be consistent with a linear

function, which is characterized with the functional form of $N(1 + \varepsilon Q)$.

V. SYSTEMATIC UNCERTAINTIES

Nine sources of systematic uncertainties are investigated: the single track and the pair cuts, the choice of the arm of the PHENIX detector (because the arrangement is not symmetric), the fit range of the correlation functions, and the Coulomb-correction method. The systematic uncertainties of the results are estimated by varying one setting at a time, while keeping the others at their default values. For the cuts this

means applying stricter or looser criteria. This is the same approach as used in Ref. [1], which one can see for details. For the pair cuts the values listed in Table I are used rather than those of Ref. [1]. For the choice of the arm of PHENIX the variation uses only one arm rather than both. The sensitivity of the results to the fit range was investigated by adding or leaving out one bin from the beginning (Q_{\min}) or the end of the fit range (Q_{\max}). The results do not vary significantly with the variation of Q_{\max} , therefore this source does not contribute to the systematic uncertainty.

As a default Coulomb-correction method, the parametrization detailed in Refs. [33,34] is used. A recent theoretical

investigation [35] suggested several alternative methods. The systematic uncertainty of the Coulomb correction is taken as the difference between our approach and the most realistic variant of the Coulomb corrections mentioned in Sec. 3.2 of Ref. [35], namely the sixth one in the enumeration. The total systematic uncertainty is estimated with a standard statistical approach. The individual contributions are summed quadratically, while both the statistical uncertainties σ_{stat} and the correlations between the uncertainties (denoted by ρ) are considered. Thus, the final systematic uncertainty (σ_{syst}) is expressed for a parameter, p (default cut denoted by p_{def} and the i th alternative cut by $p_{\text{cut},i}$), with the following form:

$$\sigma_{\text{syst,tot}}^2 = \sum_i [(p_{\text{def}} - p_{\text{cut},i})^2 - \sigma_{\text{stat}}^2(p_{\text{def}}) - \sigma_{\text{stat}}^2(p_{\text{cut},i}) + 2\rho_i \sigma_{\text{stat}}(p_{\text{def}}) \sigma_{\text{stat}}(p_{\text{cut},i})]. \quad (10)$$

The ρ_i correlation between the uncertainties can be estimated with a data-driven method by measuring the numbers of pairs yielded with the different settings

$$\rho_i = \begin{cases} \sqrt{\frac{N_{\text{def}}}{N_{\text{cut},i}}} & \text{if } N_{\text{def}} < N_{\text{cut},i}, \\ \sqrt{\frac{N_{\text{cut},i}}{N_{\text{def}}}} & \text{if } N_{\text{cut},i} < N_{\text{def}}. \end{cases} \quad (11)$$

The ρ_i correlation coefficients typically are near unity, except for the third-pad-chamber matching cut and the arm settings, for which $\rho_i \approx 0.6$.

VI. CENTRALITY AND TRANSVERSE-MASS-DEPENDENT RESULTS

The transverse-mass (m_T) and centrality dependence (expressed in terms of the average number of participants, N_{part}) of the Lévy parameters analyzed and are investigated along with their theoretically motivated combinations. The transverse mass is defined as $m_T = \sqrt{m_\pi^2 + K_T^2}$, where $K_T = \sqrt{K_x^2 + K_y^2}$ is the transverse component of the average momentum of the pair $K = 0.5(p_1 + p_2)$ and m_π is the pion mass. The number of participants N_{part} was determined via Glauber-model calculations based on Ref. [15]. The centrality dependencies of the Lévy parameters are characterized by theoretically or empirically motivated functions. The fits are found to represent the data in each of the investigated centrality-class and transverse-mass bins, the confidence levels are in the statistically acceptable $0.1\% \leq \text{CL} \leq 99.9\%$ region. The fitted correlation functions are very similar to our results obtained in the 0%–30% centrality class. For example, a Lévy fit to our Bose-Einstein correlation data was published in Ref. [1].

A. The centrality and transverse-mass dependence of the fit parameters

The dependence of the physical parameters λ , R , and α on centrality and transverse mass, m_T is determined. The parameters of the linear background in Eq. (9) are found to be $N \approx 1$ and $\varepsilon \approx 0$. In particular, the maximum of the modulus of the

coefficient of linearity is found to be $\max(|\varepsilon|) = 0.085 \text{ GeV}^{-1}$ and the average value is $\langle |\varepsilon| \rangle = 0.021 \pm 0.001 \text{ GeV}^{-1}$. The overall normalization coefficient N has a maximal deviation from unity of $\max(|N - 1|) = 0.015$, while its average deviation from unity is $\langle |N - 1| \rangle = 0.0050 \pm 0.0001$.

The results of the fits for the 6 centrality and 23 m_T bins are shown in Figs. 1–3. The intercept or the correlation-strength parameter λ has a clearly observable suppression below $m_T \leq 0.5 \text{ GeV}$ average pair transverse mass. Above that it saturates at a centrality-dependent value. This low- m_T suppression, which is characterized in more detail in Sec. VIB is observed in all centrality bins; however, in the most peripheral (50%–60%) centrality class, this observation starts to be limited by statistics. This is one of the reasons that no data are shown for the most peripheral, 60%–95% centrality class.

The saturation value of the intercept parameter could depend on background processes which are not in the scope of the present analysis. As was shown in Ref. [36], the value of the intercept parameter is lower for Gaussian, intermediate for an Edgeworth expansion [25], and higher for an exponential shape, which seems to be a systematic effect that depends on how the functional form of the Bose-Einstein correlation function extrapolates the data to the $Q = 0$ limit. This systematic effect, an overall vertical uncertainty, can however be removed by normalizing to the saturation value of $\lambda(m_T)$ at large m_T , following Refs. [1,36]. The saturation value is taken as the average value of $\lambda(m_T)$ in the interval $0.45 \leq m_T \leq 0.9 \text{ GeV}$ and is denoted by λ_{max} . This centrality-independent range is the same for each of the six centrality classes considered in this paper. However, from the similar range considered in Ref. [1] for the 0%–30% centrality selection, this range is slightly modified due to the different centrality classes considered. This range modification shifts the central values of λ_{max} slightly, but the modification is within one standard deviation, which is within the uncertainties given in Ref. [1]. However, the value of λ_{max} does depend on centrality. The resulting $\lambda(m_T)/\lambda_{\text{max}}$ ratio is shown in Fig. 4.

Radial-flow effects are known to be strongly centrality dependent, and hence are expected to significantly influence both $\lambda(m_T)$ and λ_{max} . Such an expectation is shown, for example, in Ref. [37]. If there is large radial flow, the decay

products of η' are concentrated at high m_T [compensating p_T sharing between the several daughter pions, resulting in flat $\lambda(m_T)/\lambda_{\max}(m_T)$], while if the radial flow is low (such as in the peripheral event classes), then these decay products accumulate at low m_T [resulting in a dip in $\lambda(m_T)/\lambda_{\max}$ at low m_T]. The two effects, the p_T increasing radial flow and the p_T sharing between the several pions from the $\eta' \rightarrow \eta + \pi^+ + \pi^- \rightarrow (\pi^0 + \pi^+ + \pi^-) + \pi^+ + \pi^-$ decay chain compensate one another at transverse velocity $\langle u_T \rangle \approx 0.5$, as shown in Ref. [37]. This expectation is cross-checked with the Monte Carlo simulations detailed in the Appendix, where the simulation-based centrality dependence of $\lambda(m_T)/\lambda_{\max}$ (without any in-medium η' mass modification) is also shown.

In Fig. 4, an approximate centrality independence of the characteristics of the suppression can be qualitatively observed. To test this observation quantitatively, a χ^2 test was employed to obtain conservative values of the CL using only statistical uncertainties. For each of the 15 possible pairs of centrality classes, the χ^2 is calculated under the hypothesis that their λ/λ_{\max} distributions are identical. The resulting CL values are all in the range 0.15%–89.5%, which does not reject the hypothesis.

