

1 **Title**

2 A continuum reaction-diffusion model for spread of gene silencing in chromosomal inactivation

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17

18 **Abstract**

19 Regulation of gene silencing in large regions of chromosomes is crucial for development and disease
20 progression, and there has been an increasing interest in using it for new therapeutics. One example of
21 massive gene silencing is X chromosome inactivation (XCI), a process essential for dosage compensation
22 of X-linked genes. During XCI, most genes in the X chromosome are inactivated following the transcription
23 of XIST, an X-linked long noncoding RNA. Recent experiments with transgenes showed that the spread of
24 gene silencing can be induced by XIST transcription in *cis*, but the spread is restricted in space. The
25 mechanism of controlling the spread remains unclear. In this work, we develop a continuum reaction-
26 diffusion model that elucidates chromosomal inactivation through a bistable system governed by a
27 regulatory network for XIST-mediated gene silencing. We find that the spread of XIST can be tuned by
28 known negative feedback loops regulating its synthesis and degradation, and that the spread of gene
29 silencing is controlled by a wave-pinning mechanism in which both global regulation of silencing complex
30 and local variations of histone modifications can play crucial roles. In addition, we integrate the discrete
31 three-dimensional arrangement of the X chromosome and autosomes into this continuous model. We use
32 a 3D chromosome structure inferred from experimental data and our modeling framework to show the
33 spatiotemporal regulation for spread of gene silencing. Our method enables the investigation for the
34 inactivation dynamics of large regions of chromosomes with varying degrees of the spread of gene
35 silencing. Our model provides mechanistic insights that quantitatively relate gene regulatory networks to
36 tunability and stability of chromosomal inactivation.

37 **Author Summary**

38 Precise control of gene expression is a fundamental process in biology and turning off large parts of
39 chromosomes is both common in many species and important for development and diseases. A well-
40 known example of chromosomal scale gene silencing is X chromosome inactivation (XCI), which helps
41 balancing gene activity between sexes. XCI is governed by the production of an RNA called XIST, leading to
42 most genes on one X chromosome being turned off. Experiments have shown that XIST can trigger nearby
43 genes to turn off in natural and engineered contexts, but this effect only spreads to specific regions. The
44 exact mechanism for this limit is not very well understood. In this study, the authors created a
45 mathematical model to explain how XIST spreads and controls gene activity. They found that feedback
46 systems involving XIST regulations and chromatin modifications synergize and determine how far the
47 inactivation spreads, through a process similar to a wave that gets "pinned" in place. They also included
48 the 3D shape of chromosomes in their model to better understand how gene silencing happens over time
49 and space. This model helps to quantitatively describe chromosome inactivation with gene regulatory
50 networks.

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56 **Introduction**

57 Gene silencing in large chromosomal regions is essential for the development of mammals. Particularly, X-
58 chromosome inactivation (XCI) is used to achieve dosage compensation in female cells to balance the
59 expression of X-linked genes of the two sexes (Payer and Lee, 2008). Before mammalian XCI, one of the
60 two X chromosomes in each cell of the early female embryo is randomly chosen for inactivation. During
61 XCI, the expressions of X-linked genes in the X chromosome selected to be inactivated (X_i) are repressed
62 in a progressive manner upon the transcription of XIST, an X-linked, long-noncoding RNA (lncRNA) (Heard
63 et al., 1997; Kay et al., 1993). XIST molecules are tethered to multiple chromosomal locations, where they
64 recruit polycomb repressive complexes (PRCs) for gene silencing (Sunwoo et al., 2015). The spread of gene
65 silencing was shown to be influenced by the locations of the genes in both one-dimensional (1D) genomic
66 coordinates and three-dimensional (3D) genome structures. The completion of XCI results in the stable
67 repression of most, but not all, genes in the X chromosome (Balaton and Brown, 2016). The number of X-
68 linked genes that remain to be active after the completion of XCI (i.e. the XCI escape gene, or escapees)
69 can vary across tissues and species (Berletch et al., 2015; Tukiainen et al., 2017). The mechanism
70 controlling the extent of gene silencing spread, which in turn determines the escape genes, in XCI remains
71 elusive.

72 To understand the XIST-mediated gene silencing, multiple groups integrated XIST transgenes into
73 autosomes (Lee and Jaenisch, 1997; Minks et al., 2013; Naciri et al., 2021), and these studies showed the
74 ability of the induced XIST expression in an autosome to silence genes in the adjacency of the XIST
75 integration/transcription site. However, the spread of the gene silencing seems to be limited compared to
76 that in XCI (Minks et al., 2013). Similar to the question of determinants of escape genes, factors
77 contributing to the gene silencing spread upon transgene XIST induction are unclear.

78 Previous mathematical models based on ordinary differential equations (ODEs) describing gene regulatory
79 networks have provided useful insights into the mechanisms of the decision-making process of XCI for
80 robustly selecting one X chromosome in each cell (Li et al., 2016; Mutzel et al., 2019). In addition, agent-
81 based modeling has been used to show the dynamic distribution of XIST molecules over 3D structure of
82 the X chromosome (Lappala et al., 2021). However, these modeling frameworks did not capture the
83 connection between the dynamics of XIST distribution and that of gene silencing. The latter is controlled
84 by complex regulatory networks containing feedback loops in a spatially dependent manner. These
85 feedback loops include the transcriptional and post-transcriptional self-regulations of XIST (Jachowicz et
86 al., 2022; Rodermund et al., 2021) as well as epigenetic modifications that confer gene activity switches
87 and memory (Dodd et al., 2007). The lack of quantitative description of these interconnected regulatory
88 elements has been limiting our understanding of XIST-mediated gene silencing in both natural and
89 engineered systems.

