



Searching for GEMS: Discovery and Characterization of Two Brown Dwarfs Around M Dwarfs*

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Abstract

Brown dwarfs bridge the gap between stars and planets, providing valuable insight into both planetary and stellar-formation mechanisms. Yet the census of transiting brown-dwarf companions, in particular around M-dwarf stars, remains incomplete. We report the discovery of two transiting brown dwarfs around low-mass hosts using a combination of space- and ground-based photometry along with near-infrared radial velocities. We characterize TOI-5389Ab ($68.0^{+2.2}_{-2.2} M_J$) and TOI-5610b ($40.4^{+1.0}_{-1.0} M_J$), two moderately massive brown dwarfs orbiting early M-dwarf hosts ($T_{\text{eff}} = 3569 \pm 59$ K and 3618 ± 59 K, respectively). For TOI-5389Ab, the best fitting parameters are period $P = 10.40046 \pm 0.00002$ days, radius $R_{\text{BD}} = 0.824^{+0.033}_{-0.031} R_J$, and low eccentricity $e = 0.0962^{+0.0027}_{-0.0046}$. In particular, this constitutes one of the most extreme substellar-stellar companion-to-host mass ratios of $q = 0.150$. For TOI-5610b, the best-fitting parameters are period $P = 7.95346 \pm 0.00002$ days, radius $R_{\text{BD}} = 0.887^{+0.031}_{-0.031} R_J$, and moderate eccentricity $e = 0.354^{+0.011}_{-0.012}$. Both targets are expected to have shallow, but potentially observable, occultations: $\lesssim 500$ ppm in the Johnson K band. A statistical analysis of M-dwarf/BD systems reveals for the first time that those at short orbital periods ($P < 13$ days) exhibit a dearth of $13 M_J < M_{\text{BD}} < 40 M_J$ companions ($q < 0.1$) compared to those at slightly wider separations.

Unified Astronomy Thesaurus concepts: Brown dwarfs (185); Transits (1711); Photometry (1234); Transit photometry (1709); Spectroscopy (1558); Substellar companion stars (1648)

Materials only available in the online version of record: machine-readable tables

1. Introduction

Statistically, M dwarfs have a higher occurrence rate of small planets and lower occurrence rate of large planets compared to solar-type stars (J. A. Johnson et al. 2010; X. Bonfils et al. 2013; C. D. Dressing & D. Charbonneau 2015;

G. D. Mulders et al. 2015; K. K. Hardegree-Ullman et al. 2019; D. C. Hsu et al. 2020; E. M. Bryant et al. 2023; T. Gan et al. 2023). As the most common stars in the Galaxy (T. J. Henry et al. 2006; C. Reylé et al. 2021), M dwarfs provide the greatest sample size to find rare systems in order to probe our understanding of planet formation.

Brown dwarfs (BDs) are defined as objects massive enough to burn deuterium but not massive enough that thermonuclear processes dominate their evolution ($13 \lesssim M \lesssim 80 M_J$; A. Burrows et al. 2001; A. Lecavelier des Etangs & J. J. Lissauer 2022). They are rare companions to main sequence (specifically FGK) stars, having $\lesssim 1\%$ occurrence (D. Grether & C. H. Lineweaver 2006; J. Sahlmann et al. 2011; A. Santerne et al. 2016; N. Grieves et al. 2017). The paucity of bound, short-period BDs prompted the definition of the “brown-dwarf desert” where these objects are exceedingly rare

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within $a \lesssim 3$ au of the stellar host, or equivalently $P \lesssim 5$ yr (D. Grether & C. H. Lineweaver 2006). Additionally, the census of transiting BDs around M dwarfs is exceedingly small, with only 10 confirmed (N. Grieves et al. 2021; C. I. Cañas et al. 2022; B. A. Henderson et al. 2024a). These systems in particular are useful because they probe formation mechanisms and the transition from substellar to stellar objects, and they provide a more accurate measurement of the mass as opposed to $m \sin i$.

There is some debate on classifying planets and substellar objects by formation mechanism rather than mass (i.e., core accretion versus gravitational instability; J. Maldonado & E. Villaver 2017; N. C. Santos et al. 2017; K. C. Schlaufman 2018; F. Kiefer et al. 2021). Core accretion (H. Mizuno 1980; J. B. Pollack et al. 1996; G. Chabrier et al. 2014) involves planetesimals colliding and combining until they reach a critical mass ($\sim 10\text{--}30 M_{\oplus}$), beyond which they rapidly accrete gas in order to form giant planets. Gravitational instability (A. P. Boss 1997) involves objects forming instead by a secondary gravitational collapse of the protoplanetary disk or the cloud. Thus an alternative definition would stipulate that if bodies form via core accretion they are considered planets and if they form via gravitational instability (disk instability) they are considered BDs unless they are massive enough to fuse hydrogen (G. Chabrier et al. 2014). Our targets specifically are so massive they much more likely formed via gravitational instability, since M dwarfs are observed to have lower-mass protoplanetary disks compared to solar-type stars (A. P. Boss 2006; S. M. Andrews et al. 2013; I. Pascucci et al. 2016), making it harder to form large planets via core accretion (S. Kanodia et al. 2024). The definition of BDs based on their formation mechanism blurs the lower limit of $13 M_J$. However, these more massive BDs still have unclear origins in if they formed via cloud or disk instabilities.

The Searching for Giant Exoplanets around M-dwarf Stars (GEMS) survey (S. Kanodia et al. 2024) utilizes space- and ground-based photometric and spectroscopic observations to target M dwarfs within 200 pc to find GEMS. This survey is designed to characterize 40 GEMS in order to accurately compare GEMS to the existing census of hot Jupiters around FGK stars. This is useful to constrain theories of planet formation. During the process of confirming and characterizing planets, many targets are observed to be brown dwarfs or binaries rather than planets, which is the case for the systems described in this work.

In this paper, we report the discovery and characterization of TOI-5389Ab and TOI-5610b, two BDs around M dwarfs originally flagged as objects of interest by the Transiting Exoplanet Survey Satellite (TESS; G. R. Ricker et al. 2014) science team. Section 2 presents new ground-based photometric and spectroscopic data on these targets. Section 3 describes the combination and modeling of these data sets to extract planetary, stellar, and system parameters. Section 4 provides a discussion of the new data in the context of other transiting BD systems. Finally, Section 5 summarizes and concludes our findings.

2. Observations

2.1. TESS Photometry

TOI-5389A (TIC 39143128, Gaia EDR3 7650827279 62822016) was identified via the TESS Science Processing

Operations Center (SPOC; J. M. Jenkins et al. 2016) and was observed in 1800 s cadence in Sector 22 from 2020 February 18 to 2020 March 18 and Sector 48 in 600 s cadence from 2022 January 28 to 2022 February 26. TOI-5610 (TIC 252481136, Gaia EDR3 834714838005307648) was identified via the TESS Faint Star Search (M. Kunimoto et al. 2022) and was observed in Sector 21 in 1800 s cadence from 2020 January 21 to 2020 February 18 and Sector 48 in 600 s cadence from 2022 January 28 to 2022 February 26.¹⁹ The specific TESS data analyzed can be accessed via the Mikulski Archive for Space Telescopes at the Space Telescope Science Institute.²⁰

To produce a dilution-free lightcurve, we used the TESS Gaia Light-Curve (TGLC; T. Han & T. D. Brandt 2023) python package, which deblends the lightcurve using Gaia DR3 (Gaia Collaboration et al. 2016, 2023) photometry and positions. Within the 3×3 TESS pixel aperture used by TGLC, a Gaia DR3 search using a $\sim 89''$ radius revealed that TOI-5389A has five potential nearby contaminant stars which could cause dilution in the lightcurve depending on target placement within the pixel. The collective targets have a Gaia ΔG mag range of $\sim 1\text{--}5$ mag. Similarly, TOI-5610 also has five potential contamination sources with a Gaia ΔG mag range of $\sim 1\text{--}5$ mag. Only TESS Sector 48 for TOI-5610 appeared to have appreciable dilution from another source, the others either had more favorable pixel placement or negligible dilution. TGLC corrects for most of this, and we also constrain any residual dilution using other undiluted ground-based lightcurves (Section 2.2).