The experimental result on the centrality independence of $\lambda(m_T)/\lambda_{\max}$ scaling is an unexpected, rather surprising obser-

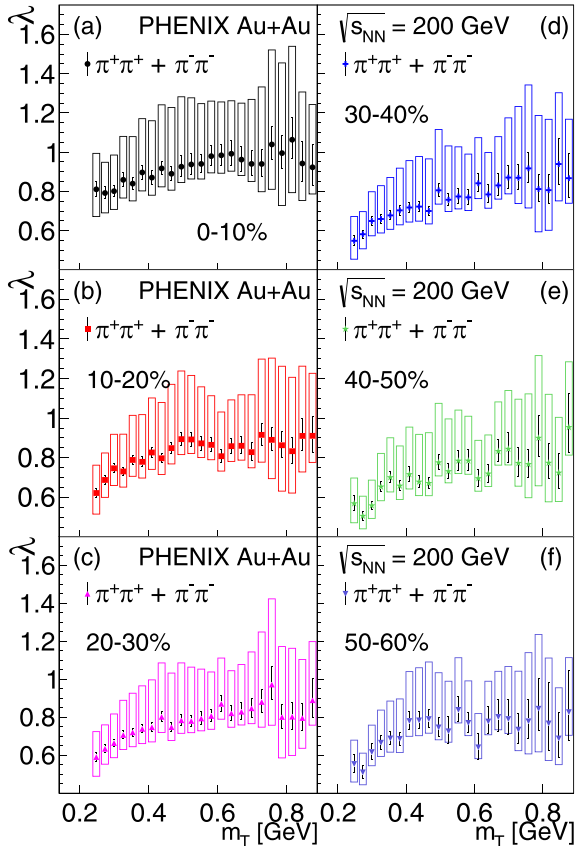


FIG. 1. The transverse-mass dependence of the correlation-strength parameter λ in six centrality bins obtained from Lévy fits with Eq. (9). The central values are shown with dots, statistical uncertainties are indicated by vertical lines, while boxes are used to illustrate the systematic uncertainties.

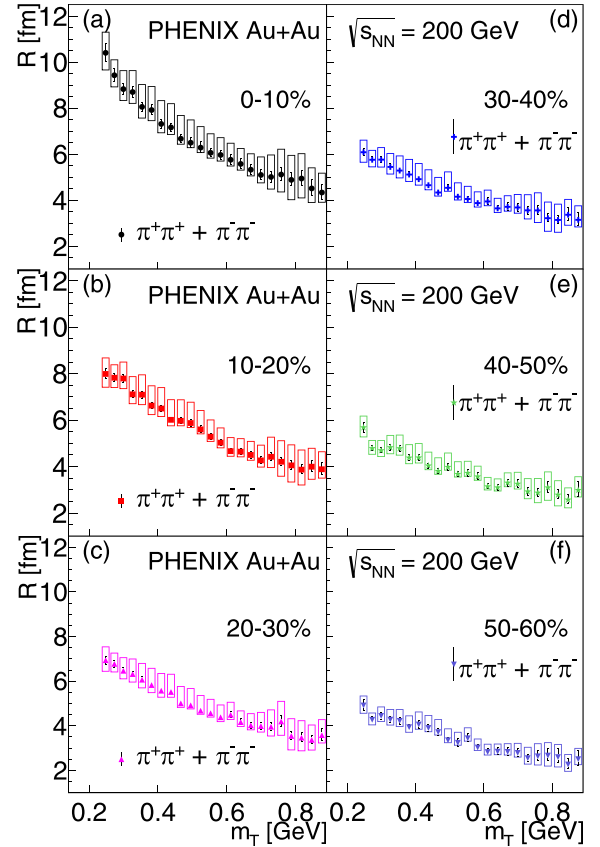


FIG. 2. The transverse-mass dependence of the Lévy-scale parameter R in six centrality bins obtained from Lévy fits with Eq. (9). The central values are shown with dots, statistical uncertainties are indicated by vertical black lines, while boxes are used to illustrate the systematic uncertainties.

vation. One possible explanation for such a scaling behavior is given in Sec. VII, based on the Monte Carlo simulations detailed in the Appendix. A quantitative analysis of these $\lambda(m_T)/\lambda_{\max}$ measurements is presented in Sec. VIB.

The Lévy-scale parameter, introduced in high-energy particle and nuclear physics as the physical-scale parameter R in Ref. [10], decreases with m_T for all centrality bins, as can be seen in Fig. 2. Analytic hydrodynamic calculations [28,39,64] predict that, in the $\alpha = 2$ special case, $R^{-2}(m_T)$ depends linearly on m_T , which can be parametrized as

$$\frac{1}{R^2} = Am_T + B \quad (12)$$

as shown in Fig. 5. The centrality dependence of the slope parameter A and the intercept parameter B is presented in Sec. VIB.

Figure 3 shows the measured values of the Lévy index α . In each centrality bin α appears to be independent of m_T . In each centrality class, the data of Fig. 3 can be well represented by the m_T -averaged values, α_0 . As indicated in Fig. 10, these values are found to depend on centrality. The measured values of α are significantly lower than the Gaussian case of $\alpha = 2$, and are significantly higher than the exponential/Cauchy

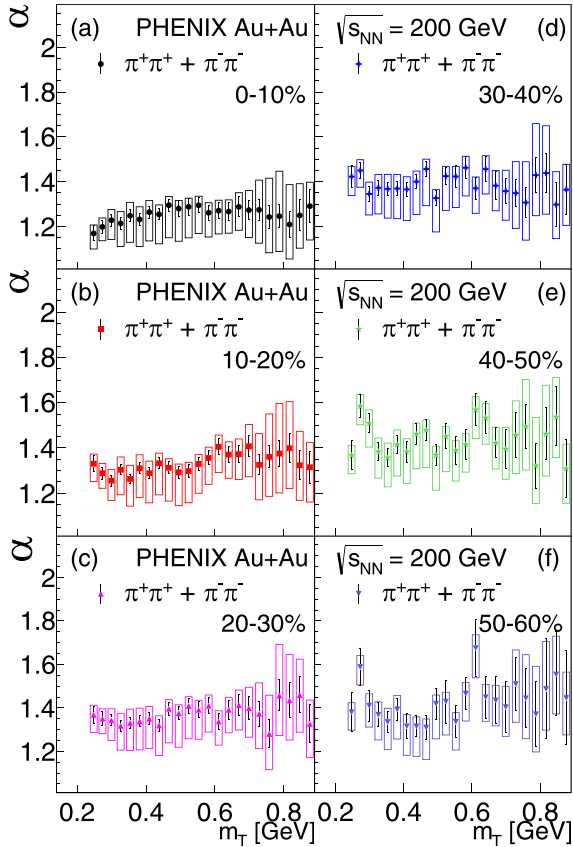


FIG. 3. The transverse-mass dependence of the Lévy-index of stability parameter α , shown in six centrality bins obtained from Lévy fits with Eq. (9). The central values are shown with dots, statistical uncertainties are indicated by vertical black lines, while boxes are used to illustrate the systematic uncertainties.

case of $\alpha = 1$. The centrality dependence of α_0 is discussed in Sec. VIB.

Similar to earlier PHENIX studies in the 0%–30% centrality-class range [1], strong correlations are observed between the parameters λ , R , and α . However, a specific combination of these three Lévy parameters, namely

$$\frac{1}{\hat{R}} = \frac{\lambda(1 + \alpha)}{R}, \quad (13)$$

shown in Fig. 6, seems to be less correlated with the direct fit of Lévy parameters λ , R , α as compared to the correlations among them, which is very similar to the correlation plots published in Ref. [1] for the 0%–30% centrality class. The correlation coefficients between the parameters λ , R , and α were determined in the Lévy fits and were considered when evaluating \hat{R}^{-1} .

In Ref. [1] \hat{R}^{-1} was found to depend approximately linearly on m_T and was parameterized in the 0%–30% centrality class as

$$\frac{1}{\hat{R}} = \hat{A}m_T + \hat{B}. \quad (14)$$

Figure 6 shows a more detailed centrality-dependent analysis. The centrality dependence of the \hat{A} slope parameter and \hat{B}

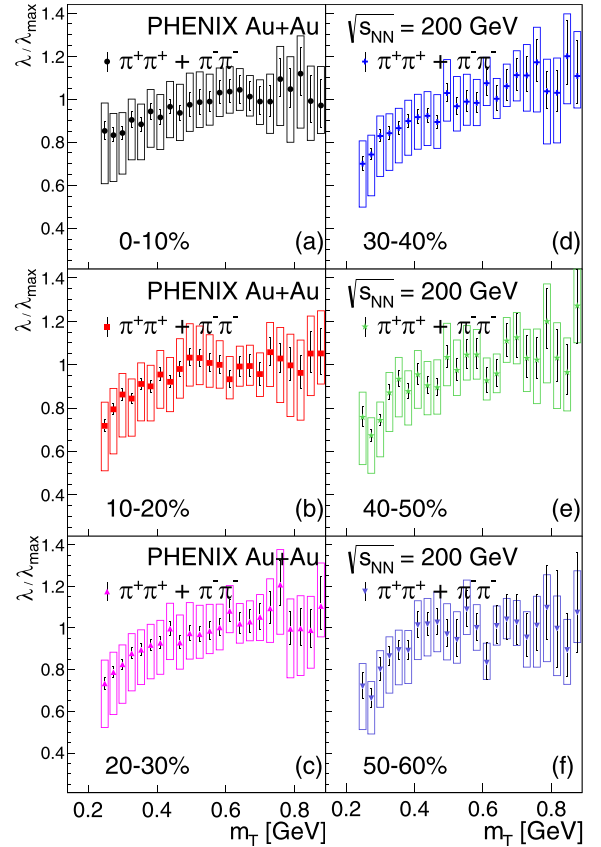


FIG. 4. The transverse-mass dependence of the normalized correlation-strength parameter λ/λ_{\max} for six centrality intervals obtained by rescaling Fig. 1 with a centrality-dependent λ_{\max} defined as the average value of $\lambda(m_T)$ in the $0.45 \leq m_T \leq 0.9$ GeV interval. The central values are shown with dots, statistical uncertainties are indicated by vertical black lines, while boxes are used to illustrate the systematic uncertainties. For each centrality bin the data are fitted with the Gaussian function of Eq. (15). This parametrization can describe the data and is discussed in detail in Sec. VIB. The fit parameters with the statistic and systematic uncertainties are shown in Fig. 7.

intercept parameter is presented in Sec. VIB. The linearity of \hat{R}^{-1} is broken at large m_T , which is likely due to the saturation of the λ parameter in that range. As noted in Ref. [1], the understanding of the approximate scaling properties of \hat{R} is another currently unsolved theoretical challenge.