90 In this work, we used a continuum reaction-diffusion framework to model XIST dynamics and gene
91 silencing induced by XIST upregulation during XCI. We show that the feedback loops of XIST's transcription
92 and degradation have significant impacts on the steady state distribution of XIST and the speed to attain
93 the distribution. We describe a wave-pinning mechanism based on chromatin-level feedback between
94 histone modifications and repressive complexes, and the proposed mechanism offers a plausible
95 explanation for controllable spread of gene silencing in the X chromosome. In addition, this generic version
96 of our model gives insights into lncRNA-mediated gene silencing in autosomes. Interestingly, both global
97 regulation of silencing complex and local variations of histone modifications can contribute to wave-
98 pinning. We extended the continuum model to capture 3D chromosome structure and showed the utility

99 of our model in explaining key observations such as escape genes under realistic spatial constraints. Overall,
100 our continuum reaction-diffusion model integrates regulatory components contributing to XCI in a systems
101 approach, and it provides new mechanistic insights into factors controlling the spread of XIST-mediated
102 gene silencing.

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105 **Results**

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107 **The roles of feedback loops controlling XIST expression on dynamics of gene silencing**

108 To understand how regulations of XIST can influence the XIST distribution, we first built a spatiotemporal
109 model with a one-dimensional (1D) domain and a reaction-diffusion system based on partial differential
110 equations (PDEs). The domain can be viewed as a chromosome region projected onto a 1D space (we will
111 consider a more realistic spatial setting later in this manuscript). A key difference between this model and
112 previous model is that our model allows the diffusion of molecules in continuous space and time, which
113 in turn permits flexibility of molecular movements and descriptions of complex gene regulatory networks.
114 The latter advantage allows us to use a simple set of differential equations to study the effects of molecular
115 concentrations and interactions on gene silencing. Specifically, we first considered a spatial temporal
116 module of regulations for XIST transcription, degradation and diffusion. We used a dimensionless PDE

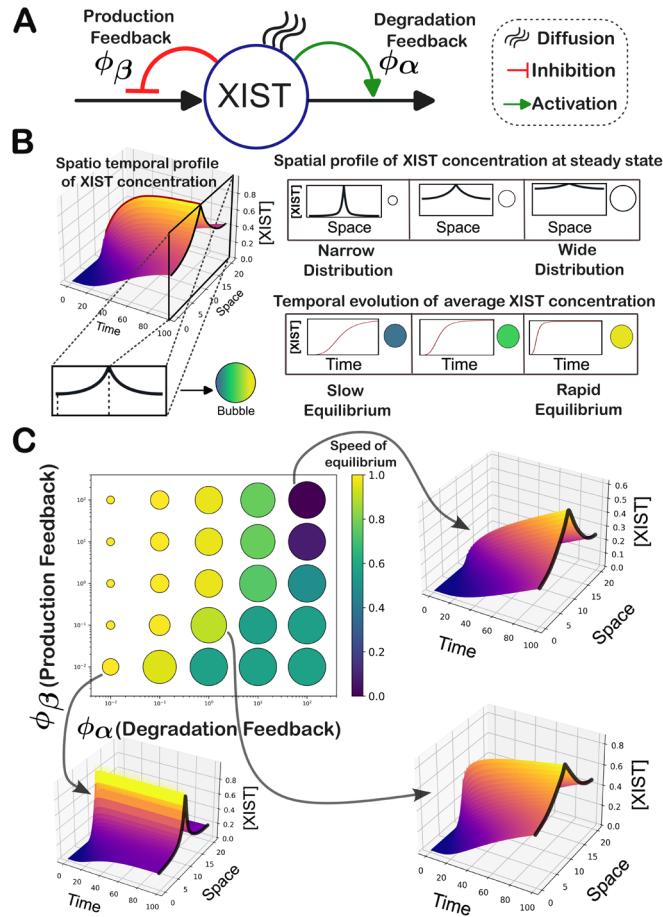
$$117 \frac{\partial [\text{XIST}]}{\partial t} = \frac{\beta_X}{1 + \left[\frac{[\text{XIST}]}{\phi_\beta} \right]^{n_\beta}} - \frac{\alpha_X \left[\frac{[\text{XIST}]}{\phi_\alpha} \right]^{n_\alpha}}{1 + \left[\frac{[\text{XIST}]}{\phi_\alpha} \right]^{n_\alpha}} [\text{XIST}] + D_X \nabla^2 [\text{XIST}] \quad (1)$$

118 to describe these regulations. Here, $[\text{XIST}]$ represents the concentration of XIST molecules, β_X is the basal
119 transcription rate of XIST, α_X is the basal degradation rate constant, D_X is the diffusion coefficient, ϕ_β is
120 the threshold of the inhibition of XIST's transcription by itself, n_β describes the nonlinearity of this
121 transcription feedback (Jachowicz *et al.*, 2022), ϕ_α is the threshold of the activation of XIST's degradation
122 by itself, and n_α describes the nonlinearity of this degradation feedback (Rodermund *et al.*, 2021). Both
123 feedback loops are experimentally observed negative feedback loops (NFLs) in which XIST can limit the
124 level of itself (**Fig 1A**). We selected biologically plausible parameter values to examine the representative
125 dynamical behaviors of the gene regulatory networks (see Supplementary Information). Importantly, the
126 roles of these feedback loops on distributions of XIST and the subsequent gene silencing were unclear due
127 to the lack of rigorous models. To examine NFLs' function in a quantitative framework, we scanned values
128 of ϕ_β and ϕ_α in a grid space with a simple assumption of all other parameters (**Fig 1B and C**). We asked
129 whether the two NFLs can determine how far XIST can reach in the 1D domain with substantial
130 concentration with respect to its level at the transcription site, and how fast XIST distribution can obtain a
131 steady state (**Fig 1B**). Overall, the range of the steady state distribution is negatively correlated with the
132 speed of equilibrium attainment (**Fig 1C**, smaller circles tend to be brighter). However, we found that
133 moderate levels of both feedback gave rise to broad distributions of XIST and relatively rapid attainment
134 of equilibrium (e.g. **Fig 1C** lower right plot. **Figures S1-3**). It should be noted that the broad distribution
135 and rapid equilibrium performance can be achieved simply by increasing the diffusion coefficient of the
136 molecule. However, this way of obtaining the performance would require substantial changes of the

137 physical properties of the molecule (e.g. size). Therefore, our results suggest a powerful approach of tuning
138 the dynamical distribution of XIST by expression regulations alone.