There are clear transit signatures at dominant periodicities of ~ 10.40 and ~ 7.95 days and approximate depths of $\sim 4\%\text{--}5\%$ and $\sim 2\%\text{--}3\%$ for TOI-5389A and TOI-5610, respectively. A Box Least Squares analysis of both targets shows no other significant periodicities.

Figure 1 shows the two full TESS Sector lightcurves and their phase-folded counterparts for TOI-5389A (black dots) and Figure 2 shows the same for TOI-5610. The blue curves shows the best-fit models, described in Section 3.

2.2. Red Buttes Observatory Photometry

The Red Buttes Observatory (RBO; D. H. Kasper et al. 2016) is a 0.6 m telescope owned by the University of Wyoming located 10 km south of Laramie, Wyoming. It is an $f/8.43$ Ritchey–Cretien by DFM Engineering, Inc. equipped with an Apogee Alta F16 camera. The Kodak KAF 16801 4096×4096 chip with $9 \mu\text{m}$ pixels binned 2×2 produces a plate scale of 0.73 pix^{-1} and a field of view of $24.9'$. All observations were conducted using the Bessell I filter and a 240 s exposure time while slightly defocused. The CCD has a gain of $1.39 e^- \text{ ADU}^{-1}$ and was operated at -15°C , which results in a significant dark current ($\sim 0.13 e^- \text{ s}^{-1} \text{ pix}^{-1}$). Exposures were calibrated using flat-field, bias, and dark corrections. The readout time was ~ 2.4 s and a typical full width at half maximum is $\sim 3''$. The typical photometric uncertainty is $0.9\%\text{--}1.2\%$. We successfully observed TOI-5389Ab on the local nights of 2023 February 28 and 2024 February 6 as well as TOI-5610b on the local nights of 2024 January 2 and 2024 February 11. Observations lasted 5–8 hr

¹⁹ While TOI-5389A does have 120 s cadence data from SPOC, it is significantly lower S/N and does not meaningfully constrain the transit.

²⁰ DOI:10.17909/zk4a-jh57.

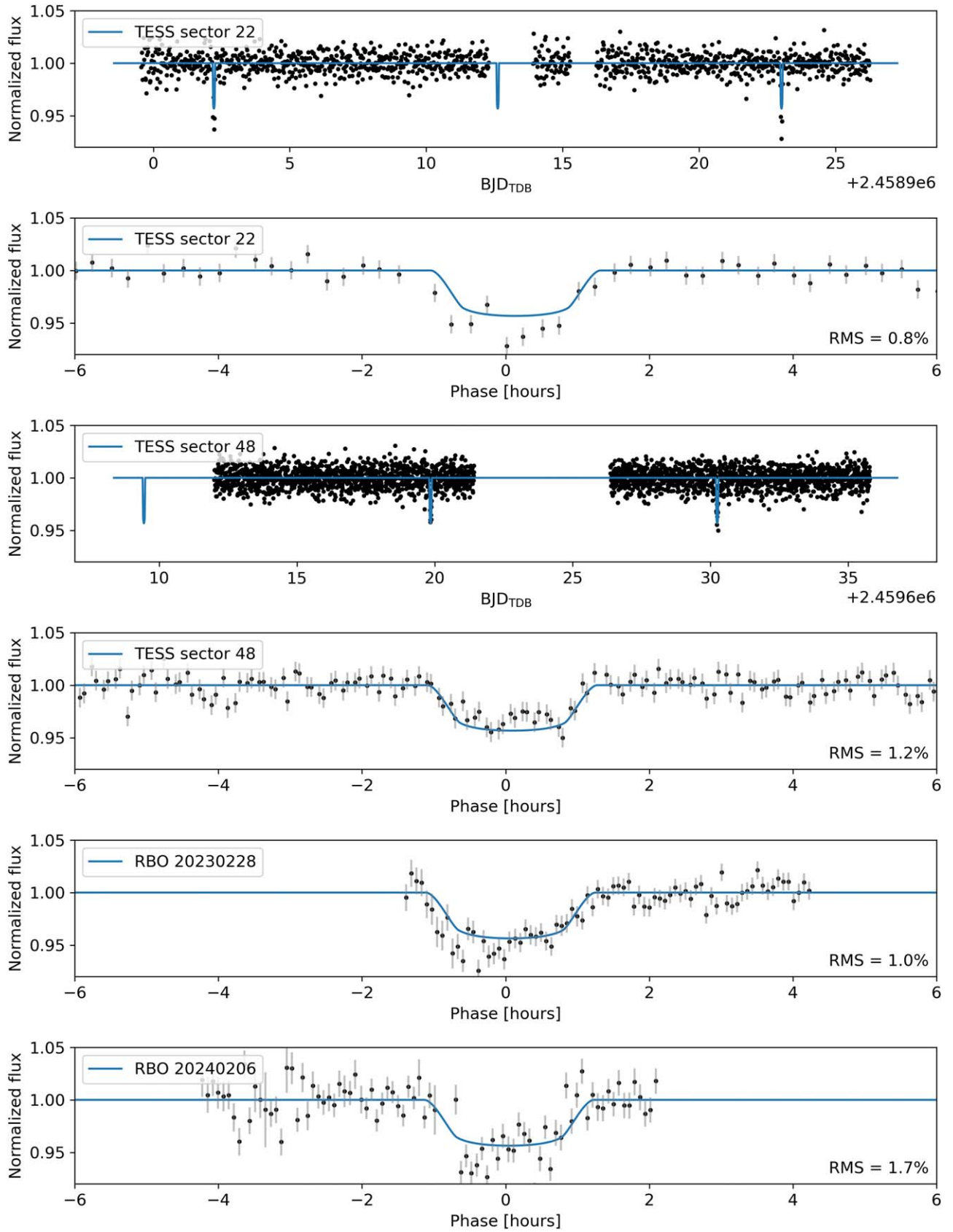


Figure 1. Observed transits for TOI-5389Ab. From first to last: TESS Sector 22 (2020 March) and its phase-folded counterpart, TESS Sector 48 (2022 February) and its phase-folded counterpart, RBO *I* band 2023 February 28, and RBO *I* band 2024 February 6.