B. The centrality dependence of the physical parameters

In this subsection, the centrality dependence of the physical parameters via theoretically or empirically motivated parametrizations are investigated. The centrality dependence of the parameters of these parameterizations is discussed here as a function of N_{part} . Quantitatively studying the suppression of $\lambda(m_T)/\lambda_{\max}$ at low m_T , uses fits of the phenomenological Gaussian parametrization, which was also used in Ref. [1],

$$\frac{\lambda}{\lambda_{\max}} = 1 - H \exp\left(-\frac{m_T^2 - m_\pi^2}{2\sigma^2}\right), \quad (15)$$

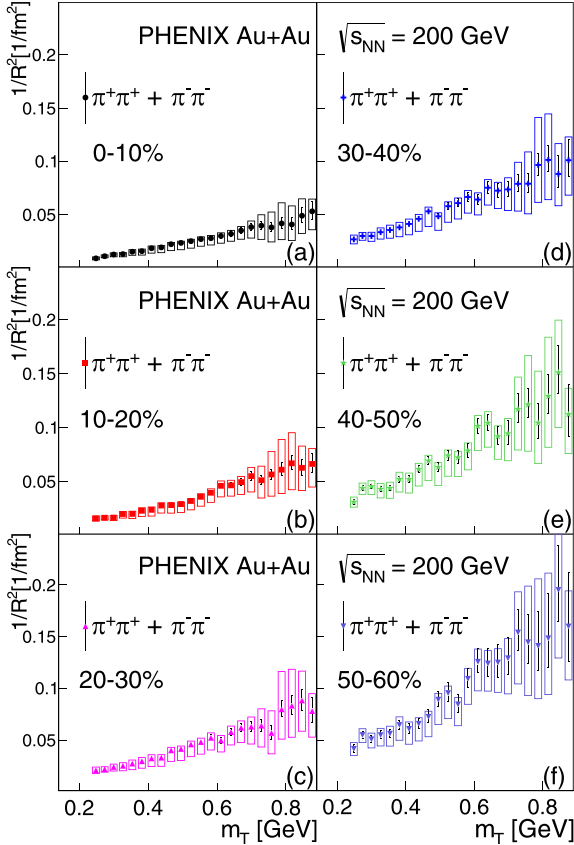


FIG. 5. The transverse-mass dependence of the inverse square of the Lévy-scale parameter R , shown in six centrality bins. The central values are shown with dots, statistical uncertainties are indicated by vertical black lines, while boxes are used to illustrate the systematic uncertainties. For each centrality bin the data are fitted with the linear function of Eq. (12). This parametrization can describe the data and is discussed in detail in Sec. VIB. The fit parameters with the statistic and systematic uncertainties are shown in Fig. 8.

where the parameter $H = H(N_{\text{part}})$ measures the depth of the suppression and parameter $\sigma = \sigma(N_{\text{part}})$ measures the width of this suppression. In principle, both H and σ could have a centrality or N_{part} dependence.

From the experimentally observed, and rather surprising, centrality-independent scaling of $\lambda(m_T)/\lambda_{\text{max}}$, the centrality independence of the H and the σ fit parameters is expected, as discussed in Sec. VIA. Indeed, both obtained parameters are consistent with the hypothesis of centrality independent values H_0 and σ_0 , as indicated in Fig. 7.

The affine linearity of the inverse square of the Lévy-scale parameter $R^{-2}(m_T)$, Eq. (12), is demonstrated in Fig. 5. The parameters of these fits as a function of N_{part} are shown in Fig. 8. The slope parameter A decreases with N_{part} . The intercept parameter B , which could be connected to the size of the source, is slightly negative within statistical uncertainties but is compatible with zero if systematic uncertainties are taken into account. Note that, in hydrodynamical calculations with $\alpha = 2$, the parameter B is dominated by the inverse of the squared geometrical size of the whole source: the smaller the

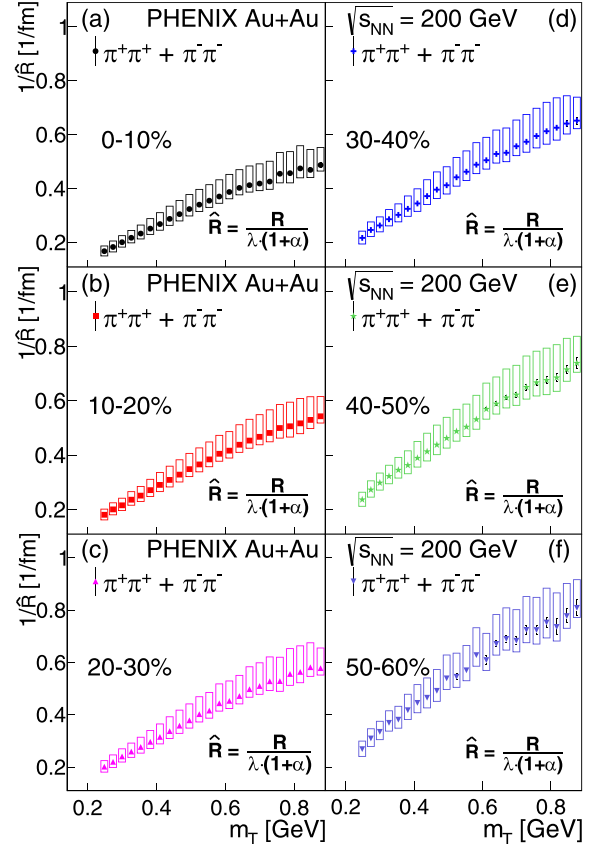


FIG. 6. The transverse-mass dependence of the inverse of the \hat{R} parameter, shown in six centrality bins. The central values are shown with dots, statistical uncertainties are indicated by vertical black lines, while boxes are used to illustrate the systematic uncertainties. For each centrality bin the data are fitted with the linear function of Eq. (14). This parametrization can describe the data and is discussed in detail in Sec. VIB. The fit parameters with the statistic and systematic uncertainties are shown in Fig. 9.

value of B , the larger the geometrical size of the source, while parameter A is dominated by radial-flow effects [28,38,39].

As mentioned in Sec. VIA, the affine linearity of $\hat{R}^{-1}(m_T)$, Eq. (14) is investigated, as a function of N_{part} . The centrality dependence of the slope and intercept parameters of these fits are shown in Fig. 9. The parameter \hat{A} exhibits a decreasing trend with increasing centrality similar to that seen for A in the parametrization of R^{-2} . Within statistical uncertainty, \hat{B} is compatible with a small positive number but both statistical and systematic uncertainties are too large to draw a strong conclusion from the present data.

As shown in Fig. 10(a), because the Lévy-exponent parameter α does not appear to depend on m_T (see Fig. 3), the m_T -averaged value, α_0 , can be used to determine the centrality dependence of α . A significant centrality dependence of α_0 is observed. The lowest value of α_0 occurs for the most central collisions. As the collisions become more peripheral α_0 saturates around ≈ 1.4 .