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141

142 **Figure 1. XIST module and its simulation results.** A. A regulatory network for controlling XIST levels. This
143 network was used to construct a PDE model to examine the time dependent spatial distribution of XIST.
144 Two feedback loops at degradation and transcription levels are controlled by two threshold parameters
145 ϕ_α and ϕ_β respectively. B. Simulation of the 1D PDE model (example result shown on left), and two
146 metrics assessing the width and the speed of XIST spread (right). C. A grid search for the strengths of the
147 feedback loops (small ϕ corresponds to lower threshold and stronger feedback) and the associated results
148 of XIST spread width and speed. In these simulations, $n_\beta = n_\alpha = 4$, $\alpha_X = \beta_X = 1$, and $D_X = 1$.

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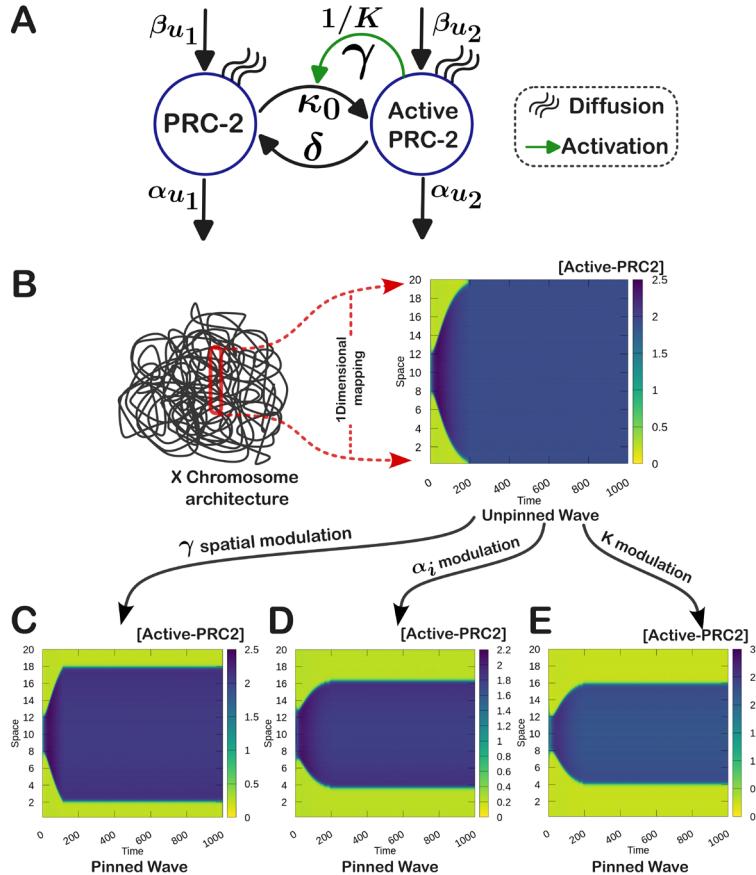
150 Tunability of the spread of gene silencing by *cis*- and *trans*-acting factors

151 Next, we focus on a module downstream of XIST that is responsible for the spread of gene silencing. Recent
152 experiments suggest that polycomb repressive complexes (PRCs) 1 and 2, both critical for XCI (Dixon-
153 McDougall and Brown, 2021; Masui et al., 2023), are involved in reciprocal regulations with histone post-
154 translational modifications. Specifically, PRC2 is activated by its enzymatic product H3K27me3 (Stafford et

155 al., 2018). Likewise, histone H2A monoubiquitination (H2Aub) binds and stimulates PRC2 (Kalb et al., 2014),
156 while H2Aub is an enzymatic product of PRC1 which in turn can be recruited by PRC2 (Fischle et al., 2003;
157 Min et al., 2003). These experiments support a positive feedback loop between repressive complexes and
158 the gene-silencing permissive post-translational modifications of histones. We use the following PDEs to
159 describe the dynamics of Polycomb repressive complex PRC2 (as a representative factor for chromatin
160 modification and a transcription repressor) in response to XIST activation:

$$\begin{aligned} 161 \quad \frac{\partial u_1}{\partial t} &= f(u_1, u_2) + \beta_{u_1} - \alpha_{u_1} u_1 + f_s([XIST]) + D_{u_1} \nabla^2 u_1 \\ 162 \quad \frac{\partial u_2}{\partial t} &= -f(u_1, u_2) + \beta_{u_2} - \alpha_{u_2} u_2 - f_s([XIST]) + D_{u_2} \nabla^2 u_2 \\ 163 \quad \frac{\partial g}{\partial t} &= \frac{\beta_g}{K^2 + u_1^n} - \alpha_g g, \end{aligned} \quad (2)$$