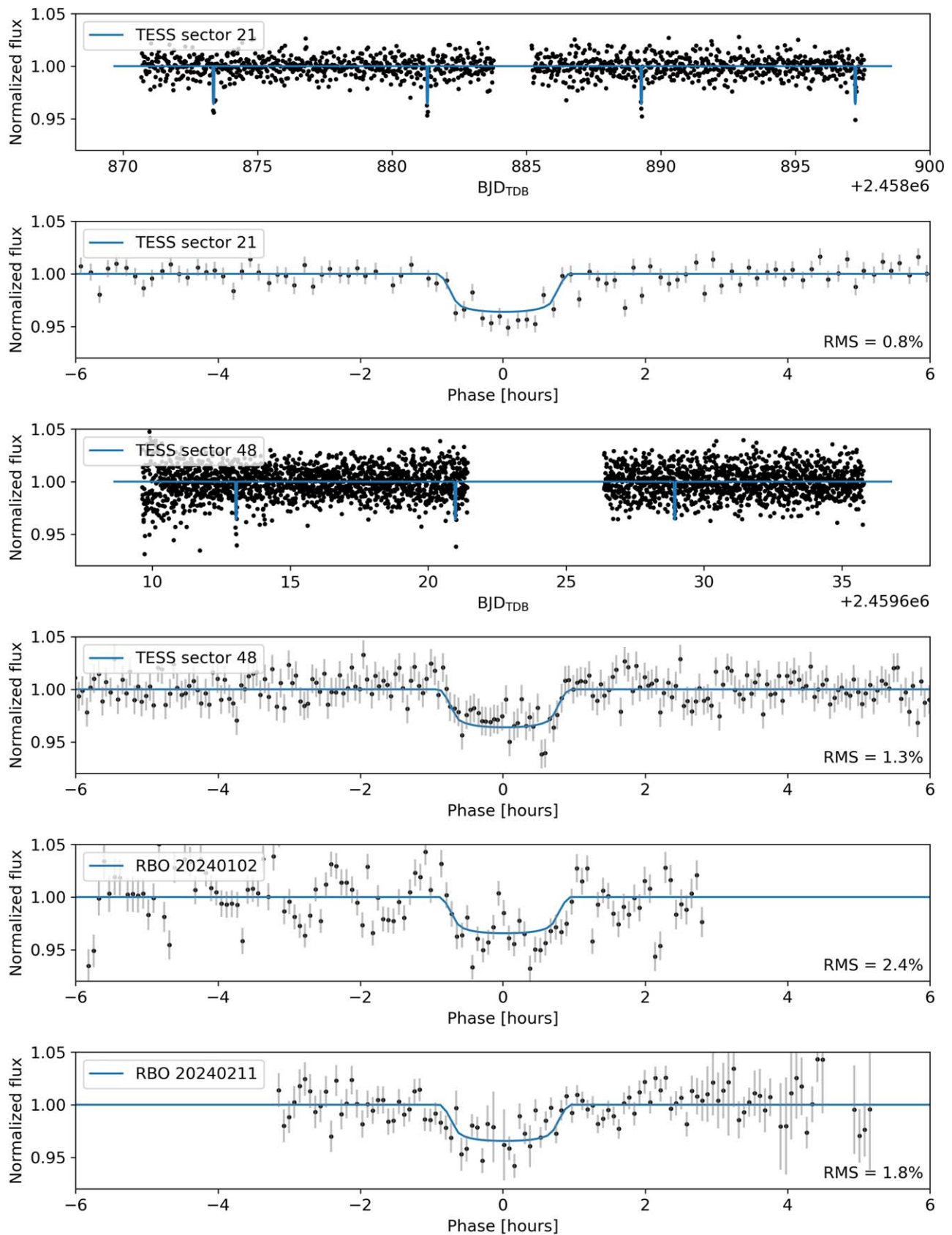


Figure 2. Observed transits for TOI-5610b. From first to last: TESS Sector 21 (2020 February) and its phase-folded counterpart, TESS Sector 48 (2022 February) and its phase-folded counterpart, RBO *I*-band 2024 January 2, and RBO *I*-band 2024 February 11.

Table 1
TOI-5389 RBO Photometric Data

BJD _{TDB}	Normalized Flux	Flux Err
2460004.605	0.995	0.013
2460004.608	1.018	0.013
2460004.611	1.011	0.013
...

(This table is available in its entirety in machine-readable form in the [online article](#).)

Table 2
TOI-5610 RBO Photometric Data

BJD _{TDB}	Normalized Flux	Flux Err
2460312.694	0.934	0.016
2460312.697	0.949	0.016
2460312.700	0.999	0.016
...

(This table is available in its entirety in machine-readable form in the [online article](#).)

each and captured two full transits for each target. Observation times had to be converted²¹ from JD_{UTC} into BJD_{TDB} as described in J. Eastman et al. (2010). The lightcurves were derived from the raw data via differential aperture photometry in AstroImageJ (K. A. Collins et al. 2017). The final two panels of Figures 1 and 2 display the normalized RBO lightcurves for each epoch. Tables 1 and 2 present the Julian dates, normalized flux, and uncertainties. The full tables are available as machine-readable files.

2.3. NESSI High-resolution Speckle Imaging

TOI-5389A and TOI-5610 were observed with the NN-Explore Exoplanet Stellar Speckle Imager (NESSI; N. J. Scott et al. 2018) on the WIYN 3.5 m telescope²² on 2023 January 28. For each target, a set of diffraction-limited frames was taken simultaneously with the red and blue NESSI cameras in the Sloan Digital Sky Survey (SDSS; J. Kollmeier et al. 2019) z' and r' filters. The data were processed following the methods described in S. B. Howell et al. (2011). Figures 3 and 4 show the reconstructed speckle images and contrast limits. For TOI-5610, these data allow us to rule out the presence of nearby sources down to magnitude limits of $\Delta r' = 4.16$ mag and $\Delta z' = 3.83$ mag at a separation of $0''.2$ and limits of $\Delta r' = 4.60$ mag and $\Delta z' = 4.67$ mag at a separation of $1''.2$. For TOI-5389A, the equivalent magnitude limits are $\Delta r' = 4.05$ mag and $\Delta z' = 3.84$ mag at a separation of $0''.2$ and $\Delta r' = 4.72$ mag, and $\Delta z' = 4.90$ mag at a separation of $1''.2$. The Gaia RUWE values of 1.147 and 0.999 for TOI-5389A and TOI-5610, respectively, are consistent with single-star astrometric solutions (RUWE > 1.4 is suggestive of binarity; V. Belokurov et al. 2020; Z. Penoyre et al. 2020). The Gaia Astrometric Excess Noise is also consistent with this.

²¹ <https://astrutils.astronomy.osu.edu/time/utc2bjd.html>

²² The WIYN Observatory is a joint facility of the NSF's National Optical-Infrared Astronomy Research Laboratory, Indiana University, the University of Wisconsin-Madison, Pennsylvania State University, Purdue University, and Princeton University.

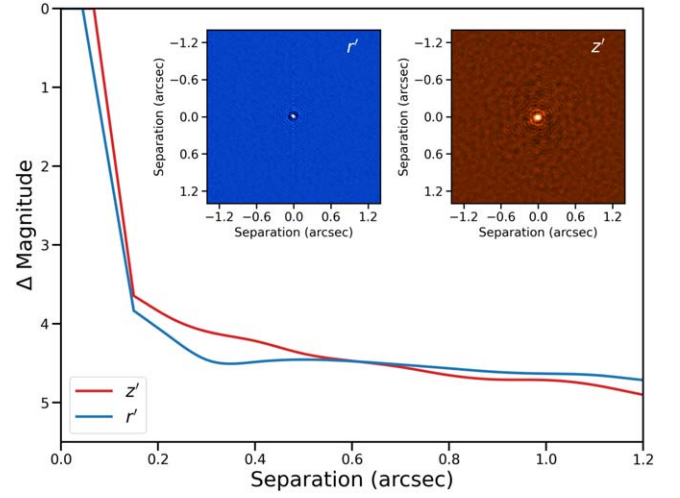


Figure 3. NESSI contrast curves for TOI-5389A. NESSI rules out nearby sources down to $\Delta r' = 4.05$ mag and $\Delta z' = 3.84$ mag within $0''.2$.

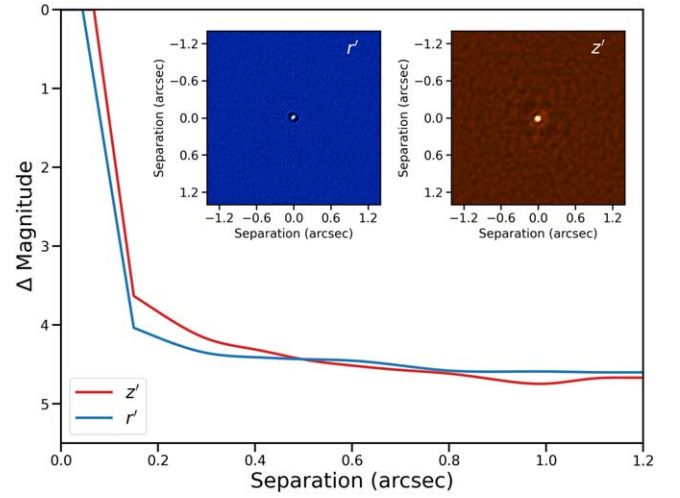


Figure 4. NESSI contrast curves for TOI-5610. NESSI rules out nearby sources down to $\Delta r' = 4.16$ mag and $\Delta z' = 3.83$ mag within $0''.2$.

2.4. The Habitable-zone Planet Finder Spectra

The Habitable-zone Planet Finder (HPF; S. Mahadevan et al. 2012, 2014) is a high-resolution, near-infrared, fiber-fed (S. Kanodia et al. 2018), stabilized (~ 1 mK; G. Stefansson et al. 2016) spectrograph with resolution $R \sim 55,000$ and wavelength coverage 8080–12780 Å. It is located on the 10 m Hobby-Eberly Telescope (L. W. Ramsey et al. 1998; G. J. Hill et al. 2021) at the McDonald Observatory in Texas. We obtained two 945 s exposures on each of 10 nights for TOI-5389A and 12 nights for TOI-5610 between 2022 November and 2023 December.