The parameter R is investigated in several m_T bins as a function of $N_{\text{part}}^{1/3}$. The linearity of this dependence, shown in

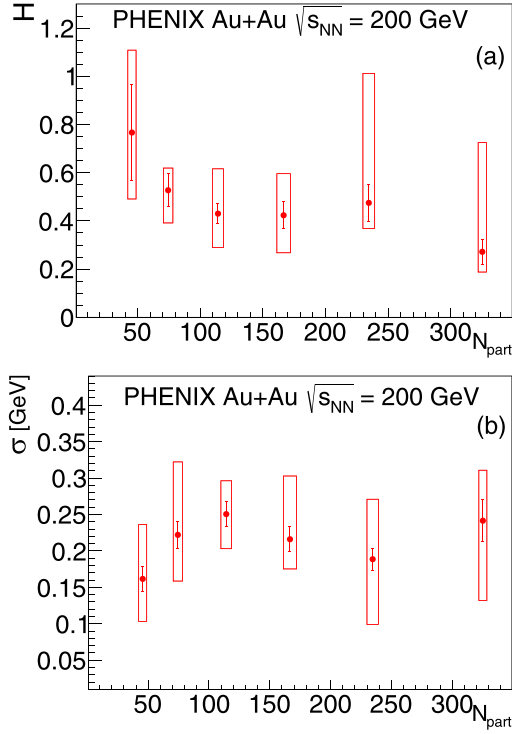


FIG. 7. The two parameters that characterize the suppression of the $\lambda/\lambda_{\max}(m_T)$ and are defined in Eq. (15), (a) the magnitude (H), and (b) the width (σ) are shown as functions of N_{part} . The central values are shown with dots, statistical uncertainties are indicated by vertical (red) lines, while boxes are used to illustrate the systematic uncertainties. Both panels (a) and (b) are consistent with a centrality-independent constant values H_0, σ_0 .

Fig. 10(b), suggests that the volume of the Lévy source, which is proportional to R^3 , is proportional to N_{part} . More detailed studies, investigating the centrality dependence in terms of constituent quark participants along the lines of Ref. [40], would be desirable, but go beyond the scope of the present analysis.

VII. COMPARISON WITH MONTE CARLO SIMULATIONS

The low- m_T decrease of the $\lambda(m_T)$ correlation-strength measurements [21,22,41] of charged pions is found in Ref. [37] to be an indirect signal of the in-medium mass reduction of the η' particles. The $\lambda(m_T)$ suppression at low transverse mass (m_T) has been investigated in Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV collisions by Monte Carlo simulations [42,43]. These calculations were also compared to the results of the previous PHENIX analysis of 0%–30% central Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV [1]. Although the measurements are performed in pair- m_T (cf. Sec. VI) ranges while the simulations are in the single-particle transverse mass $\sqrt{m^2 + p_T^2}$, this difference can be neglected as λ is a characteristic parameter of the two-particle correlation function at vanishing values of relative momentum, where $K = p$ and hence the single-particle transverse mass is equal to the pair m_T . This property is also utilized in the core-halo picture,

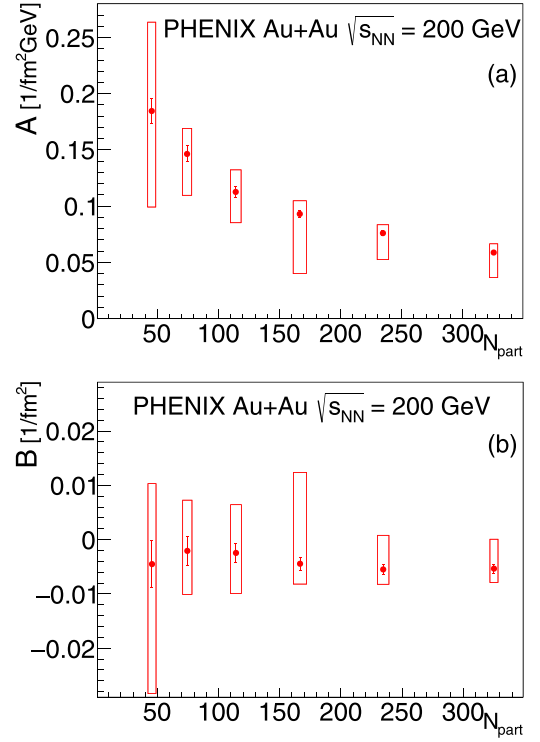


FIG. 8. The two parameters of the affine linear fit to the inverse square of the Lévy-scale parameter $1/R^2(m_T)$ are defined in Eq. (12), (a) the slope parameter (A), and (b) the intercept parameter (B) are shown as functions of N_{part} . The central values are shown with dots, statistical uncertainties are indicated by vertical (red) lines, while boxes are used to illustrate the systematic uncertainties.

where the smoothness approximation is warranted and λ is evaluated at $K = p$ [37].

The comparison of Monte Carlo resonance-model simulations of $\lambda(m_T)/\lambda_{\max}$ to the above presented data shows that, within systematic uncertainties, an in-medium mass drop of η' is not inconsistent with our measured data. Theoretically, the mass of the η' meson could be sensitive to the $U_A(1)$ symmetry restoration in hot and dense hadronic matter [44,45]. The key point being that compared to the other pseudoscalar mesons, the η' is anomalously heavy, ≈ 958 MeV, although the η' quark content is similar to that of the η meson. The η meson mass, $m_\eta \approx 548$ MeV, is closer to the mass of the charged kaons, ≈ 494 MeV while it is 410 MeV lighter than the mass of the η' meson, $m_{\eta'} \approx 958$ MeV, although the η and the η' mesons have the same quark content [46]. The extra mass is explained in the standard model in terms of the $U_A(1)$ anomaly that couples the mass of the η' to the topological properties of the quantum-chromodynamics vacuum state. If at high temperatures the structure of this vacuum changes, the extra 410 MeV mass difference may vanish and the η' mesons may return to the mass scale of the other pseudoscalar mesons, with its mass becoming similar to that of the η [45]. As the η' meson leaves the hot and dense matter, it regains its anomalously large mass at the expense of its kinetic energy and consequently this effect modifies the spectrum. See below for a discussion of how Monte Carlo calculations are used

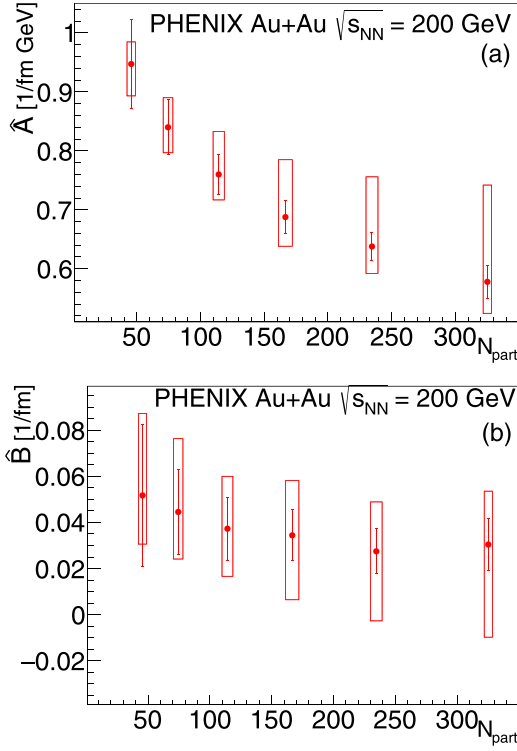


FIG. 9. The two parameters of the affine linear fit to the inverse square of the Lévy-scale parameter $1/\hat{R}(m_T)$ are defined in Eq. (14), (a) the slope parameter (\hat{A}), and (b) the intercept parameter (\hat{B}) are shown as functions of N_{part} . The central values are shown with dots, statistical uncertainties are indicated by vertical red lines, while boxes are used to illustrate the systematic uncertainties.

to investigate the effect on $\lambda(m_T)$ of this mass modification, along the lines of Refs. [1,37,42,43].

For each centrality class the SHAREV3 Monte Carlo generator [47,48] is used to evaluate the fraction of those short- and long-lived resonances that are the most important sources of pions. These fractions are used in our simulation. The input parameters to SHARE correspond to the STAR chemical freeze-out fits to the available experimental data on particle yields [49–55] and are available in the columns labeled grand-canonical-ensemble yields (GCEY) of Table VIII in Ref. [56]. The results from SHARE serve as inputs to our simulations.

The two most important parameters in the simulations are the in-medium mass of η' (denoted by $m_{\eta'}^*$) and the effective temperature of the η' condensate, the inverse-slope parameter (denoted by $B_{\eta'}^{-1}$). The in-medium mass drives the depth, while the inverse slope controls the steepness of the “dip” of the $\lambda(m_T)$ function. These two parameters are considered as fit parameters for the Monte Carlo simulations. Further details of the simulation and the estimation of its systematic uncertainties are given in the Appendix and in Refs. [1,37,42,43].

The χ^2 scans are implemented by taking small steps in the in-medium η' mass and the inverse-slope parameter $B_{\eta'}^{-1}$, comparing the results to the data. From these scans are determined the optimal (minimum χ^2) values and their uncertainties.