164 where $f(u_1, u_2) = u_2 [\kappa_0 + \gamma u_1^n / (K^2 + u_1^n)] - \delta u_1$, u_1 is concentration of the active form of PRC2, u_2 is
165 the concentration of the inactive form of PRC2. g represents the activity of any gene in a region of
166 chromosome. $f(u_1, u_2)$ represents the interconversion rates between the two forms of PRC2 (**Fig 2A**). The
167 conversion from the inactive form to the active form is assumed to be influenced by the active PRC2 via a
168 positive feedback loop, which is supported by the confirmational change of the complex induced by
169 H3K27me (Stafford et al., 2018), a histone modification catalyzed by PRC2. The feedback is characterized
170 by γ , the maximum conversion rate controlled by the feedback, K the threshold of self-activation, and n
171 the Hill exponent controlling the nonlinearity of the feedback. κ_0 is the basal conversion rate constant
172 from u_2 to u_1 , and δ is the conversion rate constant from u_1 to u_2 . α_{u_1} and α_{u_2} are the disassembly rate
173 of the active and inactive forms of PRC2 respectively. β_{u_1} and β_{u_2} are the assembly rate of the active and
174 inactive forms of PRC2 respectively. D_{u_1} and D_{u_2} are the diffusion coefficients of the active and inactive
175 forms of PRC2 respectively. The function $\beta_g / (1 + u_1)$ describes the rate of gene deactivation triggered by
176 u_1 , where β_g is the basal activity. α_g is the relaxation rate constant of the gene activity. f_s is the function
177 describing how XIST transiently influences the PRC2 activation (see details below and the rationale for
178 choosing parameter values in the Supplementary Information).



179

180 **Figure 2. Chromatin module and its simulation results.** **A.** A regulatory network for controlling Inactive
 181 and Active PRC2 levels. This network was used to construct a PDE model to examine the time dependent
 182 spatial distribution of Inactive and Active PRC2 molecules. **B.** Schematic illustration of the mapping
 183 procedure of certain chromosomal regions in one spatial dimension. The heatmap represents the time
 184 evolution of the Active PRC2 wave leading to homogeneous spatial profile (where, the blue color
 185 represents the inactive regions, and the green color represents the active regions of the chromosome). **C.**
 186 Attainment of spatial heterogeneity in terms of pinned wave by spatially modulating the more sensitive
 187 feedback coefficient γ . **D.** Visualization of the pinned wave achieved through enhancement of degradation
 188 coefficient α_X relative to the unpinned state. **E.** A pinned wave realized by increasing the value of the
 189 feedback threshold parameter K , which has a lower sensitivity towards the dynamics of XIST molecules.
 190 Parameter values for the simulations in this figure are listed in **Table S1**.

191

192 To better understand our model in a stepwise fashion, we modularized our model by separately
 193 considering the dynamics of PRC2 complex and XIST. We first focus on a *chromatin module* describing PRC2
 194 and gene activities (Eq 2) which does not include XIST explicitly (**Fig 2A**), and f_s in Eq 2 is assumed to be a
 195 transient, imposed stimulation at the beginning of the simulation (e.g. $1 < t < 4$) that only occurs in the
 196 middle of the spatial domain (see Supplementary Information for details). The dynamics of XIST will be
 197 considered more carefully in a later section when we integrate two modules (Eq 1 and Eq 2). With the
 198 chromatin module, we simulated the model at a 1D chromosome domain (**Fig 1B**), and we introduced a
 199 local perturbation in the middle of the domain (presumably triggered by tethering of XIST to chromatin)