The raw data were processed using the HxRGproc algorithms (J. P. Ninan et al. 2018) and wavelength calibrated by the method described in G. Stefansson et al. (2020). Wavelength solutions were determined from routine laser-frequency-comb calibration frames that provide a drift correction to a precision of < 30 cm s⁻¹ (G. Stefansson et al. 2020). These were not simultaneous with observations to avoid scattered light in the spectra. The extracted 1D spectra have a

Table 3
Barycentric Radial Velocity Measurements for TOI-5389A

BJD _{TDB}	RV (m s ⁻¹)	σ_{RV} (m s ⁻¹)
2459921.927	9123	23
2459921.939	9443	23
2460306.864	8455	16
2460306.876	8654	15
2459954.830	2884	16
2459954.842	2607	19
2459923.921	2424	24
2459923.933	2752	30
2460301.884	13,774	81
2460010.912	14,673	17
2460010.924	14,447	28
2460302.877	18,928	34
2460095.676	22,214	20
2460095.688	22,399	17
2459950.835	22,354	43
2459950.847	21,974	27
2460284.933	15,185	24
2460284.945	15,503	20

(This table is available in machine-readable form in the [online article](#).)

median signal-to-noise ratios (S/Ns) per pixel of 18 (TOI-5389A) and 15 (TOI-5610) in spectral order 18 at 1 μ m.

The 28 spectral orders were continuum normalized and combined into a single spectrum, which was then Doppler corrected to the barycentric frame of reference using `bar-ycorrrpy`, a python transposition of the J. T. Wright & J. D. Eastman (2014) algorithms (S. Kanodia & J. Wright 2018). Wavelengths in regions of poor atmospheric transmission were masked. Only the wavelength intervals 8600–8900 Å and 9900–10600 Å were retained for radial velocity analysis. We employed a custom python code to compute the broadening function (analogous to a cross correlation; S. Rucinski 1999) using a high-resolution PHOENIX (T. O. Husser et al. 2013) model atmosphere of the appropriate effective temperature and gravity for each target. Radial velocities were measured as the center of a Gaussian function fitted to the peak in the broadening function at each of the observed epochs. Uncertainties on the Gaussian center yield a typical precision of 20–30 m s⁻¹ via a least-squares fit to the cross-correlation function. A test of this analysis technique on HPF spectra of Barnard’s star reproduces the accepted literature velocities within 0.3 km s⁻¹, providing an overall velocity accuracy measurement (GJ699, an M dwarf having a well-established radial velocity of -110.1 km s⁻¹; P. Fouqué et al. 2018). While, in principle, the rotational velocities of the target stars are measurable from the width of the broadening function after deconvolution with the instrumental broadening, the $v \sin i$ for these spectra are all less than 2 km s⁻¹, below the instrumental limit set by HPF’s spectral resolution of $R \approx 55,000$. These slower rotational velocities are consistent with the slow rotational velocities expected of older early M dwarfs (P. Giacobbe et al. 2020).

Tables 3 and 4 list the HPF RVs, while Figures 5 and 6 show the phased RV time series and residuals. For both targets, the unphased velocity curve is consistent with uncertainties and the residuals show no long term patterns.

Table 4
Barycentric Radial Velocity Measurements for TOI-5610

BJD _{TDB}	RV (m s ⁻¹)	σ_{RV} (m s ⁻¹)
2459931.869	37,125	14
2459931.881	37,258	18
2460011.860	37,254	23
2460029.596	40,797	6
2460029.608	41,000	5
2460007.649	45,996	9
2460007.661	45,944	9
2459952.803	47,922	7
2459952.815	47,900	8
2459897.954	49,727	13
2459897.966	49,320	16
2460303.842	49,440	22
2460303.854	49,730	19
2460065.705	48,915	12
2460065.717	48,339	12
2460272.939	45,533	13
2460304.841	45,033	9
2460304.853	44,817	10
2459955.029	43,441	12
2459955.041	43,212	11
2460018.840	41,292	6
2460018.852	41,093	7

(This table is available in machine-readable form in the [online article](#).)

3. Analysis

3.1. Stellar Parameters

Table 5 lists broadband photometric measurements from the optical through mid-infrared from the Panoramic Survey Telescope and Rapid Response System (Pan-STARRS; H. A. Flewelling et al. 2020), the 2 μ m All Sky Survey (2MASS; M. F. Skrutskie et al. 2006), and the Wide-field Infrared Survey Explorer (WISE; E. L. Wright et al. 2010). Table 5 also contains stellar parameters T_{eff} , $\log g$, [Fe/H], mass (M_*), radius (R_*), parallax (Plx), R.A. proper motion (R.A. PM), decl. proper motion (decl. PM), density (ρ_*), luminosity (L), and spectral type (Sp-Type). With spectra from HPF, we used the spectral matching technique (HPF-SpecMatch) described in G. Stefansson et al. (2020) to independently obtain stellar parameters T_{eff} , $\log g$, and [Fe/H]. We selected the night with the highest S/N spectrum in the 5th spectral order ($\lambda \sim 8670$ Å–8750 Å) for the least amount of telluric contamination. HPF-SpecMatch then compares this spectrum to a library of high S/N (>100) spectra in order to derive new stellar parameters. The library of 100 stars spans $2700 \text{ K} < T_{\text{eff}} < 4500 \text{ K}$, $4.3 < \log g < 5.3$, and $-0.5 < [\text{Fe}/\text{H}] < 0.5$. The TOI-5610 metallicity estimate is close to this lower limit.

It should be noted that the metallicity estimates from HPF-SpecMatch entail some caveats. Since M dwarfs are lower temperatures and have atmospheres characterized by molecular features, their metallicities are notoriously difficult to determine (V. M. Passegger et al. 2022). In addition, HPF-SpecMatch provides only loose constraints on metallicity because the χ^2 suffers from some multimodality. The T_{eff} and $\log g$ estimates, however, are considered accurate. The errors for these quantities come from cross validation comparing values recovered from HPF-SpecMatch with library values.

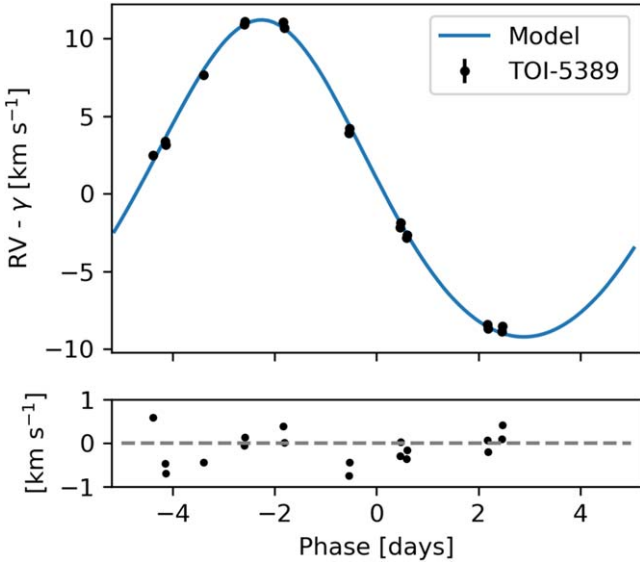


Figure 5. Phased barycentric radial velocity curve and best-fit model for TOI-5389A. Residuals are plotted below. Systemic velocity γ is subtracted. Error bars are plotted but are too small to see.

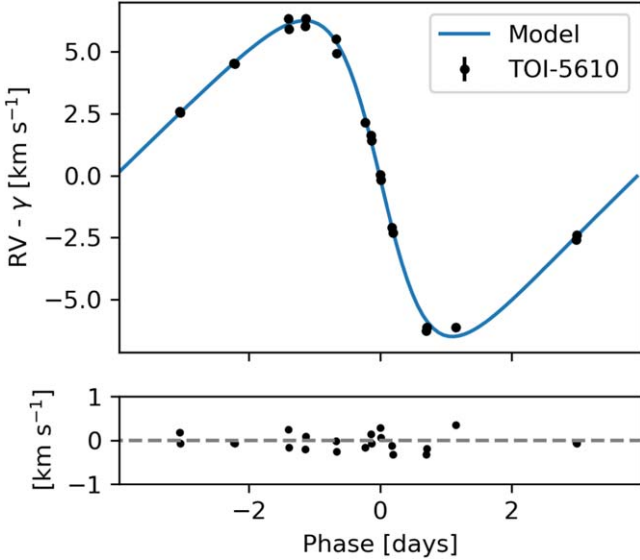


Figure 6. Phased barycentric radial velocity curve and best-fit model for TOI-5610. Residuals are plotted below. Systemic velocity γ is subtracted. Error bars are plotted but are too small to see.