Figure 11 compares the data to both the optimal fits and the case that lacks the in-medium mass modification of the η'

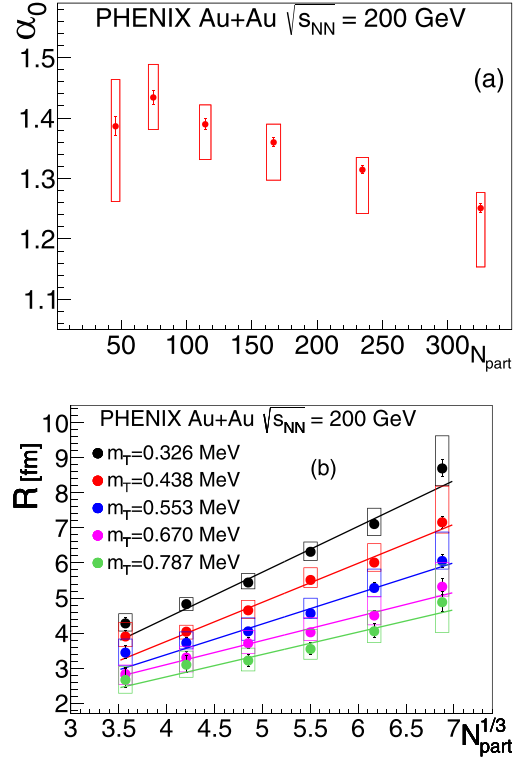


FIG. 10. The centrality dependence of (a) α_0 , the m_T -averaged value of α as a function of N_{part} and (b) the Lévy-scale parameter R as a function of $N_{\text{part}}^{1/3}$ in selected m_T bins.

meson. For $m_T \gtrsim 500$ MeV there is no significant difference between the optimal fit and the case assuming no modification. At low m_T no modification is strongly disfavored.

Our results suggest that a significant, centrality independent, in-medium mass drop of the η' is not inconsistent with the present measurements. In Figs. 12 and 13 the optimal values of $B_{\eta'}^{-1}$ and $m_{\eta'}^*$, respectively, are shown as function of N_{part} . The average value of the modified η' mass is $m_{\eta'}^* = 581_{-20}^{+12}(\text{stat})_{-91}^{+205}(\text{syst})$ MeV. In the case of the inverse-slope parameter, the point corresponding to the 0%–10% centrality bin was omitted from the constant fit, i.e., the average value indicated in Fig. 12, represents the average value of the five remaining points, that is $B_{\eta'}^{-1} = 56_{-14}^{+22}(\text{stat})_{-31}^{+190}(\text{syst})$ MeV. In Fig. 12 the solid (black) line denotes this centrality-averaged value. In at least five centrality classes, from 0%–10% to 40%–50%, the measured $\lambda(m_T)/\lambda_{\text{max}}$ functions are found to be consistently described by considering the suppressed mass of the η' . However, in the most peripheral centrality class (50%–60%) an unmodified η' mass cannot be excluded. This observation of minimal mass modification can be interpreted as being due to the lack of enough hot and dense matter. In Fig. 13, the dashed (red) line indicates the centrality-averaged value of the in-medium modified mass of the η' meson. The data are not inconsistent with Monte Carlo simulations with vanishing differences between $m_{\eta'}^*$ and $m_{\eta} \approx 548$ MeV.

These are the first, centrality-dependent experimental results that suggest the suppression of the η' meson mass in a hot and dense medium. Figure 13 compares these results to some

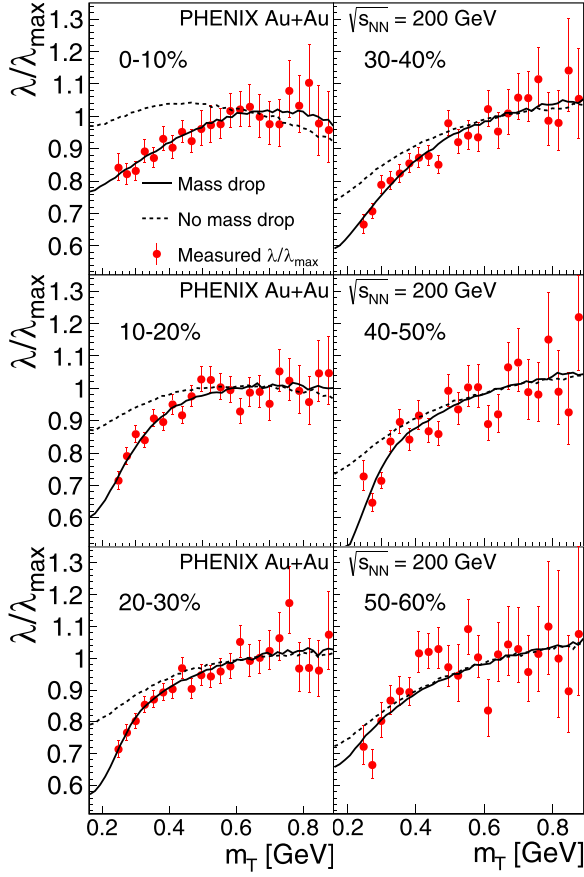


FIG. 11. Comparison of the data to the optimal fits of the six centrality classes (solid line) and the case where the mass of the η' meson in hot hadronic matter is unmodified and is equal to the Particle Data Group (PDG) value [46] (dashed line).

of the well-known theoretical predictions for the modified η' mass:

- (i) The Weinberg upper limit, $m_{\eta'}^* \leq \sqrt{3}m_\pi$, suggested in Ref. [57] clearly overestimates significantly, in each centrality bin, the possible in-medium mass drop of the η' particles.

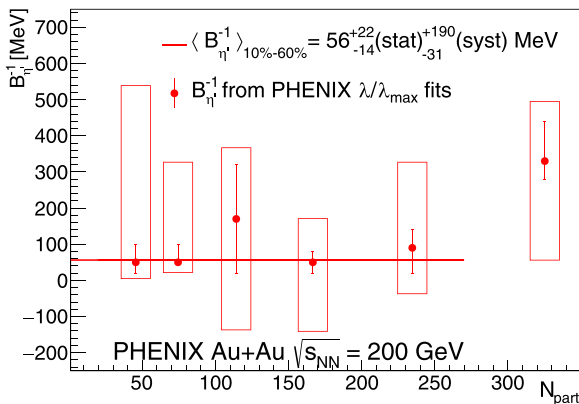


FIG. 12. The best values of $B_{\eta'}^{-1}$ spectrum of the η' condensate in the six centrality classes.

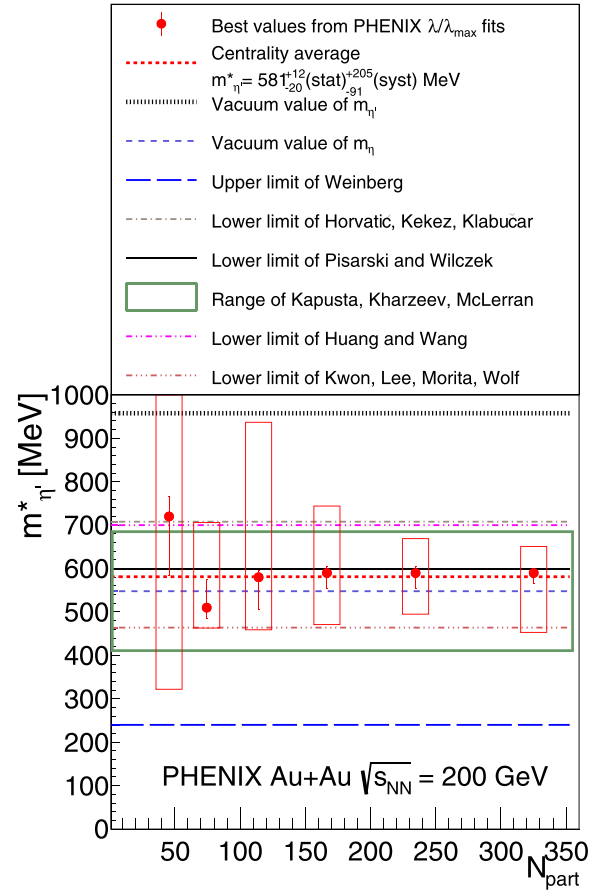


FIG. 13. The centrality dependence of the best values of the in-medium mass of the $m_{\eta'}^*$ are shown with full circle symbols, together with statistical (vertical lines) and systematic (red boxes) uncertainties. The average (using only statistical uncertainties) of the values of $m_{\eta'}^*$ found in the six centrality intervals is indicated by a dashed (red) line. The fitted values are compared to theoretical predictions of Weinberg [57], Horvatić, Kekez, and Klabučar [58], Pisarski and Wilczek [44], Kapusta, Kharzeev, and McLerran [45], Huang and Wang [59], and Kwon, Lee, Morita, and Wolf [60].

- (ii) Recent calculations by Horvatić, Kekez, and Klabučar (HKK) [58], based on the calculations of Witten-Veneziano equation and the generalization by Shore, evaluated the properties of the η and η' mesons at high temperatures, when the $U_A(1)$ and the chiral symmetry breaking are considered together. A substantial decrease in the η' mass around the chiral transition temperature was obtained in Ref. [58], but there was no decrease in the η mass. The new HKK results are an improvement on the earlier results of Refs. [42,43]. The lower limits, which are shown by the (light-green) dashed-dotted line in Fig. 13, lie above our results except in the 50%–60% centrality class, where the uncertainty on our value of $m_{\eta'}^*$ is particularly large. This is also true of the similar limit of Huang and Wang [59]. These models are in a modest tension with our results.