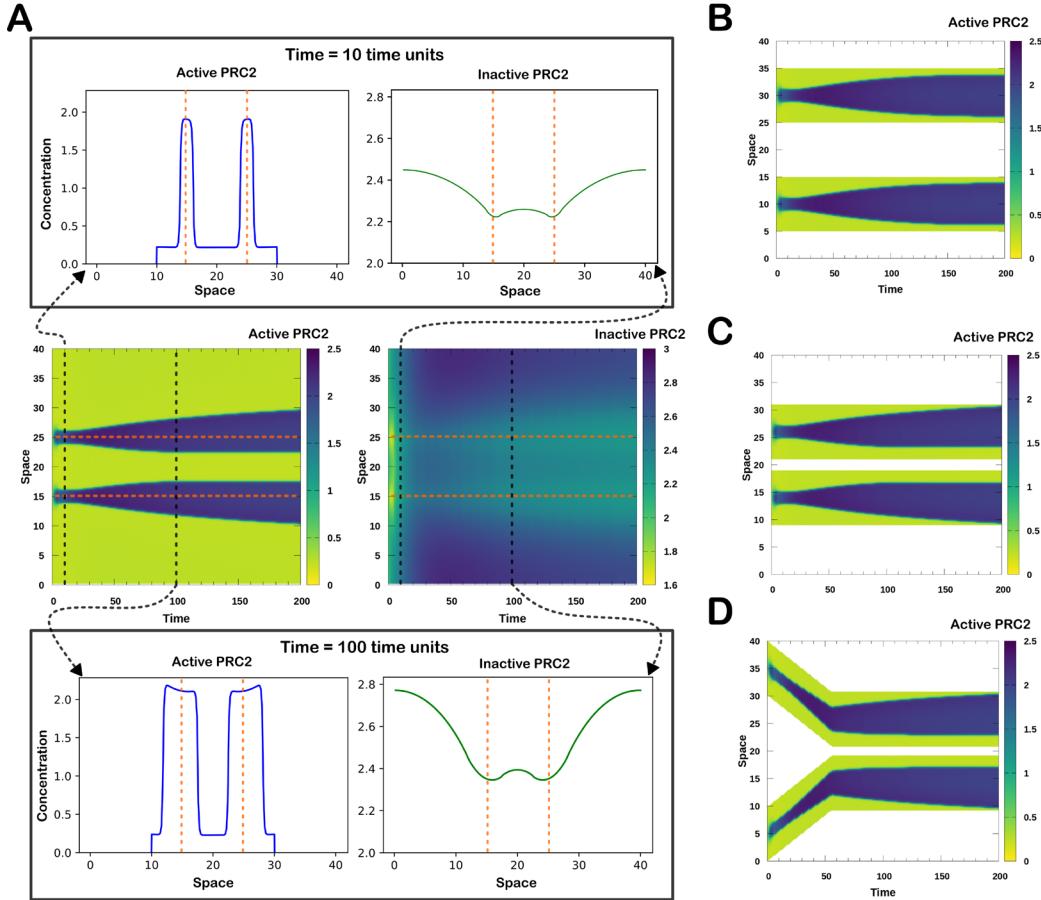
200 where the active PRC2 complex level was elevated while the rest of the region has low a concentration of
201 active PRC2 complex (i.e. a high level of gene expression). As we expected, the positive feedback loop
202 resulted in a trigger wave (also called a traveling wave) of gene silencing (**Fig 2B**), which stems from local
203 bistability of the chromatin states (**Figure S4**). This trigger wave gradually expands the active PRC2-high
204 state to the entire region. Trigger wave driven by positive feedback loop was widely studied in other
205 biological systems (Gelens et al., 2014), and our model suggests a new utility of this mechanism in
206 chromosomal inactivation. Nevertheless, we found that the gene silenced state covered the entire region
207 at steady state even when we simulated a larger domain (**Figure S5**). While this shows the robustness of
208 the trigger wave, the result is inconsistent with the observation that chromosomal inactivation only occurs
209 in one part of the X chromosome during development (Tukiainen et al., 2017), and in an even more
210 restricted region when transgene XIST is activated in an autosome (Minks et al., 2013; Naciri et al., 2021).
211 We therefore asked what can cause the restriction of the spread of gene silencing. Interestingly, reducing
212 the strength of the positive feedback loop in flanking regions of the chromosome resulted in an abrupt
213 termination of the trigger wave (**Fig 2C**). This is consistent with previous reports on sequence and
214 epigenetics determinants for XIST-mediated gene silencing (Loda et al., 2017; Tang et al., 2010).
215 Nonetheless, these local features cannot fully explain the variable extents of the gene silencing spread
216 during XCI and activation of transgene XIST (Tukiainen et al., 2017). We therefore tested a global
217 mechanism for wave pinning. By lowering the production and degradation rate constants of PRC2 (α and
218 β in Eq 2), we found that spread of gene silencing can be restricted by a conserved amount of PRC2, which
219 gave rise to a fraction of silenced chromosomal region at the steady state (**Fig 2D**). This global conservation
220 mechanism was similar to the one used for explaining robust polarized distribution of molecules at cell
221 membrane (Mori et al., 2008), and recent theoretical work showed that conservation of the total cellular
222 concentrations of histone modification enzymes is necessary for maintaining epigenetic memory during
223 cell divisions (Owen et al., 2023). We found that the strength of the positive feedback loop is important
224 for the wave pinning (**Fig 2E**), and for the position where the wave is pinned (**Figure S6**). Taken together,
225 our results provide a new mechanistic explanation, i.e. a wave-pinning model, for the spread of gene
226 silencing at the chromosomal level and we identified two distinct factors contributing to the spatial
227 restrictions of the wave. We expect that biological processes such as XCI use a combination of local and
228 global factors for the precise control of gene silencing waves.

229

230 **Spread of gene silencing from multiple initiation sites**

231 Previous studies have shown that XIST binds to multiple sites on the X chromosome, and these sites are
232 the places where gene silencing spreads start (Simon et al., 2013). To study potential interactions of
233 multiple silencing waves, we tested the wave dynamics initiated from two sites presumably bound by XIST
234 at the beginning of XCI. Note that here we are still not considering the dynamics of XIST explicitly in this
235 chromatin module. Instead, we introduced two transient stimulations in the simulated spatial domain
236 described in Eq 2. We focused on a wave pinning mechanism based on global control of polycomb
237 repressive complexes because the characteristics of wave pinning driven by local parametric changes are
238 predictable intuitively. Unsurprisingly, with two local stimulations for gene silencing, we observed two
239 silencing waves each starting to spread two both directions in the chromosome region at the beginning of
240 the simulation. However, the two waves became asymmetrical and repelled the propagation of each other
241 (**Fig 3A** left heatmap). To examine the mechanism underlying this phenomenon, we quantified the levels
242 of both active and inactive PRC2 at an early stage of the simulation (**Fig 3A** heatmaps and line plots). When
243 the active PRC2 levels were still symmetrical with respect to the location of the stimulations (**Fig 3A** orange

244 dashed vertical lines), the inactive PRC2 levels started to become asymmetrical: the region between the
245 two stimulations had a lower concentration of inactive PRC2 compared to the exterior regions due to the
246 depletion of inactive PRC2 induced by two waves, as opposed to one wave in each exterior region (**Fig 3A**,
247 green curves).



248
249 **Figure 3. Interaction among the silencing waves and its analysis:** A. The heatmaps represent
250 the spatiotemporal profile of both active and inactive PRC2 which clearly reveals the repulsive interaction
251 between silencing waves and resulting symmetry breaking. A more detailed analysis of the spatial
252 distribution of both the molecules at two different time points unravels the unequal distribution of inactive
253 PRC2 as the origin of symmetry breaking during the movement of silencing wave fronts. B-D. Illustration
254 of the effect of chromosomal movement on spatial profile of active PRC2 complex where the white, blue
255 and green regions represent the void space, silenced and active chromosomal regions respectively. The
256 results indicate that the steady state distribution of the molecules is highly dependent on the final spatial
257 orientation of the chromosome. Parameter values for the simulations in this figure are listed in **Table S2**.