Table 5 includes these best-fitting stellar parameters of $T_{\text{eff}} = 3569 \pm 59$ K, $\log g = 4.79 \pm 0.04$, and $[\text{Fe}/\text{H}] = -0.15 \pm 0.16$ for TOI-5389A, and parameters of $T_{\text{eff}} = 3618 \pm 59$ K, $\log g = 4.78 \pm 0.04$, and $[\text{Fe}/\text{H}] = -0.42 \pm 0.16$ for TOI-5610. The temperatures and gravities of both targets correspond to M1V–M2V stars as defined by T. S. Boyajian et al. (2012).

We took these spectroscopic parameters and used them as priors in a spectral energy distribution (SED) fit using EXOFASTv2 (J. D. Eastman et al. 2019) to obtain stellar mass and radius. EXOFASTv2 utilizes MIST (J. Choi et al. 2016; A. Dotter 2016) stellar-evolutionary models and tracks in conjunction with NextGen (F. Allard et al. 2012) stellar atmosphere models. In addition to stellar T_{eff} , $\log g$, $[\text{Fe}/\text{H}]$, and the Gaia parallax, we used Pan-STARRS *griz*, 2MASS *JHK*, and WISE 1 and 2 photometric measurements as listed in

Table 5 as constraints. Table 5 lists best-fit results of $M_* = 0.43 \pm 0.02 M_\odot$ and $R_* = 0.42 \pm 0.01 R_\odot$ for TOI-5389A. It also lists best-fit results of $M_* = 0.53 \pm 0.02 M_\odot$ and $R_* = 0.52 \pm 0.01 R_\odot$ for TOI-5610. As an alternative method of radius estimation we used a Python program (R. Kiman et al. 2024a) based on the R. Kiman et al. (2024b) method which uses surface brightness–color relations for the three Gaia magnitudes with Gaia DR3 parallaxes and found well-agreeing results of $0.422 \pm 0.005 R_\odot$ and $0.513 \pm 0.006 R_\odot$ for TOI-5389A and TOI-5610, respectively. We chose to use the EXOFASTv2 values since they are based on actual measurements of the star and the availability of the posteriors.

3.2. TOI-5389B: White-dwarf Companion

K. El-Badry et al. (2021) lists TOI-5389A as a wide binary with the white-dwarf star Gaia DR3 765082934121253376 (listed as spectral type DC, henceforth TOI-5389B). TOI-5389B has a *G* mag of 20.19, and the projected separation is 12.6, corresponding to 2447 au at a distance of 196 pc (TOI-5389A). Their respective Gaia DR3 parallaxes are 5.1298 ± 0.057 mas and 7.6227 ± 1.2444 mas, which are in agreement near the $\sim 2\sigma$ level. The catalog also reports that the likelihood of this being a chance alignment of proper motions is extremely small ($R \sim 10^{-9}$), implying that these objects constitute a bound system. The catalog does not contain TOI-5610, and so we assume no binarity.

Using DA and DB white-dwarf cooling curves²³ based on J. B. Holberg & P. Bergeron (2006), P. M. Kowalski & D. Saumon (2006), P. E. Tremblay et al. (2011), S. Blouin et al. (2018), and A. Bédard et al. (2020), we fit the SED of TOI-5389B using a χ^2 minimization method. We used the Gaia *G*, *BP*, *RP*, and SDSS *u*, *g*, *r*, *i*, *z* magnitudes with the TOI-5389A parallax to find the best-fitting intrinsic white-dwarf SED. The best-fit model to TOI-5389B yields a spectral type DB with $M = 0.573 M_\odot$, $T_{\text{eff}} = 6500$ K, $\log g = 7.0$, and a post-MS age of ~ 2 Gyr. Therefore, TOI-5389A must be at least 2 Gyr old.

3.3. Gyrochronology

Adopting an upper limit on $v \sin i$ of 2 km s^{-1} , corresponding to the minimum detectable rotational velocity of HPF, and using the radii of the stars, we estimate lower limits on the rotation periods of 11 and 13 days for TOI-5389A and TOI-5610, respectively. A Lomb–Scargle analysis of the TESS, Zwicky Transient Facility (ZTF DR22; E. C. Bellm et al. 2019; F. J. Masci et al. 2019), and All Sky Automated Survey for SuperNovae (ASAS-SN Sky Patrol v2.0; B. J. Shappee et al. 2014; C. S. Kochanek et al. 2017) data produced no identifiable stellar rotational period, so we proceed with our limiting estimates. S. G. Engle & E. F. Guinan (2023) presents several relations for M dwarf gyrochronology, we use the M0–2V relation for $P < 23.4933$ days. For both stars this suggests ages of $\gtrsim 0.5$ Gyr. The lack of rotational broadening and photometric rotational modulation are indicative of an inactive star.

Separately, EXOFASTv2 yields SED-based ages of $10.0^{+2.6}_{-4.0}$ and $10.0^{+2.6}_{-3.9}$ Gyr for TOI-5389A and TOI-5610, respectively.

3.4. Galactic Population Analysis

Table 6 shows galactic velocities in barycentric and local standard of rest (LSR; R. Schönrich et al. 2010) reference

²³ <http://www.astro.umontreal.ca/bergeron/CoolingModels>

Table 5
Stellar Parameters

Parameter	TOI-5389A	TOI-5610	Source
R.A.	11:15:12.4	10:24:05.48	Gaia
decl.	+39:21:32.11	+48:14:53.99	Gaia
R.A. PM (mas yr ⁻¹)	-33.0227 ± 0.0438	-27.3901 ± 0.0328	Gaia
decl. PM (mas yr ⁻¹)	-46.7001 ± 0.0404	-55.2985 ± 0.0316	Gaia
<i>g</i>	17.526 ± 0.015	17.289 ± 0.014	Pan-STARRS
<i>r</i>	16.310 ± 0.008	16.153 ± 0.010	Pan-STARRS
<i>i</i>	15.26 ± 0.007	15.290 ± 0.013	Pan-STARRS
<i>z</i>	14.796 ± 0.016	14.896 ± 0.010	Pan-STARRS
<i>J</i>	13.402 ± 0.024	13.598 ± 0.023	2MASS
<i>H</i>	12.81 ± 0.018	12.964 ± 0.024	2MASS
<i>K</i>	12.555 ± 0.023	12.781 ± 0.026	2MASS
W1	12.461 ± 0.024	12.674 ± 0.023	WISE
W2	12.323 ± 0.024	12.542 ± 0.024	WISE
W3	12.016 ± 0.327	12.01 ± 0.272	WISE
<i>T</i> _{eff} (K)	3569 ± 59	3618 ± 59	This work
log <i>g</i>	4.79 ± 0.04	4.78 ± 0.04	This work
<i>M</i> _* (<i>M</i> _⊙)	0.43 ± 0.02	0.53 ± 0.02	This work
<i>R</i> _* (<i>R</i> _⊙)	0.42 ± 0.02	0.52 ± 0.02	This work
[Fe/H] ^a	-0.15 ± 0.16	-0.42 ± 0.16	This work
Parallax (mas)	5.1298 ± 0.057	3.6509 ± 0.037	Gaia
ρ_* (g cm ⁻³)	8.19 ± 1.23	5.32 ± 0.65	This work
<i>L</i> (<i>L</i> _⊙)	0.026 ± 0.004	0.042 ± 0.005	This work
Sp-Type	M1V–M2V	M1V–M2V	This work

Note.^a These numbers entail some caveats. See Section 3.1.**Table 6**

Galactic Velocities and Probabilities of Being in the Thin Disk, Thick Disk, or Halo

Parameter	TOI-5389A	TOI-5610
<i>U</i>	-9.02	2.10
<i>U</i> (LSR)	2.08 ± 0.17	13.20 ± 0.24
<i>V</i>	-51.31	-82.81
<i>V</i> (LSR)	-39.07 ± 0.60	-70.57 ± 0.78
<i>W</i>	-14.41	-37.96
<i>W</i> (LSR)	-7.16 ± 0.10	-30.71 ± 0.09
\mathcal{P} (thin)	0.98 ± 0.02	0.69 ± 0.05
\mathcal{P} (thick)	0.02 ± 0.00003	0.30 ± 0.002
\mathcal{P} (halo)	0.00005 ± 0.0000005	0.002 ± 0.00002

frames as well as the probabilities of being associated with the thin disk, thick disk, or halo populations. Based on these results, TOI-5610 could belong to either the galactic thin or thick disk. This agrees well with our metallicity estimates, and the velocities allow for the possibility of being an old star. (J. Holmberg et al. 2009; S. Sharma et al. 2014; H.-C. Hwang & N. L. Zakamska 2020). TOI-5389A has kinematics consistent with thin disk stars, which agrees with our metallicity estimates but allows a very wide range of ages.