- (iii) The Pisarski-Wilczek lower limit of 600 MeV, as determined from Fig. 1 of Ref. [44], is shown as a solid (black) line. It is, within uncertainty, consistent, with the fitted values.
- (iv) The significant in-medium mass drop of the η' meson was related to the restoration of $U_A(1)$ symmetry and described as the return of a Goldstone-boson by Kapusta, Kharzeev, and McLerran (KKM) [45]. They have given a broad range for the possible in-medium mass of the η' meson in case of a *partial* $U_A(1)$ symmetry restoration, namely $411 \leq m_{\eta'}^* \leq 685$ MeV, shown as a (green) box in Fig. 13. Our results lie within this range.
- (v) Kwon, Lee, Morita, and Wolf (KLMW) also utilized a temperature-dependent, generalized Witten-Veneziano relation to obtain a nearly 50% decrease in the mass of the η' meson in hot and dense hadronic medium, as a consequence of the restoration of $U_A(1)$ symmetry [60]. Their 2012 prediction of $m_{\eta'}^* \geq 464$ MeV, shown as a triple dotted-dashed line in Fig. 13, is consistent with our results in each investigated centrality class.

VIII. SUMMARY AND CONCLUSION

This paper presents measurements of the two-pion BEC function in Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV using data recorded in 2010 by the PHENIX experiment. Lévy-stable distributions are utilized to characterize the data and determine the transverse-mass and centrality dependence of the Lévy parameters.

The Lévy parametrization is found to give a statistically acceptable description of the data with the Lévy exponent, α , significantly greater than 1 and less than 2. This exponent is determined in 23 m_T and in 6 centrality bins and is observed to be well described with its m_T -averaged value in every centrality bin. However, the transverse-mass-averaged values do depend on centrality.

The Lévy-scale parameter R is proportional to the HWHM of the source distribution, with a coefficient of proportionality that depends on the Lévy exponent α . The observed behavior of its inverse square as a function of m_T can be described by an affine linear fit, whose parameters are examined as a function of centrality. The parameter A , which can be related to the transverse velocity of the expansion, shows a trend similar to hydrodynamical predictions [38,39]. However, the assumption of local thermalization results in the $\alpha = 2$ special case (see, e.g., Refs. [28,38,39]), which is in significant contrast to our observations of $\alpha < 2$ in each of the investigated centrality and transverse-mass bins.

Within statistical uncertainties, the parameter B is consistent with zero or with a slightly negative value. A negative value of B is possible if the Cooper-Frye freeze-out terms are also taken into account [61]. This result may indicate a source that includes local thermalization with hydrodynamical expansion, followed by rescattering and decays of resonances. However, such nonequilibrium, scale-dependent features typically would result in deviations from the Lévy shape and from the applicability of generalized central-limit theorems.

But first-order deviations from the Lévy-stable source distributions using the expansion method of Refs. [12,13,25] were found to be consistent with zero. It is theoretically challenging to explain simultaneously the measured value of α , which is found to be significantly less than the Gaussian value of 2, and the m_T dependence of the Lévy-scale parameter R , which follows a hydrodynamically predicted scaling (see, e.g., Ref. [62]).

The connection to the initial geometry is supported by the linearity of the parameter R as a function of $N_{\text{part}}^{1/3}$ in any given m_T bin. The precise characterization of the correlation functions and the prudent handling of the Coulomb final-state interaction make it possible to determine in detail the m_T and centrality dependence of the $\lambda(m_T)$ intercept parameter, as well as its normalized form, $\lambda(m_T)/\lambda_{\text{max}}$. In the trends of the latter it can be qualitatively observed that there is a low- m_T suppression of $\lambda(m_T)/\lambda_{\text{max}}$ in every centrality bin and that the characteristics are more or less the same, i.e., the suppression is consistent with the hypothesis of centrality independence.

To quantify the suppression pattern, the Gaussian width and amplitude parameters, σ and H , were introduced. These parameters are observed to be centrality independent, except in the 50%–60% case, where the uncertainty on H increases and does not allow a statistically significant conclusion. This independence is helpful to rule out or validate models with predictions about this behavior. The pion laser model, for example, predicts a strong centrality and multiplicity dependence [61,63,64].

In Ref. [37] it was observed that the amount of suppression of $\lambda(m_T)$ and hence of $\lambda(m_T)/\lambda_{\text{max}}$ at low m_T can be an experimentally observable signal of the (partial) restoration of $U_A(1)$ symmetry and a measure of in-medium reduction of the mass of the anomalously heavy η' mass. An approximately constant trend of the intercept parameter λ as a function of m_T was observed in small systems at lower energies ($E_{\text{LAB}} = 150$ A GeV [7,20]), which is consistent with vanishing suppression ($H = 0$). These results are both qualitatively and quantitatively different from the presence of the suppression [$H_0 = 0.42 \pm 0.02(\text{stat})$], which is within uncertainties independent of the centrality in $\sqrt{s_{NN}} = 200$ GeV Au + Au collisions at RHIC. This entirely experimental-data-based, Monte Carlo independent observation suggests that the signal characterized by the value of H is dependent on the energy and/or the system size. Further detailed measurements are needed to determine the energy and system size, where H becomes nonvanishing for the first time, starting from zero in the small Be+Be collisions at the relatively low energy of $E_{\text{LAB}} = 150$ A GeV.

There could also be alternative or competing effects that could modify the correlation strength; see, e.g., Ref. [63,65]. Further theoretical studies are needed to explain these data using other methods. However, one of the possible explanations relates the observations to restoration of $U_A(1)$ symmetry in hot and dense hadronic matter as detailed in Refs. [37,42,43].

It is clear that our observed suppression of the $\lambda(m_T)/\lambda_{\text{max}}$ parameter at lower transverse mass is not inconsistent with the in-medium mass modification of the η' mesons, related to

TABLE II. The centrality-dependent, fitted values, in (MeV), of the in-medium η' mass $m_{\eta'}^*$ and the inverse slope $B_{\eta'}^{-1}$. The statistical (stat) and systematic (syst) uncertainties are also presented.

Centrality	$m_{\eta'}^*$	(stat)	(syst)	$B_{\eta'}^{-1}$	(stat)	(syst)
0%–10%	590	+5 –25	+61 –137	330	+110 –50	+165 –274
10%–20%	590	+15 –35	+79 –95	90	+50 –70	+237 –127
20%–30%	590	+15 –35	+154 –119	50	+30 –30	+121 –191
30%–40%	580	+15 –75	+357 –121	170	+150 –150	+197 –307
40%–50%	510	+65 –25	+196 –47	50	+50 –10	+277 –28
50%–60%	720	+45 –135	+508 –398	50	+50 –30	+489 –45

$U_A(1)$ symmetry restoration in hot hadronic matter. This relation was cross-checked with the help of detailed Monte Carlo simulations, and creating χ^2 and CL maps that compared simulations which allowed in-medium η' mass modification in hot and dense hadronic matter. As detailed in Fig. 13, and also in Table II, the best value of the in-medium mass of the η' meson is consistent for each of the considered centrality classes, with its centrality averaged value of $581^{+12}_{-20}(\text{stat})^{+205}_{-91}(\text{syst})$ MeV.

This mass is, within the uncertainties of this indirect measurement, the same as the Particle Data Group (PDG) value of the η meson, $m_\eta = 547.86 \pm 0.02$ MeV [46]. This observation suggests that the return of the so-called prodigal Goldstone boson [45] and the restoration of the $U_A(1)$ symmetry is not inconsistent with our measurements.

However, our measurements are inconsistent with Monte Carlo simulations that utilize the unmodified, PDG value of the η' mass, $m_{\eta'} = 957.78 \pm 0.06$ MeV, which does not allow for a description of our $\lambda(mT)/\lambda_{\text{max}}$ data, except in the most peripheral (50%–60%) centrality class. Several theoretical predictions from Refs. [44,45,57–60] are compared to the results of our χ^2 maps. These comparisons can be summarized as follows:

- (i) The Kapusta-Kharzeev-McLerran prediction [45] is in agreement with our measurements in each investigated centrality class.
- (ii) The lower limit of Kwon, Lee, Morita, and Wolf [60] is also consistent with our measurement in each investigated centrality class.
- (iii) Our measured centrality-average value of $m_{\eta'}^*$ is slightly below, but consistent with, the lower limit predicted by Pisarski and Wilczek [44].
- (iv) However, the upper limit of Weinberg [57] is several standard deviations below the central values obtained in each investigated centrality class.
- (v) The lower limit predictions of Horvatić, Kekez, and Klabučar [58] and of Huang and Wang [59] are excluded except in the 50%–60% centrality class.
- (vi) Our results also suggest that the prediction of Ref. [66] slightly underestimates the in-medium mass change of the η' .