258
259 We next asked whether the interactions of silencing waves can influence each other if they were on two
260 separate chromosome regions. Unlike the models presented earlier in this work, the two regions
261 correspond to two chromosomal segments that are not adjacent to each other in genomic coordinates
262 but are relatively close in 3D space. To simulate this in a 1D framework, we created non-chromosomal
263 space that allows diffusion of molecules but cannot be involved in any feedback-related reactions (Eq 2)

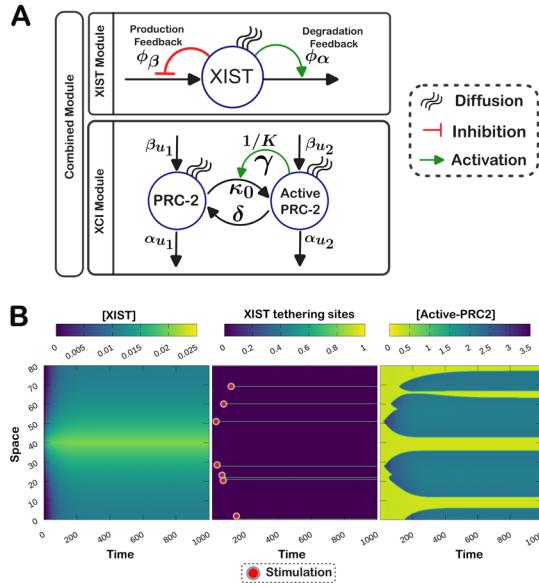
264 (white space in **Fig 3B**). When we placed the two segments, each perturbed by a wave-enabling stimulation,
265 with a long distance between them, we found that the two waves did not influence each other's dynamics
266 (**Fig 3B**). However, when the two segments were close, the two silencing waves repelled each other as we
267 saw with the one-segment model (**Fig 3C**). We next asked whether the movement of chromosomal regions
268 can influence the silencing waves. We simulated a scenario of two chromosomal regions approaching each
269 other, mimicking a chromosome undergoing compaction (**Figure S7**), and we found that the steady state
270 silencing distribution was determined by the steady state, but not transient, positions of chromosomal
271 regions (**Fig 3D**). Furthermore, the repulsion of the two silencing waves was observed for both the no-flux
272 boundary condition, which was used for all simulations presented so far, and the periodic boundary
273 condition (**Figures S8 and S9**). This shows the wide applicability of our conclusions. In summary, our results
274 show that the control of gene silencing via global levels of repressive complexes can give rise to
275 asymmetrical silencing wave propagation, and interactions of multiple waves under both static and
276 dynamic conditions of chromosomal regions' positions.

277

278 **Integration of the XIST and the chromatin modules**

279 We next combined the XIST module (Eq 1) and the chromatin module (Eq 2) to build an integrated system
280 for studying XCI-mediated gene silencing (**Fig 4A**). To bridge the gap between dynamics of XIST
281 concentration and the concentration of the repressor complex, we used an intermediate variable to
282 represent the status of the tethered XIST-protein complex (Pandya-Jones et al., 2020). Higher
283 concentration of XIST triggers the activation of the tethering sites in a probabilistic manner (see
284 Supplementary Information) (**Fig 4B** left and middle panels). The locations of the tethering sites are
285 predefined in the simulated space. The activation of the tethering sites first leads to the activation of the
286 PRC2 complex in the adjacency of the sites, which in turn initiates multiple trigger waves (**Fig 4B** right
287 panel). At the steady state, most of the simulated chromosomal region was silenced, but due to the pinning
288 of the waves, some regions remained to be active, reflecting the stable expression of the genes that have
289 escaped from the chromosomal inactivation. In this model, we assumed spatially homogeneous feedback
290 strengths to examine the effect of global regulations on the silencing wave. Interestingly, some
291 asymmetrical wave patterns were produced due to the random distribution of the tethering sites and the
292 resultant asymmetric interactions among waves that we explained in the previous section (**Fig 3** and **Fig**
293 **4**). Overall, the combined model suggests that the XIST-mediated gene silencing can be captured by a
294 dynamical system with multiple traveling waves, and such waves are initiated by both high local
295 concentration of XIST and the existence of local protein assemblies at selected sites of the chromosome.

296



297

298 **Figure 4. Integrated module and the simulation results:** A. Portrayal of the combined biochemical
299 network governing the underlying spatiotemporal dynamics of XIST, active PRC2 and inactive PRC2
300 molecules. The XIST module and X chromosome inactivation module is bridged by the tethering module
301 that incorporates the effect of XIST spatiotemporal profile to the dynamical landscape of X chromosome
302 inactivation. B. The simulation results presented in this sub panel depict the spatiotemporal profile of XIST
303 molecules (left), the tethering sites of the XIST molecules leading to stimulations (center) and the time
304 evolution of the chromosomal inactivation profile (right). A detailed description of the parameters used
305 for the simulation is provided in the Supplementary Information. Parameter values for the simulations in
306 this figure are listed in **Table S3**.

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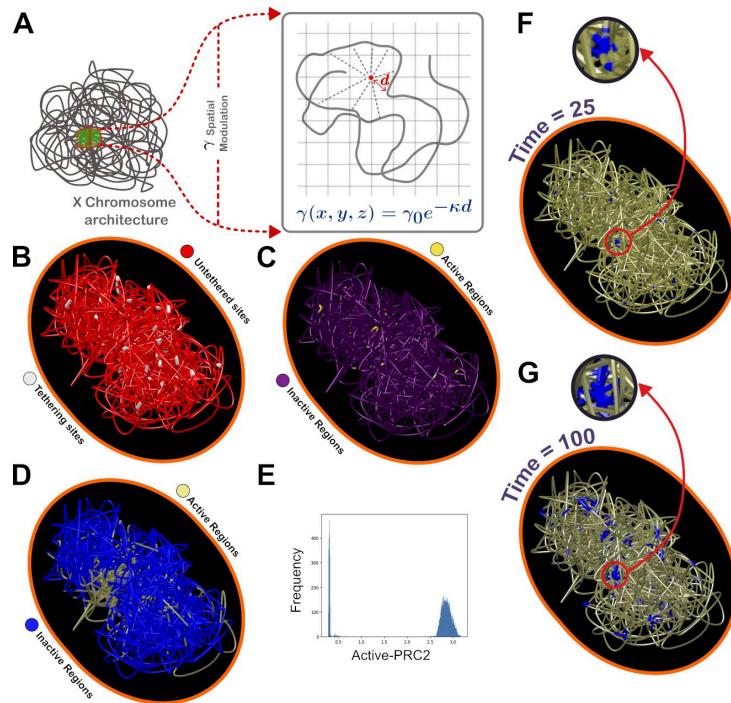
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309 **A 3D model for spatiotemporal dynamics of chromosome inactivation**