3.5. Physical and Orbital Parameters

We also used EXOFASTv2 to derive the companions' physical properties and orbital parameters by jointly modeling the four lightcurves and the radial velocity curve of each target. It should be noted that this fit also included an SED fit, which reproduced the initial values. We adopted normal stellar priors from Table 5 along with loose uniform priors on extinction,

Table 7

Physical and Orbital Parameters

Parameter	TOI-5389Ab	TOI-5610b
<i>P</i> (days)	10.40046 ± 0.00002	7.95346 ± 0.00002
<i>t</i> ₀ (BJD _{TDB})	2459609.445 ^{+0.00096} _{-0.00091}	2459628.939 ^{+0.0013} _{-0.0012}
<i>b</i> (<i>R</i> _*)	0.386 ^{+0.09} _{-0.13}	0.124 ^{+0.12} _{-0.09}
<i>b</i> _s (<i>R</i> _*)	0.396 ^{+0.08} _{-0.13}	0.26 ^{+0.24} _{-0.18}
<i>e</i>	0.096 ^{+0.003} _{-0.005}	0.354 ^{+0.011} _{-0.012}
γ (m s ⁻¹)	-11300 ⁺¹¹⁰ ₋₁₀₀	-43387 ⁺⁶⁵ ₋₆₆
<i>R</i> _{BD} (<i>R</i> _J)	0.776 ^{+0.035} _{-0.033}	0.887 ^{+0.031} _{-0.031}
<i>M</i> _{BD} (<i>M</i> _J)	68.0 ^{+2.2} _{-2.2}	40.4 ^{+1.0} _{-1.0}
log <i>g</i> _{BD}	5.394 ^{+0.030} _{-0.032}	5.105 ^{+0.029} _{-0.029}
ω (deg)	8.2 ^{+9.7} _{-9.1}	94.0 ^{+1.6} _{-1.6}
<i>a</i> (au)	0.0739 ^{+0.001} _{-0.001}	0.0647 ^{+0.0007} _{-0.0007}
<i>a</i> / <i>R</i> _*	37.90 ^{+0.91} _{-0.91}	26.76 ^{+0.59} _{-0.57}
<i>R</i> _{BD} / <i>R</i> _*	0.202 ^{+0.004} _{-0.004}	0.175 ^{+0.005} _{-0.005}
<i>M</i> _{BD} / <i>M</i> _*	0.150 ^{+0.003} _{-0.003}	0.072 ^{+0.001} _{-0.001}
<i>i</i> (deg)	89.54 ^{+0.24} _{-0.16}	89.59 ^{+0.28} _{-0.38}
<i>t</i> _s (BJD _{TDB})	2459604.867 ^{+0.024} _{-0.038}	2459632.783 ^{+0.052} _{-0.052}
<i>T</i> _{SB} (K)	1113	1049
Transit Depth (I) (ppm)	40000 ⁺¹⁸⁰⁰ ₋₁₉₀₀	35000 ⁺²⁰⁰⁰ ₋₂₀₀₀
Occultation Depth (K) (ppm)	500	200

eccentricity, equivalent evolutionary phase, systemic velocity, and system inclination. We required the default criteria for convergence given by EXOFASTv2 which is a Gelman–Rubin statistic <1.01 and *t*_z (number of independent chains) >1000. We examined the probability-distribution functions to ensure they were approximately Gaussian and that all chains converge on the same solution. Table 7 shows the best-fitting final

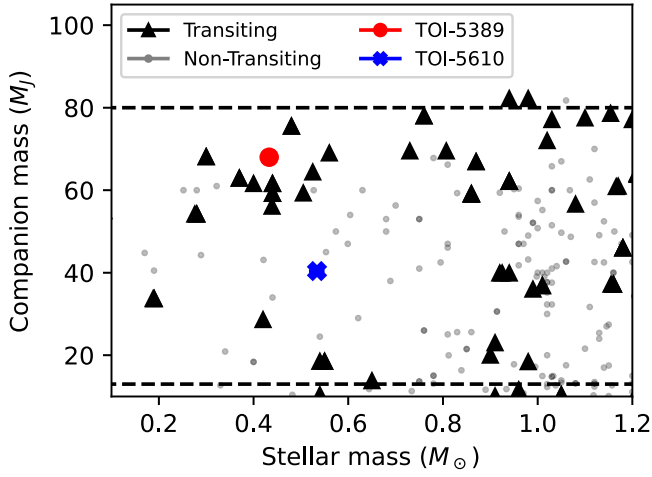


Figure 7. Host mass vs. companion mass for known transiting BDs. Black triangles denote transiting systems while gray dots denote nontransiting. The red dot denotes TOI-5389Ab and the blue \times denotes TOI-5610b. The dashed lines mark the upper and lower mass limits for BDs at 80 and 13 M_J , respectively.

parameters with uncertainties for the companions and their orbits, including period (P), time of primary eclipse (t_0), impact parameter (b), secondary impact parameter (b_s), eccentricity (e), systemic velocity (γ), BD radius (R_{BD}), BD mass (M_{BD}), BD surface gravity ($\log g_{BD}$), argument of periastron of the star (ω), semimajor axis (a), inclination (i), time of secondary eclipse (t_s), Sonora–Bobcat (M. S. Marley et al. 2021) model-based estimated BD temperature (T_{SB} , discussed in Section 4.2), transit depth in Johnson I band, occultation depth in Johnson I band, and occultation depth in the Johnson K band. For TOI-5389Ab we find a mass of $68.0^{+2.2}_{-2.2} M_J$ —near the upper mass limit for BDs, a radius of $0.824^{+0.033}_{-0.031} R_J$, and an eccentricity of $0.096^{+0.003}_{-0.005}$. For TOI-5610b we find a mass of $40.4^{+1.0}_{-1.0} M_J$, a radius of $0.887^{+0.031}_{-0.031} R_J$, and a moderate eccentricity of $0.354^{+0.011}_{-0.012}$.

4. Discussion

4.1. Population Context

We constructed a literature comparison sample of BDs, originating largely from exoplanet.eu ($m \sin i < 60 M_J$), A. T. Stevenson et al. (2023; $13 M_J < m < 80 M_J$), and S. P. Schmidt et al. (2023; $13 M_J < m < 80 M_J$), with additions/updates from E. Page et al. (2024), N. Grieves et al. (2021), P. Benni et al. (2021), and B. A. Henderson et al. (2024b). We limited the sample to BDs with main sequence (MS) host stars and masses $M > 10 M_J$, totaling 393 systems. Of these, 131 are transiting, and the remaining nontransiting systems only yield $m \sin i$, a minimum mass of the BD.