However, the lack of in-medium η' mass modification is not consistent with our measurements, except in the 50%–60% centrality class, as discussed and detailed in the

Appendix. At present, these PHENIX results provide the most detailed, centrality-dependent constraints for future theoretical studies on $U_A(1)$ symmetry restoration in hot and dense hadronic matter. These results also exhibit an unprecedented selection power by excluding certain models in certain centrality classes. Thus these indirect PHENIX measurements provide important constraints and insights to future studies of (partial) $U_A(1)$ symmetry restoration in hot and dense hadronic matter.

For future experimental studies, these results emphasize the need for direct measurements of identified η' spectra in high-energy heavy-ion collisions at RHIC and at the Large Hadron Collider, in particular in the very soft, $p_T \leq 300$ MeV kinematic domain at midrapidity. Direct experimental observations of enhanced production of soft η' mesons seem to be particularly difficult, due to the expected backgrounds. Huge backgrounds, e.g., in the $\eta' \rightarrow \gamma\gamma$ decay channel, are expected, e.g., from $\pi^0 \rightarrow \gamma\gamma$ decays. Thus a direct observation of in-medium η' mass modification seems to be experimentally challenging, but based on the indirect results summarized here, are also expected to be particularly rewarding.

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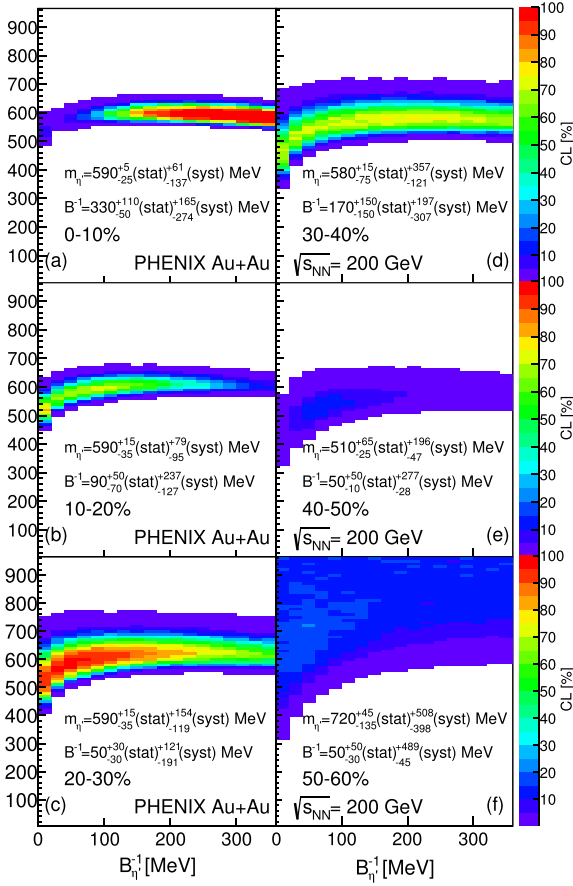


FIG. 14. The CL maps of the $m_{\eta'}^*$ and $B_{\eta'}^{-1}$ scans in the six centrality classes.

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APPENDIX: DETAILS OF THE MONTE CARLO SIMULATIONS

Monte Carlo simulations with scenarios that allow for (or exclude) a possible in-medium mass modification of the η' meson and their comparisons to data were performed in Refs. [37,42,43] and in the previous PHENIX paper [1] by utilizing standard χ^2 and CL maps. Similar but now centrality-dependent simulations are summarized in this Appendix. These comparisons result in χ^2 and CL maps that are utilized to determine the χ^2 minimum (or CL maximum) yielding the best values, as well as the statistical and the systematic uncertainties of a possible in-medium mass drop of the η' meson. Figure 14 shows an example of such a comparison, where each bin corresponds to a comparison of Monte Carlo simulations of resonance decay chains with our data on $\lambda(m_T)/\lambda_{\max}$ as a function of centrality, as detailed in Sec. VII.

Each panel of Fig. 14 contains the optimal values of $m_{\eta'}^*$ and $B_{\eta'}^{-1}$ in the scanned region. For our simulations, a unique χ^2 minimum is found, in contrast to earlier studies in Refs. [42,43]. By now, it is well known that the dominant

mechanism for soft-particle production is a thermal one in $\sqrt{s_{NN}} = 200$ GeV Au+Au collisions.

Earlier, a significantly broader class of models was considered in Refs. [42,43] that resulted in two, characteristically different minima in the 0%–30% centrality class of the same reaction. Accordingly, Refs. [42,43] evaluated the exclusion limits, focusing in particular on the smallest possible in-medium mass drop of the η' meson, which is required to describe the data with at least 0.1% CL in any of the considered Monte Carlo simulations.

Our systematic investigations include an estimation of the exclusion limits for the lowest- and largest-possible values of the in-medium mass of the η' meson. The systematic variations are consistent, i.e., all the systematic variations result in essentially the same minimum. Thus, in contrast to Refs. [42,43], our indirect observation of in-medium η' mass modification is used with the best value and its statistical and systematic uncertainties of $m_{\eta'}^*$, rather than exclusion limits.

Table II shows the centrality-dependent, fitted values of the in-medium η' mass $m_{\eta'}^*$ and the inverse slope $B_{\eta'}^{-1}$ with statistical and systematic uncertainties. The exclusion limits of these parameters were evaluated, but they do not result in statistically significant new ranges.

The η' spectrum has six parameters, which are considered as inputs to the $\lambda(m_T)$ simulations. The two most important inputs are the in-medium mass of the η' , $m_{\eta'}^*$, and the characteristic-slope parameter, $B_{\eta'}^{-1}$, of the η' mesons that are emitted with PDG mass after the decay of the condensate. The second parameter characterizes a second component, a low- p_T part of the η' spectrum and determines the slope of the suppression of the λ parameter at low- p_T [37,42,43]. These two parameters determine the shape of the observed dip in the low- m_T region of the $\lambda(m_T)$ function: a strongly reduced η' mass causes a dip in the low- m_T part of the $\lambda(m_T)$ function, while $B_{\eta'}^{-1}$ controls the slope of this dip; see Refs. [42,43]. These two parameters are considered as fit variables. The values of four nuisance parameters α_{th} , T_f , $\langle u_T \rangle$, and T_{cond} , are treated as constants in the fit, but are varied to evaluate the systematic uncertainties of the fitted values.

The χ^2 scans provide fine grids of CL in each centrality class. The in-medium mass is scanned with 10 MeV steps between 0 and 958 MeV (the unmodified mass of the η') and $B_{\eta'}^{-1}$ with 20 MeV steps between 0 and 360 MeV. The minimum value of χ^2 (or maximum value of CL) is determined, which provides the most likely values of $m_{\eta'}^*$ and $B_{\eta'}^{-1}$ and their statistical uncertainties, as well as the CL of the fits to the data, which are compared to the data in Fig. 11. The values of the four “constant” parameters, α_{th} , T_f , $\langle u_T \rangle$, and T_{cond} , are varied to evaluate the systematic uncertainties of the fitted values:

α_{th} : The centrality-dependent invariant single-particle spectra of positively and negatively charged kaons as well as protons and antiprotons of Ref. [49] are fitted with the formula

$$N(m_T) = C m_T^{\alpha_{th}} \exp\left(-\frac{m_T}{T_{eff}}\right), \quad (A1)$$

where C is a normalization constant. The polynomial exponent of the thermal spectrum is denoted by α_{th} , to distinguish it from the Lévy exponent α . Note that the parameter α_{th} of

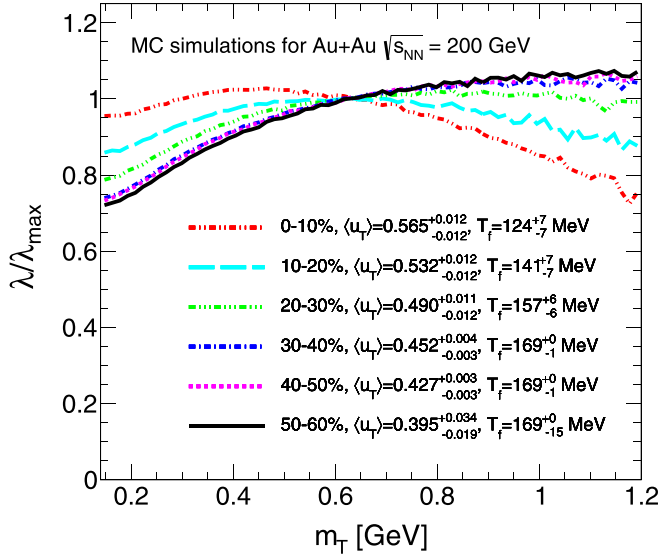


FIG. 15. Centrality dependent expectations for $\lambda(m_T)/\lambda_{\max}$, based on Monte Carlo simulations without any in-medium η' modification.