310 To illustrate the utility of our modeling framework with realistic chromosome structures and experimental
311 data, we built a proof-of-concept 3D model for X chromosome and its inactivation. We first used a
312 previously published 3D model X chromosome inferred from Hi-C sequencing data (Lappala *et al.*, 2021).
313 While the X chromosome structure is dynamic during XCI, we chose the inactive X chromosome structure
314 because we showed with the 1D model that the stationary phase of the chromosome structure is the
315 primary spatial factor determining the silencing waves (**Fig 3**). In addition, since the inferred structure
316 represents an averaged configuration across multiple cells and multiple time points at the stationary phase,
317 we introduced a continuous spatial parameter γ that determines strength of the feedback between
318 chromatin and the repressive complex, reflecting the duration or likelihood of a location in 3D space
319 occupied by chromatin (**Fig 5A**, **Figure S10**). We used the genomic location of XIST as the production site
320 of XIST RNA, and we used our model the two integrated modules simulate XCI with multiple tethering sites.
321 The location of the tethering sites were fitted to experimental data on silenced genes and escape genes
322 (Berletch *et al.*, 2015) (**Fig 5B**). We showed that the model gave rise to a reasonable agreement with
323 experimental data (**Fig 5C-E**, **Figure S11**) in that only a small fraction of X-linked genes remained active at
324 the steady state and those genes, including XIST itself, were primarily located at the exterior region of the

325 X chromosome in 3D. Among 38 experimentally identified escape genes, the model captured 12 of them
326 while maintaining a nearly 80% silencing percentage of the entire chromosome (**Fig 5E**). Furthermore, we
327 showed that, as a result of the continuous consideration of chromosome structure, silencing waves can
328 spread across multiple chromosomal regions due to their close proximity in 3D (see an example in **Fig 5F**
329 and **G**). This is consistent with a previous experiment using transgene XIST that showed the possibility of
330 gene silencing spread over 3D chromosome structure (Naciri *et al.*, 2021).

331



332

333 **Figure 5. The three-dimensional model and comparison of simulated results with experimental data: A.**
334 Schematic representation of implementing spatial dependence of the high sensitivity feedback coefficient
335 γ associated to the process of chromosomal inactivation. **B.** Spatial profile of the XIST tethering sites
336 obtained from 3D simulation marked by white beads. **C.** Spatial profile of the X chromosome inactivation
337 obtained from simulation at time unit 500 which corresponds to approximately 2 hours in physical scale
338 (details about time and length scale estimation are provided in the Supplementary Information file). **D.**
339 Visualization of the experimentally observed active and inactive genes in the X chromosome. **E.**
340 Distribution of chromosomal segments over the levels of active PRC2 in a mouse X chromosome. A total
341 of 20000 evenly spaced points on a simulated chromosome were used. **F-G.** Demonstration of the 3D
342 spatial spreading of the chromosomal inactivation through void space. (The detailed description of the
343 simulation method and utilized parameter values are provided inside the Supplementary Information file,
344 which includes **Table S4**)

345

346

347

348 **Discussion**

349 X chromosome inactivation (XCI) is a remarkable regulatory program for gene expression in which the
350 transcription of a lncRNA XIST triggers gene silencing for most of the genes in a chromosome. Substantial
351 work has been done to unravel the mechanisms of gene silencing during XCI, but many questions remain
352 open. In particular, it has been unclear how the spread of gene silencing is robustly triggered and regulated
353 in a context dependent manner. In this work, we used a mathematical model based on a continuum
354 reaction-diffusion system to show that feedback loops for XIST transcription and degradation have
355 profound an impact on both the steady state distribution of XIST and the speed to achieve the steady state.
356 We observed that the spatial range of the XIST distribution was increased by transcription-level negative
357 feedback (Jachowicz *et al.*, 2022) and decreased by degradation-level negative feedback (Rodermund *et*
358 *al.*, 2021). This suggests that these feedback loops may be used to define the boundary of XIST distribution
359 and its effect on gene silencing, and they may help to limit the XCI to a single chromosome. Previous
360 experimental and theoretical studies without considering spatial distributions showed that negative
361 feedback loop can accelerate the response of gene expression that steady state can be obtained rapidly
362 (Rosenfeld *et al.*, 2002). However, in this work, we showed that when spatial distribution of gene products
363 is considered, transcription-level negative feedback can either speed or slow down response, depending
364 on the choice of other parameters. Interestingly, these functions of negative feedback loops were obtained
365 without changing the diffusion coefficients. This may offer an alternative strategy for cells to tune
366 molecular distributions without significant changes of the physical properties of the molecules through
367 evolution.

368 The XIST-mediated gene silencing is an example of lncRNA's substoichiometric action in which a relatively
369 low expression level of a lncRNA can trigger a large effect in gene regulation. It was previously suggested
370 that mechanisms such as molecular condensate formation may explain this type of substoichiometric
371 function (Unfried and Ulitsky, 2022). Our model showed that systems-level positive feedback can amplify
372 the small initial effect induced by lncRNAs and trigger bistable switches (in the case of XCI, an on-to-off
373 switch) in a spatially sequential manner via traveling waves. This is similar to the well-known function of
374 positive feedback for long-range communication without the need of transporting molecules over a long
375 distance (Gelens *et al.*, 2014). The local functions of lncRNAs in forming condensates or scaffolding can be
376 combined with downstream positive feedback loops to give rise to large-scale effects, and this type of
377 combination may be common among lncRNA-mediated gene regulation.