Figure 7 plots companion mass versus host mass. The dashed lines denote the upper and lower mass limits for BDs at 80 and 13 M_J , respectively. TOI-5389Ab is one of the most massive BDs orbiting an M dwarf and is near the upper mass limit (80 M_J). TOI-5610b lies in the intermediate and less populated (by transits) region around 40 M_J . Among transiting BDs, there appears to be a preference toward higher masses, near the upper limit. This should not be a selection effect since radius (and therefore transit depth and transit probability) weakly scales inversely with mass for BDs; transiting low-mass brown dwarfs should be more likely to be detected. The figure

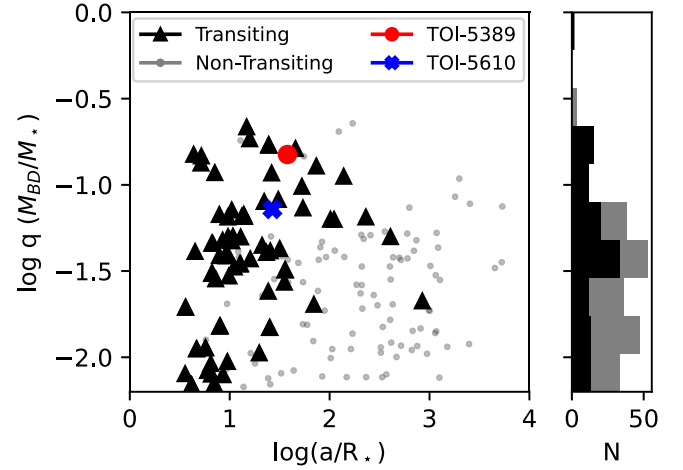


Figure 8. Comparison plot of mass ratio vs. semimajor axis. Black triangles denote transiting systems while gray dots denote nontransiting. The red dot denotes TOI-5389Ab and the blue \times denotes TOI-5610b. The histogram on the right shows the distribution in $\log q$, transiting systems in black.

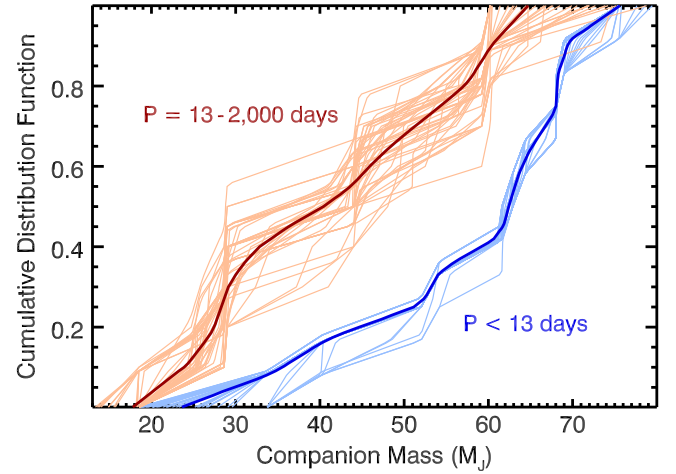


Figure 9. Cumulative distribution function for M-dwarf/BD systems divided into two populations via period. The light lines represent differing MCMC chains while the dark lines represent their averages.

also shows a paucity of low-mass BDs around late-type stars. Even low-mass BDs impart a large RV signal on the host star so that they are not likely to be underreported given modern RV sensitivities. While the comparison sample used here cannot be considered unbiased or complete, these two features considered together could be explained by an evolutionary process in which BDs accrete mass as they undergo inward migration in a gaseous disk, leading to a preponderance of close, high-mass BDs (M. R. Bate et al. 2002a, 2002b; M. R. Bate 2009; M. Moe & K. M. Kratter 2018; A. Tokovinin & M. Moe 2020).

Figure 8 plots the log of mass ratio versus the log of semimajor axis in units of the host star radius. For nontransiting systems the plotted mass ratio is a lower limit. The paucity of systems with $\log q \gtrsim -1.0$ is a consequence of sample construction, since we do not include stellar companions; furthermore, low-mass hosts comprise only a fraction of the comparison sample. The lack of transiting systems at $a/R_* > 100$ is understood as a consequence of the small separations required to detect a transit via TESS or other systematic transit finding programs. TOI-5389Ab has one of

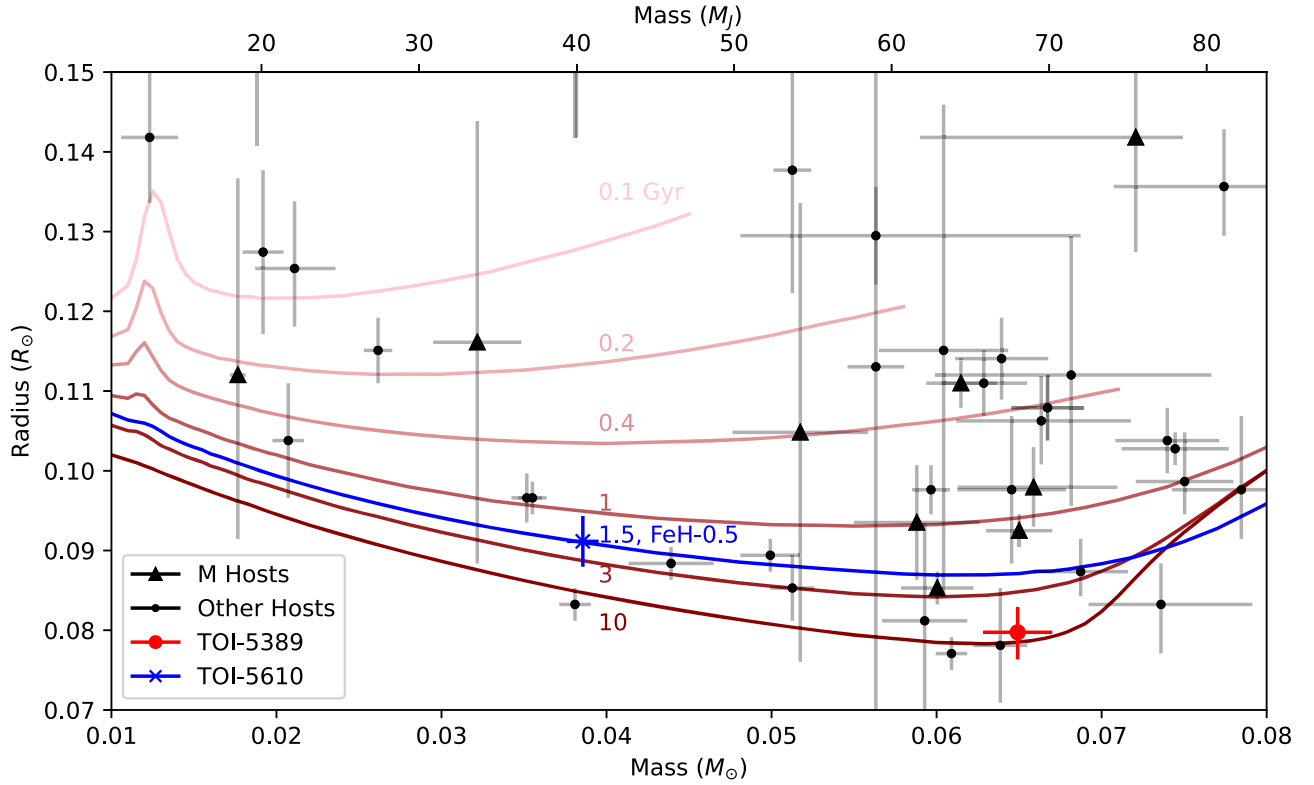


Figure 10. The Sonora–Bobcat models on a plot of companion radius vs. mass. The red lines represent isochrones of 0.1, 0.2, 0.4, 1, 3, and 10 Gyr (from light to dark), all at solar metallicity ($[\text{Fe}/\text{H}] = 0.0$). The blue line shows the nearest model to TOI-5610b which is 1.5 Gyr at $[\text{Fe}/\text{H}] = -0.5$. Only transiting systems are plotted. These models assume cloud-free atmospheres and solar C/O ratios.

the highest mass ratios for BD-MS systems, and TOI-5610 lies among the upper third. The histogram at right depicts the relative frequencies of $\log q$. There appears to be a preponderance of systems near $q \approx 0.03$ ($\log q \approx -1.5$). Given the probable biases and incompleteness of this comparison sample, we draw no further conclusions from this.