this paper is the same, as the α parameter of Refs. [37,42,43]. The exponent $\alpha_{th} = 1 - d/2$, where d is the number of spatial dimensions of the expansion; hence its allowed range is $1 \leq d \leq 3$. The value of α_{th} is fixed to $\alpha_{th} = 0$, which corresponds to an effectively two-dimensional expansion [31,38]. The same value was used in Ref. [49] by the PHENIX experiment when obtaining T_{eff} . Its value is assumed to be independent of centrality and is varied in the systematic studies between $1/2$ and $-1/2$. Also, similarly to Ref. [49], good quality exponential fits are obtained, with $CL \geq 0.1\%$, in each centrality class for charged kaons, protons, and antiprotons in the transverse-mass range of $0.1 \leq m_T - m \leq 1.0$ GeV. Similarly to Ref. [49], the mass dependence of the slopes is also well described with affine linear fits: $T_{eff}(m) = T_f + m\langle u_T \rangle^2$, but here these fits have a good CL with $CL \geq 0.1\%$ in each centrality class.

T_f : The kinetic freeze-out temperature is denoted by T_f . Significant centrality dependent results for the kinetic freeze-out temperature are obtained here. In particular, the value of T_f is significantly lower for the 0%–10% centrality class than for other centrality classes. For peripheral centrality classes, the value of T_f increases, reaching its upper limit, the value of the chemical freeze-out temperature T_{chem} . Our results for the variations of T_f and $\langle u_T \rangle$ with centrality are shown as inserted values in Fig. 15. Fixing T_f to a centrality-independent constant, results in the above affine linear-mass-dependent coefficient fits, similarly to Ref. [49], having CLs that are too small, well below the $0.1\% \leq CL$ threshold values. A possible reason for a lower kinetic freeze-out temperatures T_f in the 0%–10% centrality class is that a larger volume may cool remarkably longer and to a lower temperature, as noted by Hama and Navarra [67].

$\langle u_T \rangle$: The average radial flow is denoted by $\langle u_T \rangle$. This parameter influences the overall slope of the $\lambda(m_T)$ distribution at higher transverse mass and results in a centrality dependent

expectation for $\lambda(m_T)/\lambda_{\max}$. Note that the pion halo contains the decay products of long-lived resonances, which include the decay products of ω , η , η' and K_s^0 mesons. It is thus important to simulate these decay chains precisely and in agreement with available experimental data. The best values of T_f and $\langle u_T \rangle$ are obtained by simultaneous fits to the slope parameters of the positively and negatively charged kaons, protons, and antiprotons with the formula $T_{eff} = T_f + m\langle u_T \rangle^2$, where m is the mass of the meson or baryon. These parameters are systematically varied within the uncertainties allowed by these fits, shifting and fixing one of the $(T_f, \langle u_T \rangle)$ pair by one standard deviation and refitting the other parameter to take into account their covariation. The invariant spectra of K^\pm and those of protons and antiprotons are simulated using fits with good CL. The mass scaling of these spectra is then utilized. Nearly the same are the mass of the η meson compared to charged kaons and the mass of the η' meson compared to protons and antiprotons. This gives a good basis to describe well the null effect, which is the no in-medium mass modification scenario of the η and η' spectrum. This method has been tested and the test was published in Fig. 11 of Ref. [43]. Extrapolating the mass-scaled simulations to the PHENIX acceptance successfully reproduces the η spectrum measured in $\sqrt{s_{NN}} = 200$ GeV Au + Au collisions [68] in the $m_T - m_\eta \geq 1.75$ GeV range.

T_{cond} : The effective temperature of the in-medium η' condensate is denoted by T_{cond} , as in Refs. [37,42,43]. This parameter, together with the in-medium mass of the η' mesons, controls the thermal enhancement of the η' mesons in the medium. The smaller this number is, the smaller the η' mass drop is for a given $\lambda(m_T)/\lambda_{\max}$ measurement. Hence, the most conservative assumption is made for the default value, $T_{cond} = T_f$. As this temperature is inside the hot and dense, but already hadronic matter, part of our systematic studies is to vary T_{cond} within the allowed range of $T_{chem} \geq T_{cond} \geq T_f$. As the conservative default value lies at one of the edges of the allowed interval, this choice results in a more than usually asymmetric uncertainty distribution on the physical fit parameters, such as the in-medium mass of the η' and the slope parameter of the η' after the decay of the condensate: $m_{\eta'}^*$ and $B_{\eta'}^{-1}$. From these systematic uncertainty studies, the in-medium η' mass is frequently observed to be smaller than our quoted best values, but cannot easily be larger than the best values because the upper uncertainties from the variation of T_{cond} within the allowed range are much smaller than the lower uncertainties.

For each of the parameters α_{th} , T_f , $\langle u_T \rangle$, and T_{cond} , the χ^2 scans and the corresponding CL maps are redone for the variations. The relative difference between the default and the alternative setting is calculated. The final systematic uncertainty is the quadratic sum of these differences.

Additionally, several other cases are also investigated, including two special cases, $T_f = 140$ MeV fixed, independent of centrality, and $T_f = 177$ MeV fixed, also independent of centrality. The first case corresponds to Landau's calculation of the freeze-out temperature (equal to the pion mass, mass of the lightest neutral quanta). The second case is consistent with the PHENIX publication on charged pion, kaon, and

(anti)proton spectra [49]. The latter choice is found to be inconsistent with $T_{\text{cond}} \geq T_f$. The former gives results that are within the quoted systematic uncertainties.

Similarly to Ref. [1], the statistical uncertainty of λ_{max} is treated as a normalization uncertainty. Both this uncertainty ($\approx 1\%$) and the systematic uncertainty ($\approx 1\%$) caused by the choice of m_T range when calculating λ_{max} are negligible compared to other uncertainties, except in the 50%–60% centrality case of $B_{\eta'}^{-1}$. In that case, no significant in-medium mass modification is found. Thus, $B_{\eta'}^{-1}$ cancels from the results and cannot be precisely determined. Therefore, these three negligibly small sources of systematic errors are not included here.

The simulation does not include the PHENIX detector system; hence, the experimental systematic uncertainties are accounted for by propagating the total systematic uncertainties of the measured $\lambda/\lambda_{\text{max}}$. The centrality-dependent results for the $B_{\eta'}^{-1}$ parameters and the in-medium masses $m_{\eta'}^*$ with statistical and systematic uncertainties are shown in Figs. 12 and 13, respectively, and the corresponding numerical values are tabulated in Table II.

As mentioned in Sec. VIA, the normalized intercept parameter dependence on m_T and centrality are investigated in detail. An explanation of the dependence is given in terms of radial flow and p_T sharing among the pions arising in η' decay. The simulations suggest a $\lambda(m_T)/\lambda_{\text{max}}$ curve, that changes monotonically with centrality. These changes appear related to the monotonic decrease of radial flow $\langle u_T \rangle$, coupled to a monotonic increase of the kinetic freeze-out temperature T_f as the collisions change from most central to more and more peripheral collisions [37]. The values of T_f and $\langle u_T \rangle$

are obtained from affine linear fits with $T_{\text{eff}} = T_f + m\langle u_T \rangle^2$ to the slope parameters of the charged kaon, K^+ and K^- , as well as to the proton and the antiproton single-particle spectra. Particular attention is paid to the requirement that the single exponential fit to the single-particle spectra have acceptable CLs, $\text{CL} \geq 0.1\%$, and that the affine linear fits with $T_{\text{eff}} = T_f + m\langle u_T \rangle^2$ also reflect the slopes of the single-particle spectra T_{eff} with a $\text{CL} \geq 0.1\%$.

In addition to the systematic investigations detailed in this manuscript, several additional consistency checks were performed, such as including fit-range stability investigations and using three different methods of propagation of statistical, and systematic uncertainties. These were performed because the definition and utilization of λ_{max} in this experimental paper, which was produced using direct access to results within various experimental cuts, differ from the utilization of λ_{max} in Refs. [42,43], which reported reanalysis of already published data. In those works, only the published systematic uncertainties were available to be propagated to the final results, while in the present paper every quantity up to the final results has been evaluated within each experimental cut.

The method of propagation of statistical and systematic uncertainties was cross checked though comparison with the less direct methods of Refs. [42,43] as well as with the PHENIX method of Ref. [69]. These methods gave results that are consistent with (and typically have smaller uncertainties than) those presented in the body of this paper. Hence the central values, with statistical and systematic uncertainties, presented in the paper are obtained with the most conservative of the four different methods that were tested.

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