378 Our traveling wave model is based on the assumption of chromatin-level feedback loop between histones
379 and enzymes that catalyze histone modification reactions. This type of feedback was widely used to explain
380 epigenetic memory (Dodd *et al.*, 2007). More recently, a polymer-based model showed the importance of
381 both local attraction of epigenetic marks and global regulation of enzymes in establishing epigenetic
382 memory with perturbed 3D chromatin structure (Owen *et al.*, 2023). Our models show that imbalanced
383 distributions of polycomb repressive complexes and asymmetrical silencing waves can be a mechanism for
384 local and medium-range attraction and repulsion. The globally driven wave-pinning mechanism is
385 consistent with the finding that global regulation of key enzymes can have significant effects on epigenetic
386 memory.

387 We used a continuum reaction-diffusion framework for this study, and we the advantages of this
388 framework in describing complex regulatory networks in a spatiotemporal model. This type of network is
389 often difficult to model with agent-based strategies especially when the number of molecules is large.
390 However, a key limitation of the continuum model is that the movement of small numbers of molecules
391 and their dynamic geometries are often not described accurately. Our hybrid approach for modeling 3D
392 chromosome structure in our model is useful for integrating continuous reaction-diffusion and discrete

393 geometry. Nonetheless, we expect that future development of modeling strategies will be necessary to
394 describe the complex system of chromosome-level gene regulation with advantages of both continuum
395 and agent-based methods. Overall, our continuum reaction-diffusion model has provided new insights into
396 mechanisms for gene silencing spread mediated by lncRNAs and it will help to deepen our understanding
397 on gene regulatory networks controlling XCI in both time and space.

398

399

400 **Methods**

401

402 *Model Development*

403 We have adopted a modular approach to develop our model by dividing the process of chromosomal
404 inactivation into three key steps. First, the spatiotemporal dynamics of XIST molecules, designated as the
405 XIST module. Second, the spatiotemporal dynamics of the inactive and active PRC2 complex, termed the
406 XCI module. Third, the integration of the XIST and XCI (i.e. chromatin) modules, referred to as the
407 Tethering module. Throughout the main and Supplementary Information, we have utilized the above-
408 mentioned terminologies to refer to different modules. The XIST and chromatin modules are represented
409 by three partial differential equations (PDEs) containing production, degradation, reaction, and diffusion
410 terms. Biologically plausible parameter values were chosen to give representative results (Supplementary
411 Information), and scanning of some parameters (feedback strengths) was used to analyze the model
412 behaviors systematically (Fig 1). All simulations were performed with the no-flux boundary condition
413 unless otherwise indicated. In a few cases the terms are subject to feedback which is backed by
414 experimental evidence. For integration of the two modules, a bridging Tethering module is represented
415 by a set of conditions depending on the spatiotemporal profile of XIST molecule that invokes the
416 stimulation in XCI module. Importantly, the Tethering module introduces the stochasticity into the model
417 arising out of the limited number of transcribed copies of XIST molecules. In addition to the brief
418 description of each module provided in previous sections, we have presented a detailed description of
419 the same in the Supplementary Information.

420

421 *1D Simulations*

422 After developing and fixing various aspects of the model we proceed further and carry out the
423 phenomenological study of the system by simulating the PDEs in one dimension. We have utilized the
424 explicit forward Euler algorithm implemented in FORTRAN programming language to simulate the PDEs.
425 First, we discretized the space and time and then defined the Laplacian terms in each time step using a
426 central difference formula which is provided in Supplementary Information. A thorough account of
427 numerical simulation and parameters used are provided in the Supplementary Information. We have also
428 implemented the void space, i.e., the region without the presence of chromosome by modulating the
429 feedback coefficient γ , as the parameter has higher sensitivity towards the bistability of the system.
430 Computational details regarding the implementation of void space and spatial movement of the
431 chromosomal region are provided in the Supplementary Information. Finally, we integrated all the
432 modules and carried out the simulation for a wider spatial region and got the steady heterogeneous
433 spatial distribution which is reminiscent to the chromosomal inactivation.

434

435 *3D Simulation*

436 The outcome of the phenomenological studies performed in one spatial dimension required further
437 exploration of the system in three spatial dimensions using a realistic chromosome structure supported
438 by experimental evidence. We have carried out the three-dimensional simulations using a discrete semi-

439 empirical X chromosome structure derived from HiC data. We have also carried out necessary
440 computational maneuvers to simulate our continuum mathematical model over the template of a discrete
441 three-dimensional chromosome structure. The details of the maneuvers are provided in the
442 Supplementary Information. Furthermore, the results from one dimensional study indicated that the
443 spatial profile of the final chromosomal inactivation is explicitly dependent on the final steady geometry
444 of the chromosome. Hence, we have used the stationary structure of inactive X chromosome and
445 performed the numerical simulation in 3D.

446

447 *Data Analysis*

448 Subsequently, as a logical extension of our investigation, we have curated the available experimental gene
449 expression data for X chromosome and studied the extent of agreement of the simulated result with the
450 experimental data. Detailed descriptions of conditions to filter the state of a particular gene as on or off
451 are provided in the Supplementary Information.

452

453 **Code Availability**

454 Computer code for reproducing simulation results is available at the GitHub repository for this manuscript
455 <https://github.com/shibashispaul32/XCI>.

456

457 **Competing Interests**

458 The authors declare no competing interests.

459

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