Compared to the measured occurrence rates of stellar and planetary mass companions, solar-type FGK dwarfs exhibit a dearth of BD companions within $a < 1$ au commonly known as the BD desert (D. Grether & C. H. Lineweaver 2006; S. Csizmadia et al. 2015). The mass-ratio distribution of short-period A-dwarf host binaries also exhibits a rapid turnover below $q < 0.1$ (S. J. Murphy et al. 2018). Alternatively, BD companions at longer periods are plentiful; the companion mass distribution of solar-type binaries across intermediate separations $a = 10$ –100 au is relatively uniform above $M_2 > 13 M_J$ (E. L. Nielsen et al. 2019; K. Wagner et al. 2019), and exhibits no preference for host mass. Close companions originally fragmented on large protostellar disk or molecular core scales followed by inward migration as the binary accreted from the surrounding gaseous disk/envelope (M. R. Bate et al. 2002a, 2002b; M. R. Bate 2009; M. Moe & K. M. Kratter 2018). Only a small fraction of BD companions can migrate below $a < 1$ au through the disk without also accreting into the stellar-mass regime (A. Tokovinin & M. Moe 2020).

Our population of M-dwarf primaries with BD companions is sufficiently large to probe the nature of the BD desert for low-mass stars. We limit our sample to the 25 systems with $P < 2000$ days, $M_* = 0.08$ – $0.60 M_\odot$, and either $M_{\text{BD}} = 13$ – $80 M_J$ or $M_{\text{BD}} \sin i = 10$ – $80 M_J$. We divide this sample evenly between the 13 systems with short periods below $P <$

13 days, including our two objects TOI-5389Ab and TOI-5610b, and the 12 systems across $P = 13$ –2000 days. We display the cumulative distribution functions (CDFs) of companion masses for both subsets in Figure 9. We adopt the measured companion masses for the eclipsing and astrometric binaries. For the two and six systems with only RV measurements (only $M_{\text{BD}} \sin i$) in our short and long-period subsets, respectively, we utilize a Monte Carlo technique to simulate random orientations. We synthesize 1000 CDFs for both samples, where we draw $\cos i = U[0,1]$ from a uniform distribution for those systems without inclination measurements. We then limit each simulated population to the interval $M_{\text{BD}} = 13$ – $80 M_J$, where in most cases we remove the one object with $M_{\text{BD}} \sin i = 10.4 M_J$ and one of the two systems with $M_{\text{BD}} \sin i \geq 54 M_J$. In Figure 9, we display the first 50 Monte Carlo simulations and the average CDF for both the short and long-period subsets.

The short-period subset is weighted toward systematically more massive BD companions compared to the long-period subset, similar to the trend observed for more massive primaries. Only $2/13 = 15\%$ of the short-period systems have $M_{\text{BD}} = 13$ – $40 M_J$ whereas $6/12 = 50\%$ of the long-period systems have such low-mass BD companions. For each simulated pair of CDFs, we compute the two-sample Kolmogorov–Smirnov (KS) and Anderson–Darling (AD) statistics. The short- and long-period subsets are discrepant with each other at the $p = 0.03$ (2.2σ) and $p = 0.02$ (2.6σ) probability levels according to the average KS and AD statistics, respectively. This provides the first tentative evidence that M-dwarf primaries exhibit a dearth of $q < 0.1$ ($M_{\text{BD}} < 40 M_J$) companions at short periods compared to slightly wider systems, consistent with the BD desert observed

for AFGK primaries. This suggests that close M-dwarf/BD binaries formed via disk/core fragmentation, inward disk migration, and circumbinary accretion similar to their more massive solar-type counterparts. This split at $\sim 40 M_J$ is also close to the proposed population split at $42.5 M_J$ found in B. Ma & J. Ge (2014). However, this sample size is still small and combines different detection techniques, and so could have unknown systematics.

4.2. Age via the Sonora–Bobcat Models

Figure 10 shows the literature sample of transiting BDs and our two targets compared to the Sonora–Bobcat BD models (M. S. Marley et al. 2021) on a plot of BD radius versus mass. The red lines represent isochrones of 0.1, 0.2, 0.4, 1, 3, and 10 Gyr (from light to dark), all at solar metallicity ($[Fe/H] = 0.0$). The blue line shows the nearest model to TOI-5610b which is 1.5 Gyr at $[Fe/H] = -0.5$. Our best estimates based on these models put TOI-5389Ab at 8 Gyr and TOI-5610b at 1.5 Gyr. These two targets are among the older half of the comparison sample. These models assume a cloudless atmosphere and solar C/O ratio, which could explain discrepancies in our multiple age estimates.

4.3. Predictions of the Secondary Eclipses

Estimates of the BD temperature from the closest Sonora model corresponding to the adopted BD mass and radius (also consistent with age) indicate an effective temperature of 1113 K for TOI-5389Ab and 1049 K for TOI-5610b. Therefore, internal heating due to gravitational collapse and deuterium fusion dominate the temperature profile of these objects, reflected light would be negligible at these separations, and we use these temperatures to predict secondary depths. The orbital parameters for each brown dwarf listed in Table 7 predict²⁴ 1–2 hr duration secondary eclipses at the reference times t_s , despite the appreciable eccentricity of both targets. The estimated BD temperatures result in secondary eclipse depths of $\lesssim 10$ ppm in the Johnson *I* band for both targets. Predicted eclipses are deeper in Johnson *K* band, with estimates of 500 ppm for TOI-5389A and 200 ppm for TOI-5610. We searched the TESS photometry for secondary transits and, unsurprisingly, find none at the predicted times, given that the 1% photometric rms is far larger than the predicted depths at that bandpass. An accurate prediction for the BD secondary eclipse depth as a function of wavelength would require use of an appropriate BD atmosphere model and is beyond the scope of the nominal secondary transit estimates given here.

5. Conclusion

We have reported the characterization of two brown dwarfs orbiting M dwarfs via four observed transits from TESS and ground-based photometry in conjunction with high-resolution infrared spectroscopy. TOI-5610b has a moderately high eccentricity. TOI-5389Ab has one of the highest companion-to-host mass ratios of transiting BDs ($M_{BD}/M_* = 0.150$), near the hydrogen burning limit. Despite their eccentricities both targets do have potentially observable secondary transits, though shallow. We computed a variety of age estimates.

²⁴ We modeled the lightcurve of each system with the eclipsing binary star package PHOEBE2 v2.4 (A. Prša et al. 2016), using the adopted stellar and BD parameters tabulated above and assuming blackbody spectral energy distributions for the BDs.

The EXOFASTv2 SED fit yields ~ 10 Gyr for both targets. The WD SED fit to TOI-5389B yields a minimum post-MS age of $\gtrsim 2$ Gyr, which is likely our strongest constraint on the age. The (cloud-free) Sonora–Bobcat BD models yield 8 and 1.5 Gyr for TOI-5389Ab and TOI-5610b, respectively. Gyrochronology provide loose limits of $\gtrsim 0.5$ Gyr. Galactic kinematics for TOI-5610 imply an older star from the thin or thick disk, while kinematic ages are inconclusive for TOI-5389A. A statistical analysis of M-dwarf/BD systems reveals for the first time that those at short orbital periods ($P < 13$ days) exhibit a dearth of $13 M_J < M_{BD} < 40 M_J$ companions ($q < 0.1$) compared to those at slightly wider separations, similar to the BD desert previously observed for solar-type primaries.

These two BD characterizations add to the census of brown dwarfs with well-constrained characteristics and provide new data for differentiating formation mechanisms of substellar objects. A compilation of transiting brown-dwarf parameters from the literature suggests that high-mass BDs are more abundant than low-mass BDs—a distribution that may find an explanation in terms of BD migration within a gaseous disk during formation.

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Software: EXOFASTv2 (J. Eastman et al. 2013), barycorrpy (S. Kanodia & J. Wright 2018), HPF-SpecMatch (G. Stefansson et al. 2020), TESS-Gaia Light-Curve (T. Han & T. D. Brandt 2023), PHOEBE2 v2.4 (A. Prša et al. 2016), AstroImageJ (K. A. Collins et al. 2017), Matplotlib (J. D. Hunter 2007), pandas v1.5.3 (W. McKinney 2010; The pandas development team 2020), numpy (C. R. Harris et al. 2020), Kiman Radius-Estimation v1.0.0 (R. Kiman et al. 2024a).

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